

## Potential Energy

consider a charge  $Q$ . A test charge  $q$  experiences force  $\underline{F}_Q$ . Now let's apply a force  $\underline{F}_{ext}$  to move  $q$  from A to B at a constant speed.

$$\underline{F}_{ext} + \underline{F}_Q = 0$$

Elemental work done by  $\underline{F}_{ext}$  on  $q$  is

$$dW = \underline{F}_{ext} \cdot d\underline{l} = -\underline{F}_Q \cdot d\underline{l}$$

The total work done by  $\underline{F}_{ext}$  is found by integrating:

$$W = \int_A^B \underline{F}_{ext} \cdot d\underline{l} = - \int_A^B \underline{F}_Q \cdot d\underline{l}$$

Work done is equal to change in potential energy.

$$\Delta U = - \int_A^B \underline{F}_Q \cdot d\underline{l}$$

We know that  $\underline{F}_Q = \frac{1}{4\pi\epsilon_0} \frac{qQ}{r^2} \hat{r}$ ,  $\dots$

$$dW = - \frac{1}{4\pi\epsilon_0} \frac{qQ}{r^2} \hat{r} \cdot d\underline{l}$$

due to it being a conservative field, the path doesn't matter.

$$W = \Delta U_{AB} = - \int_A^B \frac{1}{4\pi\epsilon_0} \frac{qQ}{r^2} dr$$

$$= -\frac{qQ}{4\pi\epsilon_0} \left[ \frac{1}{r_A} - \frac{1}{r_B} \right]$$

For multiple charges, we can use the superposition principle.

$$\Delta U_{AB} = - \int_A^B \underline{F}_{Q1} \cdot d\underline{l} - \int_A^B \underline{F}_{Q2} \cdot d\underline{l} - \dots - \int_A^B \underline{F}_{QN} \cdot d\underline{l}$$

## Potential Difference

$$\Delta V_{AB} = \frac{\Delta U_{AB}}{q} = - \int_A^B \underline{E} \cdot d\underline{l}$$

path integral

$\Delta U_{AB}$  is path independent, and is independent of  $q$ . SI units of  $JC^{-1}$ . We can define where  $V=0$ , often chosen to be when  $r \rightarrow \infty$ .

The potential ( $V$ ) at point  $P$  is defined as the work needed to bring a  $+1C$  charge from  $r=\infty$  to  $P$ , at a constant speed.

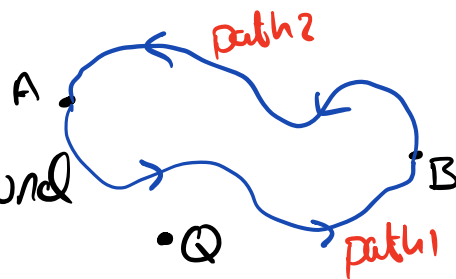
Scalar field  $\rightarrow$

$$V = - \int_{\infty}^P \underline{E} \cdot d\underline{l}$$

For a point charge  $Q$ ,  $V = \frac{Q}{4\pi\epsilon_0 r}$ .

## Circulation law

Consider the potential difference around a closed loop.



$$\Delta U_{AB} = - \int_A^B \underline{E} \cdot d\underline{l} - \int_B^A \underline{E} \cdot d\underline{l} = \int_A^B \underline{E} \cdot d\underline{l} - \int_A^B \underline{E} \cdot d\underline{l} = 0$$

$$\Rightarrow \oint \underline{E} \cdot d\underline{l} = 0$$

This is only true for electrostatics! We can use Stokes' theorem to say that

$$\oint_{\partial R} \underline{E} \cdot d\underline{l} = \iint_R (\nabla \times \underline{E}) \cdot d\underline{\Omega} = 0$$

For an infinitesimal surface element we get

$$\nabla \times \underline{E} \cdot d\underline{\Omega} = 0 \quad \Rightarrow \quad \nabla \times \underline{E} = 0$$

This equation must be satisfied at all points in the electrostatic field.