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- Postulates of quantum mechanics
 - States and wavefunctions
 - Observables and representations
 - The outcome of measurements
 - Eigenfunction expansions, including Fourier series
 - Collapse of wavefunctions and Schrödinger's cat
 - The Born interpretation: probabilities of particle positions
 - The Schrödinger equation
 - The Hamiltonian operator
 - Time-independent form by separation of variables
 - Time dependence of the wavefunction

Postulates

- This module is intended to help you achieve the following learning objectives:
 - Recall and explain the postulates of quantum mechanics
- At the end of this module, you should be able to
 - answer the following questions:
 - What does a wavefunction describe?
 - What the the commutators for position and momentum operators?
 - How do you evaluate an expectation value if the wavefunction is not an eigenfunction of an operator?
 - What happens if you perform a measurement on a superposition state?
 - What is the probability of finding a particle in a particular region?
 - What is the time-independent Schrödinger equation?

About the Postulates

- Assumptions that are the basis of quantum mechanics
- Are not intuitive
- Not provable
- Results consistent with observations

1. States and wavefunctions

- The state of a system is fully described by a function $\Psi(r_1, r_2, \ldots, t)$, known as the wavefunction
 - r_n is the position of particle n
 - particles can include electrons and nuclei
 - t is time, which is often not considered
- Over semester, we will study progressively more complex Ψ with more and more particles and dimensions

2. Observables

- Observables are represented by Hermitian operators chosen to satisfy the the commutation relations $[q,p_{q'}]=i\hbar\delta_{qq'}$, [q,q']=0, $[p_q,p_{q'}]=0$.
 - q and q' are coordinates like x, y, and z. They may be the same or different.
 - p_q is the momentum in the direction of q
- $i=\sqrt{-1}$ $\hbar=\frac{h}{2\pi}$, where $h=6.62607015\times 10^{-34}~\mathrm{J.Hz^{-1}}$ is Plank's constant
 - $\delta_{qq'}$ is a Kronecker delta function. It is one when q=q' and zero if $q \neq q'$.
- Any set of operators that satisfy this postulate may be chosen
- Many QM manipulations can actually be performed without representations

Representations

- Set of operators that represent observables
- The most popular are
 - Position representation

•
$$x \to x \times \frac{\pi}{i} \frac{\partial}{\partial x}$$

• Momentum representation

$$\begin{array}{c} x \to -\frac{\hbar}{i} \frac{\partial}{\partial p_x} \end{array}$$

- $p_x \rightarrow p_x \times$
- More generally, x may be replaced by q, which is x, y, or z.

Representations and Postulate 2

Does the position representation satisfy $[q,p_{q'}]=i\hbar\delta_{qq'}$? Here is how to show this.

$$\begin{bmatrix} q, \frac{\hbar}{i} \frac{\partial}{\partial q'} \end{bmatrix} f = q \frac{\hbar}{i} \frac{\partial f}{\partial q'} - \frac{\hbar}{i} \frac{\partial}{\partial q'} (qf)$$

$$= \frac{\hbar}{i} \left[q \frac{\partial f}{\partial q'} - \frac{\partial q}{\partial q'} f - q \frac{\partial f}{\partial q'} \right]$$

$$= -\frac{\hbar}{i} \frac{\partial q}{\partial q'} f = i\hbar \delta_{qq'} f$$

- Position representation
- Momentum representation

$$\begin{array}{c} x \to -\frac{h}{i} \frac{\partial}{\partial p_x} \end{array}$$

•
$$p_x \rightarrow p_x$$

Representations and Postulate 2

Does the momentum representation satisfy $[q,p_{q'}]=i\hbar\delta_{qq'}$? Try it.

$$\begin{split} & \left[-\frac{\hbar}{i} \frac{\partial}{\partial p_{q}}, p_{q'} \right] f \\ & = -\frac{\hbar}{i} \left[\frac{\partial}{\partial p_{q}} (p_{q'}f) - p_{q'} \frac{\partial f}{\partial p_{q}} \right] \\ & = -\frac{\hbar}{i} \frac{\partial p_{q'}}{\partial p_{q}} f = i\hbar \delta_{qq'} f \end{split}$$

- Position representation
- Momentum representation

$$\begin{array}{c} x \to -\frac{n}{i} \frac{\partial}{\partial p_x} \end{array}$$

•
$$p_x \rightarrow p_x$$

3. The outcome of measurements

The mean value of the observable in a series of measurements is equal to the expectation value of the corresponding operator,

$$\langle A \rangle = \frac{\left\langle \Psi \middle| \hat{A} \middle| \Psi \right\rangle}{\left\langle \Psi \middle| \Psi \right\rangle}$$

- This is like the expectation value of flipping a coin or rolling dice
- $\langle \Psi | \Psi \rangle$ is a normalizing constant. If Ψ is normalized (as we will assume), then $\langle \Psi | \Psi \rangle = 1$ can be omitted from the expectation value.

What if Ψ is not an eigenfunction of \hat{A} ?

. If
$$\hat{\mathbf{A}}\Psi=a\Psi$$
, then $\langle A\rangle=\left\langle \Psi\left|\hat{\mathbf{A}}\right|\Psi\right\rangle=a$.

 \bullet Otherwise, we can expand Ψ in terms of basis functions - building blocks for more complex functions

Eigenfunction Expansions

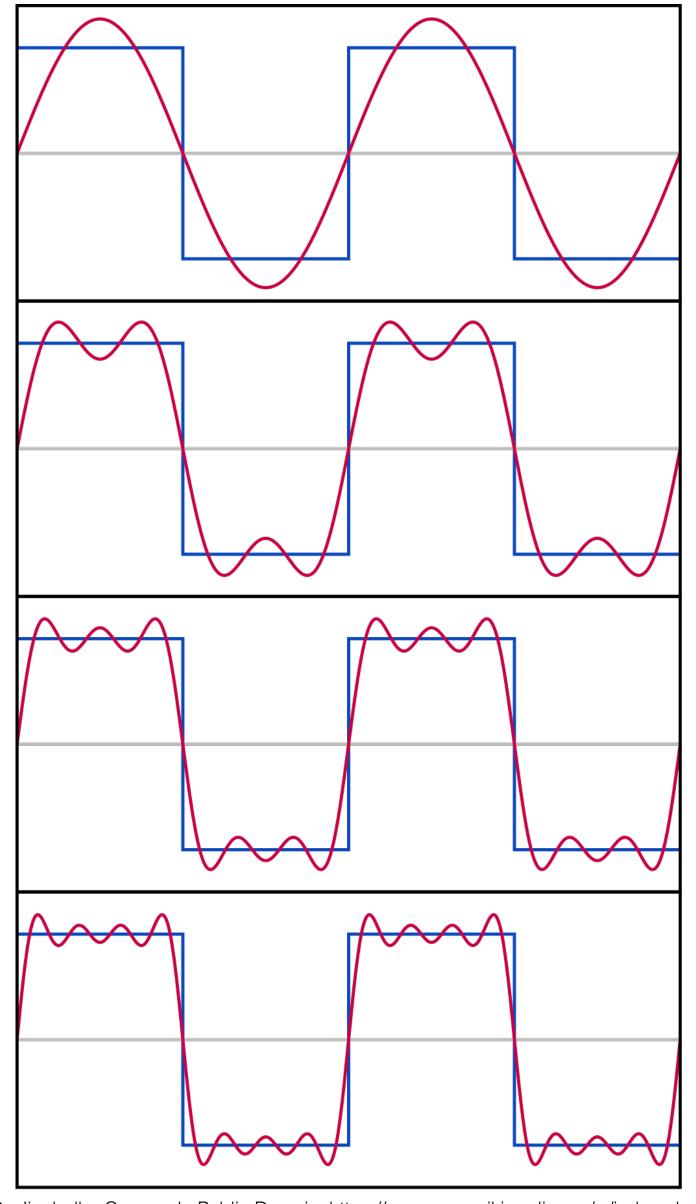
- Any function may be expressed as a linear combination of *all* the eigenfunctions of an operator, $g = \sum c_n f_n$.
 - c_n are coefficients
 - f_n are eigenfunctions. They are basis functions for g.

Example Eigenfunctions

- Any function may be expressed as a linear combination of *all* the eigenfunctions of an operator, $g = \sum c_n f_n$.
- . Consider the operator $\hat{\bf A}=\frac{d^2}{dx^2}$. What are the eigenfunctions? Hint: we talked about one in the last class period.
 - sin(ax), where a is a positive integer. The eigenvalue is $-a^2$.
 - cos(bx), where b is a positive integer. The eigenvalue is $-b^2$.
 - c, a number. The eigenvalue is 0.

Fourier Series

- An expansion of a periodic function in terms of sines and cosines
- To the right: the first four partial sums of the Fourier series for a square wave. As more terms are added, the partial sums converge to the square wave.
- For more details, examples, and expressions for coefficients, see Wolfram MathWorld.



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Mean values for linear combinations

- Given $\Psi = \sum c_m \Psi_m$, the probability of measuring a particular eigenvalue ω_n is proportional to $|c_m|^2$.
 - c_m is the coefficient
 - Ψ_m are eigenfunctions of $\hat{\mathbf{A}}$

•
$$|c_m|^2$$
 can be thought of as weights in a weighted sum,
$$\langle A \rangle = \int \left(\sum_n c_n \Psi_n\right)^* \hat{\mathbf{A}} \left(\sum_m c_m \Psi_m\right) d\tau = \sum_n \sum_m \int \left(c_n \Psi_n\right)^* \left(a_m c_m \Psi_m\right) d\tau = \sum_m a_m |c_m|^2$$

Measurement

- ullet Immediately after the *measurement*, the state of the system will be Ψ_n
- Unlike in classical mechanics, in QM, measurement alters the system. It makes the system *collapse* onto a state.
- This implies that before measurement, the system described by $\Psi = \sum c_m \Psi_m$ is in a *superposition* of states.

• Does this make sense? Schrödinger didn't think so. Yet this *Copenhagen interpretation* of QM remains.

Schrödinger's Cat

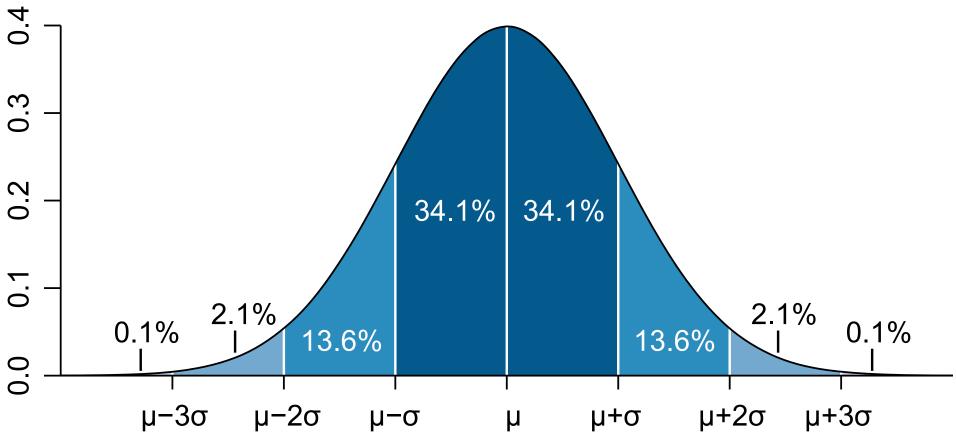


Schrödinger's cat: a cat, a flask of poison, and a radioactive source connected to a Geiger counter are placed in a sealed box. As illustrated, the objects are in a state of superposition: the cat is both alive and dead.

4. Probabilities of particle positions

- Born interpretation: the probability that a particle will be found in the volume element $d\tau$ at the point r is proportional to $|\Psi(r)|^2 d\tau$.
- The probability density function is $|\Psi(r)|^2$
 - $\int \left|\Psi\right|^2 d\tau < \infty$ because the particle must be somewhere

PDF of a normal distribution



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$$f(x) = \frac{1}{\sqrt{2\pi\sigma}} e^{-\frac{1}{2}\left(\frac{x-\mu^2}{\sigma}\right)}$$

5. The equation for the wavefunction

- The wavefunction evolves in time according to the time-dependent Schrödinger equation, $i\hbar\frac{\partial\Psi}{\partial t}=\hat{\mathbf{H}}\Psi$
 - introduced by Erwin Schrödinger in 1926
- If the potential energy is time-independent, then the equation may be simplified to the time-independent Schrödinger equation, $\hat{\mathbf{H}}\Psi=E\Psi$
 - the focus of this class
 - can be derived by separation of variables (will do so shortly)
- f H is the Hamiltonian operator, which specifies the total (kinetic + potential) energy of the system

The Hamiltonian operator

- \bullet \hat{H} is the Hamiltonian operator, which specifies the total (kinetic + potential) energy of the system
- ullet The kinetic energy of a particle with mass m and momentum p is

$$T = \frac{p^2}{2m}$$

- In the position representation, $\hat{\mathbf{T}} = \frac{\hbar^2}{2m} \nabla^2$. $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$.
- The potential energy of a particle depends on the system. Over the semester we will consider different potential energy functions

Time-Independent Schrödinger equation

using separation of variables to derive $\hat{\mathbf{H}}\Psi=E\Psi$

Separation of Variables, p. 1

- . The Schrödinger equation is $\hat{\mathbf{H}}\Psi=i\hbar\frac{\partial\Psi}{\partial t}$. In one dimension, this is $-\frac{\hbar}{2m}\frac{\partial^2\Psi}{\partial x^2}+V(x)\Psi=i\hbar\frac{\partial\Psi}{\partial t}$
- Separation of variables works by assuming that $\Psi(x,t) = \Psi(x)\Theta(t)$, which is true if the potential does not vary with time
- This yields $-\frac{\hbar^2}{2m}\Theta\frac{\partial^2\Psi}{\partial x^2} + V(x)\Psi\Theta = i\hbar\Psi\frac{\partial\Theta}{\partial t}$

Separation of Variables, p. 2

- Dividing by $\Psi\Theta$ gives, $-\frac{\hbar^2}{2m}\frac{1}{\Psi}\frac{\partial^2\Psi}{\partial x^2}+V(x)=i\hbar\frac{1}{\Theta}\frac{\partial\Theta}{\partial t}$
 - The lhs is solely a function of x and the rhs is solely a function of t; they are separable.
 - When *x* changes, the rhs does not change. When *t* changes, the lhs does not change.
 - Both sides are equal to a constant, E.
- As a self-test, write the two eigenfunction equations

$$-\frac{\hbar^2}{2m}\frac{\partial^2\Psi}{\partial x^2}+V(x)\Psi=\hat{\mathbf{H}}\Psi=E\Psi \qquad \text{Time-independent Schrödinger equation}$$

$$i\hbar\frac{\partial\Theta}{\partial t}=E\Theta$$

Time Dependence

- . One of the differential equations is $i\hbar\frac{\partial\Theta}{\partial t}=E\Theta$
- What is the solution of this differential equation?
 - $\bullet \quad \Theta = e^{-\frac{iE}{\hbar}t}$
- Thus if we have a solution for $\hat{\mathbf{H}}\Psi=E\Psi$, the complete equation is $\Psi(x,t)=\Psi(x)e^{-\frac{iE}{\hbar}t}$.
- Recall the Born interpretation: the probability that a particle will be found in the volume element $d\tau$ at the point r is proportional to $|\Psi(r)|^2 d\tau$. Does the probability of the particle position depend on time?
 - No

Summary of Postulates

- The state of a system is fully described by a function $\Psi(r_1,r_2,\ldots,t)$, known as the wavefunction
- Observables are *represented* by Hermitian operators chosen to satisfy the the commutation relations $[q,p_{q'}]=i\hbar\delta_{qq'}$, [q,q']=0, $[p_q,p_{q'}]=0$
- The mean value of the observable in a series of measurements is equal to the expectation value of the corresponding operator. For an eigenfunction expansion $\Psi = \sum c_m \Psi_m$, the
 - probability of measuring a particular eigenvalue ω_n is proportional to $|c_m|^2$. Measurement makes the system *collapse* onto a state.
- The probability that a particle will be found in the volume element $d\tau$ at the point r is proportional to $|\Psi(r)|^2 d\tau$.
- . The wavefunction evolves in time according to $i\hbar\frac{\partial\Psi}{\partial t}=\hat{\mathbf{H}}\Psi$. If the potential energy is time-independent, then $\hat{\mathbf{H}}\Psi=E\Psi$

Review

- What does a wavefunction describe?
- What the the commutators for position and momentum operators?
- How do you evaluate an expectation value if the wavefunction is not an eigenfunction of an operator?
- What happens if you perform a measurement on a superposition state?
- What is the probability of finding a particle in a particular region?
- What is the time-independent Schrödinger equation?