國立清華大學 物理系 碩士學位論文

基於 SPA-Net 的雙頂夸克 全強子衰變事件重建

Event reconstruction of all hadronic Top-quark-pair decays using SPA-Net

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Event reconstruction of all hadronic Top-quark-pair decays using SPA-Net

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Abstract

The top quarks produce by pp collision in Large Hadron Collider(LHC), have a very complicated process still can't be well-classified today. In this project, we present a novel approach to the "all hadronic decay" process of Top quarks base on the neural networks with attention mechanism, we call it "Symmetry Preserving Attention Networks" (SPA-Net). This networks identify the decay products of each quarks unambiguously and without combinatorial explosion. This approach perform a outstanding result compare to the existing state-of-the-are method. Our network can correctly assigning all hadronic decay in 93.0% of 6 jets, 87.8% of 7 jets, and 82.6% of \geq 8 jets event respectively.

摘要

在大型強子對撞機 (LHC) 實驗中,經由質子對撞所產生的頂夸克對具有非常複雜的過程以及產物,至今仍無法被非常正確的判別以及重建。在本研究中,我們提出了一個利用新穎的機器學習方法來對雙頂夸克全強子衰變過程進行重建。此方法基於 Attention mechanism,我們稱之為 Symmetry Preserving Attention Networks(SPA-Net)。這個模型架構可以在避免組合性爆炸的前提下對所有的衰變產物進行辨識以及重建。此方法對比於傳統的 χ^2 重建方式,表現出了非常巨大的差異。本方法可以在一、存在 6 jets 條件下正確的重建 93% 的事件;二、存在 7 jets 條件下正確的重建 87% 的事件;三、存在大於 8 jets 條件下正確的重建 82.6% 的事件。

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Chapter 1

Introduction

At Large Hadron Collider(LHC), two protons collide with very high energy and produce many kinds of products. A process that pp collision produces a pair of Top quark and result in the 6 jets final state is called **Full Hadronic Top-quark-pair decay**. This process has a very complicated signature due to a large number of combinations produce by the possible permutation of final state jets. These jets produced by the top quark pair is hard to tag as a specific daughter of top quarks correctly. A traditional method is to reconstruct the event using χ^2 reconstruction, but it takes a long time to compute and cannot provide enough accuracy(about 30% or less) to reconstruct an event. The importance of studying Top quark and its full hadronic decay channel is 1. Top quark is the most heaviest fundamental particle in standard model and will decay before hadronization, 2. The branching ratio of full hadonic decay is the biggest part in Top quark decay(46%).

For a problem that contains a large amount of data and highly requires computing resources, machine learning can widely provide powerful support on solving the problem and helps to reduce the CPU time. The machine learning method helps to discover physics phenomena with very outstanding effort. A remarkable discovery that helps by machine learning is the discovery of Higgs Boson. Both CMS and ATLAS groups apply machine learning methods to promote the search for the Higgs Boson. [1][2]

In this thesis, we develop a novel architechture for parton-jet assignment problem. This method is base on the state-of-the-art machine learning technology, Attention mechanism.[3] We call this novel ML model Symmetry Preserving Attention NETworks

(SPA-NET). By applying attention networks, the SPA-NET perform a outstanding performance compare to traditional method while avoiding combinatorial explosion. And thanks to the natural properties of attention network, the network reflect the permutation symmetry naturally and provide a chance to explore the model with permutation symmetry.

This work is done by the collaboration between Nation Tsing Hua University(Ta-Wei Ho), University of Washington(Shih-Chieh Hsu), and University of California Irvine(Mike Fenton, Alexander Shmakov, Daiel Whiteson, and Pierre Baldi). Mike and I focus on the physics concept, design the data format, and generate datasets. Alexander provide a techical support and machine learning network setup.

We will discuss the Top physics and the concept of machine learning in chapter 2; and explain our event generation and simulation configuration in chapter 3; then introduce how we analyze the dataset and reconstruct the event using traditional method and ML approach in chapter 4. We will discuss our work in chapter 5 and summerize in chapter 6.



Chapter 2

The Top Physics and Machine Learning

2.1 The Top Physics

Top quark, the most massive fundamental particle in Standard Model(SM), is the only quark that decays semi-weakly. Its large mass leads to a short lifetime and decay before hadronization occurs. Top quark contains so many properties that interest us, such as its mass, couplings, and cross-section, etc. Measure these properties accurately can bring us a worth understanding of fundamental interactions and provide the key to Beyond Standard Model.[4]

In Standard Model, Top quark pair produced by pp collision has three decay modes, all-hadronic channel, semi-leptonic channel, and dileptonic channel. The branch ratio of each channel, has shown in the Table 2.1. The decay width of Top quark predicted in SM is[5]:

$$\Gamma_t = \frac{G_F m_t^3}{8\pi\sqrt{2}} \left(1 - \frac{M_W^2}{m_t^2} \right)^2 \left(1 + 2\frac{M_W^2}{m_t^2} \right) \times \left[1 - \frac{2\alpha_s}{3\pi} \left(\frac{2\pi^2}{3} - \frac{5}{2} \right) \right]$$
(2.1)

Table 2.1: Top quark pair decay process[4]

Decay Channel	Process	Branch Ratio(%)
All-hadronic	$t\bar{t} \to W^+ b W^- \bar{b} \to q\bar{q}' b q'' \bar{q}''' \bar{b}$	45.7
Semi-leptonic	$t\bar{t} \to W^+ b W^- \bar{b} \to q\bar{q}' b \ell^- \bar{\nu}_\ell \bar{b} + \ell^+ \nu_\ell b q'' \bar{q}''' \bar{b}$	43.8
Dileptonic	$t\bar{t} \to W^+ b W^- \bar{b} \to \ell^+ \nu_\ell b \ell' \bar{\nu}_{\ell'} \bar{b}$	10.5

In recent study, the most precise result of Top quark mass is measured in the lepton+jets channel due to its good signal-to-background ratio and the presence of one neutrino final state. [8] Although the all-hadronic channel has the most probability to appears in the Top quark pair decay process, it couldn't provide a precise mass measurement due to its poor signal-to-background ratio. The poor signal-to-background ratio of the all-hadronic channel is due to the difficult QCD background. The CMS and ATLAS group approach a precision of Top mass measurement using the all-hadronic channel with uncertainties of 0.65% and 1.1%. [6][7]

The channel we are interested in this project is the **jet-parton assignment problem** in all hadronic decay channel. The reason that we are interested in this channel is the resolved 6 jets signature and the potential of the machine learning method apply to the ambiguous event reconstruction problem. There exist 6 jets in the final state, 2 b-jets and 4 quark jets, they can be separated into two groups (b, q, \bar{q}) and (\bar{b}, q, \bar{q}) . A schematic of the decay products is shown in 2.1.

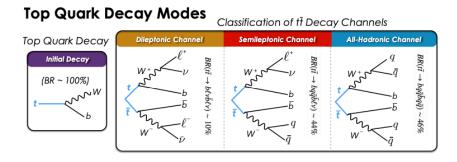


Figure 2.1: The schematic of Top quark decay channels.[8]

2.2 Machine Learning and its application on Particle Physics

Machine Learning has been applied to most of the region(e.g. Computer Vision, Solving PDE, Medical analysis, etc.) in recent age, including dose particle physics. From the search of higgs boson(neural network and BDT) to the b-tagging technology(BDT[9]), physicist already applied several kinds of machine learning method to recent research. In a nutshell, machine learning can break into several cases; it can help to do classification, regression, and clustering problems. It can not only help to accelerate the computation of well-defined problems, and also find a new path to unsolved area. We will use a state-of-the-art machine learning technology, attention mechanism.[3] The attention mechanism is a technology base on the evolution of Recurrent Neural Networks(RNN).[3] The attention mechanism will not only consider the local relationship and the sequence neighbor but also calculate the global relation base on the self-attention calculation shown in Figure 2.2. Using this novel architecture, we will train on the relationship between each jet and try to figure out the correct pair information.

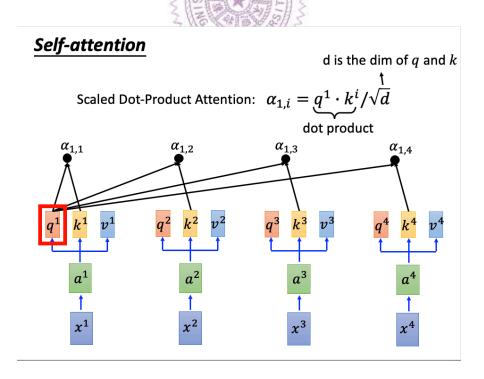


Figure 2.2: How self-attention works.[10]

Chapter 3

Event Generation

3.1 MC samples

For data preparation, we generate our dataset using a custom docker image with Mad-Graph_aMC@NLO(v2.7.2), Pythia8(v.8.2), and Delphes(v3.4.2) for showering, hadronization, and detector simulation. We apply the ATLAS parametrization during detector simulation. The data are generated at Leading order quantum chromodynamics(QCD). The top mass is configured as $m_{\rm top}=173$ GeV. We force the W quark decay hadronically into a (q,q') pair. Following is our configuration:

```
generate p p > t t~ QED=0, (t > W+ b, W+ > j j), (t~ > w- b~, w- > j j) output <file_path> launch <file_path> shower=Pythia8 detector=Delphes analysis=OFF done set nevents = 10000 set iseed = 1 Delphes/cards/delphes_card_ATLAS.tcl done exit
```

Listing 3.1: Configuration for generating samples. The "iseed" is just a place-holder, it will be changed while generating samples.

To get a more general performance, we scan the iseed value from 1 to 30000, each value

has 10 thousand events before event selection. The reason for scanning iseed value is that the iseed value is the key to the random generation. Originally, the program will choose the iseed value randomly and generate different samples. By scanning the issed value, we can make sure the seed is not reused.



Chapter 4

Data analysis and Event reconstruction

In this chapter, we will first explain the event selection, then how we apply the truth matching, traditional reconstruct method, and the machine learning approach.

4.1 Data analysis

4.1.1 Event selection

The top all hadronic decay channel has 2 b-jets and 4 quark jets, all of them in our configuration are not in the boosted region, that means the daughters of top quarks will not appears with high transverse momentum. Following the event selection used in the [8], we apply a event selection that an event should at least exists **2** b-jets and **4** quark jets satisfied p_T larger than **25** GeV and $|\eta|$ less than **2.5**. A cutflow table and figure can help us to understand how many events are killed by the selections. We may apply 5 cuts and see the evolution of survived event numbers. The rule of cuts is shown in Table 4.1, and the cutflow is shown in Figure 4.1.

Table 4.1: Rule of cuts. All the cuts require a kinematic limitation that $p_T > 25$ GeV and $|\eta| < 2.5$.

#Cut	Number of b-jets	Number of quark jets
C1	0	4
C2	0	5
C3	0	6
C4	1	6
C5	2	6

The b-tagging and jet information we used here is provided by the Delphes, a detector simulation package.[11] The b-tagging is a method of jet flavor tagging used in CMS and ATLAS.[12][13] This method base on the b-hadron properties, such as the displaced vertex from the primary vertex, large b-hadron mass, large impact parameters(d0), and semi-leptonic e/μ decay of B-hadron. This is also related to the track reconstruction and secondary vertex reconstruction. The Delphes package decides a jet is b-tagged or not base on an efficiency table and returns an array to indicate a jet is b-tagged or not. The number of survived events is an important factor in our event generation. If the survived events are very rare, it means our cuts are too tight that may also ignore the events we want. Also, the lack of survived events may slow down our data generation efficiency. In Figure 4.1, it shows that there are around $18 \sim 20\%$ events survived after C5 cut. This is an acceptable number because we obtain the evets in our dataset will at least contains 2 b-tag jets and 6 jets that passed the kinematics selection without ignoring too many events.

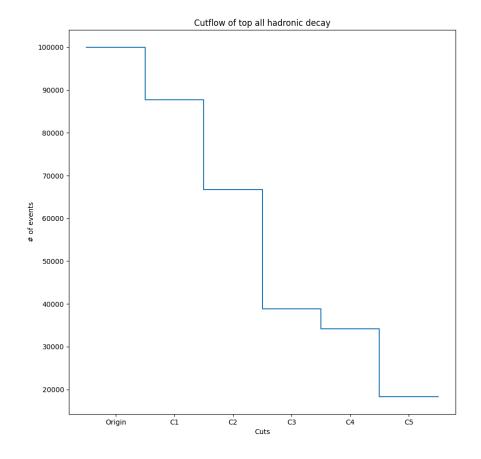


Figure 4.1: Cutflow of all hadronic top decay.

4.1.2 Truth matching

The truth matching, which is also called $\Delta \mathbf{R}$ matching, is to match the detector simulation (i.e. jet information generate by Delphes) data to truth record (i.e. Parton level information). To calculate the ΔR value, we will find the daughters of top quarks, W boson, and b quark. After the daughters of top quarks are found, we will find the daughters of W bosons. Finally, we will get six partons that come from the decay of top quark pairs. These six partons can match the jets identically by considering their distances. The formula of calculating ΔR is:

$$\Delta R = \sqrt{\Delta \eta^2 + \Delta \phi^2} \tag{4.1}$$

By using the kinematic properties provide in parton level and detector simulation information, we can calculate the ΔR value between each parton and jets. Using the result of the calculation, we may assign each parton to a specific jet.

4.1.3 Custom barcode system

To specify the relation between each parton, and the relation between mothers and daughters, we design a barcode system that helps us to declare the relationship.

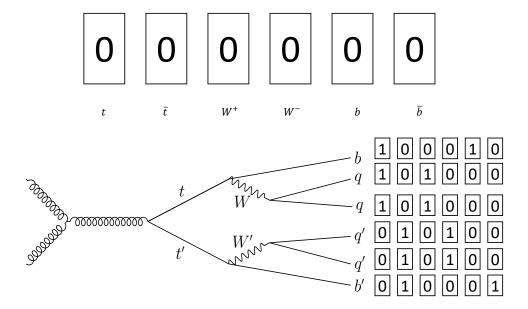


Figure 4.2: Design of barcode. These barcode can provide a pari-wide information to our network.

In Figure 4.2, we define a six-digit barcode, the first two digits are to show which top quark is the mother of this parton, the last four digits of the sequence is to declare which daughter of the top quark is the mother of parton. In case, we can use this barcode system to break six parton(jet) candidates into two subsets which contains 3 elements. The benefit of using this barcode system is not only can specify the relation without losing the information, but also provide a permutation relationship to our network. We will discuss this in the following section.

4.2 Event reconstruction

4.2.1 χ^2 minimization method

The χ^2 minimization method is a traditional method to reconstruct an event. For an event that exists 6 jets, it has about $6!/(2\times2\times2)=90$ possible combinations, the first two in the denominator is contributed by two b-tag jets, and the middle one and last one is the contributions from two pairs of quark produce by W bosons. The number of possible combinations is proportional to the number of existing jets in an event. The χ^2 minimization method will calculate all the candidates and try to find the candidate which has the smallest χ^2 value. This method base on the mass of the W boson and top quark. The origin equation of χ^2 minimization in this model is:

$$\chi^2 = \frac{(m_{bqq'} - m_t)^2}{\sigma_t^2} + \frac{(m_{\bar{b}q''q'''} - m_t)^2}{\sigma_t^2} + \frac{(m_{qq'} - m_W)^2}{\sigma_W^2} + \frac{(m_{q''q'''} - m_W)^2}{\sigma_W^2}$$
(4.2)

The equation 4.2 has four part. Each part is a "pull" that contribute by the component of observables. The parameter σ_W and σ_t is obtained by applying a fitting to the distribution of reconstructed invariant mass of W boson and top quarks. And the mass of W boson and top quark is given by the recent experiment result. To avoid the bias of top quark candidates, we may combine the first two term into one term by substituting $m_t = \frac{m_{bqq'} + m_{\bar{b}q''q'''}}{2}$, then the equation reduces to:

$$\chi^2 = \frac{(m_{bqq'} - m_{\bar{b}q''q'''})^2}{\sigma_{\Delta m_{bqq'}}^2} + \frac{(m_{qq'} - m_W)^2}{\sigma_W^2} + \frac{(m_{q''q'''} - m_W)^2}{\sigma_W^2}$$
(4.3)

Note that there are some events the three-jets invariant mass can be far away from the top quark mass but still give us a small χ^2 value. This is because we only consider the difference between two three-jets invariant mass. This can be improved by applying a

constraint to the invariant mass[8], but we didn't apply such a constraint in this report. In this project, we force the b quark candidates in 4.3 must be a b-tagged jets. This may help to reduce the number of permutations but prevent the event that contains a jet mistagged be assigned correctly.

4.2.2 Machine Learning Approach

For a machine learning model, equivariance and invariance are important properties that may affect the performance of the model. Such as a computer vision problem, the object should be invariant under translation to prevent affect the prediction. The translation, rotation, and shift of the position should not change the prediction of the model because the object remains the same object. The Convolutional Neural Network(CNN) can produce object recognition outcomes that are invariant under translations. The properties of invariance can be generalized to another geometry structure, e.g. manifolds and groups. In all hadronic top decay, we have two subsets with the same elements (b, q, q') and (\bar{b}, q'', q''') . These subsets should remain invariant under permutations of the input jets order. The reason that the order of input jets should not affect the result is because the permutation symmetry is not base on the order of jets but the pair-wise permutation.

By the invariant feature of attention architecture, rearranging the elements in a sequence leaves the attention weight unchanged. We may use this permutation symmetry present in the attention-based model to handle the reconstruction of all hadronic top decay process efficiently. In this case, the network output of the all hadronic top decay process should identify two distinct interchangeable subsets, add each contains an interchangeable qq' pair produce by the W bosons. This invariant property on the output is the unique feature of our dataset and the model should take into account.

We propose an attention-based network, called **Symmetry Preserving Attention NETwork(SPA-NET)** in this project. Its structure is shown in Figure 4.3. The input of SPA-NET is a list of unsorted jet information, with their 4-momentum (p_T, η, ϕ, M) as well as the b-tag information provided by Delphes. We take the logarithm to M and p_T and normalize all the components to have zero mean and unit variance. The input jets will be sent to the network and be embedded into

a D-dimensional latent space representation. This D-dimensional latent space is obtained by progressively increasing the latent dimensionality of the input jets up to the final dimensionality D(This operation is done by the embedding blocks in the whole stack). We target this latent space dimensionality D to a value 128 with the follow sequence: $8 \to 16 \to 32 \to 64 \to 128$. After embedding, the output will be sent to a stack of transformer encoder layers, this layer will learn the relationship between each element in the input sequence. While the encoding is finished, the encoded output will be forwarded into an important architecture in this network, a two branches structure that computes the output individually. There is a transformer encoder in each branch, this encoder layer will extract the information of top quark and the tensor attention layer will produce the top quark distribution. The output distribution predicts the top quark triplets. Figure 4.4 is a example of network outputs. Note that the attention mechanism is an architecture that allows the network to propagate the information selectively by using a "mask". By the implementation of the mask, the neural network can learn from partial information and update the parameters with selected information, and allow the network to infer the relationships between different elements in a sequence.

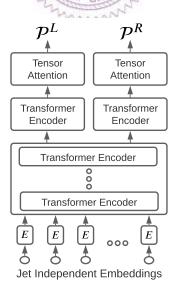


Figure 4.3: High-level structure of SPA-NET.

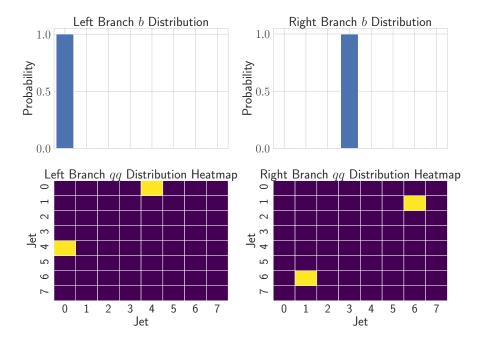


Figure 4.4: A visualization of the example single event output produce by SPA-NET. The top two plots is the projected b quark distribution, and the bottom two plots is the qq' distribution respectly.

The most important part in this network is the **Symmetry Preserving Tensor Attention**. Consider a set of weight $\theta \in \mathbb{R}^{D \times D \times D}$, this θ is not inherently symmetric at all. To make the θ become an invariant attention weighting, we apply the following transformation(eq. 4.4). This transformation will transform the θ into an auxiliary weights tensor $S^{ijk} \in R^{D \times D \times D}$. Using S^{ijk} and the embedded jets tensor $X \in \mathbb{R}^{N \times D}$ (N is the number of jets), we can calculate the dot-product attention. The dot-product attention working in flat Euclidean space and produce the output tensor O^{ijk} . The summation product tensor S^{ijk} guarantees the interchangeable of the first two dimensions of S will be symmetric and ensuring that $O^{ijk} = O^{jik}$. These properties enforce the qq' invariance.

$$S^{ijk} = \frac{1}{2} \left(\theta^{ijk} + \theta^{jik} \right)$$

$$O^{ijk} = X_n^i X_m^j X_l^k S^{nml}$$

$$(4.4)$$

To obtain the probability distributions P^L and P^R , we apply a 3-dimensional softmax on O^{ijk} to generate the joint triplet probability distribution.

$$P(i,j,k) = \frac{\exp O^{ijk}}{\sum_{ijk} \exp O^{ijk}}$$
(4.5)

We use equation 4.5 to produce the individual probability distribution of two top quarks and produce the single triplet from each by selecting the peak of these distributions.

During the training, a suitable loss function is needed to deal with the double output probability distributions. We design the loss function based on the cross-entropy between the output probability and truth distribution on the all hadronic top decay. The loss function must ensure the symmetry of the top quark pairs which are invariant concerning the permutation $tt' \leftrightarrow t't$. We create a symmetry loss function \mathcal{L} by the following function:

$$\mathcal{L} = min(\mathcal{L}_1(P^L, T_1, P^R, T_2), \mathcal{L}_1(P^L, T_2, P^R, T_1))$$
(4.6)

$$\mathcal{L}_1(P_1, T_1, P_2, L_2) = \mathcal{H}(T_1, P_1) + \mathcal{H}(T_2, P_2)$$
(4.7)

Where the \mathcal{H} is the general cross-entropy. It is possible that both branches produce the same output pairs. To make sure the network produces unique predictions, we will select the higher probability one and re-evaluate the other one, then compute the loss function. The Figure 4.4 is a example of the output produce by SPA-NET.

Chapter 5

Result and Discussion

In this chapter, we will discuss our result and compare the performance between the traditional method and machine learning approach.

5.1 Invariant mass and reconstruct efficiency

Before discussing the result, we should define the category for the reconstruction efficiency. Treating the truth matching result as the target, the prediction of χ^2 or SPA-NET may produce three kinds of results:

- Correct matched: An event that both top quarks are correctly predicted by the reconstruction method.
- Incorrect matched: An event that one or both top quarks is incorrectly predicted by the reconstruction method.
- Unmatched: An event that one of the truth match results contains an element that does not match to any jets.

Based on the category above, we can compute the reconstruction efficiency. For the reconstruction efficiency, we separate it into two parts: event-based efficiency and

quark-based efficiency. For event-based efficiency, we will calculate the efficiency base on how many **events** are matched correctly, incorrectly, and unmatched. Meanwhile, the quark-based efficiency calculates the efficiency base on how much **top quarks** are assigned correctly. The efficiency is shown in the table below:

Table 5.1: Using ϵ as the symbol of efficiency. This table performs the efficiencies of the χ^2 and SPA-NET assignments assessed by per-event efficiency ϵ^{event} and per-top efficiencies ϵ^{top} inclusively and by jet multiplicity $N_{\rm jets}$. The subscript of ϵ_1^{top} and ϵ_2^{top} is stands for the one/two reconstructable events.

	χ^2 Method		SPA-NET			
$N_{ m jets}$	$\epsilon^{\mathrm{event}}$	$\epsilon_2^{\mathrm{top}}$	$\epsilon_1^{\mathrm{top}}$	$\epsilon^{ m event}$	$\epsilon_2^{\mathrm{top}}$	$\epsilon_1^{\mathrm{top}}$
6	61.8%	65.0%	24.2%	80.7%	84.1%	56.7%
7				66.8%		
≥8	23.2%	35.5%	20.2%	52.3%	66.2%	52.9%
Inclusive	37.7%	47.0%	23.0%	63.7%	73.5%	55.2%

The χ^2 has performed a 37.7% efficiency on overall events, while SPA-NET archives an efficiency 63.7%. The χ^2 method has the best performance on the 6 jets category and has the worst effort on the category that an event contains more than 8 jets. The SPA-NET perform much better the χ^2 in all category. For the event that contains two identifiable, the χ^2 method achieves an efficiency ϵ_2^{top} 65.0%, and SPA-NET archives an ϵ_2^{top} 84.1%. Since we train the SPA-NET we the events that contain two top quarks, it is reasonable that the ϵ_1^{top} of SPA-NET is lower than ϵ_2^{top} . Also, since our definition of equation 4.2 is base on the difference between reconstructing invariant mass of 2 two quark candidates, so an event that only contains one identifiable top quark, is hard for χ^2 method to assign the jet properly. Note that in our evaluation dataset, 8.1% of events in which both tops are identifiable have at least one *b*-quark matched to non-*b*-tagged jets. These quarks, which are impossible for our χ^2 to correctly reconstruct, are reconstructed by SPA-NET with an efficiency of 29.4%.

5.1.1 Reconstructed invariant mass

Using χ^2 minimization, we may obtain the best assignment under the constraint of parameters. In this project, the parameters are configured as: $m_W = 81.3$ GeV, $\sigma_W = 12.3$ GeV, and $\sigma_{m_{bjj}} = 26.3$ GeV. The parameters σ_W and $\sigma_{m_{bjj}}$ are found by fitting the mass distribution of W boson and top quark. While computing χ^2 value, we use the b-tagging information to assign the b quark candidates.

Following is the reconstructed mass distribution of W boson and top quark using χ^2 minimization method and SPA-NET.

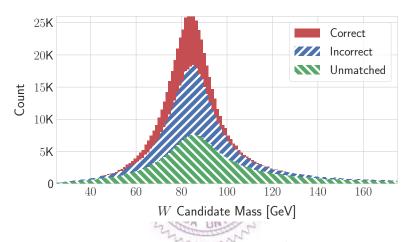


Figure 5.1: W boson mass reconstructed by χ^2 minimization method

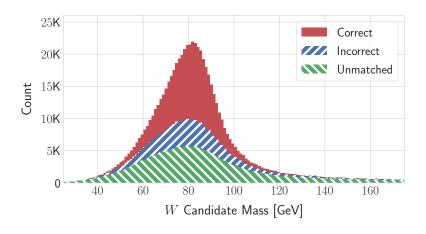


Figure 5.2: W boson mass reconstructed by SPA-NET

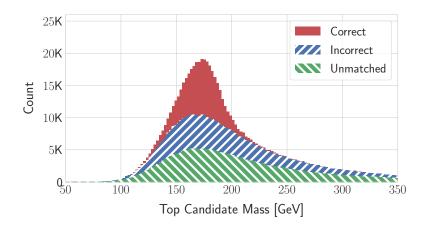


Figure 5.3: Top quark mass reconstructed by χ^2 minimization method

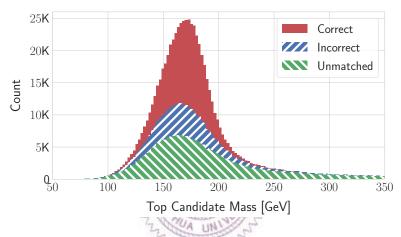


Figure 5.4: Top quark mass reconstructed by SPA-NET

Comparing the result shown in Figure 5.1 and Figure 5.2, we found the χ^2 method has a narrower peak around W boson mass than the SPA-NET. This shape comes from the incorrect and unmatched events and can be explained by the presence of m_W in equation 4.2. Another point is the Figure 5.4 and Figure 5.3 shows that the SPA-NET has the more peaked distribution compare to χ^2 method with comparable incorrect/unmatched events.

5.1.2 ROC curve

The Receiver operating characteristic (ROC) curve is a good target to estimate the performance of a machine learning model. The ROC curve of SPA-NET apply on events with one and two reconstructable top quarks is shown in Figure 5.5 and Figure 5.6. Note that the targets are defined as 1 if the prediction is correct, otherwise the 0 represents the incorrect prediction.

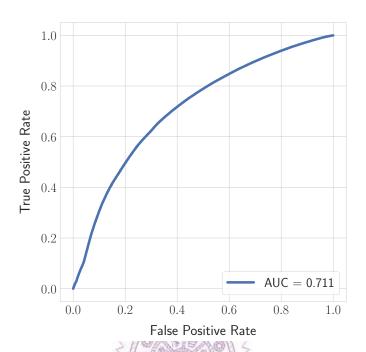


Figure 5.5: ROC curve of SPA-NET apply on events with one reconstructable top.

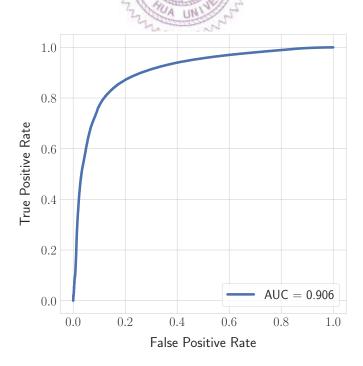


Figure 5.6: ROC curve of SPA-NET apply on events with two reconstructable top.

As the advantage we discuss in the last paragraph of section 5.1, our network archieves a lower AUC value in Figure 5.5 and perform a remarkable ROC curve in Figure 5.6. A possible solution we haven't implemented in this report is to train with the "partial" events by using the "mask".[14][15]

5.2 Reduce of time usage

The time required to compute SPA-NET is much lower than the χ^2 needed. We may compare the time they needed by considering the time complexity. The χ^2 has a time complexity scales approximately as $P(N,6) = \mathcal{O}(N^6)$ where N is the number of jets in an event. The time complexity is χ^2 proportional to the number of jets and this makes a limitation of maximum jets is indeed. Consider a 2019 DELL XPS13 computer with Core i7-1065G7 1.30GHz CPU, the SPA-NET took an average of 4.4 ms for one event. The χ^2 took an average 20 ms in 6 jets events and 369 ms in \geq 8 jets events.

5.3 Outlook

Based on the design of SPA-NET, the input is a sequence of jets information and their relationships. The network consider the permutation relation between each subset and possible permutations. In case, a physics model which has a permutation relationship might be a good item to explore with SPA-NET. For example, ttH all hadronic decay process or all hadronic four top decay process is a potential target to study with SPA-NET. Further more, not only the parton-jet assignment problem but also the clustering problem, graph matching problem, and so does other problem which contains a permutaion symmetry, is a potential item to study with SPA-NET. We plan to explore the SPA-NET with some interesting BSM problem, such as the tri-Higgs production in 2HDM model.[16]

Chapter 6

Conclusion

We perform a novel approach to parton-jet assignment using a symmetry preserving attention mechanism. This network could learn the permutation symmetry that appears in a physics process and archieves a remarkable performance. Using this architecture can lead to a great improvement of performance and reduce the cost of computing time. And we shown that a well-trained deep learning network can replace the permutaion based algorithms and avoid the combinatorial explosion by estimating symmetry-aware pair-wise similarities.

Also, this novel architecture can not only provide an improvement to all hadronic top decay parton-jet assignment problem but also can be generalized to other models which also contain the permutation symmetry.

Bibliography

- [1] G. Aad et al. [ATLAS], "Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC," Phys. Lett. B 716, 1-29 (2012) doi:10.1016/j.physletb.2012.08.020 [arXiv:1207.7214 [hep-ex]].
- [2] S. Chatrchyan et al. [CMS], "Observation of a New Boson at a Mass of 125 GeV with the CMS Experiment at the LHC," Phys. Lett. B 716, 30-61 (2012) doi:10.1016/j.physletb.2012.08.021 [arXiv:1207.7235 [hep-ex]].
- [3] A. Vaswani *et al.*, "Attention is all you need," Advances in Neural Information Processing Systems, NIPS (2017)
- [4] P. A. Zyla et al. [Particle Data Group], PTEP 2020, no.8, 083C01 (2020) doi:10.1093/ptep/ptaa104
- [5] A. Quade, "Top quark physics at hadron colliders," doi:10.1007/978-3-540-71060-8
- [6] A. M. Sirunyan *et al.* [CMS], "Measurement of the top quark mass in the all-jets final state at $\sqrt{s} = 13$ TeV and combination with the lepton+jets channel," Eur. Phys. J. C **79**, no.4, 313 (2019) doi:10.1140/epjc/s10052-019-6788-2 [arXiv:1812.10534 [hep-ex]].
- [7] M. Aaboud et al. [ATLAS], "Top-quark mass measurement in the all-hadronic $t\bar{t}$ decay channel at $\sqrt{s}=8$ TeV with the ATLAS detector," JHEP **09**, 118 (2017) doi:10.1007/JHEP09(2017)118 [arXiv:1702.07546 [hep-ex]].

- [8] T. Mccarthy, "Measurement of the Top Quark Mass in the All-Hadronic Top-Antitop Decay Channel Using Proton-Proton Collision Data from the ATLAS Experiment at a Centre-of-Mass Energy of 8 TeV," CERN-THESIS-2015-275.
- [9] M. Paganini [ATLAS], J. Phys. Conf. Ser. 1085, no.4, 042031 (2018)
 doi:10.1088/1742-6596/1085/4/042031 [arXiv:1711.08811 [hep-ex]].
- [10] Hung-Yi, Lee, "NTU course lecture note-Transformer"
- [11] J. de Favereau *et al.* [DELPHES 3], JHEP **02**, 057 (2014) doi:10.1007/ JHEP02(2014)057 [arXiv:1307.6346 [hep-ex]].
- [12] [ATLAS], ATL-PHYS-PUB-2016-012.
- [13] A. M. Sirunyan *et al.* [CMS], JINST **13**, no.05, P05011 (2018) doi:10.1088/1748-0221/13/05/P05011 [arXiv:1712.07158 [physics.ins-det]].
- [14] M. J. Fenton, A. Shmakov, T. W. Ho, S. C. Hsu, D. Whiteson and P. Baldi, [arXiv:2010.09206 [hep-ex]].
- [15] A. Shmakov, M. J. Fenton, T. W. Ho, S. C. Hsu, D. Whiteson and P. Baldi, [arXiv:2106.03898 [hep-ex]].
- [16] I. Low, N. R. Shah and X. P. Wang, [arXiv:2012.00773 [hep-ph]].