

# Simulation and Analysis of 1D Wave Propagation under Various Physical Models

Dario Liotta



UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA



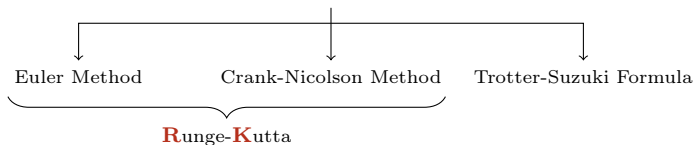
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di Fisica  
e Astronomia  
Galileo Galilei

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Academic Year 2024/2025

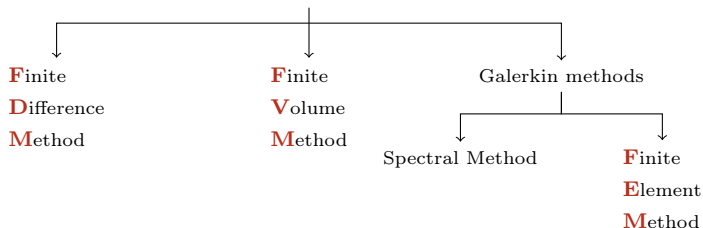
# Numerical methods for differential equations

## Ordinary Differential Equations

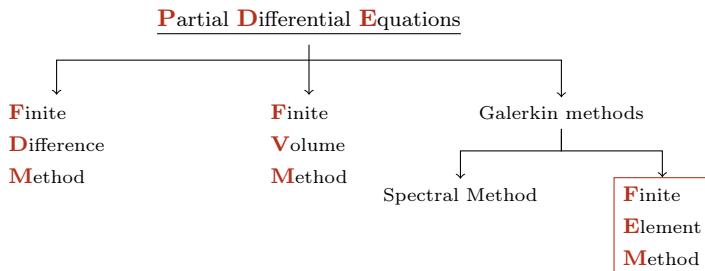
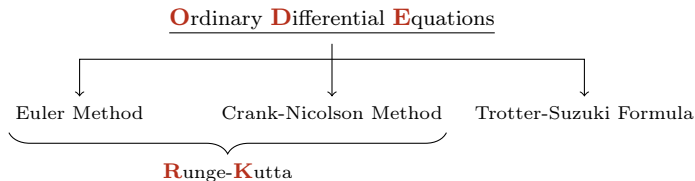



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## Partial Differential Equations



# Numerical methods for differential equations



# Introduction to the problem

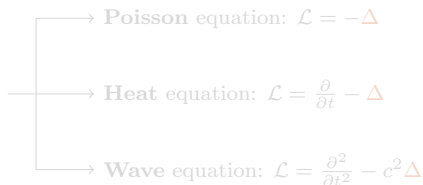
Solving a **PDE** means to find a function  $u$  such that

$$\mathcal{L}u = f$$

where  $\mathcal{L}$  is a differential operator and  $f$  is a source term.

The equation holds in a domain  $\Omega$  and is completed by prescribing **boundary conditions** on  $\partial\Omega$ .

In most physical  
applications  $\mathcal{L}$  is a  
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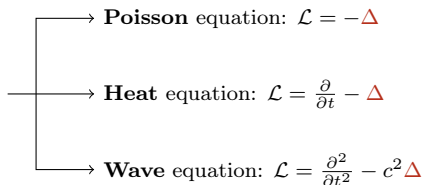
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# Weak formulation

Galerkin methods rely on a **weak formulation**

- Multiply by a **test function**  $v$  and integrate over the entire domain

$$-\int_{\Omega} (\Delta u) v d\Omega = \int_{\Omega} f v d\Omega$$

- Integrate by parts the left hand side

$$-\int_{\Omega} (\Delta u) v d\Omega = \int_{\Omega} \nabla u \cdot \nabla v d\Omega - \int_{\partial\Omega} \frac{\partial u}{\partial n} v ds$$

- Substitute and get the new expression

$$\int_{\Omega} \nabla u \cdot \nabla v d\Omega = \int_{\Omega} f v d\Omega + \int_{\partial\Omega} \frac{\partial u}{\partial n} v ds$$

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# About the test function

The test function  $v$  is introduced to check whether the PDE is satisfied on average throughout the domain.

The problem becomes to find  $u$  such that

$$a(u, v) = F(v) \quad \forall v \in V$$

where

$$a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v d\Omega \quad \text{is a bilinear form}$$

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# Benefits of the weak formulation

## Strong formulation

$$u \in C^2(\Omega)$$

Holds pointwise in  $\Omega$

Derivatives exist classically

## Weak formulation

$$u, v \in H^1(\Omega)^*$$

Holds on average on  $\Omega$

Derivatives exist in the  
distributional sense

In short: weak formulation requires **less regularity**

\*  $H^1(\Omega)$  is a **Sobolev space** of functions with square-integrable first derivatives:

$$w \in H^1(\Omega) = \left\{ w \in L^2(\Omega) \mid \nabla w \in L^2(\Omega)^d \right\}$$

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# On boundary conditions

Another difference lies in the boundary condition prescription.



$v = 0$  on  $\partial\Omega \Rightarrow$  cancels boundary term  
(no information available on  $\frac{\partial u}{\partial n}$ )

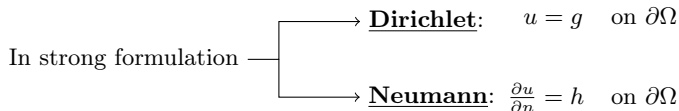
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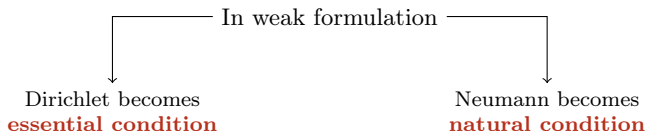
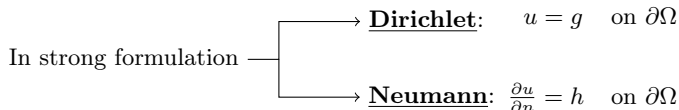
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# Shape functions

Galerkin methods allow to find an approximate solution

$$u_h \in V_h \subset H^1(\Omega) \quad \text{where } V_h \text{ is a **finite-dimensional** space}$$

In this framework, the goal is to find  $u_h$  such that

$$a(u_h, v_h) = F(v_h) \quad \forall v_h \in V_h$$

A **basis of function**  $\{\phi_i\}$  is chosen to express  $u_h$  and to use it as test:

$$u_h = \sum_{j=1}^N u_j \phi_j \implies a \left( \sum_{j=1}^N u_j \phi_j, \phi_i \right) = F(\phi_i) \quad \forall i = 1, \dots, N$$

Functions  $\phi_i$  model the solution  $\longrightarrow$  **shape functions**

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# Mesh discretization

**FEM** approach consists in the subdivision of the domain in a so-called **mesh**

This choice brings several advantages:

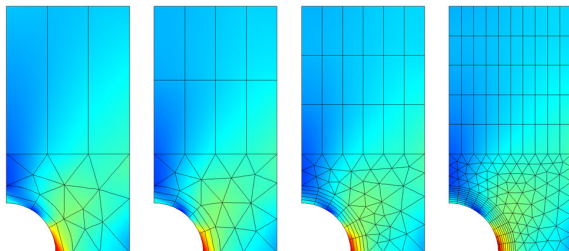
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- Better capture of **local effects**
- Possibility of **adaptive refinement**
- Natural construction of a **global solution**

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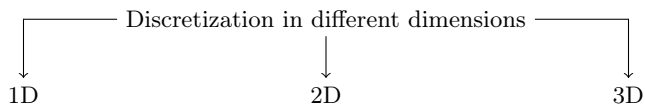
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# Elements



# Application examples

# Title

In **F**inite **E**lement **M**ethod,  $V_h$  is generated by **local basis functions** with compact support. Usually, one discretizes the domain using a so-called mesh