Prethermalization of quantum systems interacting with non-equilibrium environments

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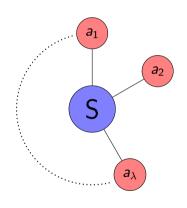
Total Hamiltonian:

$$H = H_{\rm S}$$

Components:

$$H_{\mathrm{S}}=rac{1}{2}\omega_{0}\sigma_{z}$$



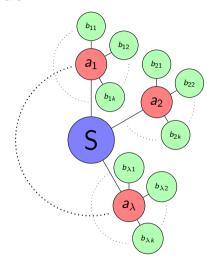


Total Hamiltonian:

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$$egin{aligned} H_{
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m RI} &= \sum_{\lambda} \omega_\lambda a_\lambda^\dagger a_\lambda \quad H_{
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► We model the spectral functions of the reservoirs with the phenomenological model:

$$J_i(\omega) = g_i \omega^{s_i} e^{-\omega/\omega_c}$$
, $i = \text{RI}, \text{RII}$

Methods

Redfield Master Equation

$$egin{aligned} rac{d}{dt}
ho_{\mathrm{S}}(t) &= -i\left[H_{\mathrm{S}},
ho_{\mathrm{S}}(t)
ight] + \left(\int_{0}^{t}d aulpha^{+}(t, au)\left[V_{ au-t}\sigma_{+}
ho_{\mathrm{S}}(t),\sigma_{-}
ight] \ &+ \int_{0}^{t}d aulpha^{-}(t, au)\left[V_{ au-t}\sigma_{-}
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ight)\,, \end{aligned}$$

where $V_t \mathcal{O} = e^{iH_{\rm S}t} \mathcal{O} e^{-iH_{\rm S}t}$ represents the free evolution of the operator.

Methods

Canonical form of the Master Equation

$$rac{d}{dt}
ho_{\mathrm{S}}(t) = -i[H(t),
ho_{\mathrm{S}}(t)] + \sum_{k=1}^{d^2-1} \gamma_k(t) \Big(L_k(t)
ho_{\mathrm{S}}(t)L_k^\dagger(t) - rac{1}{2}\{L_k^\dagger(t)L_k(t),
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with

$$\begin{split} \gamma_+(t) &= J_{\rm I}(\omega_0) n_{\rm I}(\omega_0) e^{-J_{\rm II}(\omega_0)t} + J_{\rm I}(\omega_0) n_{\rm II}(\omega_0) (1 - e^{-J_{\rm II}(\omega_0)t}) \;, \quad L_+ = \sigma_+ \\ \gamma_-(t) &= J_{\rm I}(\omega_0) (n_{\rm I}(\omega_0) + 1) e^{-J_{\rm II}(\omega_0)t} + J_{\rm I}(\omega_0) (n_{\rm II}(\omega_0) + 1) (1 - e^{-J_{\rm II}(\omega_0)t}) \;, \quad L_- = \sigma_- \end{split}$$
 where

$$n_i(\omega) = rac{1}{e^{-eta_{\mathrm{i}}\omega}-1}$$

This model exhibits an instance of prethermalization: relaxation of the open system to a state that is not the true asymptotic state

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When present, we can differentiate the following steps:

- 1. Relaxation of any initial condition to a thermal state determined by the temperature of RI.
- 2. The system remains stationary in that state
- 3. Final relaxation towards a thermal state determined by the temperature of RII.

Animations!!

Trace distance between the prethermal state and the evolved state of the system to study dependence of prethermalization time on other parameters.

$$T(\rho_{\mathcal{S}}(t), \rho_{\mathcal{S}}^{\text{th}}(\beta_{\mathcal{I}})) = \frac{1}{2} \text{Tr} \left\{ \sqrt{(\rho_{\mathcal{S}}(t) - \rho_{\mathcal{S}}^{\text{th}}(\beta_{\mathcal{I}}))^2} \right\}$$
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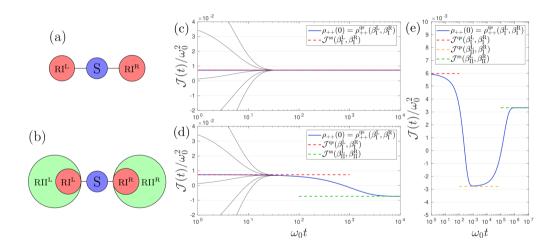
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- ► Exponential dependence with the coupling strength between environments.
- ▶ When the temperatures between reservoirs are closer, it becomes larger.
- ► Hotter RI yields longer prethermalization times, as well as colder RII.

Results - Multiple environments



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- We extended the usual methods used to study OQS to explore this more complex scenario.
- ► This model allows to indirectly control the asymptotic state of a system by modifying an environment that is not in direct contact with it
- ► Model reminiscent of the layered structure of quantum computers, with different layers that are colder close to the qubits.