Cable Project

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December 7, 2017

The temporal portion of $\psi(\xi,t) = F(\xi)G(t)$ is governed by $\frac{1}{G(t)}\frac{dG}{dt} = \lambda$. Thus, the time dependence of $\psi(\xi,t)$ is described by $G(t) = C_1 e^{\lambda t}$, where $C_1 \in \mathbb{R}$. Now, consider $\frac{d^2F(\xi)}{d\xi^2} + c\frac{dF(\xi)}{d\xi} - F(\xi)(1+\lambda) = 0$. There are three cases to consider: $\lambda = 0$, $\lambda < 0$ and $\lambda > 0$.

When $\lambda = 0$, $\frac{d^2F(\xi)}{d\xi^2} + c\frac{dF(\xi)}{d\xi} - F(\xi) = 0$ has the general solution $F(\xi) = Ae^{\frac{1}{2}\left(-c+\sqrt{c^2+4}\right)\xi} + Be^{-\frac{1}{2}\left(c+\sqrt{c^2+4}\right)\xi}$. The solution corresponding to the zero eigenvalue is $\psi(\xi,t) = A_0 e^{\frac{1}{2}\left(-c+\sqrt{c^2+4}\right)\xi} + B_0 e^{-\frac{1}{2}\left(c+\sqrt{c^2+4}\right)\xi}$.
???? Wait this doesn't depend on time so it can't decay in time????

When $\lambda < 0$, if $c^2 \geq 4(1+\lambda)$, then $F(\xi) = Ae^{\frac{1}{2}\left(-c+\sqrt{c^2+4(1+\lambda)}\right)\xi} + Be^{-\frac{1}{2}\left(c+\sqrt{c^2+4(1+\lambda)}\right)\xi}$. The boundary condition in equation **??** requires that $B = 0$ when $V > \theta$ and that $A = 0$ when $V < \theta$.

???–¿ When $\lambda < 0$, if $c^2 < |4(1+\lambda)|$, then $F(\xi) = e^{\frac{-c}{2}}\left(Ccos\left(\frac{\sqrt{|c^2+4(1+\lambda)|}}{2}\xi\right) + Dsin\left(\frac{\sqrt{|c^2+4(1+\lambda)|}}{2}\xi\right)\right)$.
$v(x,t) = V(\xi) + \epsilon\psi(\xi,t)$

0.1 Stability of the Traveling Front Solutions

Now that we have found the traveling front solutions, it is important to understand their stability using linear stability analysis. Consider the solution $v(x,t) = V(\xi) + \epsilon\psi(\xi,t)$, where $0 < \epsilon << 1$ and $\psi(\xi,t)$ represents a small perturbation to the traveling wave solution, $V(\xi)$. The goal of the linear stability analysis is to understand how $\psi(\xi,t)$ behaves over time. Recall the governing partial differential equation

$$\frac{\partial v}{\partial t} = \frac{\partial^2 v}{\partial x^2} - v + H(v-\theta) \tag{1}$$

The chain rule allows us to obtain the following partial derivatives of $v(x,t) = V(\xi) + \epsilon\psi(\xi,t)$.

$$\frac{\partial v}{\partial t} = \frac{\partial \xi}{\partial t}\frac{dV}{d\xi} + \epsilon\left(\frac{\partial \psi}{\partial \xi}\frac{\partial \xi}{\partial t} + \frac{\partial \psi}{\partial t}\right) = -c\frac{dV}{d\xi} + \epsilon\left(-c\frac{\partial \psi}{\partial \xi} + \frac{\partial \psi}{\partial t}\right)$$

$$\frac{\partial v}{\partial x} = \frac{\partial \xi}{\partial x}\frac{dV}{d\xi} + \epsilon\frac{\partial \xi}{\partial x}\frac{\partial \psi}{\partial \xi} = \frac{dV}{d\xi} + \epsilon\frac{\partial \psi}{\partial \xi}$$

$$\frac{\partial^2 v}{\partial x^2} = \frac{\partial \xi}{\partial x}\frac{d^2 V}{d\xi^2} + \epsilon \frac{\partial^2 \psi}{\partial \xi^2}\frac{\partial \xi}{\partial x} = \frac{d^2 V}{d\xi^2} + \epsilon \frac{\partial^2 \psi}{\partial \xi^2}$$

Equation (1) is rewritten as

$$-c\frac{dV(\xi)}{d\xi} + \epsilon(-c\frac{\partial \psi(\xi,t)}{\partial \xi} + \frac{\partial \psi(\xi,t)}{\partial t}) = \frac{d^2 V(\xi)}{d\xi^2} + \epsilon \frac{\partial^2 \psi(\xi,t)}{\partial \xi^2} + f(V(\xi) + \epsilon\psi(\xi,t)) \tag{2}$$

Analysis of equation (2) can be simplified by using the Taylor expansion of $f\big(V(\xi)+\epsilon\psi(\xi,t)\big)$ with respect to $V(\xi)$, and discounting $O(\epsilon^2)$ terms.

$$f\big(V+\epsilon\psi\big) \approx f(V) + \frac{\partial f(V)}{\partial V}\epsilon\psi = -V + \mathcal{H}(V-\theta) - \epsilon\psi + \delta(V-\theta)\epsilon\psi$$

Equation 2 becomes

$$-c\frac{dV}{d\xi} + \epsilon(-c\frac{\partial \psi}{\partial \xi} + \frac{\partial \psi}{\partial t}) = \frac{d^2 V}{d\xi^2} + \epsilon\frac{\partial^2 \psi}{\partial \xi^2} - V + H(V-\theta) + \epsilon\psi\big(-1+\delta(V-\theta)\big)$$

Look at the terms involving ψ and ϵ to analyze the long term behavior of the disturbance. The partial differential equation that models the behavior of the perturbation is

$$\frac{\partial \psi}{\partial t} = c\frac{\partial^2 \psi}{\partial \xi^2} + \frac{\partial \psi}{\partial \xi} - (1-\delta(V-\theta))\psi \tag{3}$$

0.2 Separation of Variables

This equation can be solved using separation of variables, where solutions have the form of $\psi(\xi,t) = S(\xi)T(t) \neq 0$.

$$\frac{1}{T}\frac{\partial T}{\partial t} = \frac{1}{S}\frac{\partial^2 S}{\partial \xi^2} + \frac{c}{S}\frac{\partial S}{\partial \xi} - 1 + \delta(V-\theta) = \lambda$$

Before proceeding with the analysis, we need to define the boundary conditions. Recall that the boundary condition $\lim_{x\to\pm\infty}\frac{\partial v(x,t)}{\partial x} = 0$ is necessary to ensure a physically realistic solution. Because we previously imposed the boundary condition that $\lim_{\xi\to\pm\infty}\frac{dV}{d\xi} = 0$, it must also be true that

$$\lim_{x\to\pm\infty}\frac{\partial\big(V(\xi)+\epsilon\psi(\xi,t)\big)}{\partial x} = \lim_{\xi\to\pm\infty}\frac{dV}{d\xi} + \lim_{\xi\to\pm\infty}\epsilon\frac{\partial\psi}{\partial\xi} = 0 \implies \lim_{\xi\to\pm\infty}\frac{d\psi}{d\xi} = 0$$

The time-dependent equation is a homogeneous ODE with the following solutions, where $a \in \mathbb{R}$.

$$\frac{1}{T}\frac{dT}{dt} = \lambda \implies \int\frac{dT}{T} = \int\lambda dt \implies T(t) = ae^{\lambda t} \tag{4}$$

For the spatial dependence, we split the domain of the equation into three regions, depending on the value of $\delta(V(\xi)-\theta)$, with the boundary condition that $\lim_{\xi\to\pm\infty}\psi = 0$:

1) $\xi < 0$ and $\delta(V(\xi)-\theta) = 0$
2) $\xi > 0$ and $\delta(V(\xi)-\theta) = 0$
3) $\xi = 0$ and $\delta(V(\xi)-\theta) = \infty$

0.2.1 Region 1 and 2

In Region 1 and 2, we solve the equation:

$$\frac{d^2S}{d\xi^2} + c\frac{dS}{d\xi} - (\lambda+1)S = 0 \tag{5}$$

The boundary conditions are in Region 1: $\lim_{\xi\to-\infty} S_1 = 0$; and in Region 2: $\lim_{\xi\to\infty} S_2 = 0$. The solutions have the form:

$$S_1(\xi) = c_1 e^{\frac{-c+\sqrt{c^2+4(\lambda+1)}}{2}\xi} S_2(\xi) = c_2 e^{\frac{-c-\sqrt{c^2+4(\lambda+1)}}{2}\xi} \tag{6}$$

0.2.2 Region 3

At Region 3, we integrate the equation at a small interval around the point of discontinuity to examine the effect $\delta(V(\xi)-\theta)$ has on the solution. The spatial equation in this region is

$$\frac{d^2S}{d\xi^2} + c\frac{dS}{d\xi} - \big[\lambda+1+\delta(V(\xi)-\theta)\big]S = 0 \tag{7}$$

Since there is a discontinuity due to the dirac delta function here at $V(\xi)=\theta$, $\xi=0$, we integrate a small interval $(+\epsilon,-\epsilon)$ around the discontinuity. Then, we taking the limit $\epsilon\to 0$ gives us information as to how the solutions in Region 1 and 2, (6), connect in Region 3.

$$\int_{-\epsilon}^{+\epsilon}\frac{d^2S}{d\xi^2}d\xi + c\int_{-\epsilon}^{+\epsilon}\frac{dS}{d\xi}d\xi - \int_{-\epsilon}^{+\epsilon}(\lambda+1)Sd\xi - \int_{-\epsilon}^{+\epsilon}\delta(V(\xi)-\theta)Sd\xi = 0$$

For the first two terms, the fundamental theorem of calculus gives us the first derivative evaluated at the endpoints and the function evaluated at the endpoints.

$$\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon)-S(-\epsilon)] - \int_{-\epsilon}^{+\epsilon}(\lambda+1)Sd\xi - \int_{-\epsilon}^{+\epsilon}\delta(V(\xi)-\theta)Sd\xi = 0$$

Now we take the limit:

$$\lim_{\xi\to0}\left[\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon)-S(-\epsilon)] - \int_{-\epsilon}^{+\epsilon}(\lambda+1)Sd\xi - \int_{-\epsilon}^{+\epsilon}\delta(V(\xi)-\theta)Sd\xi\right] = 0$$

Of the remaining two integrals, the first one goes to zero as the bounds shrink the zero since the integral is continuous, even if S is not. The second integral involving the dirac delta function evaluates to $S(0)$ by the *sampling property*, since we are specifically looking at the small region centered around where $V(\xi)=\theta$. Therefore we get:

$$\lim_{\xi\to0}\left[\frac{dS}{d\xi}(+\epsilon) - \frac{dS}{d\xi}(-\epsilon) + c[S(+\epsilon)-S(-\epsilon)\right] = S(0)$$

S must be continuous if we are assuming that the derivative $\frac{dS}{d\xi}$ exists so $S(0^+)=S(0^-)$ and the constants in front of the homogeneous solutions (6) are equal, $c_1=c_2=K$, revealing a relationship between the value of the spatial term and the discontinuity between its derivative at $\xi = 0$.

$$S(0) = \lim_{\xi \to 0} \left[K(\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} e^{\frac{-c - \sqrt{c^2+4(\lambda+1)}}{2}} - \frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} e^{\frac{-c + \sqrt{c^2+4(\lambda+1)}}{2}}) \right]$$

Taking the limit, we find the following, where $S(0) \in \mathbb{R}$ and $K \in \mathbb{R}$ since this is a physical system and we assume the solutions are real-valued.

$$S(0) = K \left(\frac{-c - \sqrt{c^2 + 4(\lambda + 1)}}{2} - \frac{-c + \sqrt{c^2 + 4(\lambda + 1)}}{2} \right) \tag{8}$$

S is the spatial part of some arbitrary perturbation function that satisfies $\lim_{\xi \to \pm\infty} S(\xi) = 0$, and its value at zero could be either non-zero, or it could be 0. Since the boundary conditions provided do not fix $S(0)$ or K, we will examine both cases. In the latter case, if $S(0) = 0$ then $K = c_1 = 0$ and we find the trivial solution of $S = 0$ which gives us $\psi(\xi, t) = 0$. In the following section, we explore the former case.

0.3 Eigenvalues

Now to examine the family of solutions that depend on λ, we start with the $\lambda = 0$ case. Assuming we can rescale $S(0)$ to unity, eq (8) simplifies to:

$$1 = K \left(\frac{-c - \sqrt{c^2 + 4)}}{2} - \frac{-c + \sqrt{c^2 + 4)}}{2} \right) = -K\sqrt{c^2 + 4}$$

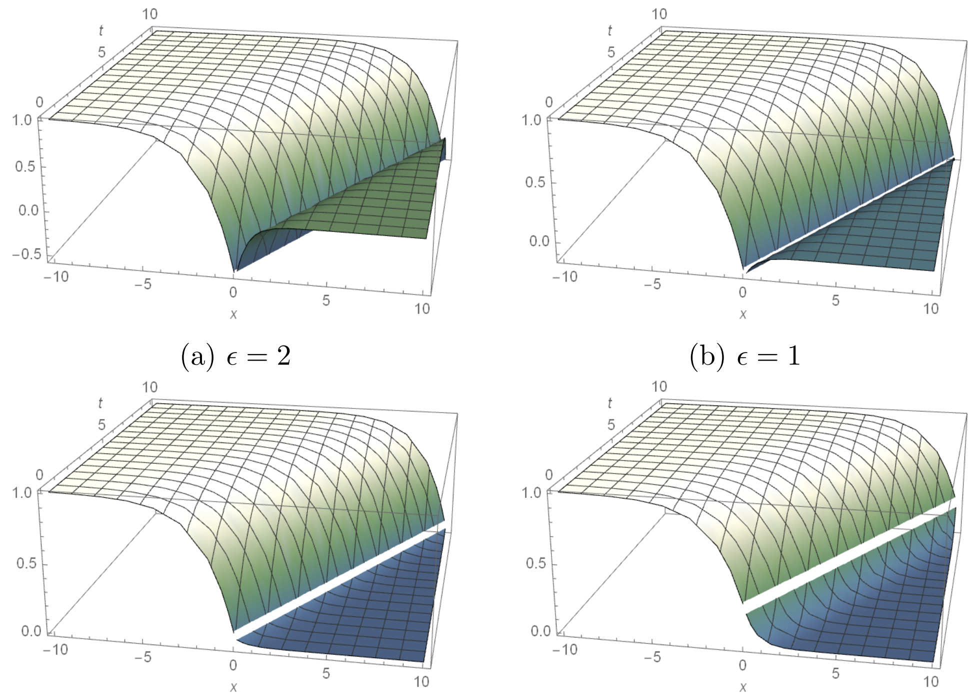
or

$$K = -\frac{1}{\sqrt{c^2 + 4}}$$

Plugging this constant back into our solutions for Region 1 and 2 (6) and multiplying it by the time-dependent solution (4), we find a solution for $\psi(\xi, t)$:

$$\psi(\xi, t) = \begin{cases} \frac{-1}{\sqrt{c^2+4}} e^{\frac{-c+\sqrt{c^2+4}}{2}}, & \xi < 0 \\ \frac{-1}{\sqrt{c^2+4}} e^{\frac{-c-\sqrt{c^2+4}}{2}}, & \xi > 0 \end{cases} \tag{9}$$

Recalling our traveling wave solution $V(\xi)$ (**??**), we note that $\psi(\xi, t) = V'(\xi)$. This solution corresponds to $v(x, t) = V(\xi) + \epsilon V'(\xi)$, a Taylor expansion linearization around the traveling wave solution. This perturbation is independent of time. We plot $v(x, t)$ with various values of ϵ to show its effect on $V(\xi)$.



(a) $\epsilon = 2$ (b) $\epsilon = 1$

(c) $\epsilon = 0.5$ (d) $\epsilon = 0.01$

Figure 1: $v(x,t) = V(\xi) + \epsilon V'(\xi)$ with various ϵ values

References

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