

ECE 371

Materials and Devices

09/03/19 - Lecture 4

Schrodinger's Equation: Meaning of the
Wave Function and Free Electron
Solution

General Information

- Homework #1 due before class today
- Homework #2 assigned and due Thursday September 12th
- Link to download Crystal Viewer (<http://crystalmaker.com/crystalviewer/index.html>) added to the website in Articles, Videos, and Additional Notes folder. See next few slides for examples of software capabilities.
- Reading for next time: 2.3.2-2.3.3

Wave Function

- The temporal and spatial evolution of a particle (e.g., electron) with one degree of freedom is given by $\psi(x, t)$
- $\psi(x, t)$ can be complex
- $\psi(x, t) \cdot \psi^*(x, t)$ is related to the probability of finding the particle within the interval $x + dx$
- The shape of the wave function is influenced by the potential energy landscape

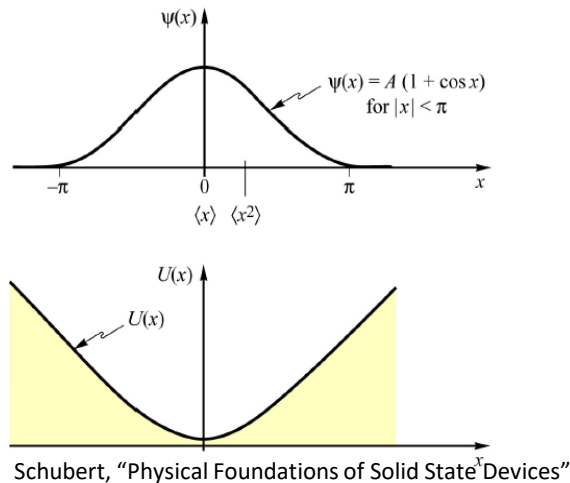


Fig. 2.1. Example for a one-dimensional wave function $\psi(x)$. Also shown is a corresponding potential function, $U(x)$. This potential function provides a driving force towards $x = 0$, that is towards minimum energy.

Schrodinger's Wave Equation

- The Schrodinger equation (SE) describes the spatial and temporal evolution of the wave function for a given potential energy landscape and set of boundary conditions

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x, t)}{\partial x^2} + V(x) \psi(x, t) = j\hbar \frac{\partial \psi(x, t)}{\partial t}$$

related to
kinetic energy

related to
potential energy

related to
total energy

- $V(x)$ is the potential energy, m is the mass of the particle, and $j = \sqrt{-1}$
↑ for crystals — use “effective mass”
- SE is a basic postulate of quantum mechanics but can be derived
- SE can be used to describe the behavior of electrons in a crystal

Time-Dependent and Time-Independent Parts of SE

- The separation of variables technique can be used to deconstruct the SE into time-dependent and time-independent parts
- We assume that the wave function can be represented as the product of a time-independent function and a time-dependent function (i.e., $\psi(x, t) = \psi(x)\phi(t)$)

Time-dependent solution
can be obtained quickly:

$$\phi(t) = e^{-j\left(\frac{E}{\hbar}\right)t} = e^{-j\omega t}$$

Sinusoidal variation with time

E is the total energy of the particle

$$E = \hbar \omega$$

Focus on
Solutions to
potential
energy (i.e.
electron in
well)

Time-independent part
of the equation:

$$\frac{d^2\psi(x)}{dx^2} + \frac{2m}{\hbar^2} (E - V(x))\psi(x) = 0$$

↑
potential

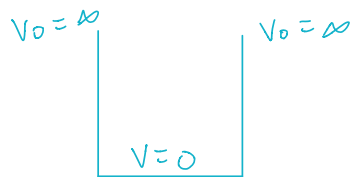
Nature of the solution for $\psi(x)$ depends upon the potential $V(x)$ and the boundary conditions

Types of solutions will be $\psi(x) = \text{sins \& cosines}$ *see in-class derivation

Solve 3 different potentials:

(1) Free electron (zero potential ($V_x = 0$))
(can treat crystals as almost free electron)

(2) Infinite potential well
($V = 0$ inside ∞ outside)



(3) Potential Step

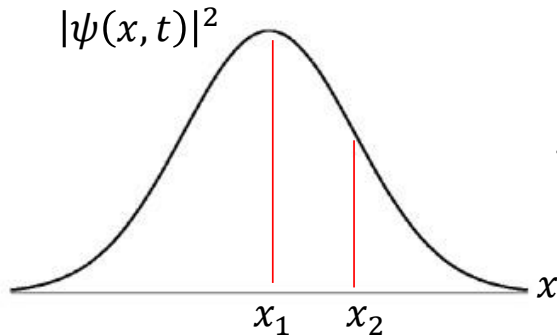


Physical Meaning of the Wave Function

- Use $\psi(x, t)$ to describe the behavior of electrons in crystals
- Wave function is complex so cannot itself represent a physical quantity
- Instead we define a probability density function that describes the probability of finding the electron between x and $x + dx$

$$|\psi(x, t)|^2 = \psi(x)\psi^*(x) = |\psi(x)|^2$$

- Probability density function of a variable describes the relative likelihood of the variable to take on a given value



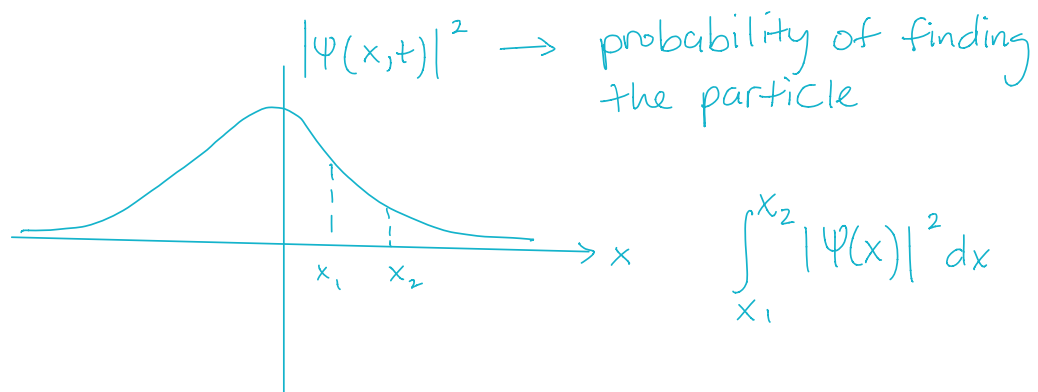
Probability that the electron is between x_1 and x_2 is

$$\int_{x_1}^{x_2} |\psi(x)|^2 dx$$

The electron must be somewhere so we can also normalize the wave function

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1$$

$$\Psi(x,t) = \Psi(x) \phi(t) = \Psi(x) e^{-j\omega t}$$



given potential \rightarrow solve for Ψ and boundary conditions

Electron in Free Space

When $V(x) = 0$ we have a free electron and the TISE becomes:

$$\frac{d^2\psi(x)}{dx^2} + \frac{2mE}{\hbar^2}\psi(x) = 0$$

completely kinetic

Time independent Schrod. eq.

The general solution is two travelling waves:

$$\psi(x, t) = A \exp[j(kx - \omega t)] + B \exp[-j(kx + \omega t)]$$

wave number:

$$k = \sqrt{\frac{2mE}{\hbar^2}}$$

For a travelling wave in the +x direction we have:

$$\psi(x, t) = A \exp[j(kx - \omega t)]$$

particle only has kinetic energy

- The probability density function is a constant (AA^*)
- Particle can be found anywhere since the momentum is well defined
- Note that the plane wave solution cannot be normalized (but a superposition of plane waves can be – wave packet)

$$\frac{d^2\psi(x)}{dx^2} + \frac{2mE}{\hbar^2}\psi(x) = 0$$

linear
2nd order
Homogeneous

$$ay'' + by' + cy = 0$$

↓

$$r^2 + r + 0 = 0$$

↓

$$r^2 + \frac{2mE}{\hbar^2}r = 0$$

$$r = \pm \sqrt{\frac{-2mE}{\hbar^2}} = \pm j \sqrt{\frac{2mE}{\hbar^2}}$$

Soln has form:

$$y = C_1 e^{r_1 x} + C_2 e^{r_2 x}$$

↓

$$\psi(x) = A e^{j \sqrt{\frac{2mE}{\hbar^2}} x} + B e^{-j \sqrt{\frac{2mE}{\hbar^2}} x}$$

↑ wave #
K (inversely proportional to wave length)

$$\psi(x) = A e^{jkx} + B e^{-jkx}$$

Adding back the time component:

$$\phi(t) = e^{-j\omega t}$$

$$\psi(x, t) = A \exp[j(kx - \omega t)] + B \exp[-j(kx + \omega t)]$$

$$\psi(x, t) = A e^{j(kx - \omega t)} + B e^{-j(kx + \omega t)}$$

What if:

Assume moving in pos x direction:

$$B = 0$$

$$\psi(x, t) = A e^{j(kx - \omega t)}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}} = \frac{p}{\hbar} = \frac{2\pi}{\lambda}$$

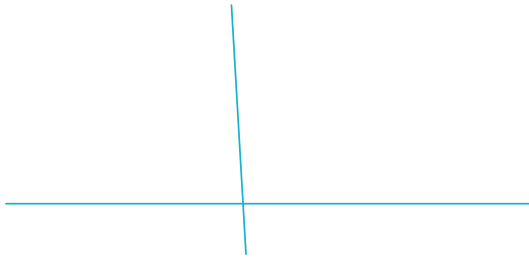
$$p = \hbar k$$

$$E = \hbar \omega = \hbar \nu$$

↓
nu?

Probability Density

$$\begin{aligned}\Psi(x,t) \cdot \Psi^*(x,t) &= A e^{j(kx - \omega t)} \cdot A e^{-j(kx - \omega t)} \\ &= A \cdot A^* \rightarrow \text{constant}\end{aligned}$$



Potential Wells

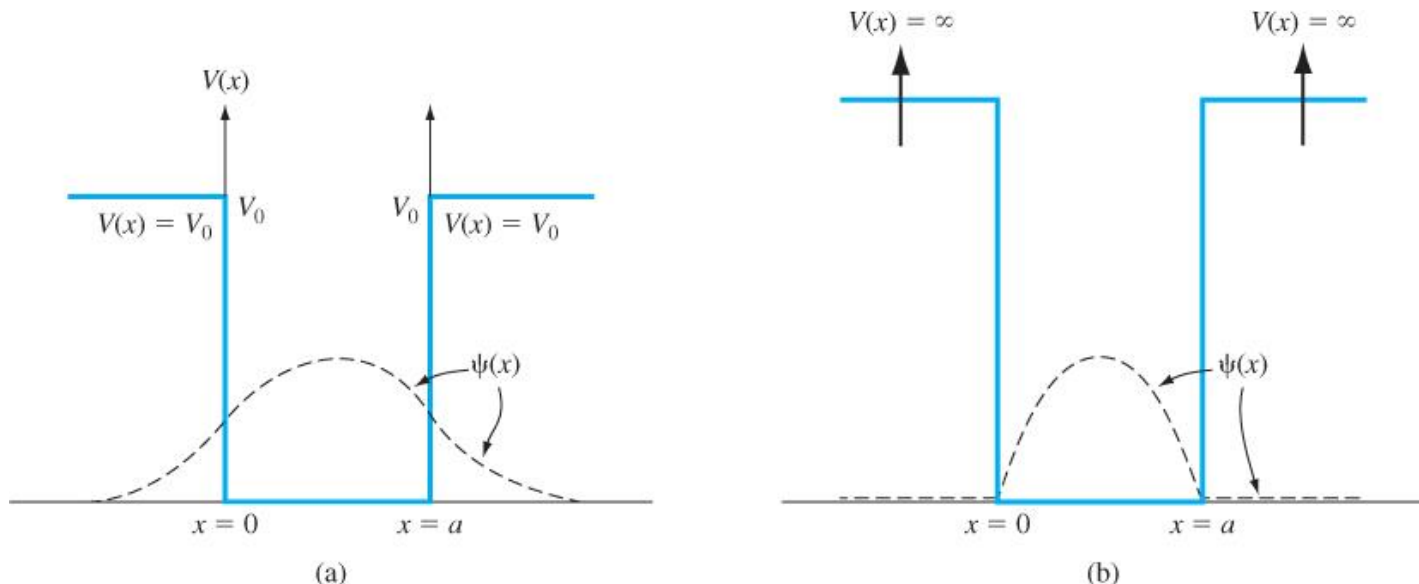


Figure 2.5 | Potential functions and corresponding wave function solutions for the case (a) when the potential function is finite everywhere and (b) when the potential function is infinite in some regions.

- Finite and infinite potential wells (bound particles)
- Potential height determines strength of confinement
- Solutions for $\psi(x)$ are sine and cosine functions for the infinite potential well