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Work supervised by F. G. Scholtz

Room 316 TALK:



PINHOLE INTERFERENCE IN 3D FUZZY SPACE

A natural quantum-to-classical transition

Stellenbosch University
August 2023

Outline



1 SETUP

- Motivation & goal
- Existing work

2 Fuzzy space FORMALISM

- States & observables
- Position measurement
- Coordinate representation

3 FREE PARTICLE solutions

- Plane & spherical waves

4 INTERFERENCE calculation

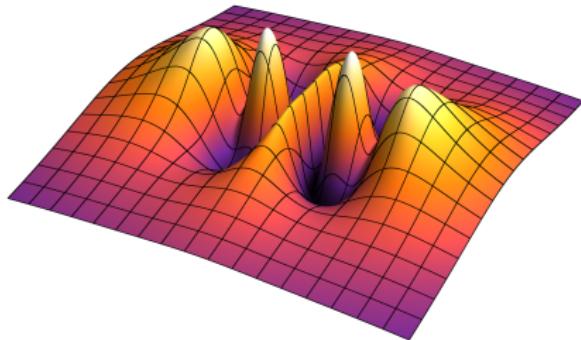
- Commutative & non-commutative

5 DISCUSSION

- Interference suppression
- Many-particles

6 SUMMARY

THE SETUP



Motivation & Goal

A lofty question



Question:

Can the structure of spacetime at **SMALLEST LENGTH SCALE** affect the physics we perceive at **LARGER LENGTH SCALES** (*i.e. classical physics*)?

Motivation & Goal

A concrete investigation



Our investigation:

Study **DOUBLE-PINHOLE SETUP** in the 3D **FUZZY SPACE** formalism of
NON-COMMUTATIVE quantum mechanics.



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- No localised states (**POSITION EIGENSTATES**)
- **MINIMUM LENGTH** scale \longleftarrow set by a parameter λ
- Otherwise **ORDINARY QUANTUM MECHANICS** !

Motivation & Goal

Why this particular investigation?



Why non-commutative geometry?

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- Doplicher et al. (1995):
 - | Non-trivial **SMALL-SCALE** spacetime structure (esp. min length)
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Why pinhole interference?

- Illustrative **TOY MODEL** ← quantum behaviour = interference
- **QUANTIFIABLE SUPPRESSION** strength
- Good setup for **EXPERIMENTAL TESTING**

Quantum-to-Classical Transition

Existing work in lower dimensions



Pittaway & Scholtz (2021)

2-SLITS in non-commutative 2D MOYAL PLANE

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- large PARTICLE NUMBER, N

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- Moyal plane interference pattern:
 - is **ASYMMETRIC** under reflection ← ∵ Moyal commutation relations break rotational symmetry
 - has **UNOBSERVABLE SUPPRESSION**



Moyal plane interference

$$P(\mathbf{D}) = 1 + \underbrace{\exp\left[-\frac{N\theta m^2}{2\hbar^2}\mathbf{v}^2(1 - \cos\alpha)\right]}_{\text{interference suppression}} \underbrace{\cos(\dots)}_{\text{interference}}$$

Quantum-to-Classical Transition

Why Moyal plane suppression is unobservable



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interference suppression interference

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- $m \sim 10^{-31} \text{ kg}$ ← mass of proton
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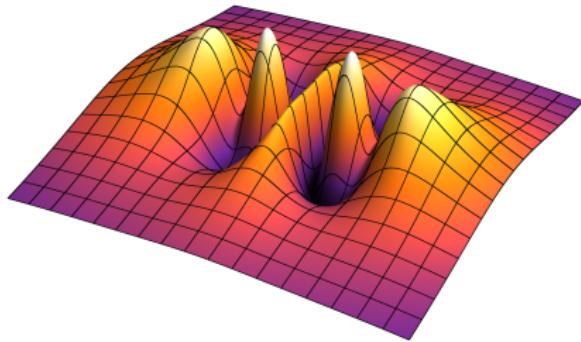
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- $\left. \begin{array}{l} \\ \\ \end{array} \right\} \implies \|\mathbf{v}\| \gtrsim 10^{16} \text{ m s}^{-1} !$

THE FORMALISM





Definition (Fuzzy space)

1 $\mathfrak{su}(2)$ ALGEBRA $\longrightarrow [\hat{x}_i, \hat{x}_j] = 2i\lambda \epsilon_{ijk} \hat{x}_k$



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Decompose into **IRREPS**:



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\implies 1 copy of each quantised radius



Definition (Quantum state space)

1 HILBERT SPACE

$$\rightarrow \quad \mathcal{H}_q := \left\{ \begin{array}{l} \text{operators on } \mathcal{H}_c \\ \text{generated by } \hat{x}_i \end{array} \right\}$$
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\uparrow
 \backslash
 \diagdown
 so $\left\| \hat{\text{Proj}}_{[1] \oplus \dots \oplus [N]} \right\|$
 $\sim \frac{4}{3}\pi [\lambda(N+1)]^3$



Definition (Observables)

HERMITIAN OPERATORS on \mathcal{H}_q



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Important observables

- POSITION \longrightarrow $\hat{X}_i |\psi\rangle := |\hat{x}_i \psi\rangle$, $\hat{R} |\psi\rangle := |\hat{r} \psi\rangle$ \longleftarrow like normal QM



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- HAMILTONIAN \longrightarrow $\hat{H} = -\frac{\hbar^2}{2m} \hat{\Delta} + V(\hat{R})$ \longleftarrow like normal QM

Physical Subspace



Physical subspace

$$\mathcal{H}_q = \bigoplus_{n \in \mathbb{N}} [\mathbf{n}] \otimes [\mathbf{n}]^* \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$$



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Projection onto \mathcal{H}_q

1 Conserved **OBSERVABLE** $\longrightarrow \hat{\Gamma} |\psi\rangle := |[\hat{n}, \psi]\rangle$

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- 2 **KERNEL** $\longrightarrow \mathcal{H}_q = \ker \hat{\Gamma}$
- 3 **PROJECTION** $\longrightarrow \hat{Q} := \frac{1}{2\pi} \int_0^{2\pi} e^{i\phi \hat{\Gamma}} d\phi$

Position Measurement

... as a POVM



Definition (Position measurement)

1 Position EIGENSTATES \longrightarrow MINIMUM-UNCERTAINTY states

$$|\mathbf{z}\rangle := e^{-\frac{1}{2}\bar{z}_\alpha z_\alpha} e^{z_\alpha a_\alpha^\dagger} |0\rangle, \text{ for}$$

$$\mathbf{z} := e^{i\gamma} \sqrt{\frac{r}{\lambda}} \begin{bmatrix} \cos(\frac{\theta}{2}) e^{-i\frac{\phi}{2}} \\ \sin(\frac{\theta}{2}) e^{i\frac{\phi}{2}} \end{bmatrix} \in \mathbb{C}^2 \quad \leftarrow \text{ encodes } \mathbf{D} = (r, \theta, \phi)$$

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2 POVM

$$\longrightarrow |z_1, z_2, n_1, n_2\rangle_{\text{ph}} := \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2\hat{r}}} |\mathbf{z}\rangle\langle n_1, n_2|$$

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3 BORN RULE $\longrightarrow P(\mathbf{D}) = \text{Tr}_{\text{q}} (\hat{\Pi}_{\mathbf{z}} \rho)$

Position Measurement

... as weak measurement



Weak measurement picture

1 $\mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$



Weak measurement picture

1 $\mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$

2 \hat{X}_i act only on \mathcal{H}_c



Weak measurement picture

- 1 $\mathcal{H}_q \subset \mathcal{H}_c \otimes \mathcal{H}_c^*$
- 2 \hat{X}_i act only on \mathcal{H}_c
- 3 \therefore POSITION measurement
 - LOCAL measurement
 - traces out “ENVIRONMENT”, \mathcal{H}_c^*

Coordinate Representation

The analogue of wavefunctions



Definition (Symbol & star product)

1 POSITION-ENCODING states $\longrightarrow |z\rangle := \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2 r}} |z\rangle\langle z|$

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ψ is indep of γ
 \therefore function on \mathbb{R}^3

Coordinate Representation

The analogue of wavefunctions



Definition (Symbol & star product)

- 1 POSITION-ENCODING states $\longrightarrow |z\rangle := \hat{Q} \frac{1}{\sqrt{4\pi\lambda^2\hat{r}}} |z\rangle\langle z|$
- 2 COORDINATE REP $\longrightarrow \psi(z) := (z|\psi) = \langle z|\sqrt{4\pi\lambda^2\hat{r}}\psi|z\rangle$
- 3 SYMBOL $\longrightarrow \langle z|\psi|z\rangle$

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4 VOROS PRODUCT $\longrightarrow \star := \exp\left[\overleftarrow{\partial}_{z_\alpha} \overrightarrow{\partial}_{\bar{z}_\alpha}\right]$



Notable properties

1 **COMPLETENESS** $\longrightarrow \int \frac{d^4z}{\pi^2} |\mathbf{z}) \bar{x}(\mathbf{z}| = \mathbf{1}_q$



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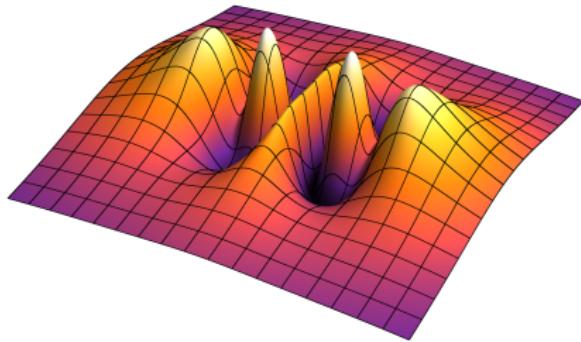
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\implies Alternate “wave-mechanics” development !



FREE PARTICLE SOLUTIONS





Non-commutative free Schrödinger equation

$$\hat{H}|\psi\rangle = -\frac{\hbar^2}{2m}\hat{\Delta}|\psi\rangle = E|\psi\rangle$$



Non-commutative free Schrödinger equation

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Types of solutions:

1 **PLANE WAVE** $\longrightarrow |k\rangle := e^{i\mathbf{k}\cdot\hat{\mathbf{x}}}$ \longleftarrow typical form

2 **SPHERICAL WAVE** $\longrightarrow |k, l, m\rangle = g_l(\hat{n} - l, k) \hat{\mathcal{Y}}_{lm}$

Plane wave solutions



1 PLANE WAVE $\longrightarrow |k\rangle := e^{ik \cdot \hat{x}}$



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$$\cos(\lambda k_3) = \cos(\lambda k_1) \cos(\lambda k_2) - \hat{k}_1 \cdot \hat{k}_2 \sin(\lambda k_1) \sin(\lambda k_2),$$

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- DISPERSION RELATION $\longrightarrow \hat{H} |k\rangle = \frac{2\hbar^2}{m\lambda^2} \sin^2\left(\frac{k\lambda}{2}\right) |k\rangle$

Spherical wave solutions



2 SPHERICAL WAVE $\longrightarrow |k, l, m\rangle = g_l(\hat{n} - l, k) \hat{\mathcal{Y}}_{lm}$

Spherical wave solutions



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Properties

- Permit 2 RADIAL SOLUTIONS $\longrightarrow g_l = A g_{J,l} + B g_{Y,l}$
c.f. spherical Bessel- & Neumann



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Spherical wave solutions



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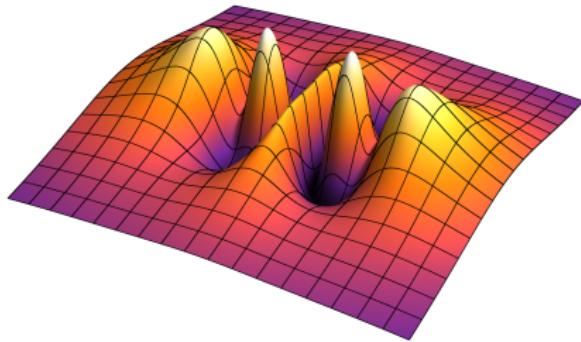
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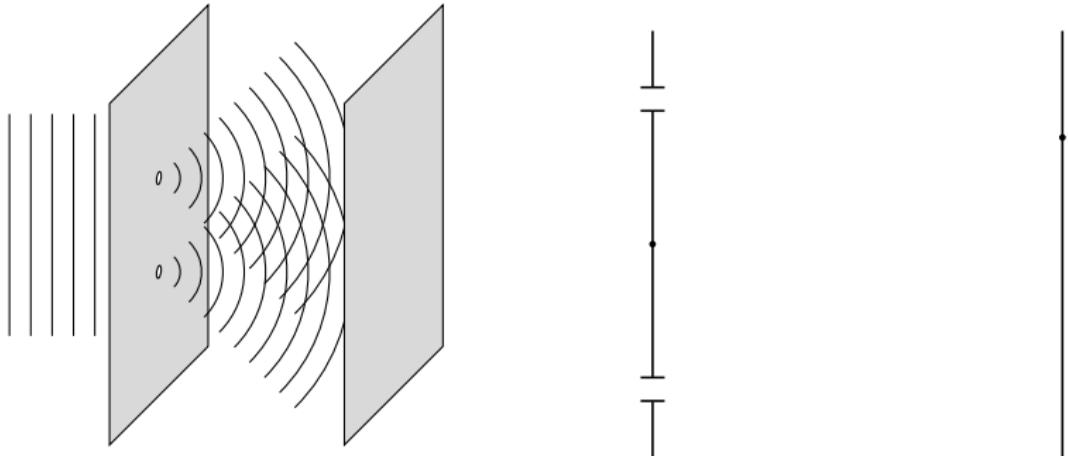
$$\therefore \langle \mathbf{z} | g_{H,l}(\hat{n}, k) | \mathbf{z} \rangle \sim \frac{e^{r(\cos(\lambda k) + i \sin(\lambda k) - 1)/\lambda}}{(ir/\lambda)^{l+1}}$$

MAIN CALCULATION



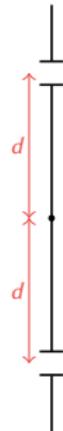
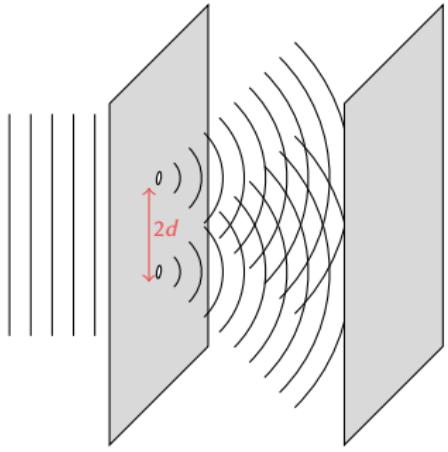
Double-Pinhole Setup

The most famous quantum experiment



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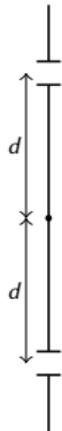
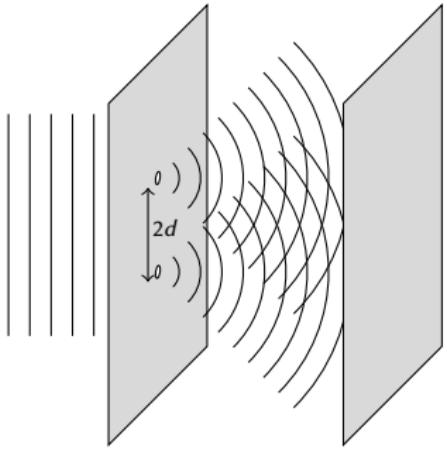
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■ PINHOLES $\longrightarrow z = \pm d$

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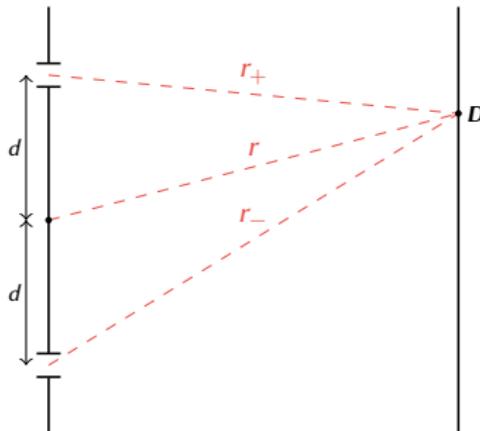
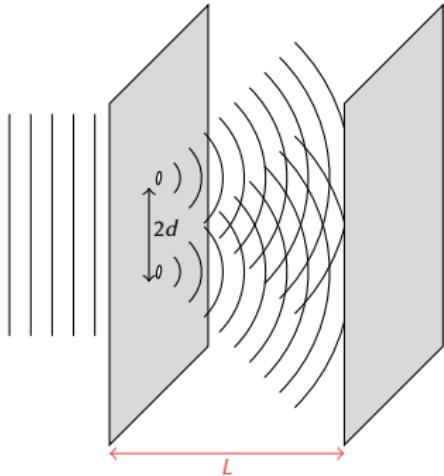
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- PINHOLES $\longrightarrow z = \pm d$
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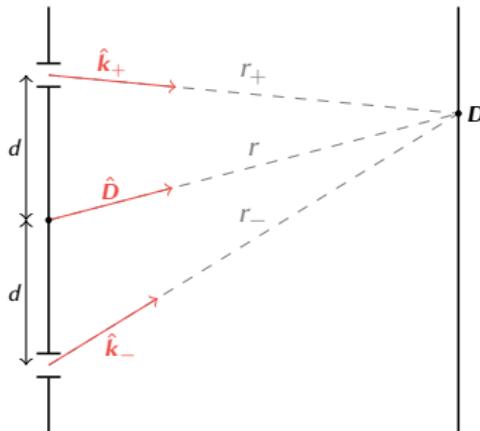
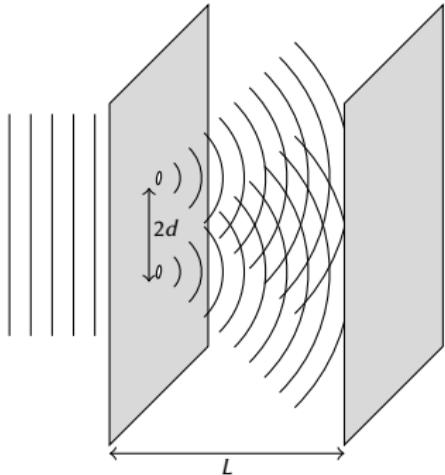
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- PINHOLES $\rightarrow z = \pm d$
- DETECTION POINT $\rightarrow D = (L, y_D, z_D) \equiv (r, \theta, \phi)$
- DISTANCES $\rightarrow L, r, r_{\pm}$ \leftarrow each $\gg d$: large separation approx

Double-Pinhole Setup

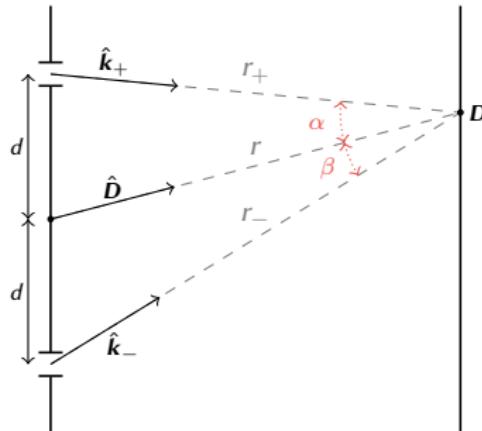
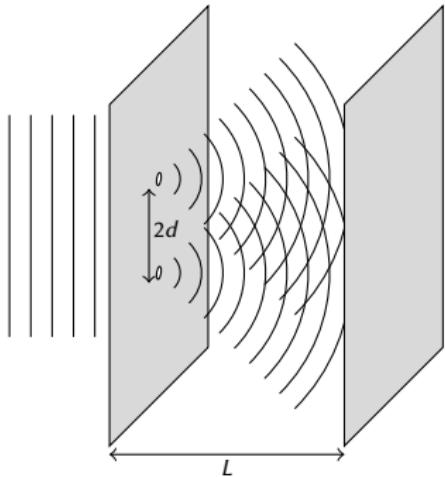
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- ANGLES $\rightarrow \alpha, \beta$

Interference Calculation

Broad strategy



Interference calculation overview

1 STATE at D $\longrightarrow \psi \sim \left\{ \begin{array}{l} \text{spherical wave} \\ \text{from } z = +d \end{array} \right\} + \left\{ \begin{array}{l} \text{spherical wave} \\ \text{from } z = -d \end{array} \right\}$



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- 3 BORN RULE $\rightarrow P(D) = \text{Tr}(\hat{\Pi}_D \rho)$

Interference Calculation

Overview of commutative treatment



Commutative interference calculation

1 STATE at D $\longrightarrow \psi(D) \sim \frac{1}{r_+} e^{ikr_+} + \frac{1}{r_-} e^{ikr_-}$ \longleftarrow asymptotic form of spherical Hankel

Interference Calculation

Overview of commutative treatment



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$$P_{\text{comm}}(\mathbf{D}) \sim \frac{1}{r_+ r_-} \left[\underbrace{\frac{2d^2}{r_+ r_-} + \cos(\alpha + \beta)}_{\text{bimodal shaping function}} + \underbrace{\cos(rk(\cos \alpha - \cos \beta))}_{\text{interference terms}} \right]$$

Interference Calculation

Overview of non-commutative treatment



Non-commutative interference calculation

1 SYMBOL at $D \longrightarrow \langle z|\psi|z\rangle \sim \langle z^+|g_k(\hat{n})|z^+\rangle + \langle z^-|g_k(\hat{n})|z^-\rangle$

Interference Calculation

Overview of non-commutative treatment



Non-commutative interference calculation

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- 2 **PARAXIAL** approximation \longrightarrow $\langle \mathbf{z}^\pm | g_k(\hat{n}) | \mathbf{z}^\pm \rangle \sim \langle \mathbf{z} | \eta_\pm e^{i\mathbf{k}_\pm \cdot \hat{\mathbf{x}}} | \mathbf{z} \rangle$

Interference Calculation

Overview of non-commutative treatment



Non-commutative interference calculation

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Interference Calculation

Overview of non-commutative treatment



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- 4 **BORN RULE** $\longrightarrow P(D) = \text{Tr}_q \left(\hat{\Pi}_{\mathbf{z}} |\psi\rangle\langle\psi| \right)$
- 5 Compute remaining **MATRIX ELEMENTS**

Interference Calculation

The main result!



$$P(\mathbf{D}) \sim \underbrace{\frac{\eta_+^2 + \eta_-^2}{2} \left(\frac{r}{\lambda} + 1 \right)}_{\text{bimodal shaping function}} + \underbrace{\eta_+ \eta_- e^{A-r/\lambda}}_{\text{interference suppression}} \times \underbrace{((A+1) \cos B - B \sin B)}_{\text{interference terms}}$$

where $\begin{cases} \eta_{\pm} := \frac{\lambda}{r_{\pm}} \exp \left[\frac{1}{\lambda} (r_{\pm} - r) (\cos(\lambda k) - 1) \right], \\ A := \frac{r}{\lambda} (\cos^2(\lambda k) + \cos(\alpha + \beta) \sin^2(\lambda k)), \\ B := \frac{r}{\lambda} \sin(\lambda k) \cos(\lambda k) (\cos \alpha - \cos \beta) \end{cases}$

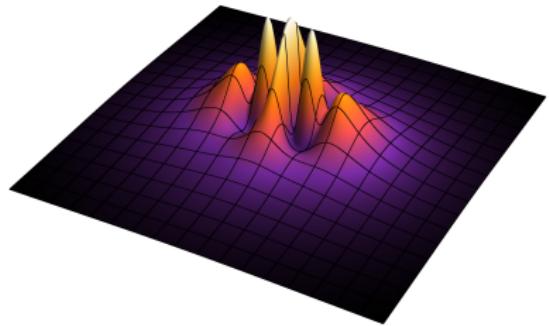
Interference Calculation

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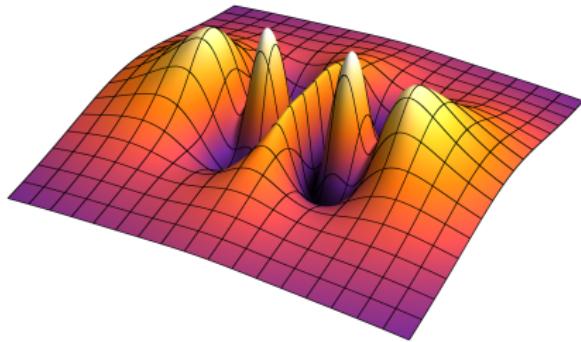


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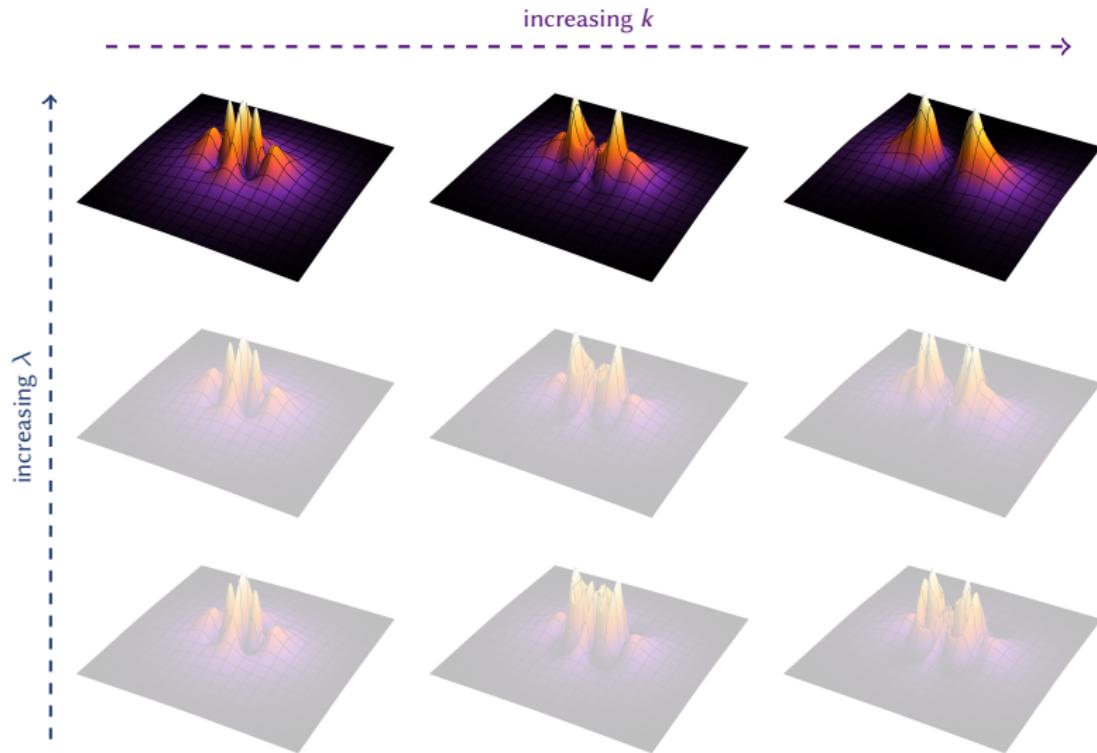


DISCUSSION OF RESULTS



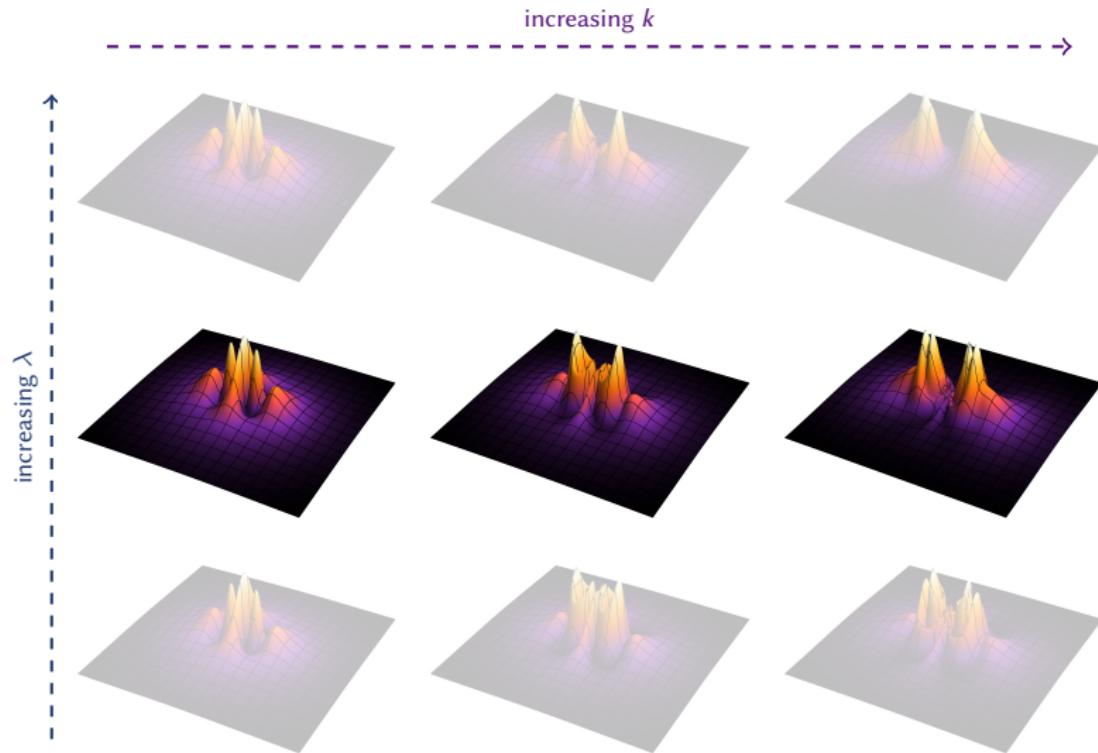
Qualitative Behaviour

& commutative limit



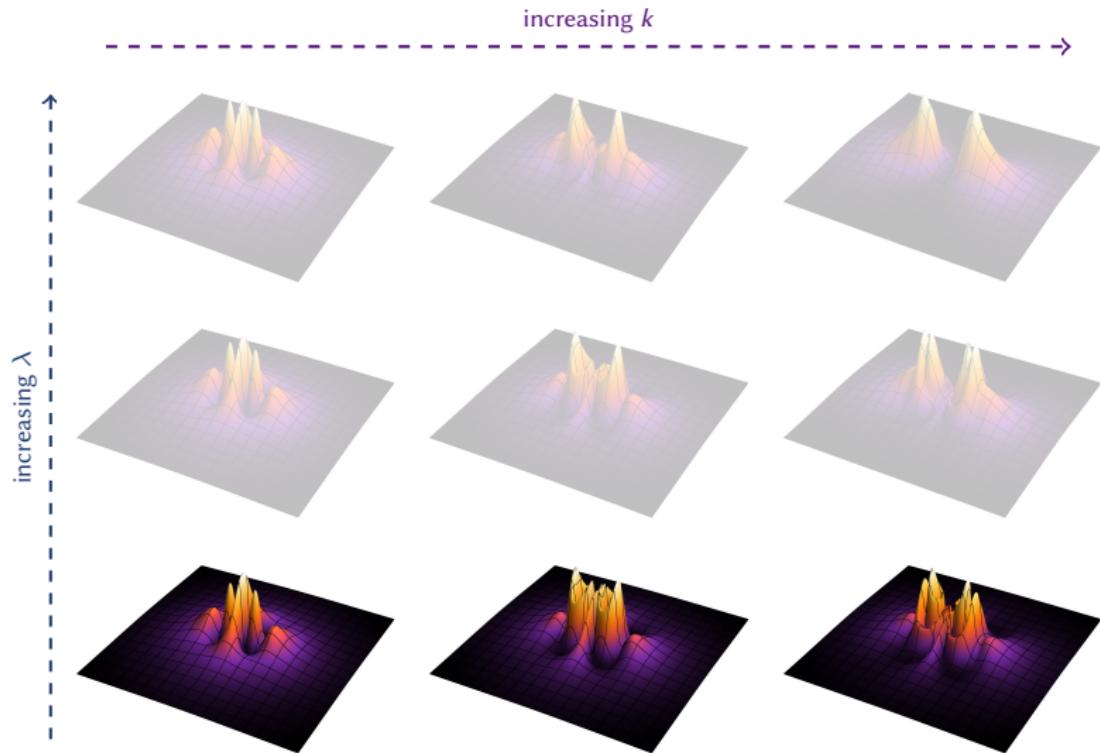
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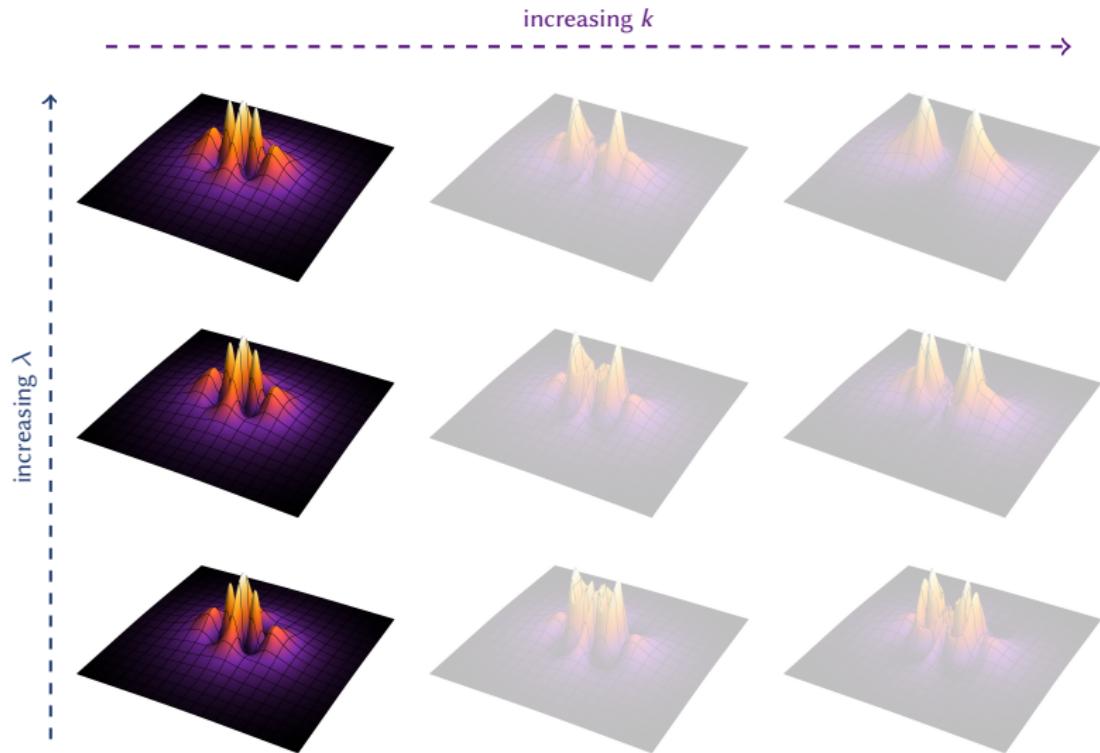
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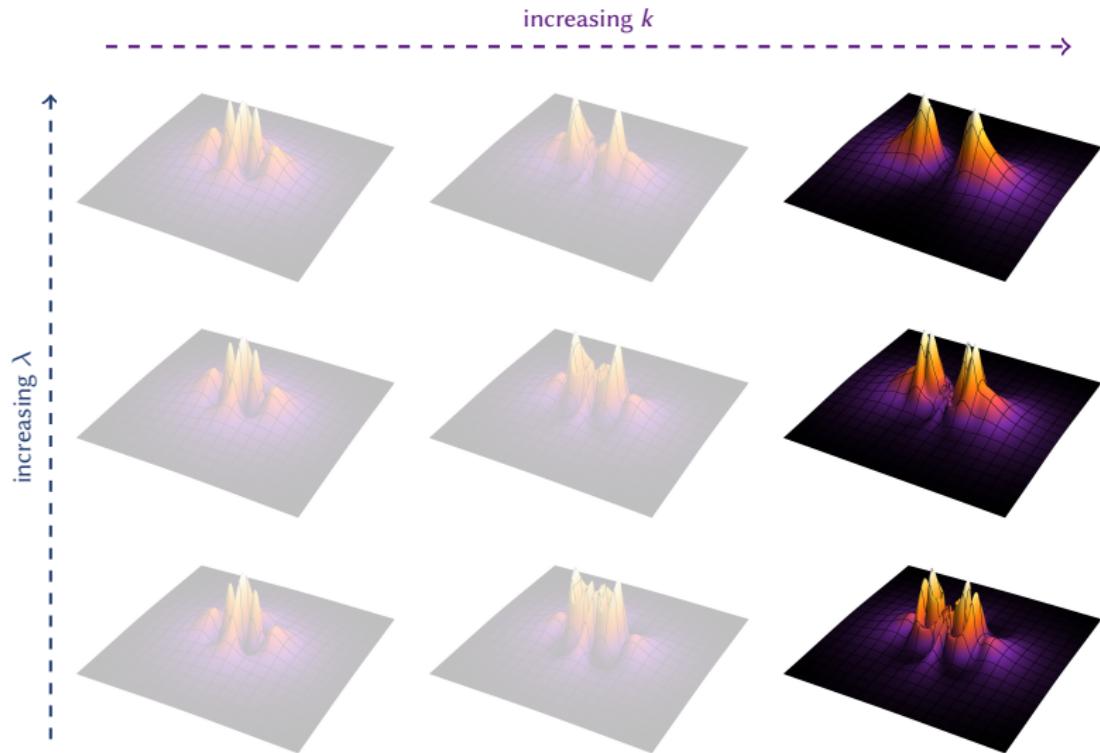
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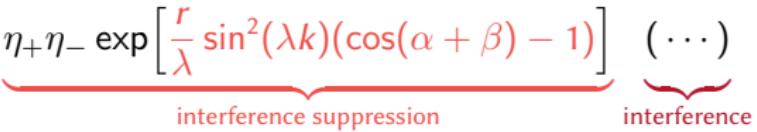
Quantum-to-Classical Transition

Is the suppression observable now?



Fuzzy space interference

$$P(\mathbf{D}) \sim \dots + \eta_+ \eta_- \exp \left[\frac{r}{\lambda} \sin^2(\lambda k) (\cos(\alpha + \beta) - 1) \right] (\dots)$$



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Observable **SUPPRESSION** \implies exponent $\sim -\frac{4\lambda d^2 m E}{r \hbar^2} \lesssim -1$

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- $m \sim 10^{-31} \text{ kg}$ ← mass of electron
- $\lambda \sim 10^{-35} \text{ m}$ ← Planck length

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 - $\lambda \sim 10^{-35} \text{ m}$ \leftarrow Planck length
- $\left. \begin{array}{l} \\ \\ \end{array} \right\} \implies r \lesssim 10^{-21} \text{ m} !$

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Fuzzy space interference

$$P(D) \sim \dots + \eta_+ \eta_- \exp \left[\underbrace{\frac{r}{\lambda} \sin^2(\lambda k) (\cos(\alpha + \beta) - 1)}_{\text{interference suppression}} \right] \underbrace{(\dots)}_{\text{interference}}$$

$$\text{Observable SUPPRESSION} \implies \text{exponent} \sim -\frac{4\lambda d^2 m E}{r \hbar^2} \lesssim -1$$

- $E \sim 1 \text{ eV}$
 - $m \sim 10^{-31} \text{ kg}$ \leftarrow mass of electron
 - $\lambda \sim 10^{-35} \text{ m}$ \leftarrow Planck length
- $\left. \begin{array}{l} \\ \\ \end{array} \right\} \implies r \lesssim 10^{-21} \text{ m} !$

Important features

- r dependence
- suppression possible at low k

Macroscopic Scaling

Extending definitions



N particles

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4 PLANE WAVES $\longrightarrow |\mathbf{k}^{(i \cdots N)}\rangle = \exp\left[i \sum_{n=1}^N \mathbf{k}^{(n)} \cdot \hat{\mathbf{x}}^{(n)}\right]$

\vdots

Macroscopic Scaling

Center-of-mass coordinates



N particles

5 CENTER-OF-MASS frame

$$\longrightarrow \quad \hat{\mathbf{x}}^{(CM)} := \frac{1}{N} \sum_{n=1}^N \hat{\mathbf{x}}^{(n)}, \quad \hat{\boldsymbol{\xi}}^{(n)} := \hat{\mathbf{x}}^{(n)} - \hat{\mathbf{x}}^{(CM)}$$

$$\mathbf{k}^{\text{tot}} := \sum_{n=1}^N \mathbf{k}^{(n)}, \quad \mathbf{q}^{(n)} := \mathbf{k}^{(n)} - \frac{1}{N} \mathbf{k}^{\text{tot}}$$

⋮

Macroscopic Scaling

Center-of-mass dynamics



N particles

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Macroscopic Scaling

Center-of-mass dynamics



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6 SPLIT Hamiltonian^{*} \longrightarrow

$$\hat{H}^{\text{tot}} = \hat{H}_{\text{free}}^{\text{tot}} + \hat{H}_{\text{interaction}}^{\text{tot}}$$

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$$\hat{H}_{\text{free}}^{\text{tot}} + \hat{H}_{\text{interaction}}^{\text{tot}} = 0$$

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7 NEGLECT corrections \longrightarrow EXPAND in λ and $T \sim \frac{\hbar^2(q^{(n)})^2}{mk_B}$

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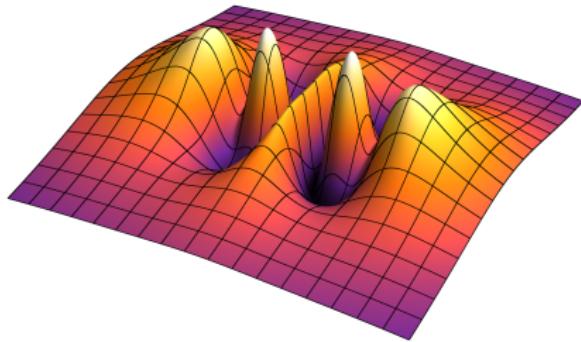
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- $\left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \implies r \lesssim 100 \text{m}$

CONCLUDING REMARKS



Experimental Prospects

How might we observe suppression practically?



Challenges

- 1 **CREATE & MANIPULATE** massive quantum superposition

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 - *Tebbenjohanns et al. (2021)* → control optically-levitated femtogram nanoparticle
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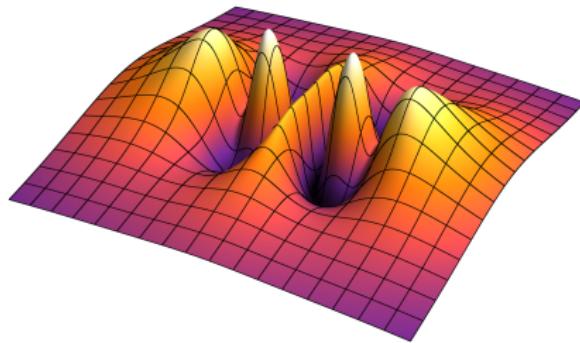
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 - *Kleckner et al. (2008)* → describe entangling photon with macroscopic cantilever



Summary

Key takeaways

- Fuzzy space → **CLASSICAL TRANSITION** without heat bath
- Quantum suppression → realistically **OBSERVABLE**
- Suppression strength → **EXTENSIVE & DISTANCE**-dependent





Future Work Proposals

1 ALTERNATE SETUP:

- Treat fuzzy-space VON NEUMANN MEASUREMENT

2 FORMALISM EXTENSIONS:

- Extend to non-commutative QFT

3 EXPERIMENTAL VERIFICATION:

- Implement proposed experiment
- Devise alternate experiment



Trinchero, D., & Scholtz, F. G. (2023, March).

Pinhole interference in three-dimensional fuzzy space.

Annals of Physics, 450, 169224.

(arXiv:2212.01449 [quant-ph])



References I

- Alekseev, A. Y., Recknagel, A., & Schomerus, V. (2000). Brane dynamics in background fluxes and non-commutative geometry. *Journal of High Energy Physics*, 2000(05), 010.
- Doplicher, S., Fredenhagen, K., & Roberts, J. E. (1995). The quantum structure of spacetime at the planck scale and quantum fields. *Communications in Mathematical Physics*, 172(1), 187–220.
- Fried, D. G., Killian, T. C., Willmann, L., Landhuis, D., Moss, S. C., Kleppner, D., & Greytak, T. J. (1998, Nov). Bose-einstein condensation of atomic hydrogen. *Phys. Rev. Lett.*, 81, 3811–3814.
- Kleckner, D., Pikovski, I., Jeffrey, E., Ament, L., Eliel, E., van den Brink, J., & Bouwmeester, D. (2008, sep). Creating and verifying a quantum superposition in a micro-optomechanical system. *New Journal of Physics*, 10(9), 095020.
- Kriel, J. N., Groenewald, H. W., & Scholtz, F. G. (2017). Scattering in a three-dimensional fuzzy space. *Physical Review D*, 95(2), 025003.



References II

- Pittaway, I. B., & Scholtz, F. G. (2021). Quantum interference on the non-commutative plane and the quantum-to-classical transition. *arXiv e-prints*, arXiv-2101.
- Scholtz, F. G., Gouba, L., Hafver, A., & Rohwer, C. M. (2009). Formulation, interpretation and application of non-commutative quantum mechanics. *Journal of Physics A: Mathematical and Theoretical*, 42(17), 175303.
- Seiberg, N., & Witten, E. (1999). String theory and noncommutative geometry. *Journal of High Energy Physics*, 1999(09), 032.
- Tebbenjohanns, F., Mattana, M. L., Rossi, M., Frimmer, M., & Novotny, L. (2021). Quantum control of a nanoparticle optically levitated in cryogenic free space. *Nature*, 595(7867), 378–382.
- van Es, J. J. P., Whitlock, S., Fernholz, T., van Amerongen, A. H., & van Druten, N. J. (2008). Longitudinal character of atom-chip-based rf-dressed potentials. *Physical Review A*, 77(6), 063623.