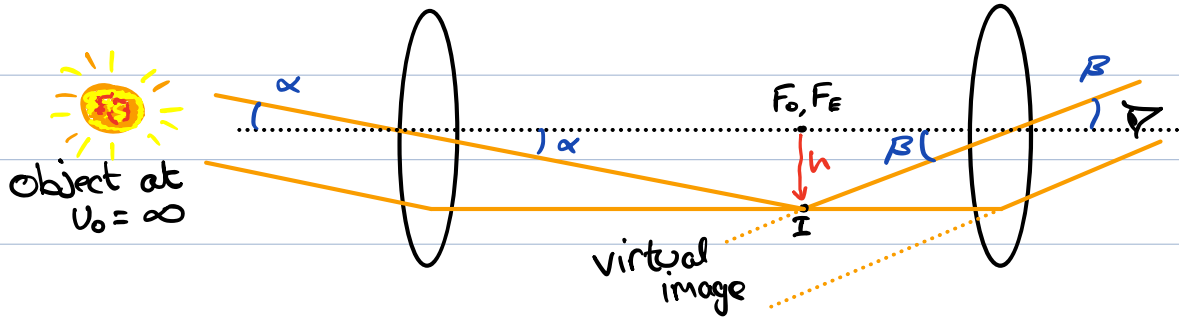


Telescopes

There are two types of telescopes: reflecting and refracting.

Keplerian Telescope (Refracting)

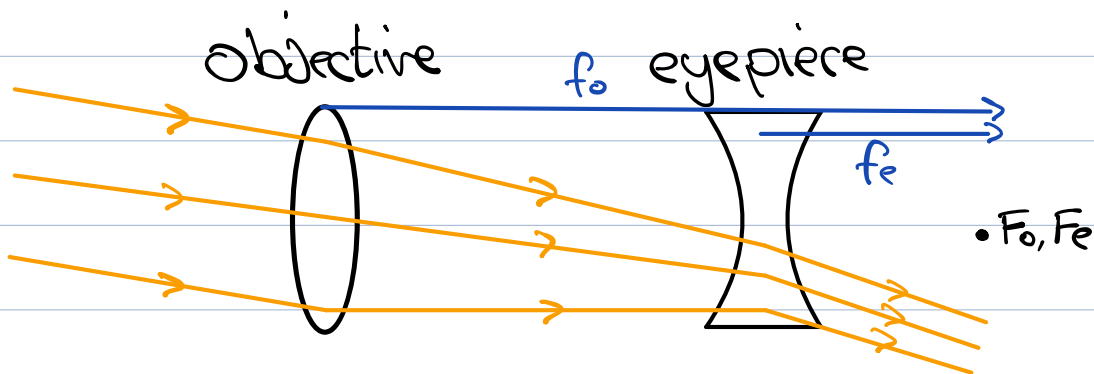
Made from two convex lenses.



The lenses must be $d = f_o + f_e$ apart. The angular magnification is $M = -\frac{\beta}{\alpha} = -\frac{f_o}{f_e}$. A high magnification telescope needs to be very long.

Galileon Telescope (Refracting)

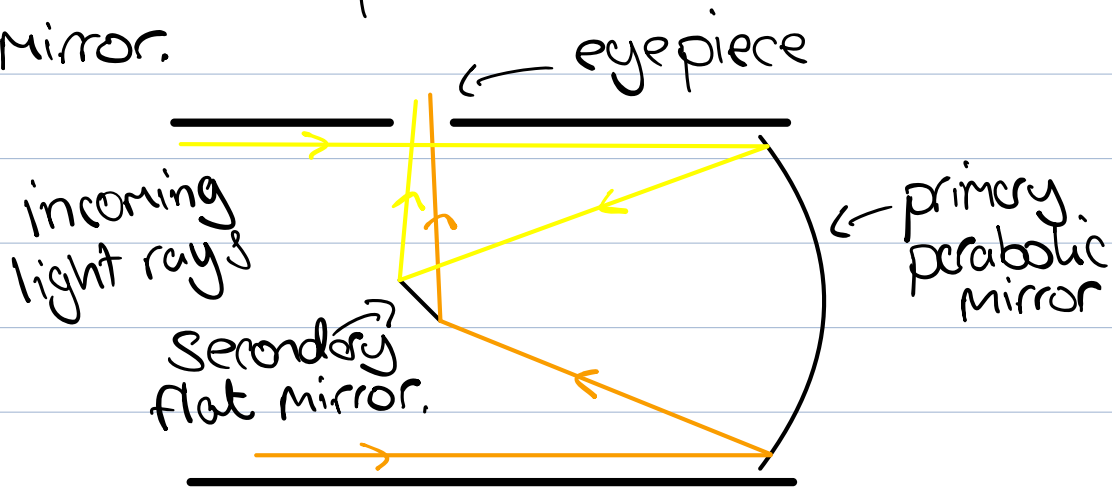
Made from one convex and one concave mirror.



The lens must be $d = f_o - |f_e|$ long - this is shorter than the keplerian telescope. Its angular magnification is $M = +\frac{\beta}{\alpha} = \frac{f_o}{|f_e|}$. It has the same mag. as keplerian but not inverted.

Newtonian Telescope (Reflecting)

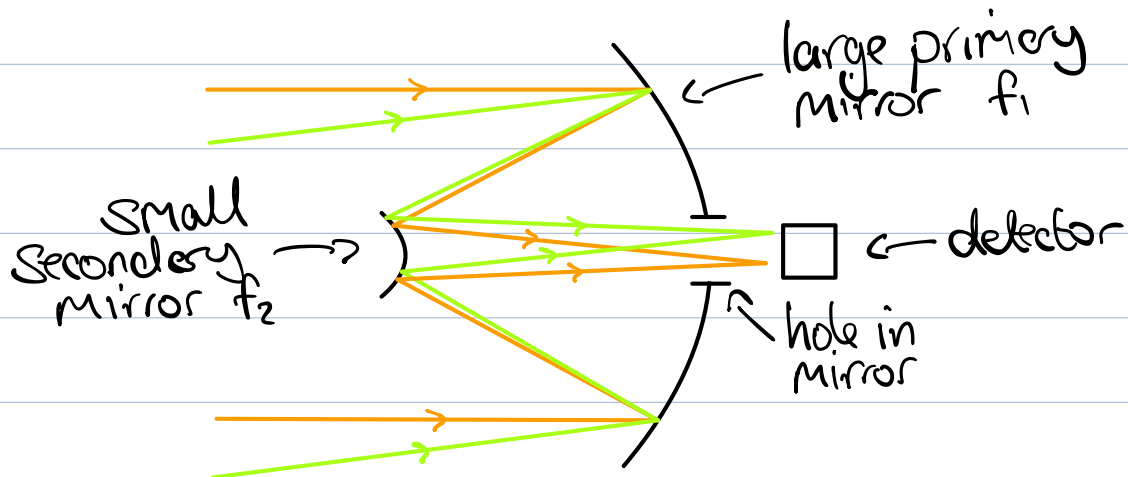
Made from one parabolic mirror and a smaller flat mirror.



Newton invented this type of telescope to solve chromatic aberration due to the lenses. Little use in present.

Cassegrain Telescope (Reflecting)

Made from one concave and one convex mirror.



The mirrors must be $d = f_1 - |f_2|$ apart. The angular magnification is $M = + \frac{f_1}{|f_2|}$. JWST

There are advantages to using reflecting telescopes.

- large diameter

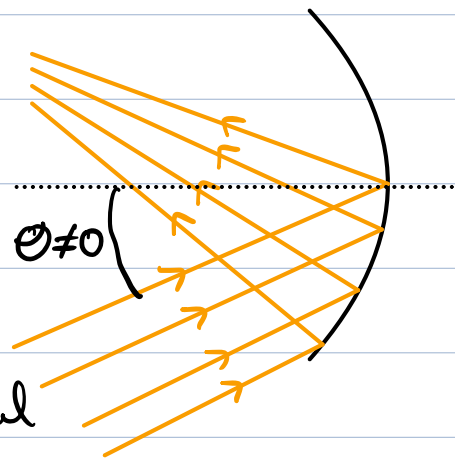
- No chromatic aberration
- can reflect visible, IR, micro + radio, lens can only do UV + visible.

Abberations

An aberration is an error in bringing together rays to an image point, the image appears blurred. It is because the OPL is not constant.

A parabolic mirror does not suffer from chromatic aberrations. However if the light entering is off-axis then we get COMA aberrations.

Central rays near the vertex ($r \approx 0$) give paraxial image point. $\theta \neq 0$



Marginal rays ($r \neq 0$) miss paraxial point with increasing error.

A Cassegrain telescope is able to reduce the COMA error.

Lens Abberations

A thin lens cannot satisfy Fermat's Principle for non-paraxial angles.

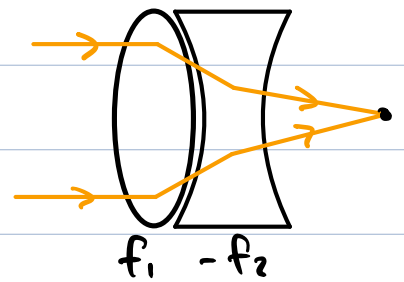
Chromatic aberrations occur since dispersion depends on refractive index $n(\lambda)$. \therefore focal length changing with wavelength.

For a thin lens, the lens' maker equation becomes

$$\frac{1}{f(\lambda)} = [n(\lambda) - 1] \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

To correct for chromatic aberration, we can use two lenses with different refractive indexes.

The two lenses have equal and opposite errors so OPL is roughly the same for all λ . Called an achromatic doublet lens.



Monochromatic Abberations

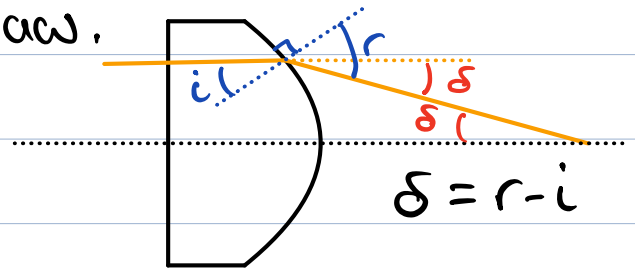
- ★ Spherical Abberation: error in imaging on-axis object for marginal rays (large r).
- ★ Coma Abberation: error in imaging off-axis object (large angle θ)
- ★ Field Curvature: relative error mapping flat object plane of points to curve of points in image plane.

Spherical Abberations

There is a significant aberration in lenses due to the nonlinearity of Snell's law.

$$n_1 \sin i = n_2 \sin r$$

$$n_1 [i - i^3/3] = n_2 [r - r^3/3]$$



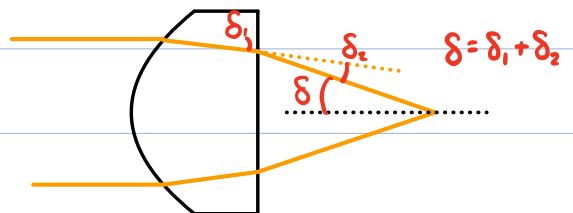
let $n_2 = 1$ & $n_1 = n$, rearranging gives

$$\delta = r - i = (n-1)i + \frac{1}{3}(r^3 - ni^3) = \delta + \Delta\delta$$

This gives the paraxial deviation $\delta = (n-1)i$ and the aberration $\Delta\delta \approx \frac{1}{3}(r^3 - ni^3) \propto \delta^3$.

This leads to spherical aberrations, even if the lens is parabolic.

To reduce spherical aberrations we should reduce the angle that light enters the lens. For a plano-convex lens, placing the convex part first reduces the error $\approx 4\times$.



$$\delta_1 \approx \delta_2 \approx \frac{\delta}{2} \quad \Delta\delta \propto 2 \cdot \left(\frac{\delta}{2}\right)^3 = \frac{\delta^3}{4}$$

More surfaces (n) is better. $\Delta\delta \propto n \cdot \left(\frac{\delta}{n}\right)^3 = \frac{\delta^3}{n^2}$
which can make the error negligible.

When the object is distant but the image is near use a planoconvex lens.

When object and image distances are similar use a biconvex lens.

A spherical ball lens is able to limit coma and spherical aberrations. It has a principle axis for all ray angles and an aperture in centre can stop marginal rays. However it doesn't stop field curvature.

A camera lens set-up also stops chromatic aberrations (different material) and a final lens solves field curvature.

In the human eye, the curvature of the retina solves field curvature

When a lens has been designed so well to eliminate all aberrations it becomes diffraction limited.