Material Motion $n^{\circ}1$:

$$\boldsymbol{\psi} = \psi_{\beta}(\boldsymbol{\chi}, t) = \begin{pmatrix} \chi_1 + \beta \sin(2\pi\chi_1) \sin(\pi\chi_2/3) \sin(\pi t/T) \\ \chi_2 + 5\beta \sin(2\pi\chi_1) \sin(\pi\chi_2/3) \sin(2\pi t/T) \\ \chi_3 \end{pmatrix}, \tag{1}$$

with T=2 and β an amplitude parameter defined for each simulation. Material Motion $n^{\circ}2$:

$$\psi = \psi_{\beta}(\chi, t) = \begin{pmatrix} \chi_1 + \beta \sin^2(2\pi\chi_1) \sin^2(\pi\chi_2/3) \sin(\pi t/T) \\ \chi_2 + 5\beta \sin^2(2\pi\chi_1) \sin^2(\pi\chi_2/3) \sin(2\pi t/T) \\ \chi_3 \end{pmatrix},$$
(2)

$$\mathbf{w} = \begin{pmatrix} \frac{\beta\pi}{T} \sin^2(2\pi\chi_1) \sin^2(2\pi\chi_1) \cos(\pi t/T) \\ \frac{10.0\beta\pi}{T} \sin^2(2\pi\chi_1) \sin^2(\pi Y/3.0) \cos(2\pi t/T) \\ 0.0 \end{pmatrix},$$
(3)

$$\mathbf{F}_{\psi} = \begin{pmatrix} 1 + 4\beta\pi \sin(2\pi X)\cos(2\pi X)\sin(\pi Y/3.0)\sin(\pi Y/3.0)\sin(\pi t/T) & 2\beta\pi/3.0\sin(2\pi X)\sin(2\pi X)\sin(2\pi$$

0.1 Conservation Laws

0.1.1 Neo Hookean model

$$p = \kappa \left(det(trueF) - 1 \right), \text{ with } trueF = FF_{\psi}^{-1}$$
 (5)

0.1.2 Set of conservation laws

$$\frac{\partial \mathbf{F}}{\partial t} = \mathrm{DIV}_{\chi} \left(\hat{\mathbf{v}} \otimes \mathbf{I} \right) \tag{6a}$$

$$\frac{\partial J}{\partial t} = \boldsymbol{H}_{\psi} : GRAD_{\chi}(\boldsymbol{w})$$
 (6b)

$$\frac{\partial \mathbf{F}_{\psi}}{\partial \mathbf{F}_{\psi}} = \mathrm{DIV}_{\chi} \left(\mathbf{w} \otimes \mathbf{I} \right) \tag{6c}$$

$$\frac{\partial \tilde{\boldsymbol{p}}}{\partial t} = \text{DIV}_{\boldsymbol{\chi}} \left(\boldsymbol{P} \boldsymbol{H}_{\psi} \right) + \text{GRAD}_{\boldsymbol{\chi}} \left(\boldsymbol{p}_{R} \right) \left(\boldsymbol{H}_{\psi}^{T} \boldsymbol{w} \right) + \boldsymbol{p}_{R} \left[\left| \boldsymbol{H}_{\psi} : \text{GRAD}_{\boldsymbol{\chi}} \left(\boldsymbol{w} \right) \right| \right]$$
(6d)

0.2 Finite Volume Method spatial discretization

0.2.1 Discretised set of conservation laws

$$\Omega_{\chi}^{a} \frac{\partial \mathbf{F}^{a}}{\partial t} = \sum_{b \in \Lambda_{a}} \hat{\mathbf{v}}^{Ave} \otimes \mathcal{C}_{\chi}^{ab} + \sum_{\gamma \in \Lambda_{a}^{B}} \hat{\mathbf{v}}_{a}^{\gamma} \otimes \mathcal{C}_{\chi}^{\gamma}$$
(7a)

$$\Omega_{\chi}^{a} \frac{\partial J^{a}}{\partial t} = \boldsymbol{H}_{\psi} : \left(\sum_{b \in \Lambda_{a}} \boldsymbol{w}^{Ave} \otimes \boldsymbol{C}_{\chi}^{ab} \right)$$
 (7b)

$$\Omega_{\chi}^{a} \frac{\partial \mathbf{F}_{\psi}^{u}}{\partial t} = \sum_{b \in \Lambda_{a}} \mathbf{w}^{Ave} \otimes \mathbf{C}_{\chi}^{ab} + \sum_{\gamma \in \Lambda_{a}^{B}} \mathbf{w}_{a}^{\gamma} \otimes \mathbf{C}_{\chi}^{\gamma}$$
(7c)

$$\Omega_{\chi}^{a} \frac{\partial \tilde{p}^{a}}{\partial t} = \sum_{b \in \Lambda_{a}} (PH_{\psi})^{Ave} \mathcal{C}_{\chi}^{ab} + \sum_{\gamma \in \Lambda_{a}^{B}} t_{\chi,a}^{\gamma} ||\mathcal{C}_{\chi}^{\gamma}|| + \sum_{b \in \Lambda_{a}} \mathcal{D}_{\chi,ab}
+ \left(\sum_{b \in \Lambda_{a}} p_{R}^{Ave} \otimes \mathcal{C}_{\chi}^{ab} \right) H_{\psi}^{T} w + \left(\sum_{\gamma \in \Lambda_{a}} p_{R,\gamma} \otimes \mathcal{C}_{\chi}^{\gamma} \right) H_{\psi}^{T} w
+ p_{R} \left[H_{\psi} : \left(\sum_{b \in \Lambda_{a}} w^{Ave} \otimes \mathcal{C}_{\chi}^{ab} \right) \right] + p_{R} \left[H_{\psi} : \left(\sum_{\gamma \in \Lambda_{a}^{B}} w_{a}^{\gamma} \otimes \mathcal{C}_{\chi}^{\gamma} \right) \right]$$
(7d)

0.2.2 Algorithms

Algorithm 1: Computation of the right hand side of the linear momentum conservation law for the first Runge-Kutta stage.

Result: rhsLm, rhsF.

1
$$oldsymbol{p}_R^n = J_\psi^{-1} oldsymbol{p}_{oldsymbol{\chi}}^n$$

$$\begin{array}{l} \mathbf{2} \ \boldsymbol{rhsLm}_a^* = \left(\sum\limits_{b \in \Lambda_a} \boldsymbol{tC}_b^n || \boldsymbol{\mathcal{C}}_{\boldsymbol{\chi}}^{ab} || \right) + \left(\sum\limits_{b \in \Lambda_a} \boldsymbol{tC2}_b^n || \boldsymbol{\mathcal{C}}_{\boldsymbol{\chi}}^{ab} || \right) \boldsymbol{H}_{\psi}^T \boldsymbol{w} + \\ \boldsymbol{p}_R^n \left(\boldsymbol{H}_{\psi} : \left[\sum\limits_{b \in \Lambda_a} \boldsymbol{tC3}_b^n || \boldsymbol{\mathcal{C}}_{\boldsymbol{\chi}}^{ab} || \right] \right) \end{array}$$

- 3 Boundary value of the first RHS term. $tC_{loc}^{\gamma} = LocalAverage(\hat{P}N_{\chi}^{\gamma})$ or apply the prescribed traction boundary condition. The nodal value is updated as $rhsLm_{\gamma}^{*} = rhsLm_{\gamma}^{n} + tC_{loc}^{\gamma}||\mathcal{C}_{\chi}^{\gamma}||$.
- 4 Boundary value of the second RHS term. $tC2_{loc}^{\gamma} = LocalAverage(\boldsymbol{p}_{R}\boldsymbol{\mathcal{N}}_{\chi}^{\gamma})$. The nodal value is updated as $\boldsymbol{rhsLm}_{\gamma}^{*} = \boldsymbol{rhsLm}_{\gamma}^{n} + \left(tC2_{loc}^{\gamma}||\boldsymbol{\mathcal{C}}_{\chi}^{\gamma}||\right)\boldsymbol{H}_{\psi,\gamma}^{T}\boldsymbol{w}_{\gamma}.$
- 5 Boundary value of the third RHS term. $tC3_{loc}^{\gamma} = LocalAverage(w\mathcal{N}_{\chi}^{\gamma})$. The nodal value is updated as $rhsLm_{\gamma}^{*} = rhsLm_{\gamma}^{n} + p_{R,\gamma}\left(H_{\psi,\gamma}: (tC3_{loc}^{\gamma}||\mathcal{C}_{\chi}^{\gamma}||)\right)$.
- 6 Volume integration: $rhsLm_a^* = rhsLm_a^*/\Omega_a$.
- 7 Update of displacements $\boldsymbol{x}^* = \boldsymbol{x}^n + \Delta t \hat{\boldsymbol{v}}^n$ and linear momentum $\boldsymbol{p}_{\chi}^* = \boldsymbol{p}_{\chi}^n + J_{\psi}J^{-1}\Delta t \ \boldsymbol{rhsLm}^*$.
- 8 Application of boundary conditions for the linear momentum.
- 9 Update of variables:

$$egin{aligned} m{true}m{F}^* &= m{F}^*m{F}_\psi^{-1,*} \ m{P}^* &= m{P}^*(m{true}m{F}^*) \ m{p}_R^* &= J_\psi^{-1}m{p}^* \ m{v}^* &= m{p}_R^*/
ho \ \hat{m{v}}^* &= m{v}^* + m{true}m{F}^*m{w}^n \ \hat{m{P}}^* &= m{P}^*m{H}_\psi \end{aligned}$$

- 10 Computation of coefficient Λ^n_H and update of U^n_p and U^n_s .
- 11 Update of the averaged values for computing the right hand sides

$$egin{aligned} oldsymbol{tC^*} &= \hat{oldsymbol{P}}^{Ave,*} \mathcal{N} + 0.5 S^* [oldsymbol{v}^*] \ oldsymbol{tC2^*} &= oldsymbol{p}_R^{Ave,*} \otimes \mathcal{N} \ oldsymbol{tC3^*} &= oldsymbol{w}^{Ave,n} \otimes \mathcal{N} \end{aligned}$$