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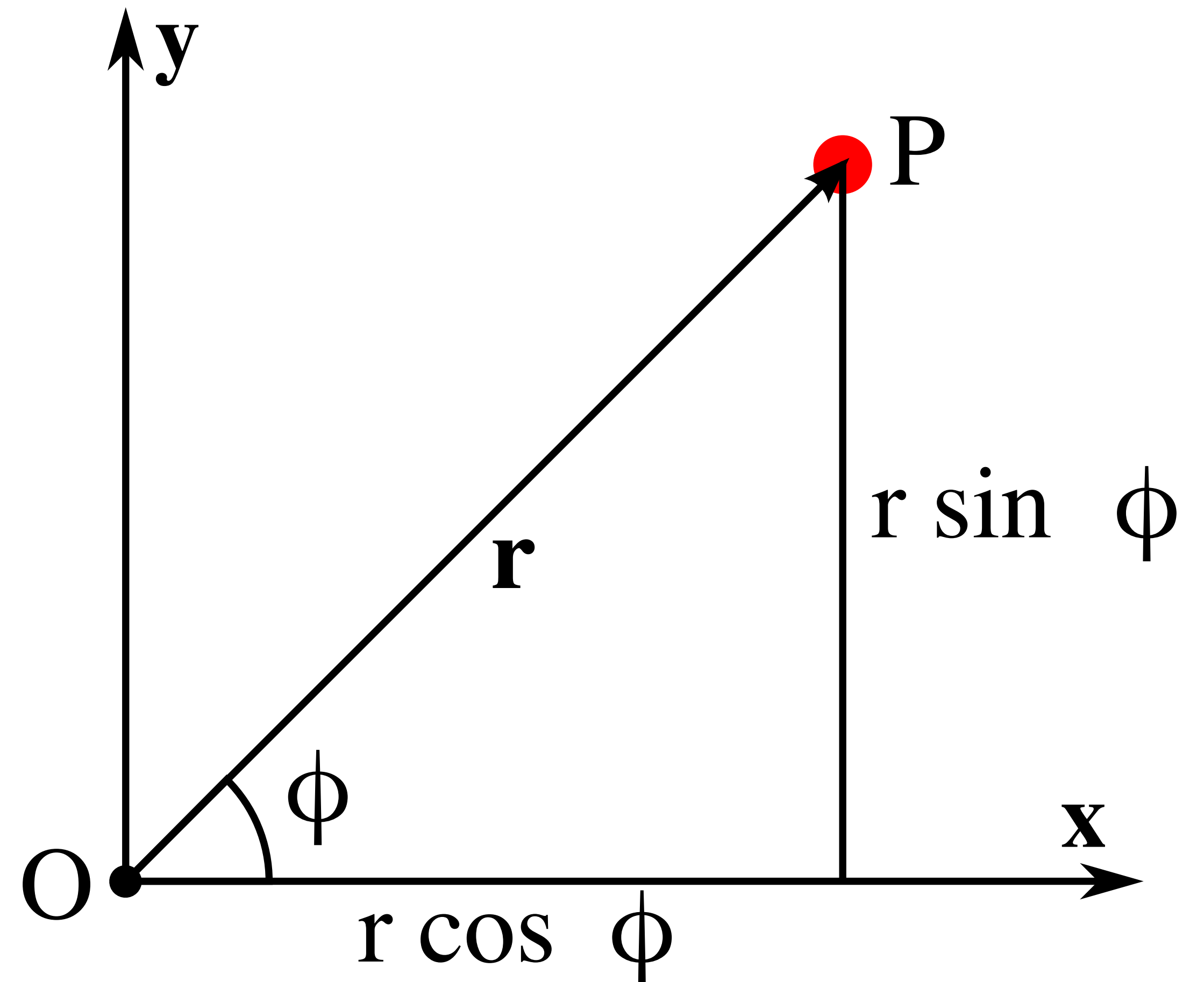
- Particle on a ring
 - Polar Coordinates
 - Laplacian and Hamiltonian operator
 - General solutions
 - Angular momentum
- Particle on a sphere
 - Spherical Coordinates
 - Laplacian and Hamiltonian operator
 - Spherical harmonics
 - Angular momentum
- This lecture is designed to help you achieve the following learning objectives
 - Obtain and interpret solutions of the Schrodinger equation for tractable systems including the particle in a box, harmonic oscillator, rigid rotor, and hydrogen atom
 - Describe how analytical models help interpret various experimental spectra

Particle on a Ring

- 1D Particle in a box, confined to a specific range of x .
 - Hamiltonian is $\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$.
- 2D Particle in a box, confined to a specific range of x and y .
 - Hamiltonian is $\hat{H} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right)$
 - Alternatively, $\hat{H} = -\frac{\hbar^2}{2m} \nabla^2$, where ∇^2 is the Laplacian operator.
- Particle on a ring, confined to a specific value of r . What is the kinetic energy?

Polar Coordinates

- $x = r \cos \phi$
- $y = r \sin \phi$
- $r = \sqrt{x^2 + y^2}$
- $\phi = \arctan \left(\frac{y}{x} \right)$



The Chain Rule

- $u(x, y)$ is a function that depends on x and y
- $\delta u(x, y) = \frac{\partial u}{\partial x} \delta x + \frac{\partial u}{\partial y} \delta y$ is differential for u
- Now let's say that $x = x(s, t)$ and $y = y(s, t)$. The chain rule says that,
$$\frac{\partial u}{\partial s} = \frac{\partial u}{\partial x} \frac{\partial x}{\partial s} + \frac{\partial u}{\partial y} \frac{\partial y}{\partial s}.$$
- What about $\frac{\partial u}{\partial t}$?
$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial u}{\partial y} \frac{\partial y}{\partial t}$$

Hamiltonian of POR

- Using the chain rule, one can obtain the Laplacian in polar coordinates as

$$\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2}$$

- $$\hat{\mathbf{H}} = \frac{-\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) = \frac{-\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} \right)$$

- For a particle on a ring, r is constant, so
$$\hat{\mathbf{H}} = \frac{-\hbar^2}{2m} \left(\frac{1}{r^2} \frac{d^2}{d\phi^2} \right)$$

Schrodinger Equation

- $\frac{-\hbar^2}{2m} \left(\frac{1}{r^2} \frac{d^2}{d\phi^2} \right) \Psi = E\Psi$ is the Schrodinger equation for the system
- $\frac{d^2\Phi}{d\phi^2} = -\frac{2IE}{\hbar^2}\Phi$ (using $I = mr^2$ and swapping notation for the wavefunction)
- The general solution is $\Phi = Ae^{im_l\phi} + Be^{-im_l\phi}$, where $m_l = \sqrt{\frac{2IE}{\hbar^2}}$.

Schrodinger Equation Solutions

- $\Phi = Ae^{im_l\phi} + Be^{-im_l\phi}$, where $m_l = \sqrt{\frac{2IE}{\hbar^2}}$.
- The boundary conditions are $\Phi(\phi + 2\pi) = \Phi(\phi)$, or $Ae^{im_l\phi}e^{2\pi im_l} + Be^{-im_l\phi}e^{2\pi im_l} = Ae^{im_l\phi} + Be^{-im_l\phi}$
- Thus, $e^{2\pi im_l} = \cos(2\pi m_l) + i \sin(2\pi m_l) = 1$ must be true.
- When is this true?
 - If m_l is an integer, $m_l = 0, \pm 1, \pm 2, \dots$
- Given $m_l = \sqrt{\frac{2IE}{\hbar^2}}$, what are the energy levels of a particle on a ring?
 - $E_{m_l} = \frac{m_l^2 \hbar^2}{2I}$
- Are there degenerate states (states with the same energy)?
 - Yes

Normalization

- If $B = 0$, then $\Phi = Ae^{im_l\phi}$
- What is A ?

$$\bullet \int_0^{2\pi} \Phi^* \Phi d\phi = |A|^2 \int_0^{2\pi} e^{-im_l\phi} e^{im_l\phi} d\phi = |A|^2 \int_0^{2\pi} d\phi = |A|^2 2\pi$$

$$\bullet |A| = \sqrt{\frac{1}{2\pi}}$$

Angular Momentum

- In classical mechanics, angular momentum is the cross product of the position \mathbf{r} and

momentum vectors \mathbf{p} ,
$$l = \mathbf{r} \times \mathbf{p} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ x & y & z \\ p_x & p_y & p_z \end{vmatrix}$$

- The z component is $l_z = xp_y - yp_x$
- Using the momentum operator in the position representation,

$$\hat{\mathbf{l}}_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

- If $\Phi = Ae^{im_l\phi}$, what is the angular momentum?

- Applying the operator to the wavefunction gives $\hat{\mathbf{l}}_z\Phi = \frac{\hbar}{i}im_l\Phi = m_l\hbar\Phi$

- Thus the momentum is $m_l\hbar$

Postulate 3 Self-Test

- $\Phi = \sqrt{\frac{1}{3}}\sqrt{\frac{1}{2\pi}} - \sqrt{\frac{2}{3}}\sqrt{\frac{1}{2\pi}}e^{-2i\phi}$

Probability

Energy

Angular
Momentum

- What are the angular momentum quantum numbers?

- What are the normalized wavefunctions ϕ_n ?

- What the the coefficients c_n ?

Expectation

Particle on a Sphere

Spherical Polar Coordinates

- The mapping from spherical polar to Cartesian coordinates is,

- $x = r \sin \theta \cos \phi$

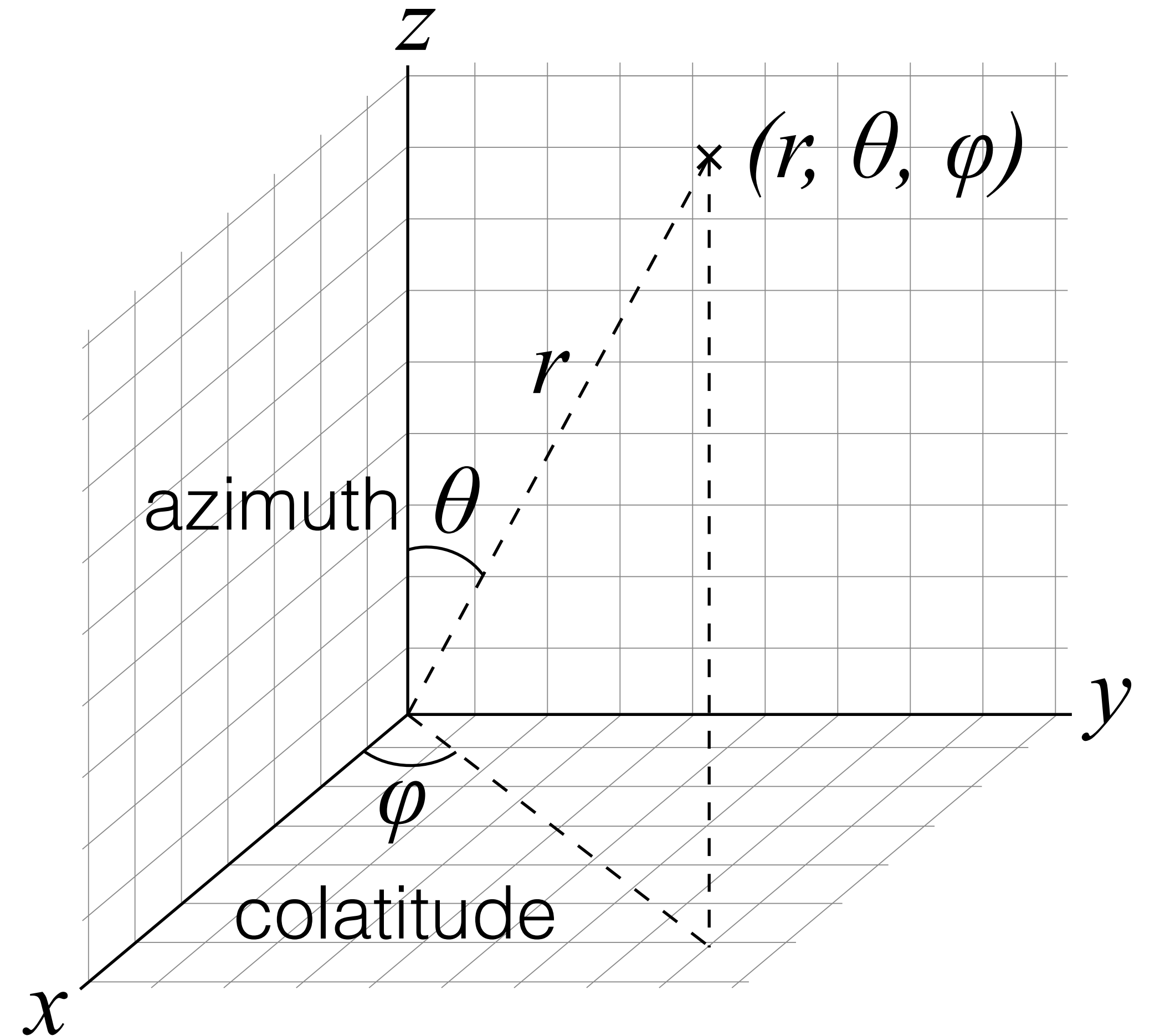
- $y = r \sin \theta \sin \phi$

- $z = r \cos \theta$

- The Laplacian is,

- $\hat{\nabla}^2 = \frac{1}{r} \frac{\partial^2}{\partial r^2} r + \frac{1}{r^2} \hat{\Lambda}^2$

- $\hat{\Lambda}^2 = \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta}$ is called the Legendrian



Schrodinger Equation

- $\hat{H} = \frac{-\hbar^2}{2m} \hat{\nabla}^2 = \frac{-\hbar^2}{2mr^2} \hat{\Lambda}^2$ is the Hamiltonian if r is constant
- $\hat{\Lambda}^2 \Psi = -\frac{2IE}{\hbar^2} \Psi$ is the Schrodinger equation
- $\hat{\Lambda}^2 = \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta}$

Separation of Variables

- Applying the Legendrian to $\Psi(\theta, \phi) = \Theta(\theta)\Phi(\phi)$ gives

- $$\hat{\Lambda}^2 \Theta \Phi = \Theta \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \Phi + \Phi \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} \Theta = -\frac{2IE}{\hbar^2} \Theta \Phi$$

- $$\frac{1}{\Phi} \frac{\partial^2}{\partial \phi^2} \Phi + \frac{1}{\Theta} \sin \theta \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} \Theta = -\frac{2IE}{\hbar^2} \sin^2 \theta \text{ by rearrangement}$$

- For Φ , this has the form $\frac{d^2}{d\phi^2} \Phi = k\Phi$, the same as the particle on the ring

- For Θ , the solutions are known as *Spherical Harmonics*

Spherical Harmonics

- The spherical harmonics, $Y_{lm_l}(\theta, \phi) = \Theta_{lm_l}(\theta)\Phi_{m_l}(\phi)$, are separable solutions to the Schrodinger equation
- For these equations, applying the Legendrian operator yields
 - $\hat{\Lambda}^2 Y_{lm_l} = -l(l+1)Y_{lm_l}$, where Y_{lm_l} are spherical harmonics
 - $l = 0, 1, 2, \dots$
 - $m_l = l, l-1, l-2, \dots, -l$

- By rearranging $\frac{2IE_{lm_l}}{\hbar^2} = -l(l+1)$, we see that the allowed energy levels are,
$$E_{lm_l} = \frac{\hbar^2 l(l+1)}{2I}.$$
 Notably, this does not depend on m_l .

Visualizing Spherical Harmonics

<http://demonstrations.wolfram.com/SphericalHarmonics/>

What is the relationship between l , m_l , and the number of angular nodes?

Angular Momentum: Components

- Recall that the z component of angular momentum is

$$\hat{\mathbf{I}}_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} + y \frac{\partial}{\partial x} \right) = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

- $\hat{\mathbf{I}}_z Y_{lm_l} = m_l \hbar Y_{lm_l}$, the same as a particle on a ring
- Due to the uncertainty principle, x and y components cannot be simultaneously determined

Angular Momentum: Magnitude

- Classical energy

- Linear momentum: $E = \frac{1}{2}mv^2$. $p = mv$. $E = \frac{p^2}{2m}$.

- Angular momentum: $E = \frac{1}{2}I\omega^2$. $L = I\omega$. $E = \frac{L^2}{2I}$.

- $\hat{\mathbf{L}}^2 = 2I\hat{\mathbf{H}} = 2I\frac{-\hbar^2}{2I}\hat{\mathbf{A}}^2 = -\hbar^2\hat{\mathbf{A}}^2$ is the square magnitude of the linear momentum

- Since $\hat{\mathbf{L}}^2 Y_{lm_l} = l(l+1)\hbar^2$, the angular momentum is quantized as $\sqrt{l(l+1)}\hbar$, where $l = 0, 1, 2, \dots$ is the angular momentum quantum number

Diatomic Rigid Rotor

- The mathematics of this model, two particles a specific distance apart, can be posed the same way as a Particle on a Sphere

- For two free particles, the Hamiltonian is $\hat{\mathbf{H}} = -\frac{\hbar^2}{2m_1} \hat{\nabla}_1^2 - \frac{\hbar^2}{2m_2} \hat{\nabla}_2^2$.

- To consider a rigid rotor, we will simplify the expression using,

$$\frac{1}{m_1} \nabla_1^2 + \frac{1}{m_2} \nabla_2^2 = \frac{1}{m} \nabla_{cm}^2 + \frac{1}{\mu} \nabla^2, \text{ where}$$

- $m = m_1 + m_2$ is the total mass of the system
- $\frac{1}{\mu} = \frac{1}{m_1} + \frac{1}{m_2}$ is the reduced mass (also seen for harmonic oscillator)
- ∇^2 pertains to relative coordinates of the two atoms.

Solutions to the RR

- $$-\frac{\hbar^2}{2m} \hat{\nabla}_{cm}^2 \Psi - \frac{\hbar^2}{2\mu} \hat{\nabla}^2 \Psi = E_{total} \Psi$$

is the Schrodinger equation

- Separation of variables $\Psi = \Psi_{cm} \Psi$ leads to

- $$-\frac{\hbar^2}{2m} \nabla_{cm}^2 \Psi_{cm} = E_{cm} \Psi_{cm}, \text{ which}$$

is a free particle

- $$-\frac{\hbar^2}{2\mu} \nabla^2 \Psi = E \Psi$$

- $$E_{total} = E + E_{cm}$$

- Since the separation between particles, a , is constant,

$$-\frac{\hbar^2}{2\mu a^2} \Lambda^2 \Psi = E \Psi. \text{ With } I = \mu a^2,$$

this is a particle on the sphere!

- Solutions are spherical harmonics
- Energy levels $E_{JM_J} = J(J+1) \frac{\hbar^2}{2I}$
- Model for microwave spectroscopy

Review

- Compare the allowed quantum numbers, energy levels, and energy gaps for a particle in a box, harmonic oscillator, particle on a ring, and particle on a sphere
- What are spherical harmonics?
- What is the relationship between l , m_l , and the number of angular nodes?
- Which analytical systems are models for infrared and microwave spectroscopy?