

# 1 Gaussian model for the $(2, 0)$ Dirac operator

## 1.1 Notation

$$D = \gamma_1 \otimes \{H_1, \cdot\} + \gamma_2 \otimes \{H_2, \cdot\}. \quad (1)$$

Define  $\gamma_{\pm} := \frac{1}{2}(\gamma_1 \mp i\gamma_2)$  and  $W := H_1 + iH_2$ .  $D$  becomes:

$$D = \gamma_+ \otimes \{W, \cdot\} + \gamma_- \otimes \{W^\dagger, \cdot\} \quad (2)$$

And the following relations hold:

$$\gamma_{\pm}^2 = 0 \quad (3)$$

$$\gamma_{\pm}\gamma_{\mp} = \frac{1}{2}(I_C \pm i\gamma_1\gamma_2) \quad (4)$$

$$(\gamma_{\pm}\gamma_{\mp})^2 = \gamma_{\pm}\gamma_{\mp} \quad (5)$$

where  $I_C$  is the identity in the space of  $\gamma$  matrices.

## 1.2 Action, $\text{Tr } D^2$ part

Because of Eq.(3), only the cross terms survive in  $D^2$ :

$$D^2 = \gamma_+\gamma_- \otimes \{W, \cdot\}\{W^\dagger, \cdot\} + \gamma_-\gamma_+ \otimes \{W^\dagger, \cdot\}\{W, \cdot\} \quad (6)$$

Taking the trace yields:

$$\text{Tr } D^2 = 2C(n \text{Tr } WW^\dagger + \text{Tr } W \text{Tr } W^\dagger) \quad (7)$$

where  $C$  is the dimension of the  $\gamma$  matrices and  $n$  the dimension of the  $W$  matrix.

## 1.3 Action, $\text{Tr } D^4$ part

Because of Eq.(3), the cross terms vanish when squaring  $D^2$ :

$$D^4 = \gamma_+\gamma_- \otimes (\{W, \cdot\}\{W^\dagger, \cdot\})^2 + \gamma_-\gamma_+ \otimes (\{W^\dagger, \cdot\}\{W, \cdot\})^2 \quad (8)$$

Taking the trace yields:

$$\begin{aligned} \text{Tr } D^4 &= 2C[n \text{Tr}(WW^\dagger)^2 + \text{Tr } W^2 \text{Tr } W^{\dagger 2} + 2(\text{Tr } WW^\dagger)^2 \\ &\quad + 2 \text{Tr } W^2 W^\dagger \text{Tr } W^\dagger + 2 \text{Tr } W^{\dagger 2} W \text{Tr } W]. \end{aligned} \quad (9)$$

## 1.4 Stationary point

Stationary points are solutions to the equation

$$\delta S = S[W + \delta W] - S[W] = 0 \quad \forall \delta W \ll 1. \quad (10)$$

To first order in  $\delta W$  the terms  $\text{Tr } D^2$  and  $\text{Tr } D^4$  give:

$$\delta \text{Tr } D^2 = 2C [n \text{Tr } W \delta W^\dagger + n \text{Tr } W^\dagger \delta W + \text{Tr } W \text{Tr } \delta W^\dagger + \text{Tr } W^\dagger \text{Tr } \delta W] \quad (11)$$

$$\begin{aligned} \delta \text{Tr } D^4 = 2C [ & 2n \text{Tr}(WW^\dagger W \delta W^\dagger) + 2n \text{Tr}(W^\dagger W W^\dagger \delta W) \\ & + 2 \text{Tr } W^2 \text{Tr } W^\dagger \delta W^\dagger + 2 \text{Tr } W^{\dagger 2} \text{Tr } W \delta W \\ & + 4 \text{Tr } WW^\dagger (\text{Tr } W \delta W^\dagger + \text{Tr } W^\dagger \delta W) \\ & + 2 \text{Tr } W^2 W^\dagger \text{Tr } \delta W^\dagger + 2 \text{Tr } W^{\dagger 2} W \text{Tr } \delta W \\ & + 2 \text{Tr } W^\dagger (\text{Tr } W^2 \delta W^\dagger + \text{Tr } WW^\dagger \delta W + \text{Tr } W^\dagger W \delta W) \\ & + 2 \text{Tr } W (\text{Tr } W^{\dagger 2} \delta W + \text{Tr } W^\dagger W \delta W^\dagger + \text{Tr } WW^\dagger \delta W^\dagger) ]. \end{aligned} \quad (12)$$

We check for solutions of the form:

$$W = I \rho e^{i\theta} \quad (13)$$

with  $\rho > 0$  and  $\theta \in [0, 2\pi)$ .

$$\text{Tr } D^2 : \quad 4Cn\rho (e^{i\theta} \text{Tr } \delta W^\dagger + e^{-i\theta} \text{Tr } \delta W) \quad (14)$$

$$\text{Tr } D^4 : \quad 4Cn\rho (8\rho^2 e^{i\theta} \text{Tr } \delta W^\dagger + 8\rho^2 e^{-i\theta} \text{Tr } \delta W) \quad (15)$$

putting the coefficient of  $\delta W$  or  $\delta W^\dagger$  to zero gives:

$$\rho^2 = 0 \quad \text{or} \quad \rho^2 = -\frac{g}{8}. \quad (16)$$

## 1.5 Quadratic action in $\delta W$

Write  $W = \rho e^{i\theta} + \epsilon V$  for small  $\epsilon$ . We will separate the contributions to the action in orders of  $\epsilon$ .

$$\begin{aligned} \text{Tr } D^2 \\ O(1) : \quad 4Cn^2 \rho^2 \end{aligned} \quad (17)$$

$$O(\epsilon) : \quad 4Cn\rho (e^{i\theta} \text{Tr } V^\dagger + e^{-i\theta} \text{Tr } V) \quad (18)$$

$$O(\epsilon^2) : \quad 2C(n \text{Tr } V V^\dagger + \text{Tr } V \text{Tr } V^\dagger) \quad (19)$$

$$\text{Tr } D^4$$

$$O(1) : \quad 16Cn^2\rho^4 \quad (20)$$

$$O(\epsilon) : \quad 24Cn\rho^3 e^{-i\theta} \text{Tr } V + \text{c.c.} \quad (21)$$

$$O(\epsilon^2) : \quad 8C\rho^2(2n \text{Tr } VV^\dagger + ne^{-2i\theta} \text{Tr } VV \\ + 2 \text{Tr } V \text{Tr } V^\dagger + e^{-2i\theta} \text{Tr } V \text{Tr } V + \text{c.c.}) \quad (22)$$

Replacing  $\text{Tr } V$  with  $\langle \text{Tr } V \rangle$  (?does this makes sense?), the action to second order in  $\epsilon$  reads:

$$S = 2Cn^2(g\rho^2 + 4\rho^4) + Cn\epsilon^2 \text{Tr } VV^\dagger(g + 16\rho^2) \\ + 8Cn\epsilon^2\rho^2 e^{-2i\theta} \text{Tr } VV + \text{c.c.} \quad (23)$$

Or:

$$S = 4Cn^2(g\rho^2 + 4\rho^4) + Cn\epsilon^2 \text{Tr}(\vec{V}^T M \vec{V}) \quad (24)$$

with  $\vec{V}^T = (V, V^\dagger)$  and:

$$M = \begin{bmatrix} 8\rho^2 e^{-2i\theta} & g + 16\rho^2 \\ g + 16\rho^2 & 8\rho^2 e^{2i\theta} \end{bmatrix} \quad (25)$$