ECE 371 Materials and Devices

09/26/19 - Lecture 10 Effective Mass, Density of States

General Information

Homework #3 assigned, due Tuesday 10/1

Midterms returned by next Thursday

Midterm solutions posted

Reading for next time: 3.5, 4.1

Effective Mass Concept

- Movement of an electron in a crystal lattice is different from that in free space
 - Influence of positively charged ions
 - Influence of other electrons

$$F_{total} = F_{ext} + F_{int}$$

- Since we don't know F_{int} very well, we fold it into an "effective mass" (m*)
- Effective mass is given in units of m₀
- Can be higher or lower than m₀
- Can be negative! (hole)
- Allows us to treat electrons in a crystal as classical particles using Newtonian mechanics
- Effective mass influences measurable device properties (current transport in transistors, efficiency of LEDs and solar cells)

The Parabolic Approximation

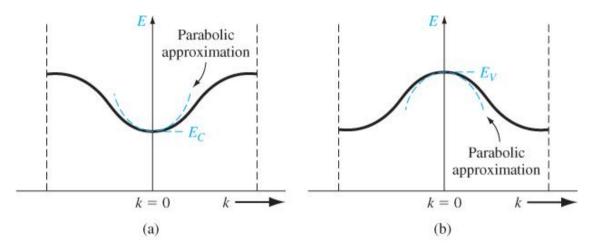


Figure 3.16 \mid (a) The conduction band in reduced k space, and the parabolic approximation. (b) The valence band in reduced k space, and the parabolic approximation.

- A free electron has a parabolic dispersion curve (E vs. k)
- The dispersion curves for electrons in crystals are roughly parabolic near the bottom (top) of the conduction (valence) bands

Concept of Hole Motion

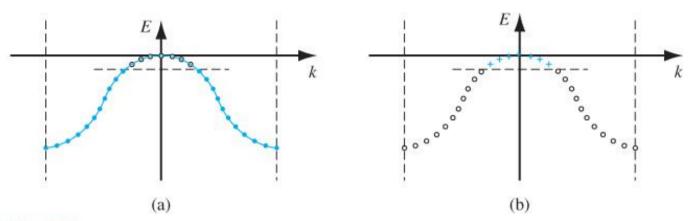


Figure 3.18 | (a) Valence band with conventional electron-filled states and empty states. (b) Concept of positive charges occupying the original empty states.

Empty electron states → Positive charges occupying empty states

*see in-class derivation

Downward parabola starting @ Ev 1 valence

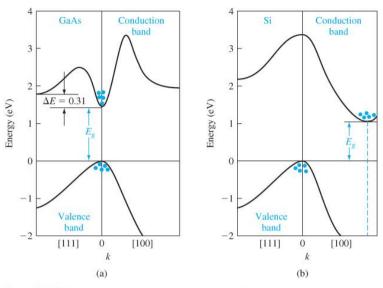
$$\frac{1}{m^*} = \frac{1}{k^2} \frac{dE}{dk^2} = \frac{-2C_2}{k^2} \Rightarrow C_2 \text{ is positive}$$

$$m^* \text{ is negative}$$

* electron moving near the top of a band has negative effective mass

(instead of using negative, use a positive effectiveness -> hole mp*

More on Effective Mass



$$m^* = \frac{\hbar^2}{\mathrm{d}^2 E / \mathrm{d}k^2}$$

Density of states effective mass

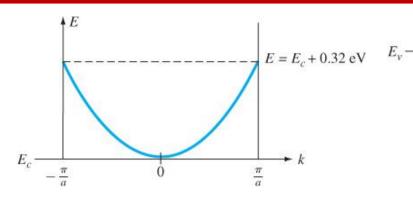
1
X

Material	m _n */m ₀	m _p */m ₀
Si	1.08	0.56
GaAs	0.067	0.48
GaN	0.2	1.2

Figure 3.25 | Energy-band structures of (a) GaAs and (b) Si. (From Sze [12].)

- m* depends upon energy (k) and direction of electron movement in the crystal
- Also, there are several different effective masses in the valence band (heavy hole (HH), light hole (LH), split-off (SO))
- Use "density of state effective mass" in this class (see Appendix F and Table B4)
- Density of states effective mass is an "average" effective mass over all directions

Test Your Understanding



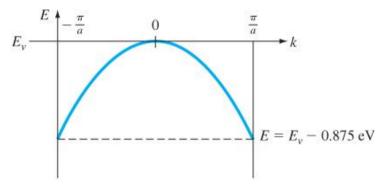


Figure 3.22 | Figure for Exercise TYU 3.3.

Figure 3.23 | Figure for Exercise TYU 3.4.

- **TYU 3.3** A simplified *E* versus *k* curve for an electron in the conduction band is given. The value of *a* is 10 Å. Determine the relative effective mass m^*/m_0 . $(SLI : I = {}^0 w/_* w : SUV)$
- **TYU 3.4** A simplified *E* versus *k* curve for a hole in the valence band is given. Assume a value of a = 12 Å. Determine the relative effective mass $|m^*/m_0|$. $(\$867.0 = |^0 w/_* w| \$ u)$

TYU 3.3 remember:
$$M^* = \frac{h^2}{\frac{d^2E}{dk^2}}$$

need $E(k) = Ck^2 + E_C$ if dealing with valence instead,

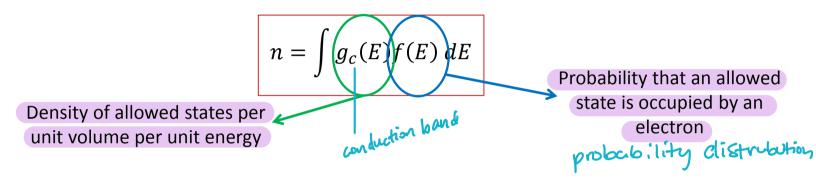
 $E(k) = E_V - ck^2$
 $E(\frac{T}{a}) = C(\frac{T}{a})^2 + E_C = \frac{T}{c} + 0.32eV$
 $C = (0.32eV) \frac{a^2}{TI^2}$
 $\frac{d^2E}{dk^2} = 2C = 2(0.32eV) \frac{a^2}{TI^2}$
 $M^* = \frac{h^2}{2(0.32eV) \frac{a^2}{TI^2}}$ ais given

 $= 1.07 \times 10^{-30} \text{ kg}$
 $M_0 = \frac{M^*}{9.11 \times 10^{-3l}} = 1.175$

written as 1.175 mo

Carrier Density and Density of States

 Goal: determine the electron (n) and hole (p) densities (#/cm³) in the semiconductor



- The density of states (DOS) describes the density of allowed states per unit volume per unit energy (g(E))
- Electron and holes are called "carriers"
- First we will calculate the DOS portion of the carrier density

Electron density

$$n = \int g_c(E) f(E) dE$$
density probability that
$$\left(\frac{\#}{\text{energ.vol}}\right) \text{ an allowed state}$$
is occupied by electron

count up the # of K states

$$V(x,y,z) = 0$$
 inside
 $V(x,y,z) = \infty$ outside
(infinite potent. box)

$$\Psi(x_1y_1,z) = A \sin\left(\frac{n_x\pi}{a}x\right) \sin\left(\frac{n_y\pi}{a}y\right) \sin\left(\frac{n_z\pi}{a}z\right)$$

$$E_{n} = \frac{K^{2}\pi^{2}}{2ma^{2}} \left(n_{x}^{2} + n_{y}^{2} + n_{z}^{2}\right)$$

(111) is the lowest possible ground state in 30

$$K_{x} = \frac{n_{x}\pi}{a}, K_{y} = \frac{n_{y}\pi}{a}, K_{z} = \frac{n_{z}\pi}{a}$$

$$K^{2} = \frac{\pi}{a} \left(n_{x}^{2} + n_{y}^{2} + n_{z}^{2}\right)$$

$$\frac{1}{2mE}$$

$$\frac{1}{4^{2}}$$

What is the area of one state?

$$K_{x+1} - K_x \rightarrow (n_x+1) \frac{\pi}{a} - n_x \frac{\pi}{a} = \frac{\pi}{a}$$

separation between states

Then area is $(\frac{\pi}{a})^2$

Volume of single state in K-space

Then D.O.S.
$$\Rightarrow$$
 $g(E) = \frac{1}{V} \frac{dNs}{dE}$

[#States

Vol • energy]

Volume (a)

Density of States (DOS)

- To determine the 3D DOS in k-space:
- 1. Calculate the number of states $(N_s(k))$ as a function of k contained within the sphere bounded by $k^2 = k_x^2 + k_y^2 + k_z^2$ by taking the volume of the sphere in k-space $(V_{sphere} = (4/3)\pi k^3)$ divided by unit volume of one state in k-space $(V_{unit-k} = (\pi/a)^3)$
- 2. Multiply by 2 for spin degeneracy and divide by 8 to keep only positive k-values
- 3. Using the parabolic approximation for a free electron, convert $N_s(k)$ to $N_s(E)$
- 4. Divide by a unit volume in real space $V_{unit} = a^3$
- 5. Differentiate with respect to E, dN_s/E , to get the DOS

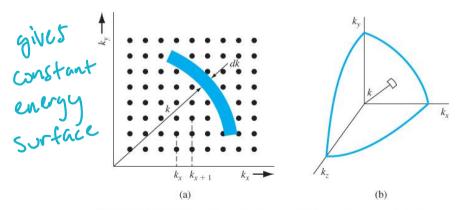


Figure 3.26 | (a) A two-dimensional array of allowed quantum states in k space. (b) The positive one-eighth of the spherical k space.

$$DOS = \frac{1}{V_{unit}} \frac{dN_s}{dE}$$

 V_{unit} = unit volume N_s = # of states E = energy

(1)
$$\frac{V_{sphere}}{V_{k}} = \frac{\frac{4}{3}\pi k^{3}}{(T_{a})^{3}}$$

(2)
$$\pm$$
 States in $\frac{1}{8}$ Sphere: $N_s(k) = \frac{a^3k^3}{3\pi^2}$

(3)
$$N_s(E) = \frac{a^3 (2mE)^{\frac{3}{2}}}{3\pi^2 h^3}$$

(4)
$$\frac{N_s(E)}{V} = \frac{(2mE)^{3/2}}{3\pi^2 k^3}$$

$$(5) \quad g(E) = \frac{1}{V} \frac{dNs}{dE}$$

$$= \frac{4\pi (2m)^{3/2}}{h^3} \sqrt{E}$$

Density is proportional to TE

Increase energy -> increase states, but only to a certain point

Translating

 $E \rightarrow E - Ec$ for conduction $E_{V} - E$ for valence

m -> mn for conduction

mp* for valence

no States in bandgap holes typically larger eff. mass than electrons

Semiconductor Density of States (3D)

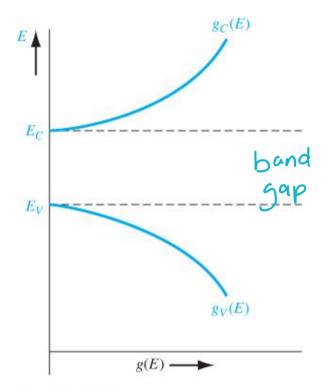


Figure 3.27 | The density of energy states in the conduction band and the density of energy states in the valence band as a function of energy.

Conduction band DOS:

$$g_c(E) = \frac{4\pi}{h^3} [2m_n^*]^{3/2} (E - E_c)^{1/2}$$

Valence band DOS:

$$g_v(E) = \frac{4\pi}{h^3} [2m_p^*]^{3/2} (E_v - E)^{1/2}$$

change E & m to m*

- Parabolic approximation
- Fewer states at lower energies
- No states in the forbidden gap
- In general, $g_c(E)$ and $g_v(E)$ are different
- 3D DOS is also called "bulk" DOS

Density of States (Lower Dimensions)

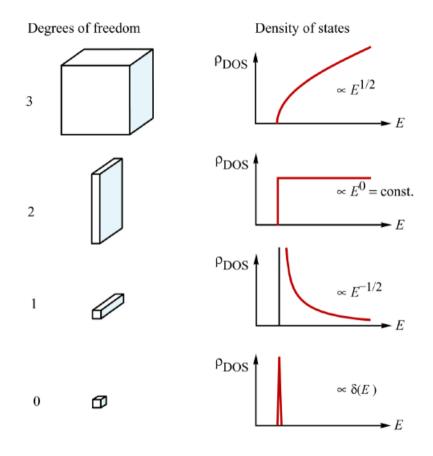


Fig. 12.7. Electronic density of states of semiconductors with 3, 2, 1, and 0 degrees of freedom for electron propagation. Systems with 2, 1, and 0 degrees of freedom are referred to as quantum wells, quantum wires, and quantum boxes, respectively.

Quantum Well Density of States

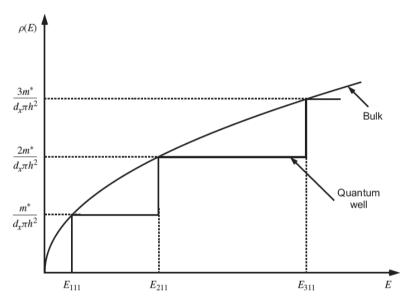


FIGURE A1.13: Density of states for an infinite-barrier quantum well and bulk material. If the barrier is not infinite, the quantum-well energies decrease slightly. If desired, the density of state plateaus can be decreased by using an effective $d_x = d_x^*$ (a different one for each state) so that the extrema continue to intersect the bulk characteristic.

QW DOS:

$$g(E) = \frac{1}{d_x} \sum_{n_x} \frac{m^*}{\pi \hbar^2} \mathcal{H}(E - E_{n_x})$$