

Exclusivity graph approach to Instrumental inequalities

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Instrumental tests have been widely used to derive causal relations between variables even in presence of unobserved latent factors, thus providing a powerful tool for different sciences ranging from xxx to xxx. Recently, much interest have been dedicated to the instrumental scenario since admits quantum violations with fewer resources respect to the Bell analogue, this property can be exploited to (drastically?) improve our current capacities to process information. Dire qualcosa della contextuality. Here, we present the instrumental inequality under a contextual approach, showing that the geometry of correlations within this framework allow us to find optimal

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INTRODUCTION

The existence of causal relationships between events is one of the basic principles that makes scientific investigation possible. Nonetheless, extracting informations about these relationships only from experimental observations is not always feasible and can be a tricky task. In this context causal models are a indispensable tool to formalize causal assumptions and to produce falsifiable statements causal relations.

A causal model is typically represented by a Directed Acyclic Graph (DAG), where the nodes represent observed events, which are described by random variables, and the edges the cause-effect relationship between them. From some causal DAGs it is possible to derive constraints on the statistical distribution of the variables which allow one to reject or accept a causal assumption[1].

It is known, since Bell theorem, that quantum mechanics allows for statistical behavior which are impossible to reproduce in a classical context, so it is reasonable to suppose that the constraints derived from classical causal assumptions could be violated by quantum systems. Indeed Bell experiment can be modeled using a causal DAG, as shown in Fig 1-b, and the Bell inequality obtained as a consequence of causal assumptions. More recently it has been proved and experimentally demonstrated[4] that the well-known “instrumental” causal structure (see Fig. 1-a) posses classical constraints that are violated by quantum mechanics.

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INSTRUMENTAL CAUSAL STRUCTURE

Among all the possible causal structures, a particular interest has been devoted to the *Instrumental causal structure* (see Fig.1-a) introduced by Pearl[2] in the 90's with applications ranging from econometrics[] to clinical trial[]. One of the most important properties of in-

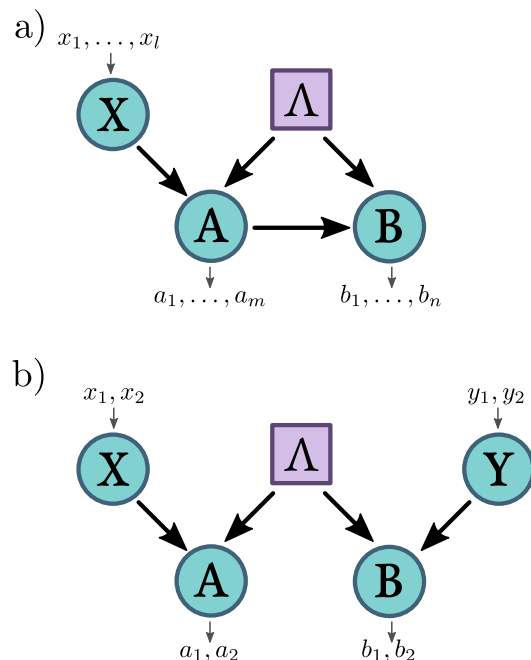


FIG. 1. ma se levamo quelle cose che stanno per dire input e ooutput nelle immagine? tipo lasiare la struttura pulita, magari il resto possiamo dirlo nella caption? Si può fare ma mi servirebbe qualcosa che ricordi il fatto che questa é una generica instrumental con l valori per X , m per A e n per b , lo metto nella caption? **a)** The DAG representing a general Instrumental scenario. **b)** The DAG of the CHSH scenario.

instrumental variables is the fact that allow us to bound causal relations between variables solely from observational data. Like in the Bell tests where the inequalities allow us to test local-realism for a particular distribution, instrumental inequalities test whether we have a valid instrument. Recently has been experimentally proved that instrumental inequalities are also violated when considering quantum resources [4], thus opening several venues for future research such as device-independent cryptography [5], random number generation [6] and basically any

application that until now has been linked with the Bell scenario. In light of the importance regarding the instrumental causal structure, we are interested in exploring new techniques in order to better address the question of whether we can obtain a quantum violation, and what is the optimal test constrained on given experimental data.

Let us consider the instrumental structure depicted in Fig.?? where X corresponds to the instrument, A and B the variables for which we are interested to bound the causal relation and Λ our quantum resource. If we allow the classical variables X, A, B to take the values l, m, n respectively the Bayesian probability distribution for the $i(j)$ -th outcome of $A(B)$ given x_k reads $P(a_i, b_j | x_k)$. Pearl showed that, under the assumption of the instrumental structure, exist the following inequalities bounding the probability distribution:

$$\sum_{j=0}^n P(a_i b_j | x_{k(i,j)}) \leq 1 \quad (1)$$

for all $i \in 1, \dots, m$ and for all the possible functions $k(i, j)$.

Davide, ho messo i "max" perche mi la ricordo così, ma tocca controllare i index. Non so se avevi levato a posta per il tuo approach boh. Si, é meglio mettere le singole disuguaglianze invece che quella col max, é carino perché ognuna corrisponde con gli edges del grafo di esclusività!

In 2001 Bonet [3] extended the results of Pearl by deriving another instrumental inequality, hence proving that the polytope of correlations is not tightly bounded by Pearl's inequalities.

For instance in the case $(l, m, n) = (3, 2, 2)$ the following inequality is also satisfied:

$$P(a_1 b_1 | x_1) + P(a_2 b_2 | x_1) + P(a_1 b_1 | x_2) + P(a_2 b_1 | x_2) + P(a_1 b_2 | x_3) \leq 2 \quad (2)$$

along with other similar inequalities obtained by permuting the labels of a_i, b_j and x_k .

These inequalities are related to the one introduced in [4]:

$$-\langle B \rangle_{x_1} + 2\langle B \rangle_{x_2} + \langle A \rangle_{x_1} - \langle AB \rangle_{x_1} + 2\langle AB \rangle_{x_3} \leq 3 \quad (3)$$

which was experimentally violated in the same work, with a value of 3.790 ± 0.013 . As one can expect, with the same data it is also possible to violate inequality (2) with a value of 2.159 ± 0.005 .

EXCLUSIVITY GRAPH METHOD

The exclusivity graph method was primarily developed to study non-contextual inequalities and obtain their quantum violation. [6] as well as for non-locality scenarios[?]. In this formalism every possible event, i.e. every possible set of outcomes $a_1 \dots, a_n$ of measurements

A_1, \dots, A_n given settings x_1, \dots, x_n , is associated to a vertex in a (undirected) graph $G = (V, E)$. Two vertices $u, v \in V$ are then connected by an edge $uv \in E$ if and only if they are exclusive, i.e. if exists a measurement that can distinguish between them.

A linear constraint can be expressed as a linear function

$$I_w(p) = \sum_{\substack{a_1, \dots, a_n \\ x_1, \dots, x_n}} w_{a_1, \dots, a_n} p(a_1, \dots, a_n | x_1, \dots, x_n) \quad (4)$$

on the probabilities for each possible event. In this framework this can be represented by weighted exclusivity graph G , and it is possible to relate classical bound and quantum violation to well-known graph invariants: the independence number $\alpha(G, w)$ and the Lovász theta $\theta(G, w)$.

Here we briefly introduce these concepts and their relationships, a more extensive and general treatment can be found in [6? , 7]

Consider a graph $G(V, E)$ with vertex weights w , and $|V| = n$. We call a *characteristic labelling* for $U \subseteq V$ a vector $x_v \in \{0, 1\}^n$ such that $x_v = 1$ if $v \in U$ and $x_v = 0$ otherwise.

An *independent set* or *stable set* is a set $S \subset V$ such that $uv \notin E$ for all $u, v \in S$. The independence number $\alpha(G, w)$ is defined as the maximum number of vertices (weighted with w) of an independent set of G . Is also customary to define the set $\text{STAB}(G)$ as the convex hull of all the characteristic labellings of stable sets:

$$\text{STAB}(G) = \{x : x \text{ is a stable labelling of } G\} \quad (5)$$

using this definition the independence number becomes:

$$\alpha(G, w) = \max\{w \cdot x : x \in \text{STAB}(G)\} \quad (6)$$

From this definitions we can see that $\alpha(G, w)$ corresponds to the classical bound of inequality, since it is exactly the maximum over the convex set defined by all the deterministic strategies respecting the exclusivity constraints.

We now call an *orthonormal labelling* of dimension d a map $a_v : V \rightarrow \mathbb{R}^d$ such that $a_v \cdot a_u = 0$ for all $uv \in E$ and $|a_v|^2 = 1$, and we define the set $\text{TH}(G)$ as:

$$\text{TH}(G) = \{x : x_v = a_v^1 \text{ where } a_v \text{ is an orthonormal labelling of } G\} \quad (7)$$

It can be proved that this set includes all the possible quantum strategies, but in general is larger as it contains also the "almost quantum correlations" described in [?]. [?] Maximizing over $\text{TH}(G)$ gives the Lovász theta:

$$\theta(G, w) = \{w \cdot x : x \in \text{TH}(G)\} \quad (8)$$

which upper-bounds the maximum quantum value.

Moreover this approach gives a useful condition to check if a given graph G admits a quantum violation.

Indeed it can be proved that $\text{TH}(G) = \text{STAB}(G)$ if and only if G does not contain a cycle C_n with $n \geq 5$ and odd, or its complement as an induced subgraph. This follows directly from the so called “sandwich theorem” and the “strong perfect graph theorem” [?].

Exclusivity graph method applied to causal models

The techniques presented in the previous section can be also employed to analyze a certain class of causal models. Suppose we have a causal scenario as in fig 2, with k observable variables, l instruments with no incoming edges, and a single unobservable latent variable Λ .

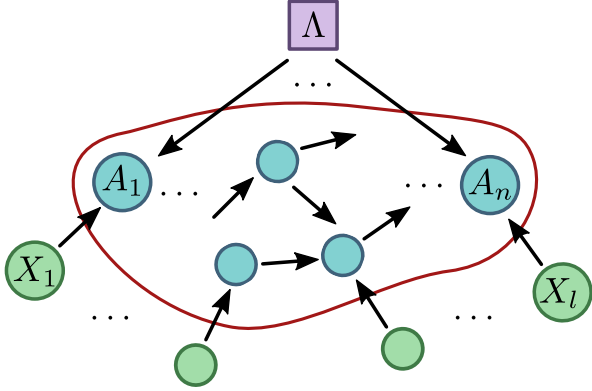


FIG. 2. A general causal graph with l instrument and a single latent variable.

In this case we can associate with such a DAG a *non-exclusivity* graph H as follows:

- Nodes are associated to events like $a|x$, where $a = (a_1, \dots, a_k)$ and $x = (x_1, \dots, x_l)$.
- Two nodes $a|x$, and $a'|x'$ are linked by an edge if and only if
 1. for all $i, j \in \{1, \dots, k\}$ if $A_i \rightarrow A_j$ then $\exists g_{ij} : g_{ij}(a_i) = a_j$ and $g_{ij}(a'_i) = a'_j$.
 2. for all $i \in \{1, \dots, l\}$ and $j \in \{1, \dots, k\}$ if $X_i \rightarrow A_j$ then $\exists f_{ij} : f_{ij}(x_i) = a_j$ and $f_{ij}(x'_i) = a'_j$.

To the complement of this graph, $G = \bar{H}$, and its subgraphs, the previous methods can be applied to obtain inequalities with their respective quantum and classical bounds.

The Instrumental exclusivity graph

Many known results about the instrumental scenario, including the inequality (2), can also be obtained using this formalism. Restricting for the moment to the

case of dichotomic observables ($n = m = 2$), we label as $p(ab|x)$ with $a, b \in \mathcal{A} = \mathcal{B} = \{0, 1\}$ and $x \in \mathcal{X} = \{0, \dots, l\}$, the probability of having outcomes a and b with setting x . As explained in the previous section we obtain the non-exclusivity graph for the instrumental scenario by linking two events $ab|x$ and $a'b'|x'$ if exist two functions $f : \mathcal{X} \rightarrow \mathcal{A}$ and $g : \mathcal{A} \rightarrow \mathcal{B}$ such that:

$$a = f(x) \quad \text{and} \quad a' = f(x') \quad (9)$$

$$b = g(a) \quad \text{and} \quad b' = g(a') \quad (10)$$

Using this method we can construct the graph for various l , as shown in figure 3, and use the methods described earlier to obtain classical and quantum bounds for several inequalities in the instrumental scenario.

First, we immediately obtain the result that there is no quantum violation in the case of $l = 2$, since the corresponding exclusivity graph (and its complement) does not contain any odd cyclic graph with more than 5 vertices.

For $l \geq 3$ instead we see that there is a violation for the inequality (2), since it forms a C_5 cyclic graph, i.e. the pentagon depicted in fig. 4. Knowing that $\alpha(C_n) = \lfloor n/2 \rfloor$ we immediately derive that the classical limit is 2. Analogously knowing that $\theta(C_n) = n \cos(\pi/n) / (1 + \cos(\pi/n))$ we get $\sqrt{5}$ as an upper bound for the quantum limit.

In this framework (2) seems analogous to the KCBS contextual inequality [8]. The difference here is that the quantum limit is not $\sqrt{5}$, which we cannot obtain in a bipartite scenario. To find a tight bound we must apply other methods to our graph, like the ones in [7], using which we find a value of $(3 + \sqrt{2})/2$.

It can be easily proved (see Methods) that no other odd anticyle beside C_5 is present for any l , so from this point of view, we can say that any quantum violation for a general $l22$ instrumental scenario has essentially the same origin: it stems from the non classicality of the C_5 graph.

This does not mean that different inequalities cannot be devised by clever choices of vertices and weights. For example the inequality for the 422 scenario presented in [?]

$$\langle AB \rangle_1 - \langle AB \rangle_2 + \langle B \rangle_0 - \langle B \rangle_3 \leq 2 \quad (11)$$

corresponds to the two sets of events whose exclusivity graphs are depicted in fig. 5. The maximum classical value for both, given by the independence number of the graphs, is 3, while their minimum is 1, which gives the maximum value of 2 as in (11) when taken with opposite sign. This inequality has a structure very similar to the CHSH (showed in fig. 6), and, as pointed out above, both arise from the non-classicality of the C_5 graph.

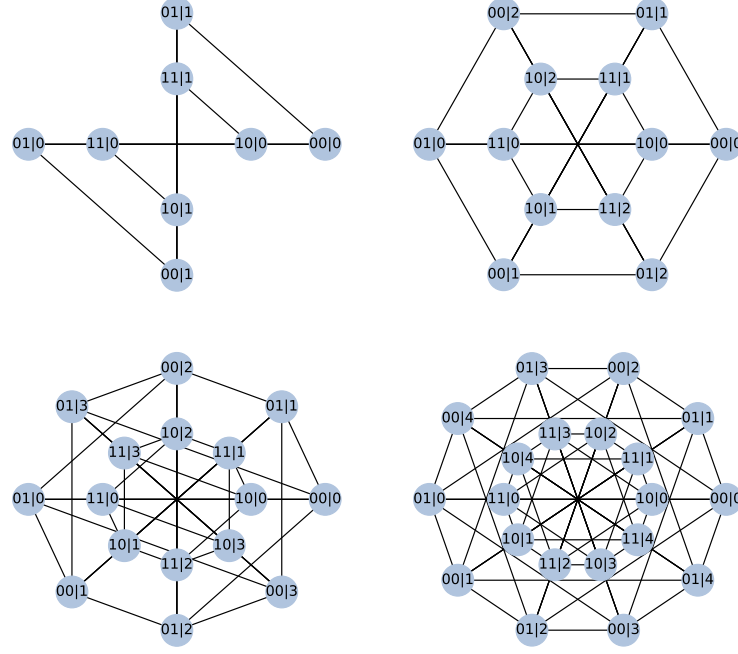


FIG. 3. Davide, sicuramente ante delle figure andranno nel SI, ma puoi comunque sistemare già la disposizione, perchè utilizzano troppo spazio, non sono segnalate con a,b,c,d etc. Magari potresti scrivere la caption per quelle che sai già che andranno nel main text. The exclusivity graph for the instrumental scenario for $l = 2, 3, 4, 5$ from top left to bottom right. In the figure the bold lines represent cliques.

Finding optimal inequalities

The presence of an odd cycle graph guarantees only that $\text{STAB}(G) \subset \text{TH}(G)$, but does not say anything on the values of $\theta(G)$ and $\alpha(G)$. Nonetheless, if that condition is satisfied there must be some optimal choice of weights w for which we can obtain a maximum quantum violation. Formally if n is the number of nodes of G , we want to find the maximum:

$$\max_{w \in [0,1]^n} \{\theta(G, w) - \alpha(G, w)\} \quad (12)$$

Using the definition (8) and (6), and remembering that $\text{STAB}(G)$ is a convex hull of a finite number of characteristic labellings we can rewrite the previous maximization as:

$$\Delta = \min\{I_y : y \text{ is a maximal stable labelling}\} \quad (13)$$

$$I_y = \max_{\substack{w \in [0,1]^n \\ x \in \text{TH}(G)}} \{w \cdot (x - y)\} \quad (14)$$

The maximization (14) for each y is efficient, being a semi-definite program, but unfortunately it has to be

run for each maximal stable set in G which are $\sim 3^{n/3}$ in the worst case. Also to enumerate all the maximal independent sets we used the Tomita version of the Bron–Kerbosch algorithm which has the same worst case complexity[11].

METHODS

There are no cycles C_n with $n \geq 7$ in the $d22$ instrumental scenario.

Here we prove that there cannot be a odd antihole with more than 5 vertices in the exclusivity graph associated to an instrumental scenario of the type $d22$.

Remind that, from the exclusivity conditions (??), given two events $ab|x$ and $a'b'|x'$, they are connected by edge if one of these two conditions is true:

1. $x = x'$.
2. $a = a'$ and $b \neq b'$.

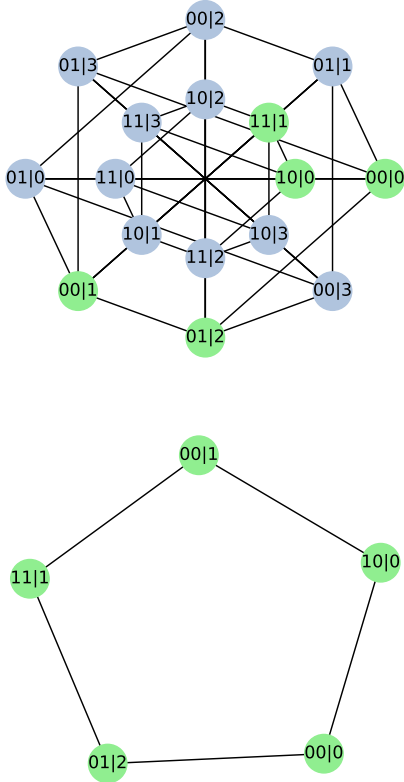


FIG. 4. The exclusivity graph of the bonet inequality.

Suppose we have a cycle C_n with $n \geq 7$, as in fig. 7, and consider node 2 which correspond to an event which we can arbitrarily identify with $00|0$. Among its neighbors 1 and 3, one will necessarily need to satisfy rule 2 (they cannot both satisfy rule 1 or the three nodes would be a clique). So without loss of generality we can assign the event $01|1$ to 3. Since nodes 5, 6, 7 must not satisfy rule 2 with both 2 and 3 they must have $a = 1$. Moreover 7 and 5 must have the same b , different from 6. In the same way 1 must not satisfy rule 2 with 6, 5 and 3, so it needs to have $a = 0$ and $b = 1$. But now, since we only have values $\{0, 1\}$ for a , we cannot avoid 4 to be linked to one of the nodes 1, 2, 6, 7.

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expressed in this publication are those of the authors and do not necessarily reflect the views of the John Templeton Foundation) **Gli acknowledgements? Ma non era un report? :D**

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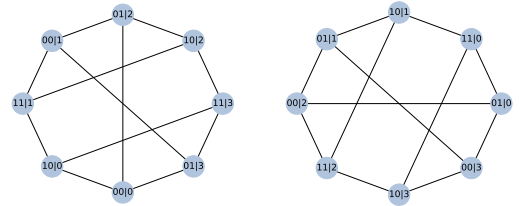


FIG. 5. The exclusivity graph for the 422 Instrumental scenario.

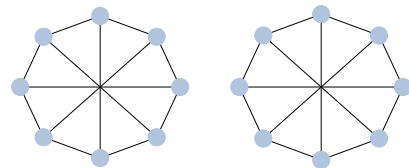


FIG. 6. The exclusivity graphs for the the CHSH scenario.

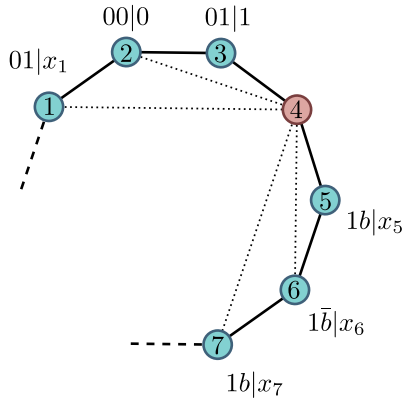


FIG. 7. Proof of the impossibility of having cycles with 7 nodes or more in the $d22$ scenario.

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