

Cable Project

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1 Perturbation

1.1 PDE for perturbation function

$$\frac{\partial \psi}{\partial t} = c \frac{\partial^2 \psi}{\partial \xi^2} + \frac{\partial \psi}{\partial \xi} - (1 - \delta(V(\xi) - \theta))\psi \quad (1)$$

We split the domain of the equation into three regions, depending on the value of $\delta(V(\xi) - \theta)$, with the boundary condition that $\lim_{\xi \rightarrow \pm\infty} \psi = 0$:

- 1) $\xi < 0$ and $\delta(V(\xi) - \theta) = 0$
- 2) $\xi > 0$ and $\delta(V(\xi) - \theta) = 0$
- 3) $\xi = 0$ and $\delta(V(\xi) - \theta) = \infty$

1.2 Separation of Variables

Looking first at region 1 and 2, we have the homogeneous PDE.

$$\frac{\partial \psi}{\partial t} = c \frac{\partial^2 \psi}{\partial \xi^2} + \frac{\partial \psi}{\partial \xi} - \psi \quad (2)$$

We assume a separation of variables, $\psi(\xi, t) = S(\xi)T(t)$, which leads to

$$\frac{1}{T} \frac{\partial T}{\partial t} = \frac{1}{S} \frac{\partial^2 S}{\partial \xi^2} + \frac{c}{S} \frac{\partial S}{\partial \xi} - 1 = \lambda$$

The general solution is

$$S(\xi) = c_1 e^{\frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} \xi} + c_2 e^{\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} \xi}$$

1.3 Region 1 and 2

In Region 1, where $\lim_{\xi \rightarrow -\infty} S_1 = 0$, and in Region 2, where $\lim_{\xi \rightarrow \infty} S_2 = 0$, we have the solutions:

$$S_1(\xi) = c_1 e^{\frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} \xi} S_2(\xi) = c_2 e^{\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} \xi} \quad (3)$$

1.4 Region 3

At Region 3, we integrate the equation at a small interval around the point of discontinuity to examine the effect $\delta(V(\xi) - \theta)$ has on the solution. After separation of variables, the inhomogeneous PDE 1 becomes

$$\frac{1}{T} \frac{\partial T}{\partial t} = \frac{1}{S} \frac{\partial^2 S}{\partial \xi^2} + \frac{c}{S} \frac{\partial S}{\partial \xi} - 1 + \delta(V(\xi) - \theta) = \lambda$$

The spatial equation is

$$\frac{d^2 S}{d\xi^2} + c \frac{dS}{d\xi} - [\lambda + 1 + \delta(V(\xi) - \theta)]S = 0 \quad (4)$$

Since there is a discontinuity here at $V(\xi) = \theta$, $\xi = 0$, we integrate a small interval $(+\epsilon, -\epsilon)$ around the discontinuity. Then, we taking the limit $\epsilon \rightarrow 0$ gives us information as to how the solutions in Region 1 and 2 (3) connect in Region 3.

$$\int_{-\epsilon}^{+\epsilon} \frac{d^2 S}{d\xi^2} d\xi + c \int_{-\epsilon}^{+\epsilon} \frac{dS}{d\xi} d\xi - \int_{-\epsilon}^{+\epsilon} (\lambda + 1)S d\xi - \int_{-\epsilon}^{+\epsilon} \delta(V(\xi) - \theta)S d\xi = 0$$

For the first two terms, the fundamental theorem of calculus gives us the first derivative evaluated at the endpoints and the function evaluated at the endpoints.

$$\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon) - S(-\epsilon)] - \int_{-\epsilon}^{+\epsilon} (\lambda + 1)S d\xi - \int_{-\epsilon}^{+\epsilon} \delta(V(\xi) - \theta)S d\xi = 0$$

Now we take the limit:

$$\lim_{\xi \rightarrow 0} \left[\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon) - S(-\epsilon)] - \int_{-\epsilon}^{+\epsilon} (\lambda + 1)S d\xi - \int_{-\epsilon}^{+\epsilon} \delta(V(\xi) - \theta)S d\xi \right] = 0$$

Of the remaining two integrals, the first one goes to zero as the bounds shrink the zero since the integral is continuous, even if S is not. The second integral involving the dirac delta function evaluates to $S(0)$ by the *sampling property*, since we are specifically looking at the small region centered around where $V(\xi) = \theta$. Therefore we get:

$$\lim_{\xi \rightarrow 0} \left[\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon) - S(-\epsilon)] \right] = S(0)$$

S must be continuous if we are assuming that the derivative $\frac{dS}{d\xi}$ exists so $S(0^+) = S(0^-)$ and the constants in front of the homogeneous solutions (3) are equal, $c_1 = c_2 = K$, revealing a relationship between the value of the spatial term and the discontinuity between its derivative at $\xi = 0$.

$$S(0) = \lim_{\xi \rightarrow 0} \left[K \left(\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} e^{\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2}} - \frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} e^{\frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2}} \right) \right]$$

Taking the limit, we find the following, where $S(0) \in \mathbb{R}$ and $K \in \mathbb{R}$ since this is a physical system and we assume the solutions are real-valued.

$$S(0) = K \left(\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} - \frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} \right) \quad (5)$$

S is the spatial part of some arbitrary perturbation function that satisfies $\lim_{\xi \rightarrow \pm\infty} S(\xi) = 0$, and its value at zero could be either non-zero, or it could be 0. Since the boundary conditions provided do not fix $S(0)$ or K , we will examine both cases. In the latter case, if $S(0) = 0$ then $K = c_1 = 0$ and we find the trivial solution of $S = 0$ which gives us $\psi(\xi, t) = 0$. In the following section, we explore the former case.

1.5 Eigenvalues

Now to examine the family of solutions that depend on λ , we start with the $\lambda = 0$ case. Assuming we can rescale $S(0)$ to unity, eq (5) simplifies to:

$$1 = K \left(\frac{-c - \sqrt{c^2 + 4}}{2} - \frac{-c + \sqrt{c^2 + 4}}{2} \right) = -K\sqrt{c^2 + 4}$$

or

$$K = -\frac{1}{\sqrt{c^2 + 4}}$$

Plugging this constant back into our