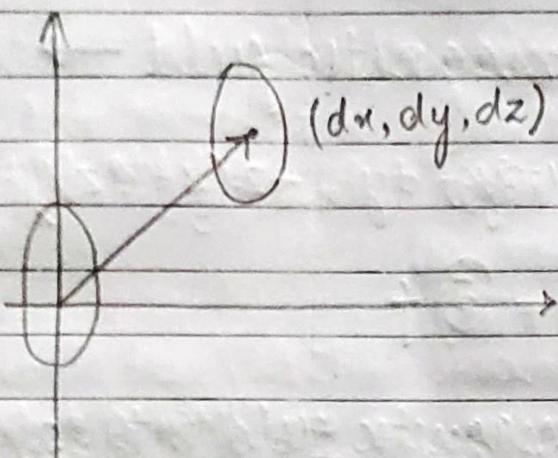


Bout CALCULATION (General Translation)



Demonstration of configuration of spheroid.

This configuration can be obtained via two mechanisms

1. Translation of spheroid.
2. Translation of applied magnetic field.

We use 2nd mechanism throughout to achieve the above specified configurations.

Using the same name of the parameters as was used in section D(d).

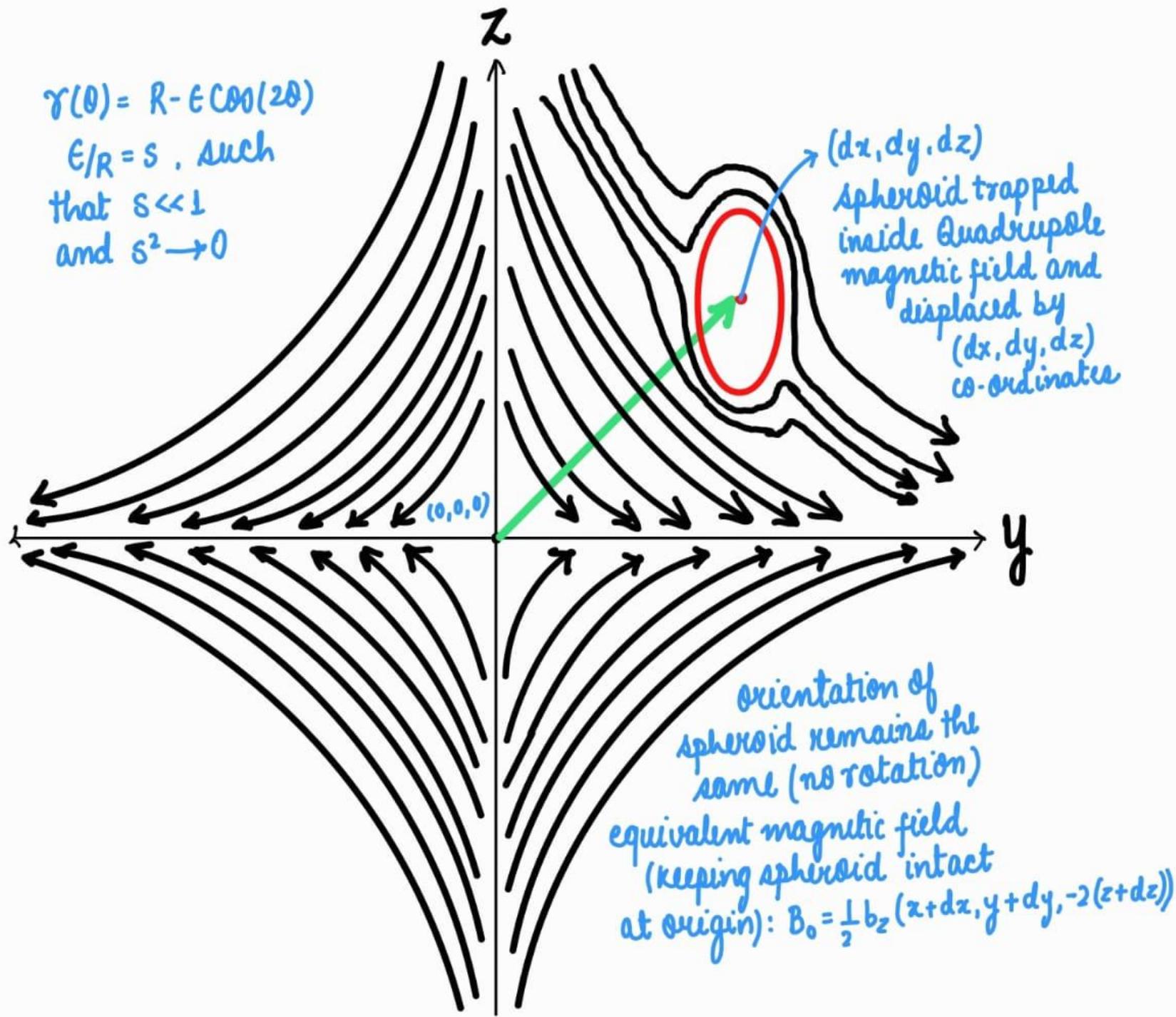
$$\boxed{r_{\text{spheroid}}} \quad r(\theta) = R - \epsilon \cos(\phi) \quad \left. \begin{array}{l} \\ = R - R_s \cos(\phi) \end{array} \right\} \rightarrow (1)_f$$

Position vector to surface of spheroid :-

$$v(\theta, \phi) = r(\theta) \hat{r} = (R - R_s \cos(\phi)) \hat{r} \quad (2)_f$$

Normal vector to surface of spheroid :-

$$\hat{n} = \frac{R - R_s \cos(\phi) \hat{r} - 4R \cos \theta \sin(\phi) \hat{s} \wedge \hat{\theta}}{\sqrt{(R - R_s \cos(\phi))^2 + 4R^2 s^2 \sin^2(\phi)}} \quad (3)_f$$



Note → Detailed explanation of eqn 1 if is given in Section A
 explanation of eqn. 2 & 3 given in section B(b)

Defining the applied magnetic field →

$$B_0(x, y, z) = \frac{1}{2} b_z (x + dx, y + dy, -2(z + dz)) \quad \text{--- (4)f}$$

$$\nabla \cdot B_0(x, y, z) = 0 \quad \text{--- (5)f}$$

converting the applied magnetic field into spherical co-ordinates system. →

→ B₀spherical

$$M_s \rightarrow c^* B_0(x, y, z) = B_0(r, \theta, \phi) \quad \text{--- (6)f}$$

$$B_0(r, \theta, \phi) = \frac{1}{2} b_z (-2dz \cos(\theta) - 2r \cos^2(\theta) + \sin\theta (dx \cos\phi + r \sin\theta \cos\phi + dy \sin\phi))$$

$$B_0(r, \theta, \phi) = b_z dz \sin\theta + \frac{1}{2} b_z \cos\theta (dx \cos\phi + r \sin\theta \cos\phi + dy \sin\phi),$$

$$B_0(r, \theta, \phi) = \frac{1}{2} b_z (dy \cos\phi - dx \sin\phi) \quad \text{--- (7)f}$$

Applying the boundary condition →

→ This part involves applying the boundary condition given in the equation eqn. 6c. Let us first focus on solving the LHS part of eqn. 6c.

LHS \rightarrow

$$B_0(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(\theta)} \longrightarrow \text{LHS}$$

$$\boxed{\text{LHSBCVsph}[\theta, \phi]} = B_0(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(\theta)} \xrightarrow{f_{\text{LHS}}(s)} \textcircled{8} f$$

We now, simplify the eqn 8f by taylor expanding w.r.t s about s=0 upto first order. This is given as,

$$\boxed{\text{AppLHSBCVsph}[\theta, \phi]} = f_{\text{LHS}}(0) + f'_{\text{LHS}}(0)s + O(s^2) \quad \textcircled{9} f$$

This eq. 9f is simplified and as we know $(B_0(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(\theta)})$ is a function of θ, ϕ and can be expressed in terms of spherical harmonic expansion.

We first verify whether eq 9f can be expressed in spherical harmonics or not.

$$\int_0^{2\pi} \int_0^\pi \text{AppLHSBCVsph}(\theta, \phi) \sin(\theta) d\theta d\phi < \infty \quad \textcircled{10} f$$

Since the integral is finite, it can be expressed in terms of spherical harmonics expansion.

$$\boxed{b_{l,k}} = \iint_0^{2\pi} \int_0^\pi Y_l^{*k}(\theta, \phi) \sin \theta (\text{AppLHSBCVsph}(\theta, \phi)) d\theta d\phi$$

$$= \iint_0^{2\pi} \int_0^\pi Y_l^{*k}(\theta, \phi) \sin \theta \underbrace{\left(B_0(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(\theta)} \right)}_{\text{Simplified version of this.}} d\theta d\phi \quad \textcircled{11} f$$

$$\left[Y_n^{-m}(0, \phi) = (-1)^m Y_n^m*(0, \phi) \right]$$

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$b_{e,k}$ are obtained from eqn 11f. All the non zero $b_{e,k}$'s are stated in the eqn 12f given below.

$$b_{0,0} = -\frac{16}{15} b_z \sqrt{\pi} R s$$

$$b_{1,0} = -\frac{2}{5} b_z dz \sqrt{\frac{\pi}{3}} (5+8s), \quad b_{1,1} = \frac{1}{5} (dx - idy) \sqrt{\frac{\pi}{6}} (4s-5)$$

$$(12)f \quad b_{2,0} = -\frac{2}{21} b_z \sqrt{\frac{\pi}{5}} R (21+11s)$$

$$b_{3,0} = \frac{16}{5} b_z dz \sqrt{\frac{\pi}{7}} s, \quad b_{3,1} = \frac{8}{5} b_z (dx - idy) \sqrt{\frac{\pi}{21}} s$$

$$b_{4,0} = \frac{48}{35} b_z \sqrt{\pi} R s$$

Note that, $b_{l,-k} = (-1)^k b_{l,k}$ — (13)f
therefore

Eqn. 13f is derived as below

$$b_{l,k} = \iint \sin \theta f(\theta, \phi) Y_l^{*k}(\theta, \phi) d\theta d\phi$$

$$b_{l,-k} = \iint \sin \theta f(\theta, \phi) Y_l^{*-k}(\theta, \phi) d\theta d\phi$$

as we know, $Y_l^{*-k}(\theta, \phi) = (-1)^k Y_l^k(\theta, \phi)$ — (14)f

$$b_{l,-k} = (-1)^k b_{l,k}^* — (15)f$$

We can even verify the all nonzero coefficients are determined or not by eqn. 16f given below

$$\sum_{l=0}^L \sum_{k=-l}^l b_{l,k} Y_l^k(\theta, \phi) - \left(B_0(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(\theta)} \right) = 0$$

— (16)f

Now, we are done evaluating the LHS part. Let us move on to calculate the RHS part and express it as sum of spherical harmonica terms.

$$\Phi_{\text{out}}(r, \theta, \phi) = \sum_n^{N_{\max}} r^{-(n+1)} \sum_{m=-n}^n g_{n,m} Y_n^m(\theta, \phi) \quad \text{--- (19f)}$$

$\boxed{\text{PhiOut}(r, \theta, \phi, N_{\max})}$

Taking the gradient of the eqn 19f we obtain the eqn 5c.

Note → The RHS part of the eqn. always remains same as we are only shifting B_0 and keeping other things const.

RHS consists of $\nabla \Phi \cdot \hat{n}$ which always remains same.

However, for completion of calculation, I am stating all the calculations explicitly.

$$\boxed{\text{gradph}(r, \theta, \phi, N_{\max})} = \nabla \Phi_{\text{out}}(r, \theta, \phi) \quad \text{--- (20f/5c)}$$

RHS part of eqn 6c can be written as,

$$\nabla \Phi_{\text{out}}(r, \theta, \phi) \cdot \hat{n}(\theta, \phi) \Big|_{r=r(0)} = \quad \text{--- (21f/19d)} \\ + \quad \text{--- (20d)}$$

Eqn. 19d & 20d are illustrated in three steps in mathematica file

$\boxed{\text{rhsBCVSPH}}$

\downarrow
 $\boxed{\text{Simplified RHSBC}}$

\downarrow
 $\boxed{\text{more simplified RHSBC}}$

Since, we know the coefficients of LHS, we can predict which components of RHS will be zero.

From eqn. 12 f we can say that only non zero coefficients are $m=0$ & $m=\pm 1$ for other m 's we do not have any non zero coefficient. Hence RHS should also contain only $m=0$ & $m=\pm 1$ terms.

We expand RHS (more simplified RHSBC) upto $N_{max} = 10$, although we knew that $N_{max} = 8$ is enough to encounter the expression & coefficients to be correct.

Then we Taylor expand this above evaluated expression w.r.t s about $s=0$ up to first order.

$$\boxed{\text{AppRHSBC}(0, \varphi)} = \underbrace{(\nabla \Phi_{out}(0, 0, \varphi) \cdot \hat{n}(0, \varphi)|_{r=r(0)})}_{\text{its Taylor series expanded form}} - ②f$$

Also note that while expanding the eqn. 22 f we removed all the terms other than $m=0$ and $m=\pm 1$ since we already knew that their contribution is going to be zero to equate with LHS terms.

We evaluate the coefficients of RHS which are written as

$\begin{matrix} 2 \\ 1 \\ 1 \end{matrix}$

$$\boxed{q_{l,k}} = \iint_0^{\pi} \sin \theta Y_l^k(0, \varphi) (\text{AppRHSBC}) d\theta d\varphi - ②f$$

$$\text{RHS} = \sum q_{l,k} Y_l^k(0, \varphi) - ②f$$

The coefficients are given in Mathematica file.

$\boxed{\text{tablegd}}$ = Table of the coefficients of RHS. - (25)f

Now we compare the corresponding coefficients of LHS & RHS.

$$\sum_{l,k} b_{l,k} Y_l^k(0, \varphi) = \sum_{l',k'} 99 l' k' Y_l^{k'}(0, \varphi)$$

$$[b_{l,k} = 99 l' k'] \quad (\text{for } l=l' \& k=k')$$

$\boxed{b\text{coeff LHS}}$

$\boxed{g\text{coeff RHS}}$

(26)f → It is divided
in several parts.

first of all we compare the coefficients with $k=0$.

As we have already encountered a problem in Section for solving this.

$$b_{0,0} = h_0(g_{0,0}, g_{2,0})$$

$$b_{1,0} = h_1(g_{1,0}, g_{3,0})$$

$$b_{2,0} = h_2(g_{0,0}, g_{2,0}, g_{4,0})$$

$$b_{3,0} = h_3(g_{1,0}, g_{3,0}, g_{5,0})$$

$$b_{4,0} = h_4(g_{2,0}, g_{4,0}, g_{6,0})$$

$$b_{5,0} = h_5(g_{3,0}, g_{5,0}, g_{7,0})$$

$$b_{6,0} = h_6(g_{2,0}, g_{6,0}, g_{8,0})$$

- (27)f

(Divided
in parts)

Ans

we can see that there are 7 eqns and 9 variables again. To rectify this problem. We ignore $g_{7,0}$ & $g_{8,0}$ because those terms contained S/R^9 & S/R^{10} terms as their coefficients. We consider these terms to be $S/R^9 \rightarrow 0$, $S/R^{10} \rightarrow 0$ hence, now we have

7 eqns and 7 variables which gives the solution for each g 's.

Similarly, $K=1$,

$$b_{1,1} = P_1(g_{1,1}, g_{3,1})$$

$$b_{2,1} = P_2(g_{2,1}, g_{4,1})$$

$$b_{3,1} = P_3(g_{1,1}, g_{3,1}, g_{5,1})$$

$$b_{4,1} = P_4(g_{2,1}, g_{4,1}, g_{6,1})$$

$$b_{5,1} = P_5(g_{3,1}, g_{5,1}, g_{7,1})$$

$$b_{6,1} = P_6(g_{4,1}, g_{6,1}, g_{8,1})$$

} — (28) f

(Divided in parts)

again we ignore $g_{7,1}$ & $g_{8,1}$ because of their coefficients containing S/R^9 & S/R^{10} term which tends to zero.

glist → list of non simplified g's — (29) f

g0 → list of g's with $k=0$ and non zero value. — (30) f

g1 → list of nonzero g's with $k=1$. — (31) f

gg → list of all non zero g's after simplification — (32) f

The non zero simplified coefficients of Φ_{out} term are →

$$g_{1,-1} = -\frac{1}{10} b_2 (dx + idy) \sqrt{\frac{\pi}{6}} R^3 (5+3s)$$

$$g_{1,0} = \frac{1}{5} b_2 dz \sqrt{\frac{\pi}{3}} R^3 (5+9s)$$

$$g_{1,1} = \frac{1}{10} b_2 (dx - idy) \sqrt{\frac{\pi}{6}} R^3 (5+3s)$$

$$g_{2,0} = \frac{2}{21} b_2 \sqrt{\frac{\pi}{5}} R^5 (7+5s)$$

$$g_{3,-1} = \frac{2}{5} b_2 (dx + idy) \sqrt{\frac{3\pi}{7}} R^5 s$$

$$g_{3,0} = -\frac{12}{5} b_2 dz \sqrt{\frac{\pi}{7}} R^5 s$$

$$g_{3,1} = -\frac{2}{5} b_2 (dx - idy) \sqrt{\frac{3\pi}{7}} R^5 s$$

$$g_{4,0} = -\frac{16}{21} b_2 \sqrt{\pi} R^7 s$$

(33)f

Calculating B_{out} from the evaluated coefficients →

$$\Phi_{out}(r, \theta, \phi) = \sum_{n=0}^N r^{-(n+1)} \sum_{m=-n}^n g_{n,m} Y_n^m(\theta, \phi) \quad (34)f$$

$$\nabla \Phi_{out}(r, \theta, \phi) \quad (35)f$$

$$B_{out} = B_0(r, \theta, \phi) - \nabla \Phi_{out}(r, \theta, \phi) \quad (36)f$$

final simplified equation of $B_{out}(0, \phi)$ at $r=r(0)$ is given in eq. 38f in the mathematica file.

This equation is further used for calculating force acting of spheroid using maxwell stress tensor.