# Global existence for nonlinear heat equation and harmonic maps between Riemannian manifolds

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#### **Contents**

Estimate of higher derivatives.

2

#### Global existence for nonlinear heat equation.

Let M be a compact Riemannian manifold. We want to solve the following nonlinear heat equation where  $F: M \longrightarrow M' \subset B \subset V = \mathbb{R}^N$ :

$$\frac{dF_t}{dt} = -\Delta F_t + \Gamma(F_t)(\nabla F_t)^2$$

We have proved that the solution exists in short-time and is smooth whenever it exists. We will now establish long-time existence using continuity method: we will show that if the solution exists on  $[\alpha, \omega_n]$  where  $\omega_n$  is an increasing sequence to  $\omega$ , then the solution exists on  $[\alpha, \omega]$ . We then apply short-time existence to gain a small open interval where solution still exists. We then conclude that the solution exists globally on  $[\alpha, +\infty)$  since this interval is connected.

The crucial step to prove that the solution can be extended on  $[\alpha, \omega]$  is to uniformly bound all of its derivatives in time of evolution  $[\alpha, \omega)$ . These estimates will also be useful to justify the convergence of  $F_t$  in  $C^{\infty}(M)$  to a smooth function  $F_{\infty}$  which will eventually be a harmonic map from M to M'.

Recall that we proved in Corollary ??, under the hypothesis of negative curvature, the boundedness of  $||F_t||_{W^{2,2}(M)}$  by a constant C depending only on curvatures of M, M' and the initial total energies. Since  $\frac{dF_t}{dt}$  relates to spatial derivatives of F by the nonlinear heat equation, it is easy to see that  $||F_t||_{W^{2,2}(M\times[\tau,\tau+\delta])}$  is bounded by a constant independent of  $\tau$ . We will denote  $W^{k,p}(M \times [\beta, \gamma])$  by  $W^{k,p}([\beta, \gamma])$ .

**Theorem 1** ( $W^{2,2}$ -boundedness). Suppose Riem(M')  $\leq 0$ . There exists a constant C depending only on  $\delta$ , the metrics and initial total energies such that

$$||F||_{W^{2,2}(\tau,\tau+\delta)} \le C \quad \text{for all } \alpha \le \tau < \omega - \delta.$$

Proof. Since

$$\|F\|_{W^{2,2}([\tau,\tau+\delta])}^2 \leq \int_{\tau}^{\tau+\delta} \|F_t\|_{W^{2,2}(M)}^2 dt + 2 \int_{\tau}^{\tau+\delta} \|\Delta F_t\|_{L^2}^2 dt + 2 \int_{\tau}^{\tau+\delta} \left\|\Gamma(F_t)(\nabla F_t)^2\right\|_{L^2}^2 dt$$

The first term and the second term are bounded by  $C^2\delta$ , the third one, since  $\Gamma(F_t)$  is bounded, by  $C^2\delta$  where C is a constant only depending on the metrics and initial total energies.

The estimates of higher derivatives of F will be established in the same strategy as the bootstrap: first in  $W^{2,p}$  for all p then in  $W^{k,p}$  for all k, p, then in  $C^{\infty}$ .

### 1 Estimate of higher derivatives.

**Lemma 2** ( $W^{2,p}$ -boundedness). Suppose Riem(M')  $\leq 0$ . For all  $p \in (1, +\infty)$ , there exists a constant C > 0 depending only on  $\delta$ , p, the metrics and initial energies such that for all  $\alpha + \delta \leq \tau \leq \omega - \delta$ :

$$||F||_{W^{2,p}([\tau,\tau+\delta])} \le C$$

*Proof.* Applying Gårding Inequality to the parabolic equation  $AF = \Gamma(F)(\nabla F)^2$  where  $A := \frac{\partial}{\partial t} + \Delta$  is the heat operator, one has

$$||F||_{W^{2,p}([\tau,\tau+\delta])} \le C \left( ||\Gamma(F)(\nabla F)^2||_{L^p([\tau-\frac{\delta}{3},\tau+\delta])} + ||F||_{W^{2,2}([\tau-\frac{\delta}{3},\tau+\delta])} \right)$$

The second term of RHS is already bounded by applying Theorem 1 to  $\frac{4\delta}{3}$ . For the first term:

$$\|\Gamma(F)(\nabla F)^2\|_{L^p([\tau-\frac{\delta}{2},\tau+\delta])} \leq C(M')\||\nabla F|^2\|_{L^p([\tau-\frac{\delta}{2},\tau+\delta])} = C(M')\|e(F)\|_{L^p([\tau-\frac{\delta}{2},\tau+\delta])}.$$

Recall that by Theorem ??, the potential density satisfies  $\frac{de}{dt} + \Delta e - Ce \leq 0$  for certain constant C depending only on the metric of M. By Maximum principle (Theorem ??), one has  $e \leq \psi_{\tau}$  where  $\psi_{\tau}$  is the solution of

$$\begin{cases} \frac{d}{dt}\psi_{\tau} + \Delta\psi_{\tau} - C\psi_{\tau} = 0\\ \psi_{\tau}\big|_{\tau - \frac{\delta}{2}} = e\big|_{\tau - \frac{\delta}{2}} \end{cases}$$

We apply Gårding Inequality again for  $\psi_{\tau}$  and obtain

$$\|e(F)\|_{L^{p}([\tau-\frac{\delta}{2},\tau+\delta])} \le \|\psi_{\tau}\|_{L^{p}([\tau-\frac{\delta}{2},\tau+\delta])} \le C\|\psi_{\tau}\|_{L^{1}([\tau-\frac{\delta}{2},\tau+\delta])}. \tag{1}$$

Now apply  $L^1$ -Comparison Theorem ?? to  $\psi_{\tau}$ , one has

$$\|\psi_{\tau}\|_{L^{1}([\tau-\frac{\delta}{2},\tau+\delta])} \leq \int_{0}^{3\delta/2} \|\psi_{\tau}|_{\tau-\frac{\delta}{2}}\|_{L^{1}} e^{Bt} dt \leq \int_{0}^{3\delta/2} e^{Bt} dt. \|e|_{\tau-\frac{\delta}{2}}\|_{L^{1}} \leq C. \tag{2}$$

The lemma follows from (1) and (2).

We can now estimate higher order derivatives.

**Theorem 3** ( $W^{k,p}$ -boundedness). Suppose Riem(M')  $\leq 0$ . For all  $p \in (1, +\infty)$  and  $k < +\infty$ , there exists C depending only on k, p, the metrics and initial energies such that

$$||F||_{W^{k,p}([\tau,\tau+\delta])} \leq C$$

for all  $\alpha + \delta \leq \tau \leq \omega - \delta$ .

*Proof.* Applying Gårding Inequality to the equation  $\frac{dF}{dt} + \Delta F_t = \Gamma(F)(\nabla F)^2$  then Regularity estimate for the quadratic term (Theorem ??), one has for  $\epsilon \ll \delta$ :

$$||F||_{W^{k,p}([\tau,\tau+\delta])} \le C_{\epsilon} \left( ||F||_{W^{2,p}([\tau-\epsilon,\tau+\delta])} + ||\Gamma(F)(\nabla F)^{2}||_{W^{k-2,p}([\tau-\epsilon,\tau+\delta])} \right)$$

$$\le C_{\epsilon} \left( 1 + C \left( 1 + ||F||_{W^{s,q}([\tau-\epsilon,\tau+\delta])} \right)^{q/p} \right)$$

as long as k-1 < s and  $\frac{1}{p} > \frac{k}{s} \cdot \frac{1}{q}$ . Therefore if  $||F||_{W^{s,q}([\tau,\tau+\delta])} \le C(\delta,s,q)$  for all  $\beta \le \tau \le \omega - \delta$  and  $q \in (1,+\infty)$ , we just proved that

$$||F||_{W^{k,p}([\tau,\tau+\delta])} \le C(\epsilon,k,p)$$

$$\text{for all } \begin{cases} \beta + \epsilon \leq \tau \leq \omega - \delta \\ k < s + 1, p \in (1, +\infty) \end{cases} \quad \text{since } \|F\|_{W^{s,q}([\tau - \epsilon, \tau + \delta])} \leq 2C(\delta, s, q).$$

One can then conclude by induction on k, with step  $\frac{1}{2}$ , starting with k=2 and  $\epsilon=\frac{\delta}{2}$  divided by 2 after each induction step.

## 2 Global existence for nonlinear heat equation.

**Theorem 4** (Global existence). *Suppose* Riem(M')  $\leq 0$ . *The solution of nonlinear heat equation* 

$$\frac{dF}{dt} = -\Delta F + \Gamma(F)(\nabla F)^2 \tag{3}$$

with smooth initial condition exists globally for all time  $t > \alpha$ .

*Proof.* Let  $F_n$  be a sequence of solution of (3) on  $[\alpha, \omega_n]$  with  $\omega_n$  increasing to  $\omega$  then they coincide by uniqueness of the solution. As discussed in the beginning of this part, it is sufficient to prove that the solution extends to  $[\alpha, \omega]$ . Let F be the solution on  $[\alpha, \omega)$  such that  $F|_{[\alpha, \omega_n]} = F_n$ , then by Theorem 3, for all  $\tau \in [\alpha, \omega - \delta)$ :

$$\|D_t^u D_x^v F\|_{L^{\infty}(M \times [\tau, \tau + \delta])} \le C_{\text{Sobolev}} \|D_t^u D_x^v F\|_{W^{k, p}(M \times [\tau, \tau + \delta])} \le C_{\text{Sobolev}} C(k, p, \delta)$$

where, if we choose k sufficiently large,  $C_{\text{Sobolev}}$  is the constant of Sobolev imbedding  $W^{k,p}(M \times [0,\delta]) \hookrightarrow C(M \times [0,\delta])$  and  $C(k,p,\delta)$  is the constant provided by Theorem 3.

So all partial derivatives of F are uniformly bounded on  $[\alpha, \omega)$ . This proves that F extends to a solution on  $[\alpha, \omega]$ . In fact  $F|_{\tau} := F|_{M \times \{\tau\}}$  converges in  $C^{\infty}(M)$  as  $\tau \to \omega$ , since

$$\|D^{\alpha}F|_{\tau}-D^{\alpha}F|_{\tau'}\|_{L^{\infty}}\leq \max_{\|\beta\|=\|\alpha\|+1}\|D^{\beta}F\|_{L^{\infty}}|\tau-\tau'|.$$

We have just proved the first part of the following theorem.

**Theorem 5** (Eells-Sampson). 1. Let M, M' be compact Riemannian manifolds with Riem $(M') \le 0$ . Then for every smooth map  $f_0: M \longrightarrow M' \subset B \subset \mathbb{R}^N$ , the nonlinear heat equation

$$\begin{cases} \frac{df_t}{dt} = \tau(f_t), & \text{for all } t \ge 0 \\ f\big|_{t=0} = f_0, \end{cases}$$

admits a globally defined smooth solution  $f_t$ . Moreover, all derivatives  $D^{\alpha} f_t$  remain bounded as  $t \to +\infty$ .

2. For a suitable sequence  $\{t_n\}$  increasing to  $+\infty$  the sequence  $\{f_{t_n}\}$  converges in  $C^{\infty}(M)$  to a function  $f_{\infty}$  with  $\tau(f_{\infty})=0$ . Therefore any map  $f_0: M \longrightarrow M'$  is homotopic to a harmonic map.

*Proof.* For any sequence  $\{t_n\}$ , one can extract from  $\{f_{t_n}\}$ , since their derivatives are uniformly bounded, a subsequence  $\{f_{t_{n_i}}\}$  converging in  $C^k(M,\mathbb{R}^N)$ . By a diagonal argument, one can extract from any sequence  $\{f_{t_n}\}$  a subsequence converging in  $C^{\infty}(M,\mathbb{R}^N)$  to  $f_{\infty}$ . Abusively denote this subsequence by  $\{f_{t_n}\}$ , by Theorem ??

$$\lim_{n\to\infty} K(f_{t_n}) = \lim_{n\to\infty} \int_M |\tau(f_{t_n})|^2 = 0$$

Therefore  $\tau(f_{t_n}) \to 0$  in  $L^2(M)^{\oplus N}$ . But also  $\tau(f_{t_n}) \to \tau(f_{\infty})$  in  $C^{\infty}(M, \mathbb{R}^N)$ , one has  $\tau(f_{\infty}) = 0$ .