

MuMaX3 Tutorial Presentation

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Continuum magnetization field: $\mathbf{M}(\mathbf{r}, t) = M_s \mathbf{m}(\mathbf{r}, t)$

Energy

$$E = \int_V dV \left\{ A \sum_{\alpha, \beta \in \{x, y, z\}} (\partial_\alpha m_\beta)^2 \right. \quad \text{Exchange} \quad (1)$$

$$+ D(\mathcal{L}_{yz}^{(x)} + \dots) \quad \text{DMI} \quad (2)$$

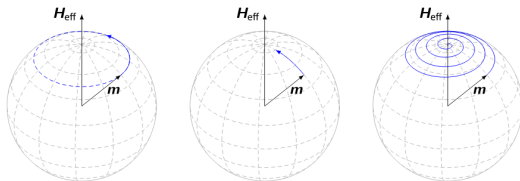
$$- K(\mathbf{m} \cdot \hat{u})^2 + K_c(m_x^2 m_y^2 + \dots) \quad \text{Anisotropies} \quad (3)$$

$$- \mu_0 \mathbf{M} \cdot \mathbf{H} \quad \text{Zeeman} \quad (4)$$

$$\left. - \frac{\mu_0}{2} \mathbf{M} \cdot \mathbf{H}_d \right\} \quad \text{Demag} \quad (5)$$

Effective field

$$\mathbf{H}_{\text{eff}}(\mathbf{r}) = -\frac{1}{\mu_0 M_s} \frac{\delta E(\mathbf{r})}{\delta \mathbf{m}} \quad (6)$$



LLG Precession, Damping and Full form¹.

□ Dynamics (spin-wave, vortex dynamics, domain-wall motion, FMR, etc.):

Landau-Lifshitz-Gilbert equation

$$\frac{\partial \mathbf{M}}{\partial t} = -\frac{\gamma}{1 + \alpha^2} \left(\mathbf{M} \times \mathbf{H}_{\text{eff}} + \frac{\alpha}{M_s} \mathbf{M} \times \mathbf{M} \times \mathbf{H}_{\text{eff}} \right) \quad (7)$$

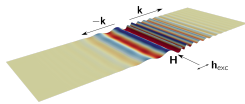
¹Image taken from <http://micromagnetics.org/micromagnetism/>. Based on: C. Abert, Discrete Mathematical Concepts in Micromagnetic Computations. 2013.

- Static (hysteresis, local energy minima, phase diagrams, etc.):
- Energy minimization via gradient algorithms: steepest descent

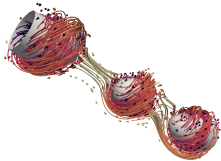
Exl et al. JAP 115, 17D118 (2014)

Minimizer stopping criteria

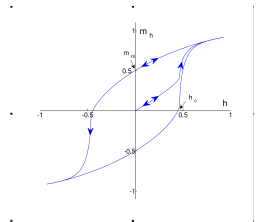
The minimization procedure stops when $\max(dM) < \text{"MinimizerStop"}$ for the last $\text{"MinimizerSamples"}$ steps.
The default values are $1e-6$ for MinimizerStop and 10 for minimizerSamples .



Spin waves

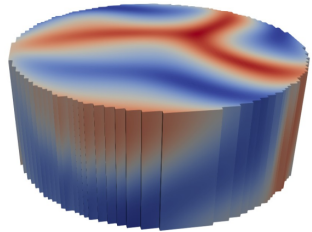
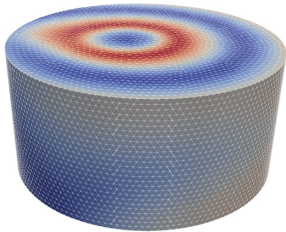


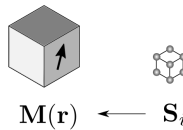
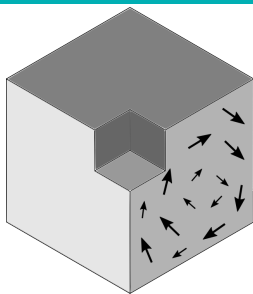
Local energy minimum states



Hysteresis

Numerical discretization





$$\ell_{\text{ex}} = \sqrt{\frac{2A}{\mu_0 M_s^2}}$$

Exchange length (8)

$$\delta_{\text{DW}} = \sqrt{\frac{A}{K}}$$

Domain Wall width (9)

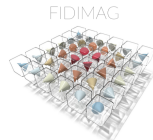
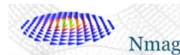
(10)



- Micromagnetic



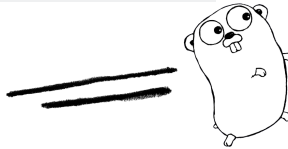
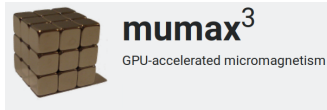
mumax³
GPU-accelerated micromagnetism
[Home](#) [Download](#) [Examples](#) [API](#)



FIDIMAG

- Atomistic





Golang

<https://mumax.github.io/>

<https://mumax.github.io/examples.html>

Script language based on **Go**: <https://mumax.github.io/api.html>

<https://mumax.github.io/api.html>

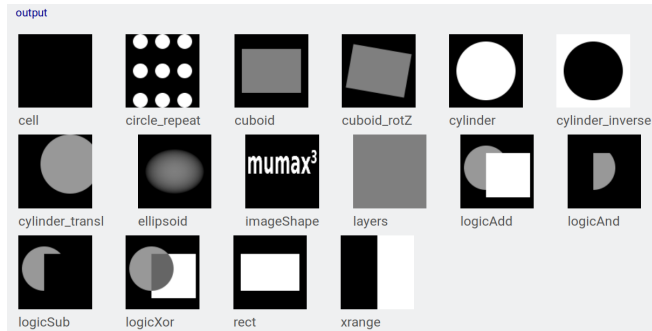
Aex

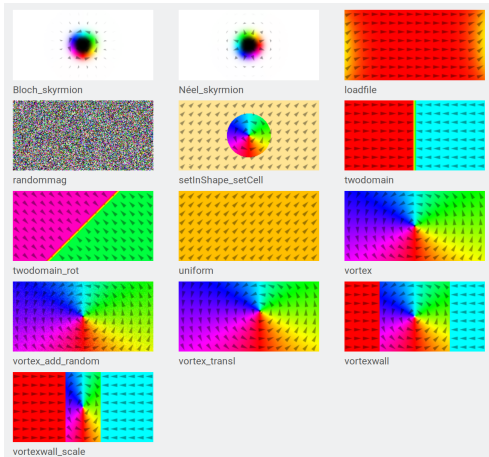
Exchange stiffness (J/m)

```
methods: Average( ) EvalTo( Slice ) GetRegion( int ) IsUniform( )  
MSlice( ) Region( int ) Set( float64 ) SetRegion( int ScalarFunction )  
SetRegionFuncGo( int func() float64 ) SetRegionValueGo( int float64 )
```

```
examples: [1] [2] [3] [6] [7] [8] [9] [10] [11] [12] [13] [14] [15]
```

<https://mumax.github.io/examples.html>





```

1 // Standard problem 2
2
3 lex := 5e-9
4 d := 30 * lex
5
6 lx := 5 * d
7 ly := 1 * d
8 lz := 0.1 * d
9
10 Ms := 1.0 / mu0 // 1 Tesla
11 Msat = Ms // Or: Msat.Set(Ms)
12 Aex = 0.5 * mu0 * lex * lex * pow(Ms, 2)
13
14 nx := pow(2, ilogb(lx / (0.5 * lex)) + 1)
15 ny := pow(2, ilogb(ly / (0.5 * lex)) + 1)
16 Print(nx, ny)
17 SetGridSize(nx, ny, 1)
18 SetCellSize(lx / nx, ly / ny, lz / 1.)
19
20 // alpha = 0.02
21 m = uniform(1., 1., 0.9)
22 SaveAs(m, "m_initial")
23
24 // ...

```

This standard micromagnetic problem includes both magnetostatic and exchange energies, but has the advantage of only one scaled parameter. If crystalline anisotropy is neglected and the geometry is fixed, scaling of the static micromagnetic equations (Brown's equations) yield a hysteresis loop which depends only on the scaled geometry to the exchange length when expressed as M/M_s versus H/H_m , where $H_m = M_s$ (SI) or $4\pi M_s$ (cgs emu). The exchange length is $lex = (A/K_m)^{1/2}$, where A is the exchange stiffness constant and K_m is a magnetostatic energy density, $K_m = 1/2\mu_0 M_s^2$ (SI) or $2\pi M_s^2$ (cgs emu).

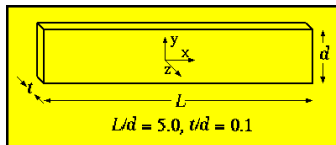
The field should be applied in the $[1,1,1]$ direction, approximately 54.74 degrees ($\arccos[1/\sqrt{3}]$) from each of the coordinate axes. This field orientation is chosen to avoid potential symmetry breaking problems.

Geometry:

Let us take a thin film of thickness t , width d , and length L . We suggest to make the problem virtually 2D by choosing $t/d = 0.1$, and to obtain interesting non-uniform reversal modes, $L/d = 5$.

Material parameters:

The magnetostatic exchange length, lex
Zero magnetocrystalline anisotropy



Desired output for comparison:

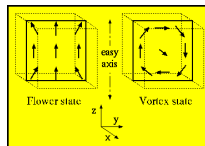
Calculated as a function of d/l_{ex} , with aspect ratios held constant at $t/d = 0.1$ and $L/d = 5.0$,

- Coercivity (H_c/H_m , the magnitude of the field at which the projection of the magnetization along the field, $M_x + M_y + M_z$, is zero.)
- Remanence (M_x/M_s , M_y/M_s , M_z/M_s , at $H = 0$)

```

1 // Standard problem 3
2
3 N := 64
4
5 SetGridSize(N, N, N)
6 // Setting the initial state
7
8 Lmin := 8.0
9 Lmax := 9.0
10 Lstep := 0.05
11 L := Lmin
12 setcellsize(L/N, L/N, L/N)
13
14 // Initial magnetization (xy-plane)
15 m = vortex(1, -1)
16
17 // Parameters in reduced units -> lex = 1.0
18 Msat = sqrt(2 / mu0)
19 Aex = 1.0
20 Ku1 = 0.1
21 anisU = vector(1, 0, 0)
22
23 tableadd(E_total)
24 tableaddvar(L, "L", "")
25
26 for L = Lmin ; L <= Lmax; L += Lstep {
27     ...
28 }

```



This problem is to calculate the single domain limit of a cubic magnetic particle. This is the size L of equal energy for the so-called flower state (which one may also call a splayed state or a modified single-domain state) on the one hand, and the vortex or curling state on the other hand. Material parameters:

Geometry:

A cube with edge length, L , expressed in units of the intrinsic length scale, $lex = (A/Km)^{1/2}$, where Km is a magnetostatic energy density, $Km = 1/2\mu_0 Ms^2$ (SI).

Material parameters:

Uniaxial anisotropy Ku with $Ku = 0.1 Km$, and with the easy axis directed parallel to a principal axis of the cube.

Desired output for comparison:

- ☐ The indicated "single domain limit," L where the energy of the flower state is equal to the energy of the vortex state.
- ☐ The partial energies (exchange, stray field, anisotropy) all in units of Km .
- ☐ The average magnetization along the three axes.

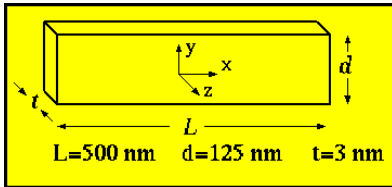
Normalize micromagnetic equations:

$$x \rightarrow \ell_{\text{ex}} x \quad A \rightarrow 1 \quad K \rightarrow K/K_m \quad M_s \rightarrow \sqrt{2/\mu_0}$$

$$E = \int_V dV \left\{ \begin{aligned} & A \sum_{\alpha, \beta \in \{x, y, z\}} (\partial_\alpha m_\beta)^2 && \text{Exchange} \\ & - K(\mathbf{m}_z)^2 && \text{Anisotropies} \\ & - \mu_0 \mathbf{M} \cdot \mathbf{H} && \text{Zeeman} \\ & - \frac{\mu_0}{2} M_s^2 \mathbf{m} \cdot \int \mathbf{N}(\mathbf{r} - \mathbf{r}') \cdot \mathbf{m}(\mathbf{r}') dV' \end{aligned} \right\} \quad \text{Demag}$$

```

1 SetGridsize(128, 32, 1)
2 SetCellsize(500e-9 / 128,
3             125e-9 / 32,
4             3e-9 / 1)
5
6 Msat = 800e3
7 Aex = 13e-12
8 alpha = 0.02
9
10 m = uniform(1, .1, 0)
11 relax()
12 save(m) // relaxed state
13
14 autosave(m, 200e-12)
15 tableautosave(10e-12)
16
17 B_ext = vector(-24.6E-3, 4.3E-3, 0)
18 run(1e-9)
    
```



Standard problem #4 is focused on dynamic aspects of micromagnetic computations. The initial state is an **equilibrium s-state** such as is obtained after applying and slowly reducing a **saturation field along the [1,1,1] direction to zero**.

Geometry: A film of thickness, $t=3$ nm, length, $L=500$ nm and width, $d=125$ nm will be used.

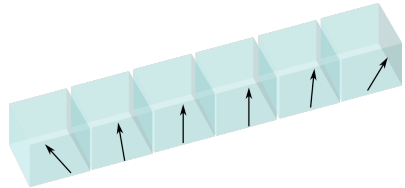
Material parameters similar to Permalloy: $A = 1.3e-11$ J/m, $M_s = 8.0e5$ A/m, $K = 0.0$, $\alpha = 0.02$, $\gamma_0 = 2.21e5$ (m / A s)

Applied Fields: Two switching events will be calculated using fields applied in the x-y plane of different magnitude and direction.

Field 1: $\mu_0 H_x = -24.6$ mT, $\mu_0 H_y = 4.3$ mT, $\mu_0 H_z = 0.0$ mT which is a field approximately 25 mT, directed 170 degrees counterclockwise from the positive x axis

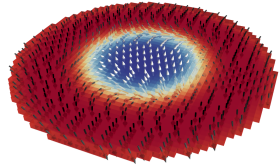
Field 2: $\mu_0 H_x = -35.5$ mT, $\mu_0 H_y = -6.3$ mT, $\mu_0 H_z = 0.0$ mT which is a field approximately 36 mT, directed 190 degrees counterclockwise from the positive x axis

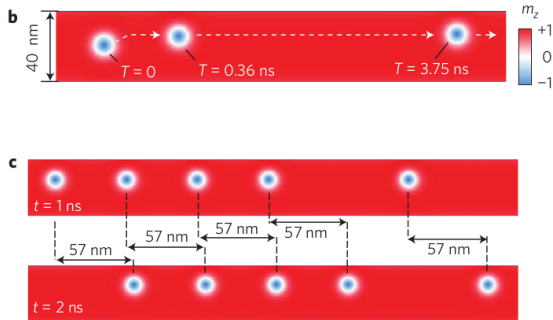
```
1 SetMesh(100, 1, 1, 1e-09, 1e-09, 1e-9,  
2         0, 0, 0)  
3  
4 Msat      = 0.86e6  
5 Aex       = 13.0e-12  
6 Ku1       = 0.4e6  
7 anisU     = vector(0, 0, 1)  
8 Dbulk     = 3.0e-3  
9 NoDemagSpins = 1  
10  
11 // Initial state  
12 m.setRegion(0, uniform(0, 0, 1))  
13  
14 OutputFormat = OV2F2_TEXT  
15  
16 // Relax with specific torque:  
17 // alpha      = 0.9  
18 // RunWhile(MaxTorque > 1e-4)  
19  
20 // Relax with conjugate gradient:  
21 minimize();  
22 SaveAs(m, "one_dim_bulk")
```



Standard Problem DMI paper

```
1 // Disk radius
2 R := 50e-9
3
4 // Number of cells
5 nx := 50
6 ny := 50
7 nz := 1
8 // Mesh discretisation
9 dx := 2e-9
10 dy := 2e-9
11 dz := 2e-9
12 // Set mesh and disk geometry
13 SetGridSize(nx, ny, nz)
14 SetCellSize(dx, dy, dz)
15 SetGeom(Circle(2 * R))
16 DefRegion(1, circle(R / 2))
17
18 Msat      = 860e3
19 Aex       = 13e-12
20 Ku1       = 0.4e6
21 anisU     = vector(0, 0, 1)
22 Dind      = -3e-3
23 // No Demag:
24 NoDemagSpins = 1
25
26 // Initial state
27 m.setRegion(0, uniform(0, 0, -1))
28 m.setRegion(1, uniform(0, 0, 1))
29
30 OutputFormat = OV2F2_TEXT
31 // relax()
32 minimize()
33 SaveAs(m, "Skyrmion-Inter")
34 Snapshot(m)
```





Sampaio et al. Nucleation, stability and current-induced motion of isolated magnetic skyrmions in nanostructures

$$\boldsymbol{\tau}_{\text{ZL}} = \frac{(1 + \alpha\xi)}{1 + \alpha^2} \mathbf{m} \times (\mathbf{m} \times (\mathbf{u} \cdot \nabla) \mathbf{m}) + \frac{(\xi - \alpha)}{1 + \alpha^2} [\mathbf{m} \times (\mathbf{u} \cdot \nabla) \mathbf{m}] \quad (11)$$

$$\mathbf{u} = \frac{\mu_B P}{2e\gamma_0 M_s (1 + \xi^2)} \mathbf{j} \quad (12)$$

Notice that for the current density \mathbf{j} the electron velocity \mathbf{u} is in the opposite direction.

The design and verification of MuMax3. AIP Advances 4, 107133 (2014).

```
1 // Track length
2 Lx := 400e-9
3 Ly := 40e-9
4 Lz := 0.4e-9
5
6 // Number of cells (power of 2) -> discretization of 1.25
7 nx := pow(2, 6) * 5
8 ny := 32
9 nz := 1
10
11 // Set mesh and disk geometry
12 SetGridSize(nx, ny, nz)
13 SetCellSize(dx, dy, dz)
14
15 Msat      = 580e3
16 Aex       = 15e-12
17 Ku1       = 0.8e6
18 anisU     = vector(0, 0, 1)
19 Dind      = 3e-3
20 // No Demag:
21 // NoDemagSpins = 1
```

```
22 // Initial state
23 m = NeelSkyrmion(1, -1).scale(3, 3, 1).translate()
24
25 OutputFormat = OVf2_BINARY
26 minimize()
27 SaveAs(m, "sk_track_relax")
28 SnapshotAs(m, "sk_track_relax.png")
29
30 alpha = 0.3
31 // Now set the current motion: in-plane -> Zhang-Li model
32 // (see eq in Landau-Lifshitz form)
33 // non-adiabaticity:: (0.15, 0.30 and 0.60)
34 pol = 0.4
35 j = vector(-0.1e12, 0, 0) // 10 MA cm^-2 in the x-direction
36 xi = 0.15
37
38 autosave(m, 5e-10)
39 tableAutosave(1e-11)
40 tableAdd(ext_bubblepos)
41
42 run(10e-9)
```

Sampaio et al. Nucleation, stability and current-induced motion of isolated magnetic skyrmions in nanostructures

Non-adiabaticity parameter β (0.15, 0.30 and 0.60)

Current densities: 0.01 - 100 MA cm⁻²

Polarization $P = 0.4$

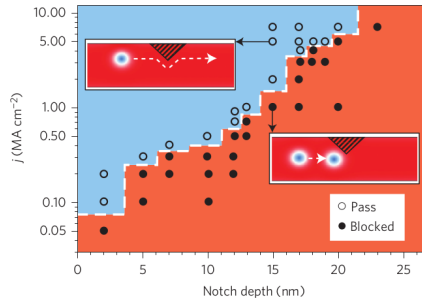


Figure 4 | Pinning of the current-induced motion of skyrmions by defects in a 40-nm-wide nanotrack. The defect (notch) is a triangular region of higher perpendicular anisotropy ($K = 1.2 \text{ MJ m}^{-3}$ instead of 0.8 MJ m^{-3} for the rest of the track), and the motion has been simulated for different notch

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