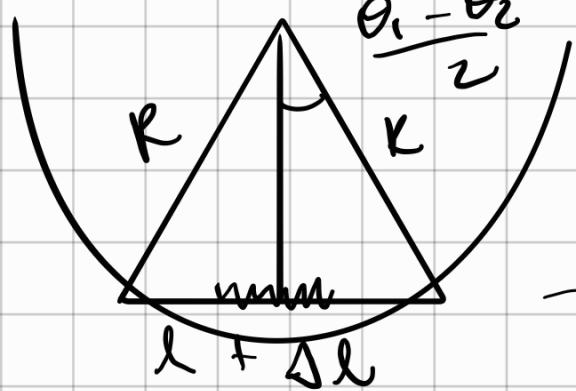


110 parte

3)



$$2K \sin\left(\frac{\theta_1 - \theta_2}{2}\right) = l + \Delta l$$

$$V = -\frac{1}{2}K\left(2K \sin\left(\frac{\theta_1 - \theta_2}{2}\right) - l\right)^2 - m_1 g R \sin \theta_1 - m_2 g R \sin \theta_2.$$

$$T = \frac{1}{2}m_1(R\dot{\theta}_1^2) + \frac{1}{2}m_2(R\dot{\theta}_2^2)$$

$$L = \frac{1}{2}R^2(m_1\dot{\theta}_1^2 + m_2\dot{\theta}_2^2)$$

$$+ \frac{1}{2}K\left(2R \sin\left(\frac{\theta_1 - \theta_2}{2}\right) - l\right)^2$$

$$+ gR(m_1 \sin \theta_1 - m_2 \sin \theta_2)$$

en este caso los coordenados generalizados son θ_1 y θ_2 a)

b.) $\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{\theta}_1}\right) = \frac{dL}{d\theta_1}$ para θ_1

$$\rightarrow R^2 m_1 \ddot{\theta}_1 = K \left(2R \sin\left(\frac{\theta_1 - \theta_2}{2}\right) - l\right)^\circ$$

$$K \cos\left(\frac{\theta_1 - \theta_2}{2}\right)$$

$$+ g R m_1 \cos \theta_1$$

para θ_2

$$\begin{aligned} R^2 m_2 \ddot{\theta}_2 &= K \left(2 R \sin\left(\frac{\theta_1 - \theta_2}{2}\right) - l \right) \\ &\quad \cdot \left(l - R \cos\left(\frac{\theta_1 - \theta_2}{2}\right) \right) \\ &\quad + g R m_2 \cos \theta_2 \end{aligned}$$

C.

d.)

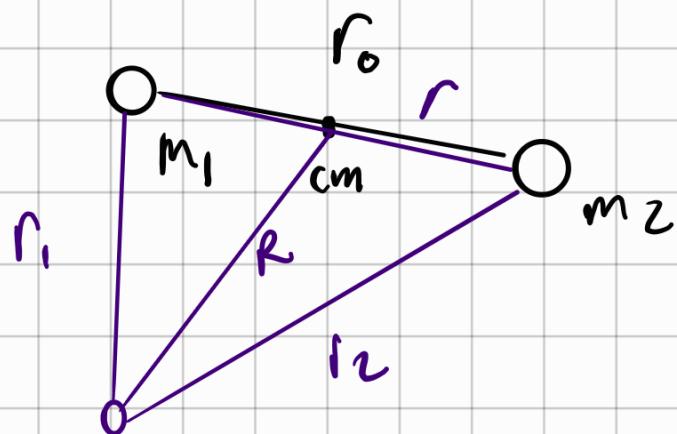
$$\frac{dV}{d\theta_1} = K \left[2R \sin\left(\frac{\theta_1 - \theta_2}{z}\right) - l \right] \left[K \cos\left(\frac{\theta_1 - \theta_2}{z}\right) \right] \\ + m_1 g R \cos \theta_1 = 0$$

$$\frac{dV}{d\theta_2} = K \left[(2K \sin\left(\frac{\theta_1 - \theta_2}{z}\right) - l) \right] \left[-K \cos\left(\frac{\theta_1 - \theta_2}{z}\right) \right] \\ + m_2 g R \cos \theta_2 = 0$$

1)

2.)

$$m_1 + m_2 = M$$



$$r = r_1 - r_2$$

$$R = \frac{m_1 r_1 + m_2 r_2}{M}$$

$$\vec{r}_1' = \vec{r}_1 - \vec{R} = -\frac{m_2}{M} \vec{r}$$

$$\vec{r}_2' = \vec{r}_2 - \vec{R} = -\frac{m_1}{M} \vec{r}$$

$$\vec{v}_1' = -\frac{m_2}{M} \dot{\vec{r}}$$

$$\vec{v}_2' = \frac{m_1}{M} \dot{\vec{r}}$$

$$L = \frac{1}{2} (m_1 + m_2) \dot{R}^2 + \frac{1}{2} \mu \dot{r}^2 + \underbrace{6 \frac{m_1 m_2}{r}}_{V(r)}$$

T

Ya que $\dot{R} = \text{cte}$ ya que R es coordenada cíclica podemos tomar el cm como origen del sistema $\rightarrow R = 0$

$$L = \frac{1}{2} \mu \dot{r}_2^2 + \frac{6 m_1 m_2}{r}$$

La energía se conserva ya que + es coordenada cíclica.

$$E_0 = E$$

$$-\frac{6m_1m_2}{r_0} = \frac{1}{2}\mu\dot{r}^2 - \frac{6m_1m_2}{r}$$

$$\frac{1}{2}\mu\dot{r}^2 = 6m_1m_2 \left(\frac{1}{r} - \frac{1}{r_0} \right)$$

$$\dot{r} = \sqrt{26m \left(\frac{1}{r} - \frac{1}{r_0} \right)}$$

m

$$v_i = m_2 \sqrt{26m \left(\frac{1}{r} - \frac{1}{r_0} \right)}$$

$$v_i = m_1 \sqrt{26m \left(\frac{1}{r} - \frac{1}{r_0} \right)}$$

$$V(r) = mgR_{cm} + \frac{1}{2}k(r-l)^2$$

2.)

$$F = -\frac{\partial V}{\partial r}$$

$$\dot{L} = \frac{1}{2}\mu(\dot{r}^2 + r^2\dot{\theta}^2) - V(r)$$

θ es cíclico

$$L^2 = \mu^2 r^4 \dot{\theta}^2$$

$$\frac{dL}{d\theta} = \mu r^2 \dot{\theta} = L = \text{cte}$$

$$V(r) = \int_0^r \left(\frac{K}{r^2} + \frac{L}{r^3} \right) dr$$

$$V(r) = -\frac{K}{r} - \frac{L}{2r^2}$$

Si E se conserva $E = T + V = \text{cte}$

$$\frac{1}{2}\mu\dot{r}^2 + \frac{L^2}{2\mu r^2} - \frac{K}{r} - \frac{L}{2r^2} = \text{cte}$$

V_{ef}

2do parte:

$$1. V(r) = -\frac{k}{r} + \frac{\epsilon}{r^2} = -k u + \epsilon u^2$$

$$\frac{dV}{du} = -k + 2\epsilon u \quad \varphi = u \text{ modo reducido}$$

sustituyéndolo en la ecuación de binet

$$\frac{du}{d\theta^2} + u = -\frac{\varphi}{l^2} (-k + 2\epsilon u)$$

$$\frac{d^2u}{d\theta^2} + \left(1 + \frac{2\epsilon l}{l^2}\right)u = \frac{\varphi k}{l^2}$$

la sol a la ecuación

$$u(\theta) = A \cos \left(\sqrt{\frac{1 + 2\epsilon \varphi}{l^2}} \theta + \phi \right) + \frac{k\varphi}{l^2 \left(1 + \frac{2\epsilon \varphi}{l^2}\right)}$$

el término φ indica que la órbita no es cerrada

el ángulo de precesión

$$\Delta\theta = 2\pi \left(\frac{1}{\alpha} - 1 \right) \quad (\text{cuando } \epsilon \text{ es pequeño})$$
$$d \approx 1 + \frac{\epsilon \ell}{l^2}$$

$$\Delta\theta \approx -\frac{2\pi\epsilon\ell}{l^2}$$

2. $F(r) = -\frac{\partial V}{\partial r} \rightarrow V = -\frac{K}{r} - \frac{\lambda}{er^2}$

$$V = -Ku - \frac{\lambda}{2} u^2$$

de la ec de Binet $\ell = M$ modo

$$\frac{\partial^2 u}{\partial \theta^2} + u = -\frac{\ell}{l^2} \frac{\partial V}{\partial u}$$

reducido

$$\frac{\partial^2 u}{\partial \theta^2} + u = \frac{\ell}{l^2} (K + \lambda u)$$

$$\frac{\partial^2 u}{\partial \theta^2} + \left(1 - \frac{\lambda \ell}{l^2} \right) u = \underbrace{\frac{\ell}{l^2} K}_{b}$$

$$U(\theta) = \frac{b}{\lambda} + C_2 \sin(\sqrt{\lambda} \theta) + C_3 \cos(\sqrt{\lambda} \theta)$$

$$\frac{1}{r} = \frac{b}{\lambda} + C_2 \sin(\sqrt{\lambda} \theta) + C_3 \cos(\sqrt{\lambda} \theta)$$

elipse

Si $\lambda = \frac{l}{\mu^2}$ $\lambda = 0$ no hay precesion

Si $\lambda > \frac{l}{\mu^2}$ $\lambda < 0$ orbita es inestable

Si $\lambda < \frac{l}{\mu^2}$ $\lambda > 0$ elipse con precesion

$$3. \quad l = \frac{1}{2} \mu \dot{\theta}^2 + \frac{K}{r} - \frac{1}{2} \frac{l^2}{\mu r^2}$$

E y l se conservan

en el punto r_{\min} $\dot{r} = 0 \rightarrow E = V_{\text{ef}}$

y la E total en el infinito $= \frac{1}{2} m v_0^2$

Sabemos que $l = r \times p = b \mu v_0$

$$E_i = E_f$$

$$\frac{1}{2} \mu V_0^2 = -\frac{k}{r_{min}} - \frac{1}{2} \frac{(UV_0b)^2}{\mu r_{min}^2}$$

$$\frac{1}{2} \mu V_0^2 r_{min}^2 = -kr_{min} - \frac{1}{2} \mu V_0^2 b^2$$

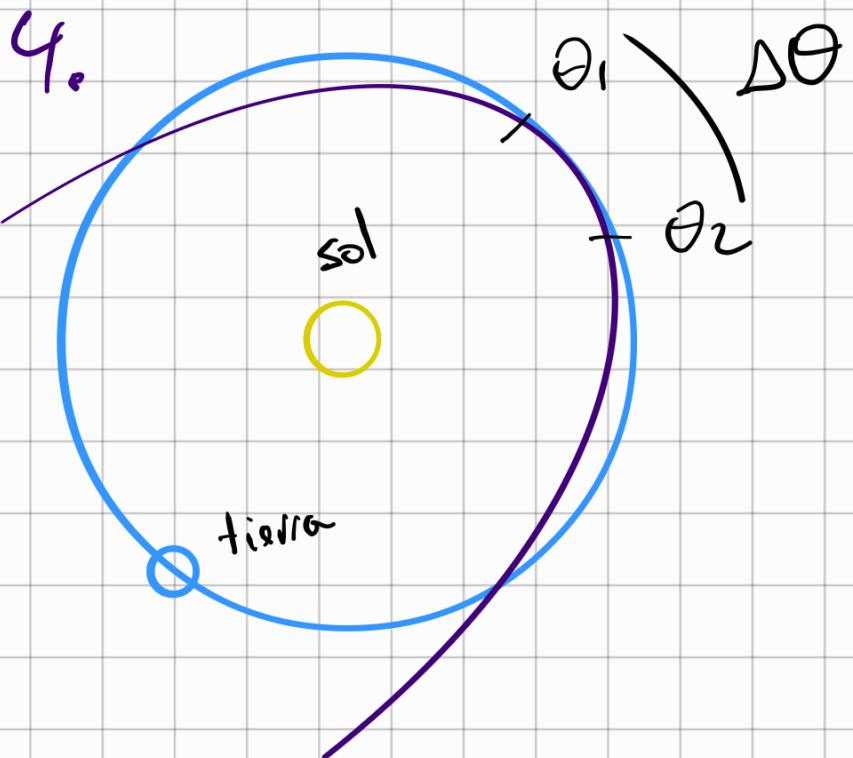
$$\mu V_0^2 r_{min}^2 + 2kr_{min} + \mu V_0^2 b^2 = 0$$

$$r_{min} = \frac{-k \pm \sqrt{k^2 + \mu^2 V_0^4 b^2}}{\mu V_0^2}$$

el angulo de de dispersion χ

$$\chi = \pi - 2 \arcsin \left(\frac{1}{e} \right)$$

$$\chi = \pi - 2 \arcsin \left(\frac{1}{\sqrt{1 + \left(\frac{V_0 b}{k} \right)^2}} \right)$$



igualamos las trayectorias

trayectoria para el cometa

$$\frac{q_c}{r} = 1 + e_c \cos \theta$$

$$r = q_c + q_c \cos \theta$$

para una parábola la excentricidad es 1

trayectoria de la tierra

$$\frac{q_t}{r} = 1 \quad \text{la excentricidad es 0}$$

$$r = \frac{1}{q_t}$$

$$q_c + q_c \cos\theta = \frac{1}{q_f} \rightarrow \cos\theta = \frac{1}{q_f} - 1$$

$$q = \frac{l^2}{\mu K}$$

para el cometa $l \approx m$
para la tierra $m = M$

$$\theta = \arccos \left[\frac{1}{\left(\frac{l_c}{m+K} \right) \left(\frac{l^2}{mK} \right)} \right] - 1$$

$$\theta = \arccos \left[\frac{m+MK^2}{l^2 l_c^2} \right]$$

el área que barre en la órbita terrestre es:

$$A = \frac{1}{2} r_+^2 \Delta\theta$$

r_+ = radio de la órbita de la tierra

$$A = \frac{l_c}{2M} \Delta t$$

$$\rightarrow \Delta t = \frac{Mr_+^2}{l_c} \Delta\theta$$

$$5. d) r = a(1 + \cos\theta)$$

$$\bar{F} = -\frac{\partial V}{\partial r}$$

entonces de la ec de Binet

$$F(r) = -\frac{v^2}{mr^2} \left(\frac{\partial^2 u}{\partial \theta^2} + u \right)$$

$$u = \frac{1}{r} = \frac{1}{a(1 + \cos\theta)}$$

$$\frac{\partial^2 u}{\partial \theta^2} = \frac{\cos\theta + \cos^2\theta + 2\sin^2\theta}{a(1 + \cos\theta)^3}$$

$$= \frac{\cos\theta + \cos^2\theta + 2 - 2\cos^2\theta}{a(1 + \cos\theta)^3}$$

$$= \frac{\cos\theta - \cos^2\theta + 2}{a(1 + \cos\theta)^3}$$

$$F(r) = \left(\frac{\cos\theta - \cos^2\theta + 2}{a(1+\cos\theta)^3} + \frac{1}{a(1+\cos\theta)} \right)$$

$\bullet -\frac{l^2}{mr^2}$

$$\cos^2\theta + 2\cos\theta + 1$$



$$= \left(\frac{2 + \cos\theta - \cos^2\theta + (1 + \cos\theta)^2}{a(1+\cos\theta)^3} \right) \cdot -\frac{l^2}{mr^2}$$

$$= \left(\frac{3 + 3\cos\theta}{a(1+\cos\theta)^3} \right) \cdot -\frac{l^2}{mr^2} = \frac{3(1+\cos\theta)}{a(1+\cos\theta)^3}$$

$$\bullet \frac{l^2}{mr^2} = -\frac{l^2}{mr^2} \left(\frac{3}{r^2} \right) = -\frac{3l^2}{mr^4}$$

b.) el área de la órbita

$$A = \int_0^{2\pi} \frac{1}{2} a^2 (1+\cos\theta)^2 d\theta = \frac{3\pi a^2}{2}$$

lo lasso de barriob $\dot{A} = \frac{l}{2m}$

el periodo $T = \frac{A}{\dot{A}} = \frac{3\pi a^2 m}{\ell}$

$$V(r) = - \int F(r) dr = - \int \frac{3\ell^3}{mr^4} = - \frac{\ell^2}{mr^3}$$

para que la partícula escape $E = 0$

$$V_{ef} = - \frac{\ell^2}{mr^3} + \frac{\ell^2}{mr^2}$$

