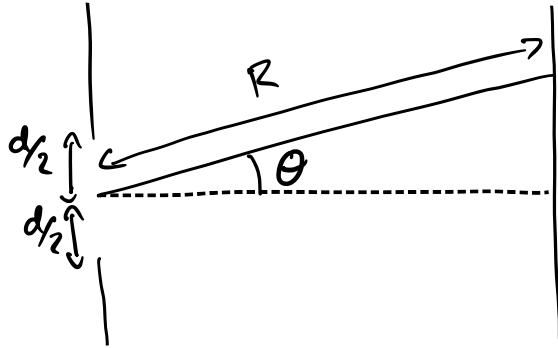


Fraunhofer Diffraction

light passing through a small aperture forms a diffraction pattern.



The plane wave arriving at x over the infinitesimal area dx will contribute $d\psi(\theta)$ to the amplitude radiation field.

$$d\psi(\theta) = a(x) e^{i \left[\underbrace{\frac{2\pi}{\lambda} R}_{\text{path length}} + \underbrace{\frac{2\pi}{\lambda} x \sin \theta}_{\text{path difference}} \right]} dx$$

relative amplitude

This works in the approximation $\lambda \ll d \ll R \gg d$. We can now drop the common phase, as we only need the pattern on the screen.

$$d\psi(\theta) \sim a(x) e^{i \left[\frac{2\pi}{\lambda} x \sin \theta \right]} dx$$

Assuming small angles ($l \sim \sin \theta$)

$$d\psi(l) \sim a(x) e^{i \left[\frac{2\pi}{\lambda} lx \right]} dx$$

We can now integrate to get the total contribution

$$\Psi(l) = A \int_{-d/2}^{d/2} a(x) e^{i \frac{2\pi}{\lambda} l x} dx$$

↑ aperture function

where A is the overall normalisation.

Let $\underline{k} = \frac{2\pi l}{\lambda}$ we can see that

$$\Psi(l) = A \int_{-d/2}^{d/2} a(x) e^{i k x} dx$$

Which is just a fourier transform of our aperture function $a(x)$.

- ★ A very narrow slit can be represented by the delta function so the diffraction pattern is smoothly spread in all directions.
- ★ A very wide slit will produce a delta function as the diffraction pattern.
- ★ To get the diffraction pattern for two slits we just fourier transform a single slit and add it to the same transform shifted by an amount.

This is how X-ray crystallography works. (Bragg).