Stoke's Theorem

COn sider two adjacent in finterinal square loops, not necessarily in the same place

$$(\nabla \times \mathbb{R}) \cdot \hat{n} = \lim_{A \to 0} \frac{\int_{\mathbb{R}} \mathbb{R} \cdot d\mathbf{r}}{A}$$
  $(d\mathbf{s} = A\hat{n})$ 

... (V×B).du = & B.dr for an infinitesimal loop

$$(000) 1 \qquad (000) 2$$

$$4 \sqrt{2} \qquad (000) 2$$

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$$\sum_{i=1}^{4} \underline{B} \cdot \underline{d} \underline{r} = \sum_{i=1}^{4} \omega_{i}' = (\nabla_{x} \underline{B}_{i}) \cdot \underline{d} \underline{s}$$

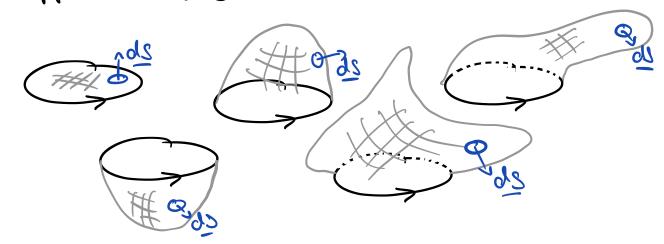
$$\sum_{i=1}^{4} \underline{B} \cdot \underline{d} \underline{r} = \sum_{i=1}^{4} \omega_{i}^{2} = (\nabla_{x} \underline{B}_{i}) \cdot \underline{d} \underline{s}$$

now less join the two loops together

$$\sum_{i+2} B \cdot dx = \sum_{i=1}^{4} \omega_{i}^{i} + \sum_{i=1}^{4} \omega_{i}^{2} - \omega_{i}^{1} - \omega_{i}^{2}$$

now, lets keep adding loops to form a macrosopric Surface attached to a large loop. closed of B. dr = SI(7x13)-ds

This applies to any surface attached to the loop.



To be clear on orientation of <u>ds</u>, "collopse" Jurfae on to the loop.

Why? Sometimes we wont a surface integral but the line integral is easier. Or, you chosse a different surface on the same loop.

Example B= zî-yî-xx, Shupe: hemisphere 2>0

 $dS = a^2 sin \theta d\theta d\phi \hat{\rho}$ 

$$\oint_{\mathbb{R}} \vec{G} \cdot \vec{q} \cdot \vec{q} = \iint_{\mathbb{R}} \Delta^* \vec{G} \cdot \vec{q} \cdot \vec{q}$$

& B. dr = & Bady + Bydy + Bzdzo = Zdzo-ydy

we now have to convert into cyclindrical so that our start e end limits are different.

$$x = a\cos \phi \qquad dx = -a\sin \phi d\phi$$

$$\partial \vec{S} \cdot d\vec{r} = \left[ \frac{2}{3} - \sigma_s 2i u k \cos \alpha q \alpha = -\sigma_s \left[ \frac{5}{7} 8i u_s \alpha \right]_{su}^0 = 0 \right]$$

Now let's calculate using curl. 
$$ds = \alpha^2 \sin \theta d\phi d\theta^2$$
  

$$\int \int \nabla x \, \mathbf{E} \cdot ds = \int \int 23 \cdot (\alpha^2 \sin \theta d\theta dx) \hat{r}$$
The probability of the probability

$$=\int_{0}^{\infty}\int_{0}^{\infty}2a^{2}\sin\theta\cos\phi d\theta d\phi = 0$$

even quicker: Choose a surface in the x-y place.
$$\frac{dS}{dS} = dS \hat{\kappa}$$

$$(77 \times B) \cdot dS = 23 \cdot dS \hat{\kappa} = 0$$

Vector Identities There are many vector identities. 3 are partially important!

I) 
$$\nabla x(\nabla \Omega) = 0$$
  
if  $B = \nabla \Omega$ , the  $B$  is conservable/irrotational

II) 
$$\nabla \cdot (\nabla \times \underline{V}) = 0$$

$$= \frac{\partial}{\partial x} (\frac{\partial v_z}{\partial y} - \frac{\partial v_z}{\partial z}) + \frac{\partial}{\partial y} (\frac{\partial v_z}{\partial z} - \frac{\partial v_z}{\partial y})$$

$$= 0$$

$$= if \quad \underline{G} = \nabla \times \underline{V} , \quad \underline{B} \quad \text{is Subsolidad'} \quad \underline{V} \quad \text{is collect} \quad \text{the Yeator potential} \quad \text{.}$$

III) 
$$\nabla \times (\nabla \times \mathcal{B}) = \nabla (\nabla \cdot \mathcal{B}) - \nabla^2 \mathcal{B}$$
  
USing in ESM to derive wave equation,  

$$\nabla^2 \mathcal{B} = \nabla^2 \mathcal{B}_x \hat{c} + \nabla^2 \mathcal{B}_y \hat{o} + \nabla^2 \mathcal{B}_{\bar{z}} \hat{\kappa}$$

 $\int \int (\nabla \pi \mathbf{B}) \cdot d\mathbf{b} = \mathbf{Sock}$ Sock
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