

1 **Relocating a Cluster of Earthquakes Using**
2 **a Single Seismic Station**

3 David J. Robinson, Malcolm Sambridge, Roel Snieder, Juerg Hauser

David Robinson: Earth Monitoring and Hazard Group
4 Geoscience Australia
GPO Box 383
Canberra ACT 2601 Australia
E-mail: david.robinson@ga.gov.au

Abstract

Coda waves arise from scattering to form the later arriving components of a seismogram. Coda wave interferometry is an emerging tool for constraining earthquake source properties from the interference pattern of coda waves between nearby events. A new earthquake location algorithm is derived which relies on coda wave based probabilistic estimates of earthquake separation. The algorithm can be used with coda waves alone or in tandem with travel time data. Synthetic examples in 2D and 3D and real earthquakes on the Calaveras Fault, California are used to demonstrate the potential of coda waves for locating poorly recorded earthquakes. It is demonstrated that coda wave interferometry: (a) outperforms traditional earthquake location techniques when the number of stations is small; (b) is self-consistent across a broad range of station situations; and (c) can be used with a single station to locate earthquakes.

Introduction

Accurate earthquake location is important for many applications. Locations are required for: magnitude determination (*Richter, 1935; Gutenberg, 1945*); computing moment tensors (*Sipkin, 2002*); seismological studies of the Earth's interior (*Spencer and Gubbins, 1980; Kennett et al., 1995; Curtis and Snieder, 2002; Kennett et al., 2004*); understanding strong motion and seismic attenuation (*Toro et al., 1997; Campbell, 2003*) and modeling earthquake hazard or risk (*Frankel et al., 2000; Stirling et al., 2002; Robinson et al., 2006*). The accuracy required in earthquake location depends on the application. For example, imaging the structure of

27 a fracture system from microseismicity requires greater detail than determining
28 whether a $M_w = 7.5$ earthquake occurs offshore for tsunami warning. This pa-
29 per focuses on reducing location uncertainty for a cluster of events when they are
30 recorded by a small number of stations.

31 Absolute location describes the location of an earthquake with respect to a
32 global reference such as latitude, longitude (or easting/northing) and depth. Un-
33 certainties associated with absolute locations are influenced by source to station
34 distances, the number of stations and their geometry, signal-to-noise, clarity of
35 onsets and accuracy of the velocity model used in computing travel times. Uncer-
36 tainties in absolute location are typically of the order of several kilometers because
37 they are susceptible to uncertainty in the velocity structure along the entire path
38 between the source and receiver. For example, *Shearer* (1999) states that location
39 uncertainty in the ISC (International Seismological Centre) and PDE (National
40 Earthquake Information Center) catalogues are generally around 25 km horizon-
41 tally and at least 25 km in depth (Here the depth uncertainties of 25 km assume the
42 use of depth dependent phases such as pP . Without such phases the uncertainty
43 is higher). *Bondár et al.* (2004) demonstrate that at the local scale, absolute loca-
44 tions are accurate to within 5 km with a 95% confidence level when local networks
45 meet a number of station related criteria. Such errors are too large for many appli-
46 cations, particularly those focussed on imaging rupture surfaces from aftershock
47 sequences.

48 Relative earthquake location involves locating a group of earthquakes with
49 respect to one another and was first introduced by *Douglas* (1967) who devel-
50 oped the technique commonly known as joint hypocenter determination (*Douglas*
51 (1967) originally used the term joint epicentre determination. However, he was

52 solving for hypocentre). In principle, relative locations can be computed by differ-
53 encing absolute locations. However, *Pavlis* (1992) shows that inadequate knowl-
54 edge of velocity structure leads to systematic biases when relative positions are
55 computed in this way. To reduce errors from unknown velocity structure, rela-
56 tive location techniques compute locations directly from travel time differences
57 between two waveforms (*Ito*, 1985; *Got et al.*, 1994; *Nadeau and McEvilly*, 1997;
58 *Waldhauser et al.*, 1999). By doing so, they remove errors associated with ve-
59 locity variations outside the local region, because such variations influence all
60 waveforms in the same manner (*Shearer*, 1999).

61 Reported location uncertainties from relative techniques are around 15 to 75 m
62 in local settings with good station coverage (*Ito*, 1985; *Got et al.*, 1994; *Wald-*
63 *hauser et al.*, 1999; *Waldhauser and Schaff*, 2008). Here, ‘good coverage’ implies
64 multiple stations distributed across a broad range of azimuthal directions. Rela-
65 tive location techniques have been used to image active fault planes (*Deichmann*
66 *and Garcia-Fernandez*, 1992; *Got et al.*, 1994; *Waldhauser et al.*, 1999; *Wald-*
67 *hauser and Ellsworth*, 2002; *Shearer et al.*, 2005); study rupture mechanics (*Ru-*
68 *bin et al.*, 1999; *Rubin*, 2002); interpret magma movement in volcanoes (*Frèmont*
69 *and Malone*, 1987); and monitor pumping-induced seismicity (*Lees*, 1998; *Ake*
70 *et al.*, 2005).

71 In traditional approaches to absolute and relative location only early onset
72 body waves, typically *P* and/or *S* waves, are used. The data utilised may be the
73 direct arrival times; travel time difference computed between picked arrivals of
74 two waveforms; or time differences inferred from time-lagged cross correlation
75 of relatively small windows around the body wave arrivals. In all three cases,
76 the majority of the waveform is discarded. Furthermore, obtaining high accuracy

77 with these techniques requires multiple stations with good azimuthal coverage.
78 In this paper we demonstrate that it is possible to significantly reduce location
79 uncertainty when few stations are available by using more of the waveform.

80 Coda refers to later arriving waves in the seismogram that arise from scattering
81 (*Aki, 1969; Snieder, 1999, 2006*). Coda waves are ignored in most seismological
82 applications due to the complexity involved in constraining complex heterogeneous
83 velocity models in real settings. In this paper we develop an approach for locating
84 earthquakes using coda waves. *Snieder and Vrijlandt (2005)* demonstrate that the
85 coda of two earthquakes can be used to estimate the separation between them.
86 Their technique, known as coda wave interferometry (CWI), is based on the inter-
87 ference pattern between the coda waves. Unlike travel time based location tech-
88 niques, CWI does not require multiple stations or good azimuthal coverage. In
89 fact, it is possible to obtain estimates of separation using a single station (*Robin-*
90 *son et al., 2007a*). This makes CWI particularly interesting for regions where
91 station density is low such as intraplate settings. In this paper we demonstrate
92 how CWI separation estimates can be used to constrain location with data from a
93 single station. Our technique can be used on coda waves alone or in combination
94 with travel times. We begin by introducing the theory of CWI based earthquake
95 location. This is followed by a demonstration of capability using synthetic exam-
96 ples and application to earthquakes on the Calaveras fault, California using CWI
97 alone and CWI in combination with travel time constraints.

98 Theory

99 *Snieder and Vrijlandt (2005)* introduce a CWI based estimator of source sep-

100 aration δ_{CWI} between two earthquakes

$$101 \quad \delta_{CWI}^2 = g(\alpha, \beta) \sigma_\tau^2, \quad (1)$$

102 where σ_τ is the standard deviation of the travel time perturbation between the coda
 103 waves of two earthquakes, and α and β are the near-source P and S wave veloc-
 104 ities, respectively. The function g depends on the type of excitation (explosion,
 105 point force, double couple) and on the direction of source displacement relative
 106 to the point force or double couple. For example, for two double couple sources
 107 displaced in the fault plane,

$$108 \quad g(\alpha, \beta) = 7 \frac{\left(\frac{2}{\alpha^6} + \frac{3}{\beta^6}\right)}{\left(\frac{6}{\alpha^8} + \frac{7}{\beta^8}\right)}, \quad (2)$$

109 whereas, for two point sources in a 2D acoustic medium

$$110 \quad g(\alpha, \beta) = 2\alpha^2 \quad (3)$$

111 (*Snieder and Vrijlandt, 2005*). In this paper we use equation 2 which assumes that
 112 the source mechanism of both events are identical, an assumption likely to be true
 113 for events in the same fault plane. *Robinson et al. (2007b)* explore the impact of a
 114 change in mechanism.

115 The σ_τ in equation (1) is related to the maximum of the cross correlation be-
 116 tween the coda of the two waveforms, R_{max} , and hence can be computed di-
 117 rectly from the recorded data. The original formulation by *Snieder and Vrijlandt*
 118 (2005) used a second-order Taylor series expansion of the waveform autocorrela-
 119 tion function to relate σ_τ and R_{max} by

$$120 \quad R_{max}^{(t, t_w)} = 1 - \frac{1}{2} \overline{\omega^2} \sigma_\tau^2, \quad (4)$$

121 In this paper we use the autocorrelation approach of *Robinson et al.* (2011) to
 122 relate the parameters directly and we apply a restricted time lag search when eval-
 123 uating R_{max} . These extensions to the original technique of *Snieder and Vrijlandt*
 124 (2005) increase the range of applicability of CWI by 50% (i.e. from 300 to 450 m
 125 separation for 1 to 5 Hz filtered coda waves).

126 *Robinson et al.* (2011) show that CWI leads to probabilistic constraints on
 127 source separation and introduce a Bayesian approach for describing the probab-
 128 ility of true separation given the CWI data. Their approach is summarised by

$$129 \quad P(\tilde{\delta}_t | \tilde{\delta}_{CWIN}) \propto P(\tilde{\delta}_{CWIN} | \tilde{\delta}_t) \times P(\tilde{\delta}_t) \quad (5)$$

130 where $P(\tilde{\delta}_t | \tilde{\delta}_{CWIN})$ is the posterior function indicating the probability of true sep-
 131 aration $\tilde{\delta}_t$ given the noisy CWI separation estimates $\tilde{\delta}_{CWIN}$; $P(\tilde{\delta}_{CWIN} | \tilde{\delta}_t)$ is the
 132 likelihood function (or forward model) giving the probability that the separation
 133 estimates $\tilde{\delta}_{CWIN}$ would be observed if the true separation was $\tilde{\delta}_t$; and $P(\tilde{\delta}_t)$ is the
 134 prior PDF accounting for all a-priori information. The tilde above the separation
 135 parameters in equation (5) indicates the use of a wavelength normalised separation
 136 parameter

$$137 \quad \tilde{\delta} = \frac{\delta}{\lambda_d}, \quad (6)$$

138 which measures separation ($\delta = \delta_{CWIN}$ or δ_t) with respect to dominant wavelength
 139 λ_d . In this paper we consider a uniform prior over appropriate bounds to ensure
 140 that the posterior function is dominated by the recorded data. The procedure for
 141 computing the likelihood $P(\tilde{\delta}_{CWIN} | \tilde{\delta}_t)$ is derived by *Robinson et al.* (2011) and
 142 summarised in Appendix . With these two pieces in place we can compute the
 143 posterior $P(\tilde{\delta}_t | \tilde{\delta}_{CWIN})$ (or PDF) for the separation between any pair of events
 144 directly from their coda waves.

145 We seek a probability density function (PDF) which links individual pairwise
 146 posteriors $P(\tilde{\delta}_t|\tilde{\delta}_{CWIN})$ to describe the location of multiple events whose maxi-
 147 mum corresponds to the most probable combination of locations. More impor-
 148 tantly however, the PDF shall quantify location uncertainty and provide informa-
 149 tion on the degree to which individual events are constrained by the data. For
 150 convenience, we begin with three earthquakes having locations \mathbf{e}_1 , \mathbf{e}_2 and \mathbf{e}_3 . Us-
 151 ing a Bayesian formulation we write

$$152 \quad P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\mathbf{d}) \propto P(\mathbf{d}|\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3) \times P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3), \quad (7)$$

153 where $P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\mathbf{d})$, $P(\mathbf{d}|\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ and $P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ are the posterior, likeli-
 154 hood and prior functions, respectively. In equation (7) \mathbf{d} represents observations
 155 that constrain the locations. They can be any combination of travel times, geode-
 156 tic information or CWI separations. For example, if coda waves are used we have
 157 $P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\tilde{\delta}_{CWIN})$ and $P(\tilde{\delta}_{CWIN}|\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$, where $\tilde{\delta}_{CWIN}$ are the wavelength
 158 normalised separation estimates. Alternatively, if we use CWI and travel time data
 159 we may write $P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\tilde{\delta}_{CWIN}, \Delta_{TT})$ and $P(\tilde{\delta}_{CWIN}, \Delta_{TT}|\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3)$ where
 160 Δ_{TT} represent travel time differences. In the following derivation and in Syn-
 161 thetic Experiments and Relocating Earthquakes on the Calaveras Fault we focus
 162 on the constraints imposed by coda waves, whereas in Combining Travel Time and
 163 CWI Constraints we demonstrate how CWI and travel time data can be combined.

164 For three earthquakes we have likelihoods; $P(\tilde{\delta}_{CWIN,12}|\mathbf{e}_1, \mathbf{e}_2)$, $P(\tilde{\delta}_{CWIN,13}|\mathbf{e}_1, \mathbf{e}_3)$
 165 and $P(\tilde{\delta}_{CWIN,23}|\mathbf{e}_2, \mathbf{e}_3)$. In writing these likelihoods we have replaced the condi-
 166 tional term on separation $\tilde{\delta}_t$ with the locations (e.g. \mathbf{e}_1 and \mathbf{e}_2). This can be done
 167 because knowledge of location translates to separation. Note that the reverse is
 168 not true. That is, knowledge of separation between a single event pair does not
 169 uniquely translate to location but rather places a non-unique constraint on location.

170 Furthermore, since the pairwise functions are independent the joint likelihood be-
 171 comes

$$172 \quad P(\tilde{\delta}_{CWIN}|\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3) = P(\tilde{\delta}_{CWIN,12}|\mathbf{e}_1, \mathbf{e}_2) \quad (8)$$

$$\times P(\tilde{\delta}_{CWIN,13}|\mathbf{e}_1, \mathbf{e}_3) \times P(\tilde{\delta}_{CWIN,23}|\mathbf{e}_2, \mathbf{e}_3).$$

173 Similarly, the earthquake locations are independent and the joint prior becomes

$$174 \quad P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3) = P(\mathbf{e}_1) \times P(\mathbf{e}_2) \times P(\mathbf{e}_3). \quad (9)$$

175 Combining equations (8) and (9) gives the joint posterior function

$$176 \quad P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\tilde{\delta}_{CWIN}) = c \prod_{i=1}^3 P(\mathbf{e}_i) \quad (10)$$

$$\times \prod_{i=1}^2 \prod_{j=i+1}^3 P(\tilde{\delta}_{CWIN,ij}|\mathbf{e}_i, \mathbf{e}_j)$$

177 for three events.

178 A detailed understanding of the location of a single event (e.g. \mathbf{e}_2) is obtained
 179 by computing the marginal

$$180 \quad P(\mathbf{e}_2|\delta_{CWIN}) = \int \int P(\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3|\tilde{\delta}_{CWIN}) d\mathbf{e}_1 d\mathbf{e}_3, \quad (11)$$

181 where the intergral is taken over all plausible locations for \mathbf{e}_1 and \mathbf{e}_3 . Alterna-
 182 tively, we can compute the marginal for a single event coordinate by integrating
 183 the posterior over all events and remaining coordinates for the chosen earthquake.
 184 Evaluation of the normalizing constant c in equation (10) involves finding the inte-
 185 gral of the posterior function over all plausible locations. In many applications the
 186 constant of proportionality c can be ignored. For example, it is not required when
 187 seeking the combination of locations which maximise the posterior function, nor
 188 in Bayesian sampling algorithms such as Markov-chain Monte-Carlo techniques
 189 which only require evaluation of a function proportional to the PDF.

190 Extending to n events we get the posterior function

$$\begin{aligned}
 P(\mathbf{e}_1, \dots, \mathbf{e}_n | \tilde{\delta}_{CWIN}) &= c \prod_{i=1}^n P(\mathbf{e}_i) \\
 &\times \prod_{i=1}^{n-1} \prod_{j=i+1}^n P(\tilde{\delta}_{CWIN,ij} | \mathbf{e}_i, \mathbf{e}_j).
 \end{aligned}
 \tag{12}$$

192 When evaluating equation (12) over a range of locations it is necessary to compute
 193 and multiply many numbers close to zero. This is because the PDFs tend to zero
 194 as the locations get less likely (i.e. near the boundaries of the plausible region).
 195 Such calculations are prone to truncation errors and so we work with the negative
 196 logarithm

$$L(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n) = -\ln \left[P(\mathbf{e}_1, \dots, \mathbf{e}_n | \tilde{\delta}_{CWIN}) \right] \tag{13}$$

198 OR

$$\begin{aligned}
 L(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n) &= -\ln [c] - \sum_{i=1}^n \ln [P(\mathbf{e}_i)] \\
 &- \sum_{i=1}^{n-1} \sum_{j=i+1}^n \ln \left[P(\tilde{\delta}_{CWIN,ij} | \mathbf{e}_i, \mathbf{e}_j) \right].
 \end{aligned}
 \tag{14}$$

200 The logarithm improves numerical stability by replacing products with summa-
 201 tions. The negative facilitates the use of optimisation algorithms that are designed
 202 to minimise an objective function.

203 The event locations $\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n$ are defined by coordinates \hat{x} , \hat{y} and \hat{z} where
 204 the hat indicates use of a local coordinate system. We choose a local coordinate
 205 system which removes ambiguity associated with transformations of the coordi-
 206 nate system. It is necessary to do this because the distances between events are
 207 invariant for rotations, reflections and translations of the seismicity pattern and
 208 hence cannot be resolved from CWI alone. In defining this coordinate system we

209 fix the first event at the origin

$$210 \quad \mathbf{e}_1 = (0, 0, 0), \quad (15)$$

211 the second event on the positive \hat{x} -axis

$$212 \quad \mathbf{e}_2 = (\hat{x}_2, 0, 0), \hat{x}_2 > 0 \quad (16)$$

213 the third on the $\hat{x} - \hat{y}$ plane

$$214 \quad \mathbf{e}_3 = (\hat{x}_3, \hat{y}_3, 0), \hat{y}_3 > 0 \quad (17)$$

215 and the fourth to

$$216 \quad \mathbf{e}_4 = (\hat{x}_4, \hat{y}_4, \hat{z}_4), \hat{z}_4 > 0. \quad (18)$$

217 This coordinate system reduces translational (equation 15) and rotational (equa-
218 tions 16 to 18) non-uniqueness without loss of generality. It is necessary to work
219 with a local coordinate system when using coda waves alone because the CWI
220 technique constrains only event separation between earthquakes. The inclusion of
221 travel times in Combining Travel Time and CWI Constraints allows us to move to
222 a global reference system.

223 In summary, the posterior $P(\mathbf{e}_1, \dots, \mathbf{e}_n | \tilde{\delta}_{CWIN})$ and its negative logarithm L
224 describe the joint probability of multiple event locations given the observed coda
225 waves. The most likely set of locations is given by the minimum of L . In this
226 paper we use the Polak-Ribiere technique (*Press et al.*, 1987), a conjugate gradient
227 method, to minimize L . It uses the derivatives of L , derived in Appendix , to
228 guide the optimization procedure. Note that when optimizing equation 14 the
229 values of $\ln [c]$ and $\ln [P(e_i)]$ can be ignored because they are constant ($\ln [P(e_i)]$
230 is constant because we consider a uniform prior).

Synthetic experiments

We use synthetic examples in 2D and 3D with 50 earthquakes to test the performance of the optimization routine. In these examples the synthetic earthquakes are located randomly and CWI data generated according to the event separation. It is not necessary to generate synthetic waveforms and compute CWI estimates directly because we are testing the performance of the optimization routine only. The ability of CWI to estimate event separation has been demonstrated already (*Snieder and Vrijlandt, 2005; Robinson et al., 2007a, 2011*). We undertake a complete coda wave location experiment, including the calculation of CWI separation estimates, for recorded earthquakes in Relocating Earthquakes on the Calaveras Fault and in Combining Travel Time and CWI Constraints.

Examples 1 and 2 - 2D synthetic experiments

We design a 2D synthetic acoustic experiment (example 1) by randomly selecting \hat{x} - and \hat{y} -coordinates such that $-50 \leq \hat{x}, \hat{y} \leq 50$ m. These are indicated with triangles in Figure 1. We assume a local velocity of $\alpha = 3,300 \text{ ms}^{-1}$ between all event pairs and a dominant frequency of 2.5 Hz to represent waveform data filtered between 1 and 5 Hz. The CWI data are defined by the dominant wavelength normalized positive bounded Gaussian PDF with statistics $\bar{\mu}_N$ and $\bar{\sigma}_N$ (*Robinson et al., 2011*). A hypothetical CWI mean is created by setting

$$\bar{\mu}_N = \mu_1 \left(\tilde{\delta}_t \right) \quad (19)$$

using equation (41). This assumption ensures that the sample mean of hypothetical separation estimates is consistent with known CWI biases (*Robinson et al., 2011*). In example 1 we use $\bar{\sigma}_N = 0.02$ between all event pairs. Application of our

254 optimization procedure on the hypothetical CWI data yields the circles in Figure
 255 1. The optimization does not lead to the exact solution due to the addition of noise
 256 ($\bar{\sigma}_N = 0.02$) on the hypothetical CWI data. The average coordinate error is 2.0 m
 257 which is small compared to the noise of $\bar{\sigma}_N = 0.02$ which for $v_s = 3300 \text{ ms}^{-1}$
 258 and $f_{dom} = 2.5 \text{ Hz}$ corresponds to roughly 25 m.

259 *Robinson et al.* (2011) demonstrates that the noise on CWI estimates is often
 260 larger than 0.02 and that it increases with event separation. Consequently, example
 261 1 is simplistic because we fix $\bar{\sigma}_N = 0.02$ for all pairs. In example 2 we increase
 262 the uncertainty and introduce a distance dependance into the hypothetical $\bar{\sigma}_N$ by
 263 defining $\bar{\sigma}_N = \epsilon(\delta_t)$, where $\epsilon(\delta_t)$ is the half-width of the errorbars for a synthetic
 264 acoustic experiment with filtering between 1 and 5 Hz (see Fig. 4(b) of *Robinson*
 265 *et al.*, 2011). Repeating the optimization leads to the circles in Figure 2 which
 266 have an average coordinate error of 2.8 m.

267 Conjugate gradient based optimization techniques are susceptible to the pres-
 268 ence of local minima. This is because they use the slope of the target function to
 269 explore the solution space. We explore the impact of local minima for our CWI
 270 location problem by beginning the optimization from 25 randomly chosen starting
 271 positions. We observe no differences in the solution for either example.

272 Three observations can be drawn from the error structure in Figures 1 and 2.
 273 Firstly, the location errors depicted by gray bars increase between examples 1 and
 274 2 with the introduction of larger noise. Secondly, the errors are larger for events
 275 at greater distances from the center. This is because events near the center of
 276 the cluster are constrained by links from all angles, whereas those on the outside
 277 are moderated by links from a limited number of directions. This observation is
 278 analogous to problems associated with poor azimuthal coverage in triangulation

279 problems such as individual earthquake location from limited travel time data, or
280 GPS positioning with few satellites. Our third observation is that the location er-
281 rors form a pattern of circular rotation, despite our attempt to correct for rotational
282 non-uniqueness with the local coordinate system.

283 The local coordinate system works by constraining the location of the first
284 three earthquakes. Earthquake 1 is fixed at the origin, earthquake 2 on the pos-
285 itive \hat{x} -axis and earthquake 3 has $\hat{y} > 0$. As the number of events increase the
286 strength of these constraints on later events weakens allowing small rotations of
287 events with respect to each other. That is, even though the rotational freedom
288 of the cluster is in principal removed by the constraints imposed on the events
289 (see equations (15) to (17 - equation 18 is needed in 3D only) we observe that in
290 practice the presence of noise allows the rotational non-uniqueness to reappear.
291 This is because errors align themselves in directions least constrained by data.
292 For the CWI technique this amounts to rotations in 2D. The same phenomena
293 is observed in linear inversion where noise creates large spurious model changes
294 in directions of the eigenvectors with the smallest singular values (*Aster et al.*,
295 2005). Fortunately however, combining coda waves with measurements of travel
296 times alleviates this problem and facilitate the removal of a local coordinate sys-
297 tem altogether (see Combining Travel Time and CWI Constraints). On balance
298 however, we gain confidence in the optimization procedure due to its stability for
299 different starting locations and because of the small average coordinate errors of
300 2.0 m and 2.8 m for examples 1 and 2, respectively.

301 **Example 3 - The impact of incomplete event pairs in 2D**

302 Synthetic examples 1 and 2 use 100% direct linkage between event pairs. That
303 is, there is a constraint between each earthquake and all other events. In reality,
304 we might expect that the separation between some pairs will not be constrained
305 by CWI data due to poor signal to noise ratio in the coda for common stations.
306 Obviously, the fewer stations that record an event the more likely it is that links
307 between it and other events will be broken. In such cases the probabilistic distance
308 constraint between a pair of events may only exist indirectly through multiple
309 pairs. In this section we consider the impact of reduced linkage between event
310 pairs. In example 3, we repeat example 2 using 90%, 80%, ..., 10% of the links.
311 As with the above examples, we undertake the optimization with 25 randomly
312 chosen starting locations.

313 Figures 3(a) and (b) illustrates the maximum Δ_{max} (top) and mean Δ_{μ} (mid-
314 dle) of the coordinate error as a function of percentage of earthquake pairs that
315 are directly linked by a separation estimate. We show the statistics for the ‘best’
316 optimization solution (black) and for the solution space when all 25 optimizations
317 are considered (gray). In the former case the best solution is determined by the
318 set of event locations which lead to the smallest value of L . The error in the best
319 solution is consistent when 30% or more of the branches are used. The errors
320 increase when only 10% or 20% of the constraints are included. Interestingly, this
321 breakdown around 20% to 30% coincides with the point where the average num-
322 ber of branches required to link an event pair reaches 2 (see Fig.3 (bottom)). Since
323 the average number of branches can be computed in advance it can be used as an
324 indication of the inversion stability prior to optimization. A higher breakdown is
325 observed when all 25 solutions are considered collectively. For example, the max-

imum coordinate error Δ_{max} exceeds that for the best solution for linkage $\leq 60\%$ confirming that the optimization is susceptible to local minima and that a range of starting points should be considered. Some optimizations fail to converge after 1200 iterations when the linkage is 60% or lower. All optimizations fail when the linkage is 20% or lower. Despite their failure to converge, the locations at final iteration are close to the actual solution.

The derivatives used in the conjugate gradient method depend on events connected by CWI measurements. Consequently, earthquakes that are only connected via other events do not ‘communicate’ with each other directly. To some extent, this should be addressed during the iterative process where location information can spread to events which have no direct links. However, the lack of direct connection through the gradient could prevent convergence in extreme cases, or more likely slow the procedure down. This could explain why some examples do not converge after 1200 iterations. *VanDecar and Snieder (1994)* show that derivative based regularization acts slowly through iterative least-squares, because every cell in one iteration communicates only with its neighbours, and they demonstrate that this can be fixed with preconditioning in some cases. Their findings suggest that it may be possible to improve the convergence (stability and/or speed) of the CWI optimization by preconditioning.

Example 4 - The impact of incomplete event pairs in 3D

In Example 4 we expand the optimization routine to 3D by randomly picking a set of actual event locations for 50 earthquakes with $-50 \text{ m} \leq \hat{x}, \hat{y}, \hat{z} \leq 50 \text{ m}$. As in the 2D case we assume a local velocity of $v = 3,300 \text{ ms}^{-1}$ between all event pairs and a dominant frequency of 2.5 Hz to represent waveform data filtered between

350 1 and 5 Hz. The hypothetical CWI mean is created using equation (19) which en-
351 sures consistency between the sample mean of hypothetical separation estimates
352 and CWI biases. We use a standard deviation for the noisy CWI estimates of
353 $\bar{\sigma}_N = \epsilon$ and perform the optimization using 10%, 20%, ..., 100% of the direct
354 links. In each case we repeat the optimization 25 times using randomly chosen
355 starting locations. The results are summarised in Figure 4.

356 When 70% of the direct constraints are considered all optimization results
357 (gray) are consistent with the best solution obtained from all 25 starting locations
358 (black). The best solution constrains the event locations down to 30% of the direct
359 links. There is one notable difference between the 3D and 2D results. In 2D the
360 final iteration was close to the actual solution when the optimization failed to
361 converge. Conversely, in 3D the optimization appears to converge to the correct
362 solution or fail completely, leading to a set of locations at final iteration which
363 do not resemble the actual solution. This is depicted in Figure 4 by the absence
364 of the gray and black lines below 60% and 30% of the constraints, respectively.
365 The reason for this difference may be due to the increased number of degrees of
366 freedom in 3D requiring a greater number of iterations to converge. Nevertheless,
367 the accurate convergence of the best solution for cases with 30% linkage or higher
368 is encouraging for the potential of coda wave optimization to constrain earthquake
369 location.

370 **Summary of synthetic experiments**

371 In summary, the synthetic examples demonstrate the ability of coda wave data
372 to constrain relative event location using optimization. The optimization error is
373 influenced by the noise on CWI estimates with greater $\bar{\sigma}_N$ leading to larger errors

374 in the solutions. When 70% or more of the direct branches are used the optimizer
375 is stable with no observable difference in the solution for 25 randomly chosen
376 starting locations. As the direct linkage reduces to 50% the optimization becomes
377 less stable and the best solution from 25 random starting locations is required to
378 find the optimal solution. All optimisations fail to converge as the number of links
379 decrease below 30%.

380 **Relocating Earthquakes on the Calaveras Fault**

381 In this section we relocate 68 earthquakes from the Calaveras Fault, California.
382 The 68 earthquakes are selected from the 308 earthquake Calaveras example re-
383 leased with the open source Double Difference algorithm or hypoDD (*Waldhauser*
384 *and Ellsworth*, 2000; *Waldhauser*, 2001) [See also Data and Resources]. These
385 events are chosen for four reasons. Firstly, they are recorded by a large number
386 of stations (Fig. 5) and therefore lend themselves to accurate travel time loca-
387 tion. This makes them ideal for assessing the performance of a new location tech-
388 nique. Secondly, they are distributed with separations from near zero to hundreds
389 of meters making them ideal for application of CWI. Thirdly, Calaveras earth-
390 quakes have been well researched with several studies having relocated events in
391 the region (*Waldhauser*, 2001; *Schaff et al.*, 2002; *Waldhauser and Schaff*, 2008).
392 Finally, the hypoDD locations for these 68 earthquakes align in a streak increas-
393 ing the likelihood that they have near identical source mechanisms, a necessary
394 assumption for the application of equation 2. The relocations in this paper are
395 sorted into four examples as summarised in Table 1.

396 **Example 5 - comparison of CWI, catalogue and hypoDD loca-** 397 **tions**

398 Figure 6 illustrates three sets of locations for the Calaveras earthquakes. The first
399 column shows the original catalogue locations for all 308 earthquakes. That is,
400 each event is located individually using all available travel time arrivals and a
401 regional velocity model. The 68 earthquakes of interest in this study are differen-
402 tiated in black. Catalogue locations suggest that the 68 earthquakes of interest are
403 spatially widely distributed on the scale of Figure 6.

404 To apply CWI we download available waveforms from the Northern Califor-
405 nia Earthquake Data Center (See Data and Resources). Unsuitable waveforms are
406 removed using the conditions summarised in Table 2. Remaining waveforms are
407 filtered between 1 and 5 Hz and aligned to P arrivals at 0 s. CWI estimates are ob-
408 tained from 5 s wide non-overlapping time windows between $2.5 \leq t \leq 20$ s and
409 used to create probabilistic constraints on event separation. We utilize the local
410 coordinate system introduced in Theory and find the optimum relative locations
411 using Polak-Ribiere optimization.

412 CWI locations for the 68 events are illustrated in column two of Figure 6.
413 Catalogue locations (gray) are shown for the remaining 240 earthquakes and are
414 included to ease comparison. The third column of Figure 6 illustrates the locations
415 given by hypoDD with Singular Value Decomposition (SVD), absolute arrival
416 times and cross correlation computed travel time differences.

417 Absolute locations cannot be found by CWI alone. This is because of the
418 non-uniqueness associated with translation, rotation and reflection. For the sake
419 of comparison, we arbitrarily choose a ‘master’ event and translate our relative
420 locations to align with the hypoDD location for the same event. This arbitrary

translation does not change the relative locations. We return to this issue of relative versus absolute location in Example 7 by introducing a combined travel time and coda wave inversion.

The spatial distribution of the CWI locations is clearly tighter than the catalogue locations of column 1. That is, CWI provides an independent indication of clustering for the 68 events and to first order, similar locations to those from hypoDD (column 3). There is a small second order difference between the CWI and hypoDD based locations. In particular, the lineation is less clear in the CWI locations (column 2) than the hypoDD locations (column 3). Our experience suggest that the coda are less supportive of the presence of streaks although a complete understanding of these differences is left for future work. Our attention now is devoted towards understanding how both techniques perform with fewer stations (Example 6) and exploring how CWI and travel times can be combined (Examples 7 and 8).

Example 6 - Dependence on the number of stations

Accurate location of the Calaveras events is possible using arrival phases because of the excellent recording situation in California with many stations and strong azimuthal coverage (see Fig. 5). In contrast, a small number of stations and poor azimuthal coverage are common limitations when trying to locate intraplate clusters. For example, there are only four network seismic stations in the South West Seismic Zone of Western Australia, a region similar in size to that hosting 805 stations in Figure 5.

We explore the impact of poorer recording situations in example 6 by re-locating the 68 Calaveras events using hypoDD and coda waves with a reduced

445 number of stations. We begin with 10 stations and repeat the process removing
446 one at a time until a single station remains. The 10 stations considered are shown
447 in Figure 7 and the order of removal explained in Table 3.

448 CWI locations are illustrated in Figure 8 for the inversions with seven, five,
449 four, three, two and one station. We observe a high level of consistency between
450 these 6 inversions and the locations shown in Figure 6 (column 2) when all stations
451 are considered. That is, the coda wave approach is self-consistent regardless of
452 the number of stations available, reinforcing our hypothesis that coda waves can
453 constrain location in what would normally be regarded as a poor station network.

454 Figure 9 illustrates the hypoDD inversion results for seven, five and four sta-
455 tions. The travel time problem is ill-posed for fewer than four stations so it is not
456 possible to apply hypoDD with SVD for three or fewer stations. The hypoDD
457 locations are not self-consistent as the number of stations is reduced. We observe
458 a general increase in scatter and a higher number of stray events outside the clus-
459 ter when less stations are used with hypoDD. Even with seven stations the linear
460 geometry of Figure 6 (column 3) is less evident.

461 As the number of stations are reduced both the CWI and hypoDD techniques
462 are not able to re-locate all events. To use the coda waves we need at least one pair-
463 wise separation constraint to be formed from the available stations. This means
464 that for every event there must be at least one station that records it and at least
465 one other earthquake sufficiently well to apply CWI. Fortunately, we can make
466 an assessment of this prior to starting the inversion. The top panel of Figure 10
467 demonstrates that when five or more stations are used, CWI can constrain the lo-
468 cation of all 68 earthquakes. When less than five stations are used the coda waves
469 constrain a decreasing number of events until at one station it is only possible to

470 locate 55 of the 68 events. The hypoDD algorithm also fails to locate all events
471 as the number of stations is reduced. In the case of hypoDD an event can be
472 identified as unlocatable in one of two stages. Firstly, the data are analyzed to
473 ensure that there exists travel time differences for each event and at least one other
474 earthquake. This is analogous to the situation for the coda wave technique. The
475 hypoDD program also has a secondary identification phase in which events that
476 can not be located sufficiently are rejected during the inversion. This process is
477 related to the iterative removal of outliers described by *Waldhauser and Ellsworth*
478 (2000). The top panel of Figure 10 shows that the number of events re-located by
479 hypoDD fluctuates between 63 and 28 earthquakes for ten to four stations and it
480 demonstrates that the number of events located by hypoDD is less than or equal
481 to the number located by CWI.

482 The remaining panels of Figure 10 illustrate a statistical comparison of the
483 CWI and hypoDD reduced station locations to those using hypoDD with all avail-
484 able data. For the CWI inversions the mean and maximum coordinate difference
485 is consistent regardless of the number of stations considered. In contrast, the
486 hypoDD mean and maximum coordinate error fluctuate above those for CWI con-
487 firming that the hypoDD inversion is less stable than CWI with fewer stations.

488 **Combining Travel Time and CWI Constraints**

489 In Examples 5 and 6 we compare the location of the Calaveras earthquakes us-
490 ing coda wave and arrival time based constraints independently. Since the arrival
491 time (direct or difference) and coda wave data come from different sections of the

492 waveform they provide independent constraints on the locations. In this section
 493 we devise a location algorithm which incorporates both CWI and travel time data.

494 We do not propose a new technique for earthquake location using travel time
 495 differences. Rather, we exploit the information created by hypoDD with SVD to
 496 define a probability density (or posterior) function

$$497 \quad P(\mathbf{e}_p | \Delta_{TT}) \frac{1}{(2\pi)^{3/2} \sqrt{|\Sigma|}} \times \exp \left(-\frac{1}{2} ([\mathbf{e}_p - \mu_{\mathbf{e}_p}]^T \Sigma^{-1} [\mathbf{e}_p - \mu_{\mathbf{e}_p}]) \right), \quad (20)$$

498 where

$$499 \quad \mathbf{e}_p = (x_p, y_p, z_p)^T \quad (21)$$

500 is the location of event p ,

$$501 \quad \mu_{\mathbf{e}_p} = (\mu_{x_p}, \mu_{y_p}, \mu_{z_p})^T \quad (22)$$

502 is the most likely location as determined using the travel time data, and

$$503 \quad \Sigma = \begin{pmatrix} \sigma_{x_p}^2 & 0 & 0 \\ 0 & \sigma_{y_p}^2 & 0 \\ 0 & 0 & \sigma_{z_p}^2 \end{pmatrix} \quad (23)$$

504 is the covariance matrix. In this paper we define the mean location $\mu_{\mathbf{e}_p}$ and covari-
 505 ance matrix by the hypoDD optimum solution and its uncertainties. It is important
 506 to note that hypoDD must be used with SVD to obtain useful estimates of σ_{x_p} , σ_{y_p}
 507 and σ_{z_p} because the errors reported by conjugate gradient methods (LSQR) are
 508 grossly underestimated in hypoDD (Waldhauser, 2001).

509 We pose the location problem using the negative log likelihood

$$510 \quad L(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_1, \mathbf{e}_n) = -\sum_{i=1}^n \ln [P(\mathbf{e}_i | \Delta_{TT})] - \sum_{i=1}^{n-1} \sum_{j=i+1}^n \ln [P(\delta_{CWIN} | \mathbf{e}_i, \mathbf{e}_j)], \quad (24)$$

511 where $(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n)$ is the joint location,

$$512 \quad \sum_{i=1}^n \ln [P(\mathbf{e}_i | \Delta_{TT})] \quad (25)$$

513 incorporates the travel time constraints and

$$514 \quad \sum_{i=1}^{n-1} \sum_{j=i+1}^n \ln [P(\delta_{CWIN} | \mathbf{e}_i, \mathbf{e}_j)] \quad (26)$$

515 the coda waves.

516 We must differentiate L to use the Polak-Ribiere conjugate gradient technique
517 of *Press et al.* (1987). The derivative of $L(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_n)$ with respect to x_p is given
518 by

$$519 \quad \frac{\partial L}{\partial x_p} = -\frac{\partial \ln [P(\mathbf{e}_p | t_{DD})]}{\partial x_p} - \sum_{i=p+1}^N \frac{\partial \ln [P(\delta_{CWIN} | \mathbf{e}_p, \mathbf{e}_i)]}{\partial x_p} \\ - \sum_{j=1}^{p-1} \frac{\partial \ln [P(\delta_{CWIN} | \mathbf{e}_j, \mathbf{e}_p)]}{\partial x_p} \quad (27)$$

520 where

$$521 \quad \sum_{i=p+1}^N \frac{\partial \ln [P(\delta_{CWIN} | \mathbf{e}_p, \mathbf{e}_i)]}{\partial x_p} \quad (28)$$

522 and

$$523 \quad \sum_{j=1}^{p-1} \frac{\partial \ln [P(\delta_{CWIN} | \mathbf{e}_j, \mathbf{e}_p)]}{\partial x_p} \quad (29)$$

524 are defined in Appendix and

$$525 \quad \frac{\partial \ln [P(\mathbf{e}_p | t_{DD})]}{\partial x_p} = -\frac{1}{2} [1, 0, 0]^T \Sigma^{-1} [\mathbf{e}_p - \mu_{\mathbf{e}_p}] \\ - \frac{1}{2} [\mathbf{e}_p - \mu_{\mathbf{e}_p}]^T \Sigma^{-1} [1, 0, 0]. \quad (30)$$

526 Similarly, for the derivatives with respect to y_p and z_p we have

$$527 \quad \frac{\partial \ln [P(\mathbf{e}_p | t_{DD})]}{\partial y_p} = -\frac{1}{2} [0, 1, 0]^T \Sigma^{-1} [\mathbf{e}_p - \mu_{\mathbf{e}_p}] \\ - \frac{1}{2} [\mathbf{e}_p - \mu_{\mathbf{e}_p}]^T \Sigma^{-1} [0, 1, 0] \quad (31)$$

528 and

$$\begin{aligned} \frac{\partial \ln[P(\mathbf{e}_p|t_{DD})]}{\partial z_p} &= -\frac{1}{2}[0, 0, 1]^T \Sigma^{-1}[\mathbf{e}_p - \mu_{\mathbf{e}_p}] \\ &\quad -\frac{1}{2}[\mathbf{e}_p - \mu_{\mathbf{e}_p}]^T \Sigma^{-1}[0, 0, 1]. \end{aligned} \quad (32)$$

529

530 Combining the travel time and coda wave data offers two advantages. Firstly,
531 it combines independent constraints on the event locations offering further confi-
532 dence in the resulting solution. Secondly, the travel time constraints in the form
533 of equation (25) resolve the inherent non-uniqueness of the CWI inversion that
534 is associated with translation, rotation and reflection around a global coordinate
535 system. This means that it is no longer necessary to use a local coordinate system
536 and we can solve directly for location with respect to a global reference. Collec-
537 tively, these advantages improve the behavior of the Polak-Ribiere optimization
538 leading to faster and more stable convergence. Consequently, we no longer have
539 to consider multiple randomly chosen starting locations.

540 **Example 7 - Combining travel time and CWI constraints**

541 Figure 11 illustrates the earthquake locations obtained when we combine the travel
542 time and coda wave data using all data (left) and five stations (right). The linear
543 features observed in the original hypoDD inversions (see Fig. 6) are evident in
544 both cases. However, the coda waves introduce a scatter around these streaks.
545 That is, the locations in figure 11 result from a trade-off between hypoDD's desire
546 to place the events on linear features and the coda waves voracity to push them
547 away from streaks. When all stations are used the hypoDD constraints are strong
548 and little off-streak scatter is introduced. As we reduce hypoDD's leverage by
549 decreasing the number of stations to five, we observe an increase in off-streak
550 scatter resulting from the enhanced influence of the coda.

551 **Example 8 - Combining CWI and travel times when the travel**
552 **times constrain a limited number of events**

553 In intraplate regions such as Australia it is common to deploy temporary seis-
554 mometers to monitor aftershocks for significant events (*Bowman et al.*, 1990;
555 *Leonard*, 2002). Traditionally, these deployments facilitate a higher accuracy of
556 location for events occurring during the deployment period. Using our combined
557 inversion it is possible to re-locate all events by employing the detailed travel
558 time data when the temporary network is in-situ and using coda waves from net-
559 work stations when the deployment is absent. The hypothesis, to be tested in this
560 section, is that conducting such a combined inversion will improve the location
561 accuracy of events outside the deployment period.

562 An estimate of the cumulative number of aftershocks $N(t)$ after t days is
563 given by the modified Omori formula

$$564 \quad N(t) = K \frac{c^{1-p} + (t + c)^{1-p}}{p - 1} \quad (33)$$

565 (*Utsu et al.*, 1995). The empirically derived constants, K , C and p vary between
566 tectonic settings. For example, using recorded aftershocks with $M \geq 3.2$ of the
567 Hokkaido-Nansei-Oki, Japan $M_s = 7.8$ earthquake of 12 July 1993, *Utsu et al.*
568 (1995) obtained maximum likelihood estimates for K , p and c of 906.5, 1.256
569 and 1.433, respectively. With these empirically derived values an array deployed
570 within 4 days and left for 150 days will record roughly one half of the aftershocks
571 occurring within the first 1000 days. That is,

$$572 \quad \frac{N(150 + 4) - N(4)}{N(1000)} = \frac{2257 - 934}{2626} \approx 0.5. \quad (34)$$

573 This idea is illustrated in Figure 12 which shows the best fitting Omori Formula

574 separated into segments before (gray), during (black) and after (gray) the pseudo
575 temporary deployment.

576 With this idea of a temporary deployment in mind we have another attempt at
577 relocating the Calaveras earthquakes. In Example 8 we consider the travel time
578 constraints on half (34) of the earthquakes and incorporate coda wave data from a
579 single station for all 68 earthquakes. The combined inversion is shown in column
580 1 of Figure 13. The inversion result is similar to the combined inversion when
581 all travel time data is incorporated (see Fig. 11). The slight increase in scatter
582 observed here can be explained by the events with no travel time constraints and
583 the tendency of the coda to push events away from streaks.

584 Remarkably, the combined coda wave and travel time inversion locates all 68
585 earthquakes to an accuracy similar to the inversions with all data. In contrast when
586 travel time data is used alone it is only possible to locate the 34 events recorded
587 by the pseudo temporary deployment. This ability of coda waves to constrain
588 the location of events recorded by a single station creates new opportunities for
589 understanding earthquakes in regions with limited station coverage.

590 **Discussion and Conclusions**

591 Coda wave interferometry is an emerging technique for constraining earth-
592 quake location. The technique relies on the interference between coda waves of
593 closely located events and is hence useful for studying earthquake clusters and/or
594 aftershock sequences. Coda wave constraints are independent of travel times and
595 can be used in isolation or combination with early onset body waves. The strength

596 of coda is that it is possible to constrain earthquake location from a single station,
597 an outcome demonstrated most clearly by Figures 8 and 13.

598 Coda wave interferometry offers a new technique for understanding earth-
599 quakes in intraplate areas with sparse networks and poor azimuthal coverage. In
600 particular, the ability to combine coda wave constraints with travel times makes it
601 possible to link well constrained events from a temporary deployment with those
602 recorded outside the deployment period. All that is required to achieve this is at
603 least one network station which has recorded sufficient events from both periods.
604 CWI facilitates the location of poorly recorded events to an accuracy approach-
605 ing those recorded during the temporary deployment and therefore opens new
606 avenues for imaging intraplate fault structures and improving our understanding
607 of intraplate seismicity and earthquake hazard.

608 Another potential application of CWI is in the area of hydraulic fracturing
609 such as hot rock geothermal projects, petroleum reservoir engineering, tight gas
610 extraction, CO₂ geosequestration and/or underground brine injection. Monitoring
611 pumping-induced micro earthquakes is a key step in understanding the migration
612 of fluids in such reservoirs. There is a trade-off in the ability of surface deployed
613 networks to locate events which are small and/or deep. Downhole seismic mon-
614 itoring is likely to play increasingly important roles in deep reservoir projects.
615 CWI creates new possibilities to monitor pumping induced micro earthquakes
616 from fewer boreholes and hence dramatically reduce the costs of reservoir mon-
617 itoring at large depths. It may also be possible to utilize coda for understanding
618 hazard in tunneled mining operations where the location of deep tunnels prohibits
619 azimuthal coverage of induced events.

Data and Resources

We thank the Northern California Earthquake Data Center (NCEDC) for providing the Calaveras data and the Northern California Seismic Network (NCSN), U.S. Geological Survey, Menlo Park and Berkeley Seismological Laboratory, University of California, Berkeley for contributing it to the NCEDC. The waveforms can be downloaded from <http://www.ncedc.org/ncedc/access.html> (last accessed August 2012). We also acknowledge Felix Waldhauser and William Ellsworth, the authors of the openly available Double Difference location algorithm, hypoDD which can be downloaded from <http://www.ldeo.columbia.edu/felixw/hypoDD.html> (last accessed August 2012). The International Seismological Centre can be found at <http://www.isc.ac.uk/> (last accessed December 2012). The National Earthquake Information Center catalogue can be accessed from <http://earthquake.usgs.gov/earthquakes/eqarchives/epic/> (last accessed December 2012).

Acknowledgments

Geoscience Australia, the Research School of Earth Sciences at The Australian National University, and the Center for Wave Phenomena at the Colorado School of Mines, are acknowledged for supporting this research. The paper is published with permission of the CEO Geoscience Australia. Work was conducted as part of an Australian Research Council Discovery Project (DP0665111). This paper has benefited significantly from: reviews by Trevor Allen and Clive Collins at Geoscience Australia; as well as comments and advice received from Felix

642 Waldhauser on this and earlier research on the application of CWI to earthquake
643 location.

644 **References**

645 Ake, J., D. O’Connell, and L. Block (2005), Deep-injection and closely monitored
646 induced seismicity at Paradox Valley, Colorado, *Bull. Seism. Soc. Am.*, 95(2),
647 664–683.

648 Aki, K. (1969), Analysis of the seismic coda of local earthquakes as scattered
649 waves, *J. Geophys. Res.*, 74(2), 615–631.

650 Aster, R. C., B. Borchers, and C. H. Thurber (2005), *Parameter estimation and*
651 *inverse problems*, *International Geophysics Series*, vol. 90, Elsevier Academic
652 Press, USA.

653 Bondár, I., S. C. Myers, E. R. Engdahl, and E. A. Bergman (2004), Epicentre
654 accuracy based on seismic network criteria, *Geophys. J. Int.*, 156, 483–496.

655 Bowman, J. R., G. Gibson, and T. Jones (1990), Aftershocks of the 1988 January
656 22 Tennant Creek, Australia intraplate earthquakes: evidence for a complex
657 thrust-fault geometry, *Geophys. J. Int.*, 100, 87–97.

658 Campbell, K. W. (2003), Strong motion attenuation, in *International Handbook*
659 *of Earthquake and Engineering Seismology*, vol. B, edited by W. H. K. Lee,
660 H. Kanamori, P. C. Jennings, and C. Kisslinger, chap. 60, pp. 1003–1012, Aca-
661 demic Press, London.

- 662 Curtis, A., and R. Snieder (2002), Probing the Earth's interior with seismic to-
 663 mography, in *International Handbook of Earthquake Engineering Seismology*,
 664 vol. A, edited by W. H. Lee, H. Kanamori, P. C. Jennings, and C. Kisslinger,
 665 chap. 52, pp. 861–874, Academic Press, London.
- 666 Deichmann, N., and M. Garcia-Fernandez (1992), Rupture geometry from high-
 667 precision relative hypocentre locations of microearthquake clusters, *Geophys.*
 668 *J. Int.*, *110*, 501–517.
- 669 Douglas, A. (1967), Joint epicentre determination, *Nature*, *215*, 47–48.
- 670 Frankel, A. D., C. S. Mueller, T. P. Barnhard, E. V. Leyendecker, R. L. Wesson,
 671 S. C. Harmsen, F. W. Klein, D. M. Perkins, N. C. Dickman, S. L. Hanson,
 672 and M. G. Hopper (2000), USGS National seismic hazard maps, *Earthquake*
 673 *Spectra*, *16*(1), 1–19.
- 674 Frèmont, M.-J., and S. D. Malone (1987), High precision relative locations of
 675 earthquakes at Mount St. Helens, *J. Geophys. Res.*, *92*(B10), 10,223–10,236.
- 676 Got, J.-L., J. Frèchet, and F. W. Klein (1994), Deep fault plane geometry inferred
 677 from multiplet relative relocation beneath the south flank of Kilauea, *J. Geo-*
 678 *phys. Res.*, *99*(B8), 15,375–15,386.
- 679 Gutenberg, B. (1945), Amplitudes of surface waves and magnitudes of shallow
 680 earthquakes, *Bull. Seism. Soc. Am.*, *35*, 3–12.
- 681 Ito, A. (1985), High resolution relative hypocenters of similar earthquakes by
 682 cross-spectral analysis method, *J. Phys. Earth*, *33*, 279–294.

- 683 Kennett, B. L. N., E. R. Engdahl, and R. Buland (1995), Constraints on seismic
684 velocities in the Earth from traveltimes, *Geophys. J. Int.*, 122, 108–124.
- 685 Kennett, B. L. N., S. Fishwick, and M. Heintz (2004), Lithospheric structure in
686 the Australian region - a synthesis of surface wave and body wave studies, *Ex-
687 ploration Geophysics*, 35, 242–250.
- 688 Lees, J. M. (1998), Multiplet analysis at Coso Geothermal, *Bull. Seism. Soc. Am.*,
689 88(5), 1127–1143.
- 690 Leonard, M. (2002), The Burakin WA earthquake sequence Sept 2000 – June
691 2002, in *Total Risk Management in the Privatised Era, Australian Earthquake
692 Engineering Society Conference*, vol. 10th, edited by M. Griffith, D. Love,
693 P. McBean, A. McDougall, and B. Butler, pp. 22(1)–22(5), AEES, University
694 of Adelaide.
- 695 Nadeau, R. M., and T. V. McEvilly (1997), Seismological studies at Parkfield V:
696 Characteristic microearthquake sequences as fault–zone drilling targets, *Bull.
697 Seism. Soc. Am.*, 87(6), 1463–1472.
- 698 Pavlis, G. L. (1992), Appraising relative earthquake location errors, *Bull. Seism.
699 Soc. Am.*, 82(2), 836–859.
- 700 Press, W. H., B. P. Flannery, S. A. Teukolsky, and W. T. Vetterling (1987), *Nu-
701 merical Recipes: The Art of Scientific Computing*, Cambridge University Press,
702 USA.
- 703 Richter, C. F. (1935), An instrumental earthquake magnitude scale, *Bull. Seism.
704 Soc. Am.*, 25(1), 1–32.

- 705 Robinson, D., T. Dhu, and J. Schneider (2006), Practical probabilistic seismic risk
706 analysis: A demonstration of capability, *Seism. Res. Let.*, 77(4), 452–458.
- 707 Robinson, D. J., M. Sambridge, and R. Snieder (2007a), Constraints on coda wave
708 interferometry estimates of source separation: The 2.5d acoustic case, *Explora-
709 tion Geophysics*, 38(3), 189–199.
- 710 Robinson, D. J., R. Snieder, and M. Sambridge (2007b), Using coda wave in-
711 terferometry for estimating the variation in source mechanism between double
712 couple events, *J. Geophys. Res.*, 112, b12302, doi:10.1029/2007JB004925.
- 713 Robinson, D. J., M. Sambridge, and R. Snieder (2011), A probabilistic approach
714 for estimating the separation between a pair of earthquakes directly from their
715 coda waves, *J. Geophys. Res.*, B04309, 1–17.
- 716 Rubin, A. M. (2002), Aftershocks of microearthquakes as probes of the mechanics
717 of rupture, *J. Geophys. Res.*, 107(B7,2142), 10.1029/2001JB000,496.
- 718 Rubin, A. M., D. Gillard, and J.-L. Got (1999), Streaks of microearthquakes along
719 creeping faults, *Nature*, 400, 635–641.
- 720 Schaff, D. P., G. H. R. Bokelmann, and G. C. Beroza (2002), High-resolution
721 image of Calaveras Fault seismicity, *J. Geophys. Res.*, 107(B9), 2186,
722 doi:10.1029/2001JB000,633.
- 723 Shearer, P., E. Hauksson, and G. Lin (2005), Southern California hypocenter relo-
724 cation with waveform cross-correlation, Part 2: Results using source-specific
725 station terms and cluster analysis, *Bull. Seism. Soc. Am.*, 95(3), 904–915.
726 doi:10.1785/0120040,168.

- 727 Shearer, P. M. (1999), *Introduction to Seismology*, Cambridge University Press,
728 USA, 260pp.
- 729 Sipkin, S. A. (2002), USGS earthquake moment tensor catalog, in *International*
730 *Handbook of Earthquake Engineering Seismology*, vol. A, edited by W. H. Lee,
731 H. Kanamori, P. C. Jennings, and C. Kisslinger, chap. 50, pp. 823–825, Aca-
732 demic Press, London.
- 733 Snieder, R. (1999), Imaging and averaging in complex media, in *Diffuse waves in*
734 *complex media, NATO Science Series C*, vol. 531, edited by J. P. Fouque, pp.
735 405–454, Kluwer Academic Publishers.
- 736 Snieder, R. (2006), The theory of coda wave interferometry, *Pure Appl. Geophys.*,
737 *163*, 455–473.
- 738 Snieder, R., and M. Vrijlandt (2005), Constraining the source separation
739 with coda wave interferometry: Theory and application to earthquake dou-
740 blets in the Hayward Fault, California, *J. Geophys. Res.*, *110*(B04301),
741 doi:10.1029/2004JB003317.
- 742 Spencer, C., and D. Gubbins (1980), Travel-time inversion for simultaneous earth-
743 quake location and velocity structure determination in laterally varying media,
744 *Geophys. J. R. Astr. Soc.*, *63*, 95–116.
- 745 Stirling, M. W., G. H. McVerry, and K. R. Berryman (2002), A new seismic
746 hazard model for New Zealand, *Bull. Seism. Soc. Am.*, *92*(5), 1878–1903.
- 747 Toro, G. R., N. A. Abrahamson, and J. F. Schneider (1997), Model of strong
748 ground motions from earthquakes in Central and Eastern North America: Best
749 estimates and uncertainties, *Seism. Res. Let.*, *68*(1), 41–57.

- 750 Utsu, T., Y. Ogata, and R. S. Matsu'ura (1995), The Centenary of the Omori
751 Formula for a decay law of aftershock activity, *J. Phys. Earth*, *43*, 1–33.
- 752 VanDecar, J. C., and R. Snieder (1994), Obtaining smooth solutions to large linear
753 inverse problems, *Geophysics*, *59*, 818–829.
- 754 Waldhauser, F. (2001), hypoDD – a program to compute double-difference
755 hypocenter locations (hypoDD version 1.0 - 03/2001), *Open file report 01-113*,
756 United States Geological Survey, Menlo Park, California.
- 757 Waldhauser, F., and W. L. Ellsworth (2000), A double–difference earthquake lo-
758 cation algorithm: method and application to the northern Hayward Fault, Cali-
759 fornia, *Bull. Seism. Soc. Am.*, *90*(6), 1353–1368.
- 760 Waldhauser, F., and W. L. Ellsworth (2002), Fault structure and mechanics of
761 the Hayward Fault, California, from double–difference earthquake locations, *J.*
762 *Geophys. Res.*, *107*(B3), 10.1029/2000JB000,084.
- 763 Waldhauser, F., and D. P. Schaff (2008), Large-scale relocation of two decades
764 of Northern California seismicity using cross-correlation and double-difference
765 methods, *J. Geophys. Res.*, *133*, B08311, doi10.1029/2007JB005479.
- 766 Waldhauser, F., W. L. Ellsworth, and A. Cole (1999), Slip-parallel lineations on
767 the Northern Hayward Fault, California, *Geophys. Res. Lett.*, *26*(23), 3525–
768 3528.

Authors' Affiliations, Addresses

David Robinson: Earth Monitoring and Hazard Group
 Geoscience Australia
 GPO Box 383
 Canberra ACT 2601 Australia
 E-mail: david.robinson@ga.gov.au

Also at: Research School of Earth Sciences
 Australian National University
 Canberra ACT 0200, Australia
 E-mail: david.robinson@anu.edu.au

Malcolm Sambridge: Research School of Earth Sciences
 Australian National University
 Canberra ACT 0200, Australia

Roel Snieder: Center for Wave Phenomena and Department of
 Geophysics
 Colorado School of Mines
 Golden CO 80401-1887, USA

Juerg Hauser: CSIRO Earth Science and Resource Engineering
 Dick Perry Av
 Kensington WA 6151 Australia

Table 1: Location examples for the 68 Calaveras earthquakes.

Example 5	Comparison of CWI, catalogue and hypoDD locations (using all available data).
Example 6	Exploration of station dependance for CWI and hypoDD (using a subset of data).
Example 7	Combined use of CWI and travel time data with all and a reduced number of stations.
Example 8	Combined use of CWI and travel time data when travel times constrain only 50% of the events.

Table 2: Conditions used to identify unsuitable waveforms before applying CWI
(Originally published as Table 5 *Robinson et al.*, 2011)

	condition
1	waveform is clearly corrupted
2	waveform indicates recording of more then one event
3	signal to noise ratio is obviously low
4	there is insufficient coda recorded after the first arrivals
5	there is insufficient recording before the arrivals (needed for accurate noise energy estimate)

Table 3: Stations considered when exploring the impact of reduced station coverage.

Number of Stations	Station Names
10	CCO, JCB, JST, CMH, HSP, JAL, CSC, JST, CAD, JHL, JRR
9	CCO, JCB, JST, CMH, HSP, JAL, CSC, JST, CAD, JHL
8	CCO, JCB, JST, CMH, HSP, JAL, CSC, JST, CAD
7	CCO, JCB, JST, CMH, HSP, JAL, CSC
6	CCO, JCB, JST, CMH, HSP, JAL
5	CCO, JCB, JST, CMH, HSP
4	CCO, JCB, JST, CMH
3	CCO, JCB, JST
2	CCO, JCB
1	CCO

Figure 1: Example 1 - Synthetic relocation of 50 earthquakes in 2D using all constraints with noise $\bar{\sigma}_N = 0.02$. Actual and optimization event locations are shown in triangles and circles, respectively.

Figure 2: Example 2 - Synthetic relocation of 50 earthquakes in 2D using all constraints with noise $\bar{\sigma}_N = 2\epsilon(\delta_t)$. Actual and optimization event locations are shown in triangles and circles, respectively.

Figure 3: Example 3 - Statistical measures of error in the solutions for the 2D synthetic cases when all and best optimization results are considered. The statistics Δ_{max} and Δ_μ are the maximum and mean coordinate error, respectively. The bottom subplot shows the average minimum number of branches required to link the 2450 pairs.

Figure 4: Example 4 - Statistical measures of error in the optimization solutions for the 3D synthetic cases when all (gray) and best (black) results are considered. The statistics Δ_{max} and Δ_μ are the maximum and mean coordinate error, respectively. The absence of the gray and black lines below 60% and 30% indicates a breakdown in the solutions when all or best optimization result(s) are considered, respectively.

Figure 5: Map showing location of the Calaveras cluster (white star) and 805 seismic stations (black triangles).

Figure 6: Example 5 - Comparison of relative earthquake locations using three different methods: catalogue location (column 1), CWI (column 2) and hypoDD (column 3). Note that in the case of the hypoDD and CWI inversions we consider only the 68 earthquakes in black, the gray catalogue locations for the remaining 240 (308-68) earthquakes are shown for the purpose of orientation only.

Figure 7: Location of the 10 stations (triangles) used to relocate the Calaveras events in Example 6 to 8. Stations are removed one at a time according to the order in Table 3 and the events relocated. The events are indicated with black circles.

Figure 8: Example 6 - CWI relative locations with reduced stations.

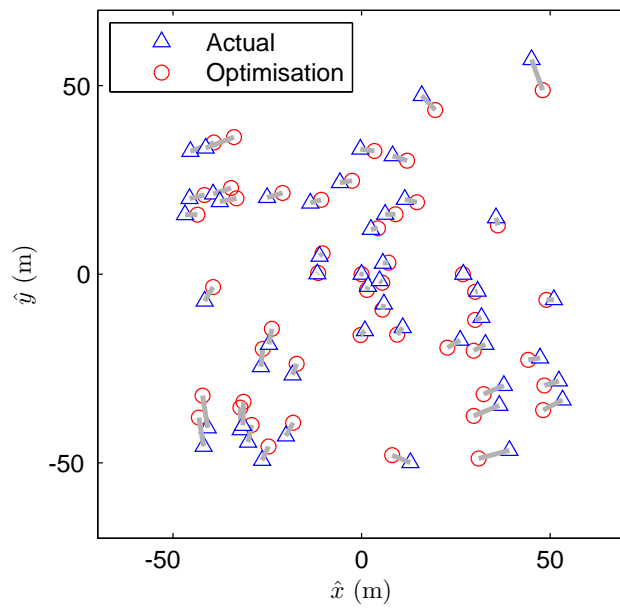
Figure 9: Example 6 - HypoDD (SVD) relative locations with reduced stations.

Figure 10: Example 6 - Number of constrainable events n_E in the CWI and hypoDD inversions as a function of the stations considered (top). Mean (middle) and maximum (bottom) of the difference computed between the reduced station inversion results (CWI and hypoDD) and the complete hypoDD locations for all 308 events.

Figure 11: Example 7 - Combined HypoDD (SVD) and CWI relative locations using data from all stations (left) and 5 stations (right).

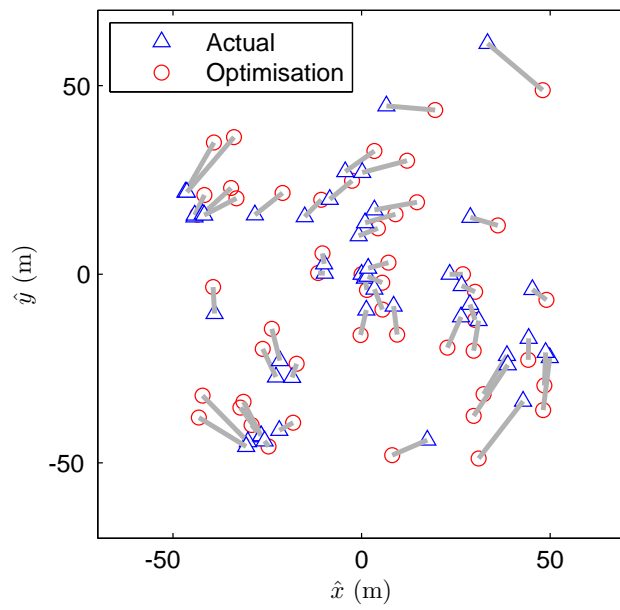
Figure 12: Cumulative number of aftershocks for the Hokkaido-Nansei-Oki, Japan $M_s = 7.8$ earthquake of 12 July 1993 using equation (33). The leftmost red, middle blue and rightmost red lines signify aftershocks occurring before, during and after the deployment of a pseudo temporary array installed 4 days after the main shock and left for 150 days. A temporary deployment of this kind will record roughly 50% of the aftershocks in the 1000 days following the mainshock.

Figure 13: Example 8 - Mimicking the deployment of a temporary network by ignoring data from all but station CCO for 50% (or 34) of the events. Relative locations are shown for the combined CWI and travel time inversion (left) and the inversion with travel times only (right). Only by combining the data is it possible to locate all 68 events. Furthermore, combining the data leads to a solution more consistent with Figure 6.



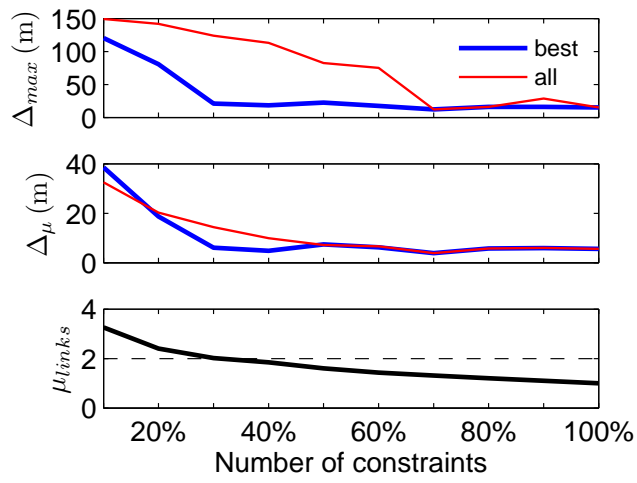
771

772 **Figure 1:**



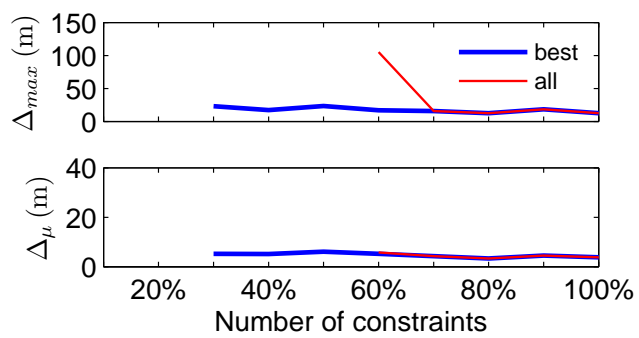
773

774 **Figure 2:**



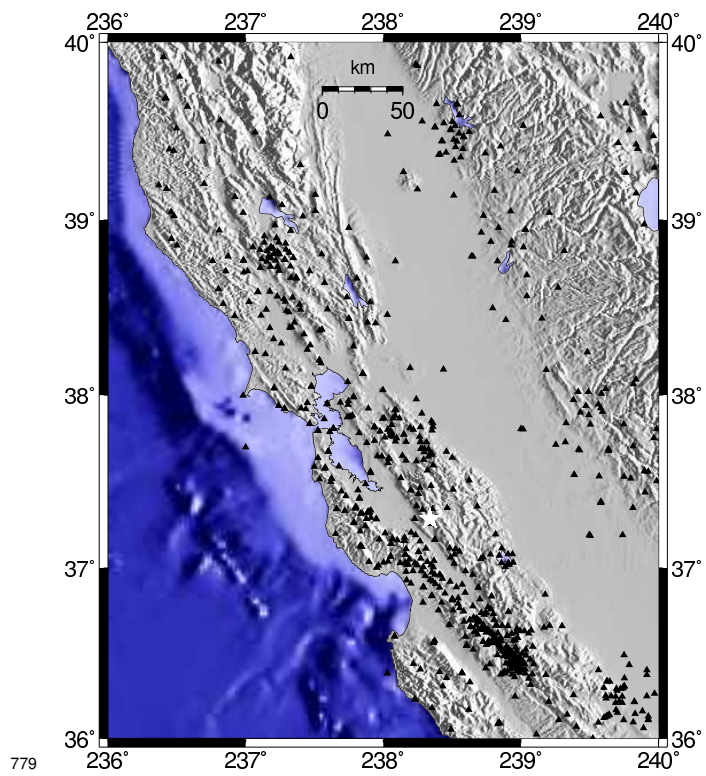
775

776 **Figure 3:**

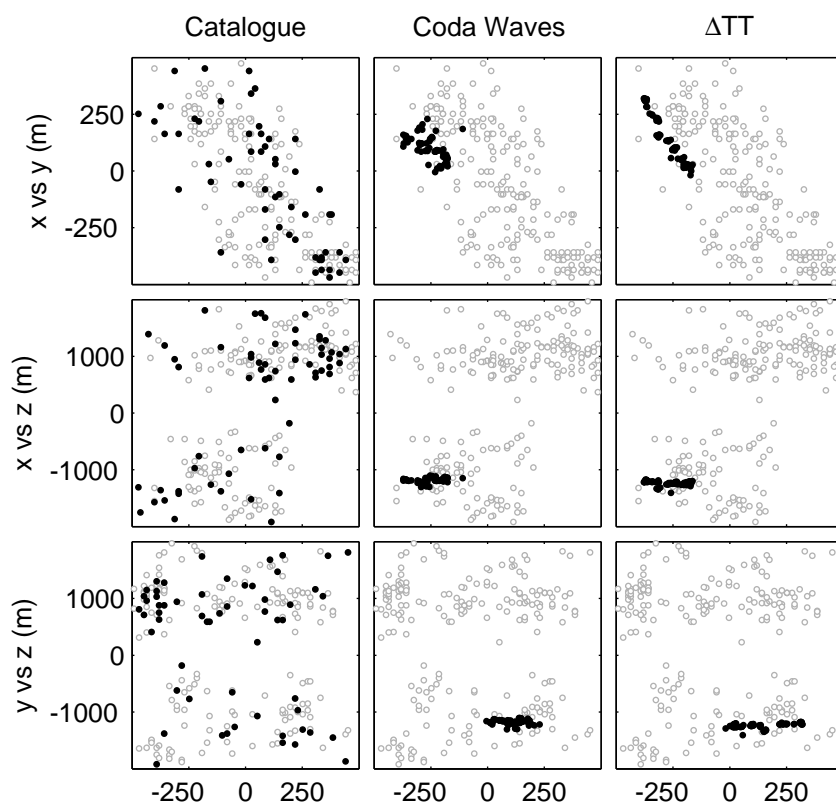


777

778 **Figure 4:**

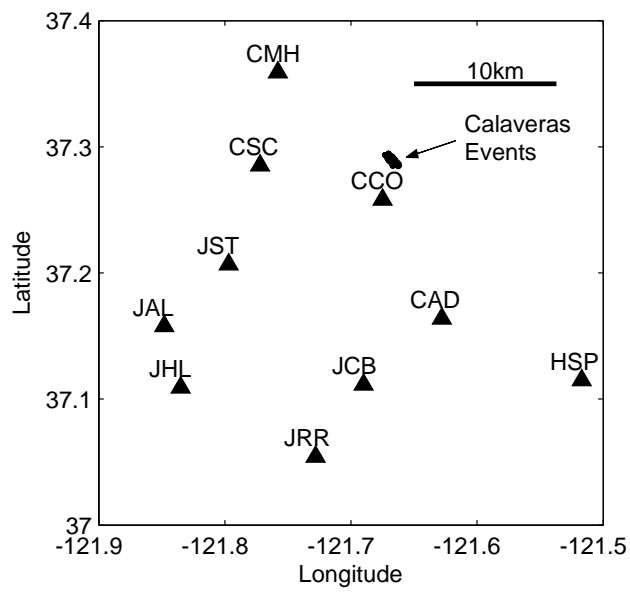


779 **Figure 5:**
780



781

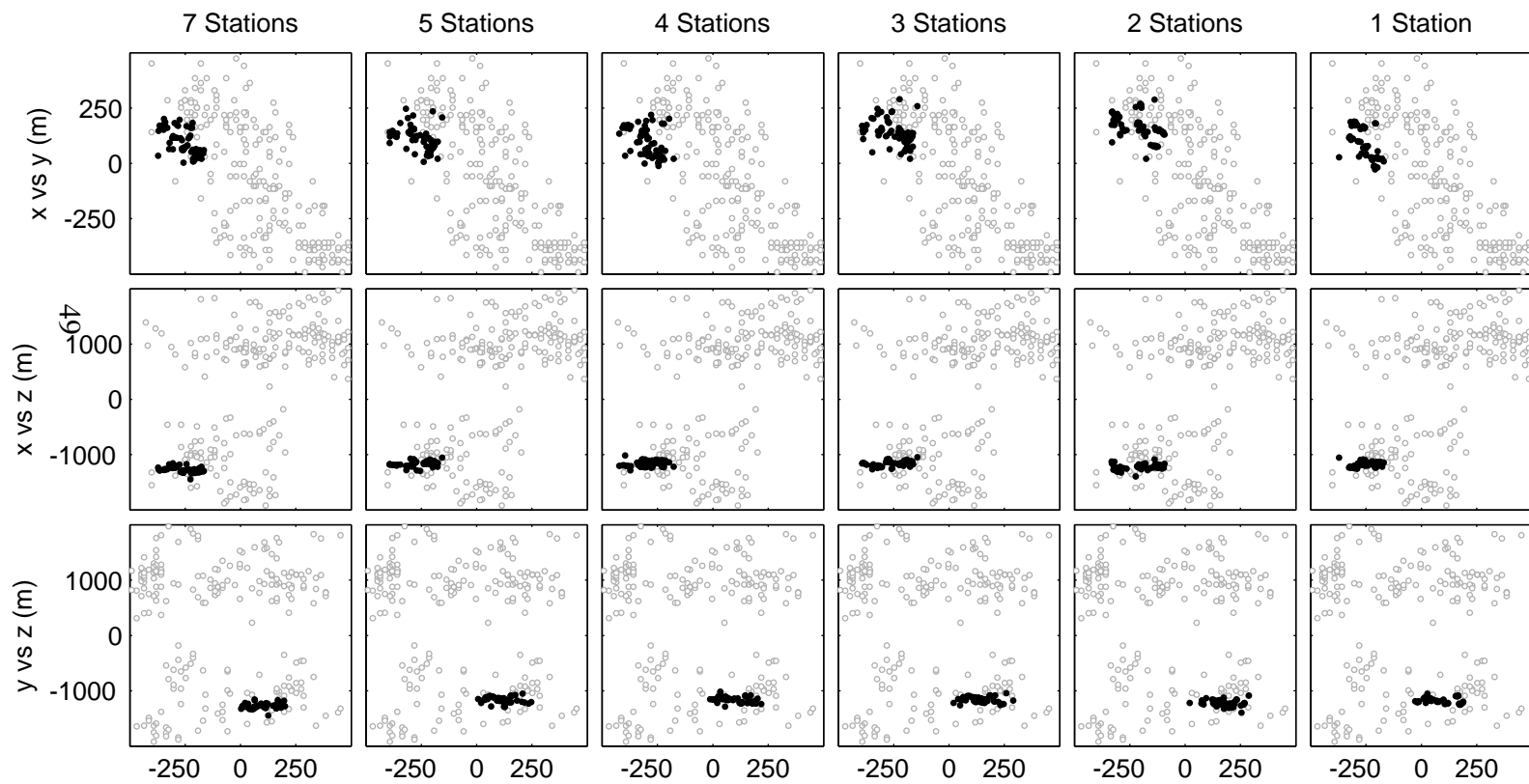
782 **Figure 6:**



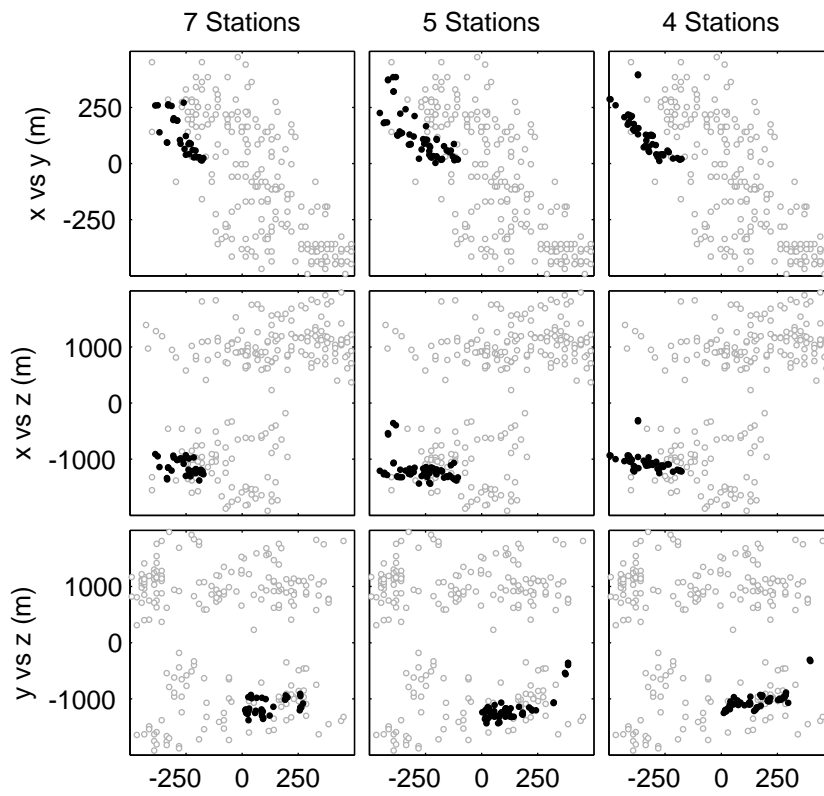
783

784 **Figure 7:**

Figure 8:

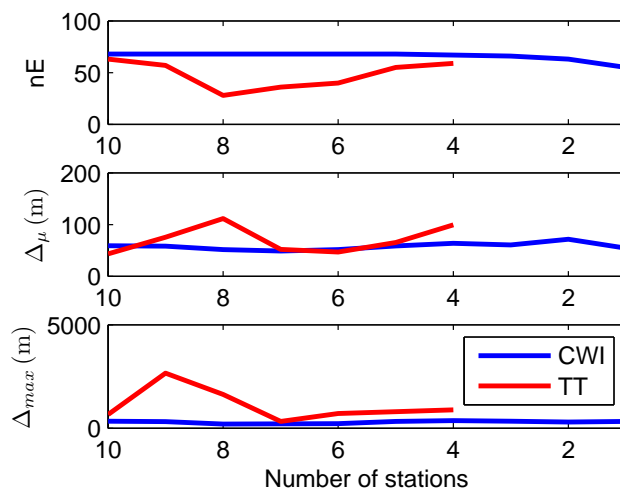


786



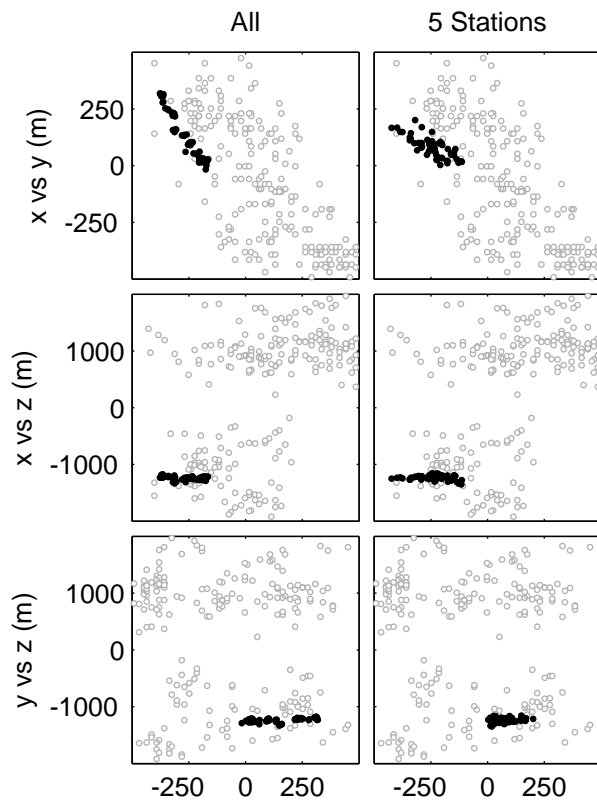
787

788 **Figure 9:**



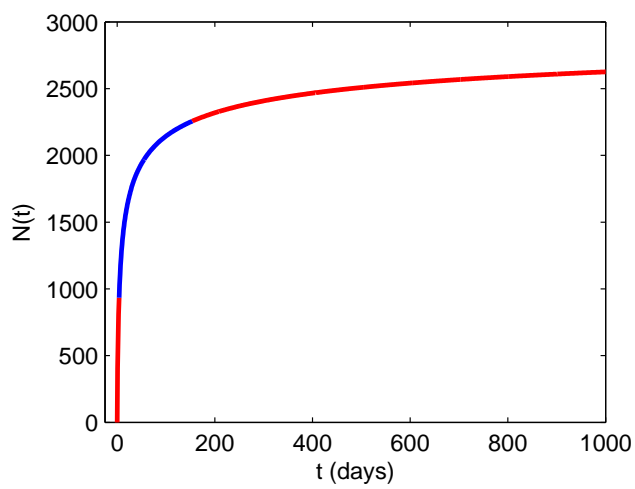
789

790 **Figure 10:**



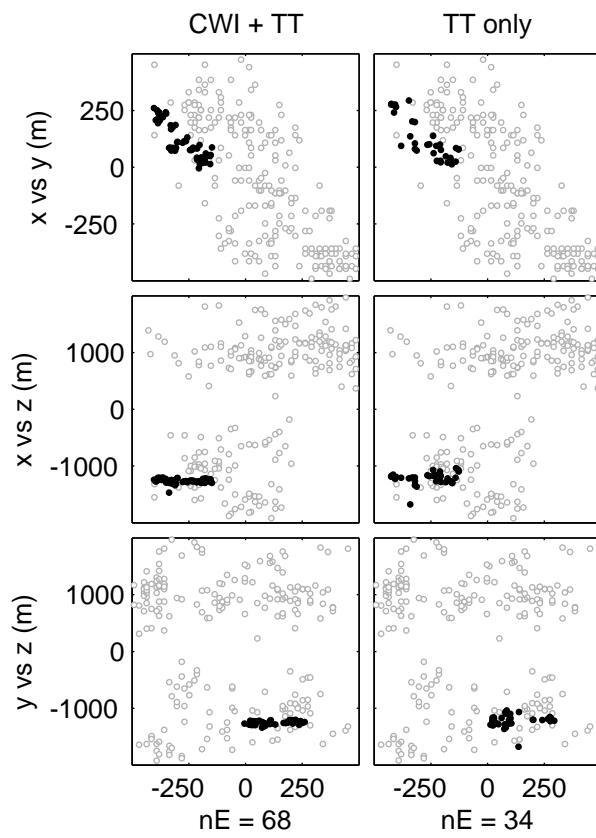
791

792 **Figure 11:**



793

794 **Figure 12:**



795

796 **Figure 13:**

The Likelihood

The likelihood $P(\tilde{\delta}_{CWIN}|\tilde{\delta}_t)$ used in equation (5) is given by

$$P(\tilde{\delta}_{CWIN}|\tilde{\delta}_t) = A(\tilde{\delta}_t)C(\bar{\mu}_N, \bar{\sigma}_N) \times \int_0^\infty B(\tilde{\delta}_t, \tilde{\delta}_{CWI})D(\tilde{\delta}_{CWI}, \bar{\sigma}_N, \bar{\mu}_N)d\tilde{\delta}_{CWI} \quad (35)$$

where $\tilde{\delta}_{CWI}$ is an estimate of CWI separation in the absence of noise,

$$A(\tilde{\delta}_t) = \frac{1}{(1 - \Phi_{\mu_1, \sigma_1}(0))\sigma_1\sqrt{2\pi}}, \quad (36)$$

$$B(\tilde{\delta}_t, \tilde{\delta}_{CWI}) = e^{\frac{-(\tilde{\delta}_{CWI} - \mu_1)^2}{2\sigma_1^2}}, \quad (37)$$

$$C(\bar{\mu}_N, \bar{\sigma}_N) = \frac{1}{(1 - \Phi_{\bar{\mu}_N, \bar{\sigma}_N}(0))\sigma_N\sqrt{2\pi}}, \quad (38)$$

$$D(\tilde{\delta}_{CWI}, \bar{\sigma}_N, \bar{\mu}_N) = e^{\frac{-(\tilde{\delta}_{CWI} - \bar{\mu}_N)^2}{2\bar{\sigma}_N^2}} \quad (39)$$

and $\Phi_{\mu, \sigma}(x)$ is the cumulative Gaussian distribution function

$$\Phi_{\mu, \sigma}(x) = \frac{1}{\sigma\sqrt{2\pi}} \int_{-\infty}^x e^{\frac{-(s-\mu)^2}{2\sigma^2}} ds \quad (40)$$

(Robinson *et al.*, 2011). The parameters μ_1 and σ_1 used in equation (36) are defined by the expressions

$$\mu_1(\tilde{\delta}_t) = a_1 \frac{a_2 \tilde{\delta}_t^{a_4} + a_3 \tilde{\delta}_t^{a_5}}{a_2 \tilde{\delta}_t^{a_4} + a_3 \tilde{\delta}_t^{a_5} + 1} \quad (41)$$

and

$$\sigma_1(\tilde{\delta}_t) = c + a_1 \frac{a_2 \tilde{\delta}_t^{a_4} + a_3 \tilde{\delta}_t^{a_5}}{a_2 \tilde{\delta}_t^{a_4} + a_3 \tilde{\delta}_t^{a_5} + 1} \quad (42)$$

with coefficients a_1 to a_5 and c defined in Table 4. The parameters $\bar{\mu}_N$ and $\bar{\sigma}_N$ used in equation (38) are obtained by finding the values which minimize the difference

Table 4: Coefficients for equations (41) and (42).

$\mu_1(\tilde{\delta}_t)$	$\sigma_1(\tilde{\delta}_t)$
$a1 = 0.4661$	$a1 = 0.1441$
$a2 = 48.9697$	$a2 = 101.0376$
$a3 = 2.4693$	$a3 = 120.3864$
$a4 = 4.2467$	$a4 = 2.8430$
$a5 = 1.1619$	$a5 = 6.0823$
	$c = 0.017$

in a least squares sense between the noisy CWI estimates $\tilde{\delta}_{CWIN}$ computed from the waveforms and the positively bounded Gaussian density function

$$P(\tilde{\delta}_{CWIN}|\tilde{\delta}_t, \tilde{\delta}_{CWI}) = \frac{1}{(1-\Phi_{\tilde{\mu}_N, \tilde{\sigma}_N}(0))\tilde{\sigma}_N\sqrt{2\pi}} e^{\frac{-(\tilde{\delta}_{CWIN}-\tilde{\mu}_N)^2}{2\tilde{\sigma}_N^2}} \quad (43)$$

with $\tilde{\delta}_{CWIN} \geq 0$.

Derivatives

The derivatives of $L(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_N)$

$$\frac{\partial L}{\partial \hat{x}_1}, \frac{\partial L}{\partial \hat{y}_1}, \frac{\partial L}{\partial \hat{z}_1}, \frac{\partial L}{\partial \hat{x}_2}, \frac{\partial L}{\partial \hat{y}_2}, \frac{\partial L}{\partial \hat{z}_2}, \dots, \frac{\partial L}{\partial \hat{x}_N}, \frac{\partial L}{\partial \hat{y}_N}, \frac{\partial L}{\partial \hat{z}_N} \quad (44)$$

are required by the Polak-Ribiere algorithm. These are used to guide the optimization procedure towards the values of $(\mathbf{e}_1, \mathbf{e}_2, \dots, \mathbf{e}_N)$ which minimize L .

The equations for the derivatives are convoluted so we build them gradually.

We start with an expression for δ_t , the wavelength normalized separation between

828 two events $\mathbf{e}_p = (\hat{x}_p, \hat{y}_p, \hat{z}_p)$ and $\mathbf{e}_q = (\hat{x}_q, \hat{y}_q, \hat{z}_q)$

$$829 \quad \delta_t = \frac{f_{dom}}{v_s} \sqrt{(\hat{x}_p - \hat{x}_q)^2 + (\hat{y}_p - \hat{y}_q)^2 + (\hat{z}_p - \hat{z}_q)^2}, \quad (45)$$

830 where f_{dom} is the dominant frequency of the waveforms and v_s is the velocity
831 between the events. Expression 45 has derivatives

$$832 \quad \begin{aligned} \frac{\partial \tilde{\delta}_t}{\partial \hat{x}_p} &= \frac{f_{dom}^2 (\hat{x}_p - \hat{x}_q)}{v_s^2 \tilde{\delta}_t}, \quad \frac{\partial \tilde{\delta}_t}{\partial \hat{y}_p} = \frac{f_{dom}^2 (\hat{y}_p - \hat{y}_q)}{v_s^2 \tilde{\delta}_t}, \\ \frac{\partial \tilde{\delta}_t}{\partial \hat{z}_p} &= \frac{f_{dom}^2 (\hat{z}_p - \hat{z}_q)}{v_s^2 \tilde{\delta}_t}, \quad \frac{\partial \tilde{\delta}_t}{\partial \hat{x}_q} = \frac{f_{dom}^2 (\hat{x}_q - \hat{x}_p)}{v_s^2 \tilde{\delta}_t}, \\ \frac{\partial \tilde{\delta}_t}{\partial \hat{y}_q} &= \frac{f_{dom}^2 (\hat{y}_q - \hat{y}_p)}{v_s^2 \tilde{\delta}_t}, \quad \frac{\partial \tilde{\delta}_t}{\partial \hat{z}_q} = \frac{f_{dom}^2 (\hat{z}_q - \hat{z}_p)}{v_s^2 \tilde{\delta}_t}. \end{aligned} \quad (46)$$

833 For brevity we focus the following derivation in terms of \hat{x}_p . The remaining terms
834 for \mathbf{e}_p (i.e. \hat{y}_p and \hat{z}_p) can be computed by following the same procedure. The
835 derivatives for \mathbf{e}_q can be attained by exploiting the symmetry

$$836 \quad \frac{\partial \tilde{\delta}_t}{\partial \hat{x}_q} = -\frac{\partial \tilde{\delta}_t}{\partial \hat{x}_p}. \quad (47)$$

837 The chain rule gives

$$838 \quad \frac{\partial \mu_1}{\partial \hat{x}_p} = \frac{\partial \mu_1}{\partial \tilde{\delta}_t} \frac{\partial \tilde{\delta}_t}{\partial \hat{x}_p} \quad (48)$$

839 where differentiating equation (41) gives

$$840 \quad \frac{\partial \mu_1}{\partial \tilde{\delta}_t} = a_1 \frac{a_2 a_4 \tilde{\delta}_t^{a_4-1} + a_3 a_5 \tilde{\delta}_t^{a_5-1}}{(a_2 \tilde{\delta}_t^{a_4} + a_3 \tilde{\delta}_t^{a_5} + 1)^2}. \quad (49)$$

841 Similarly, we have

$$842 \quad \frac{\partial \sigma_1}{\partial \hat{x}_p} = \frac{\partial \sigma_1}{\partial \tilde{\delta}_t} \frac{\partial \tilde{\delta}_t}{\partial \hat{x}_p} \quad (50)$$

843 where $\frac{\partial \sigma_1}{\partial \tilde{\delta}_t}$ has the identical form as 49 with different constants a_1, a_2, \dots, a_5 (see
844 table 4).

845 The cumulative Gaussian distribution function 40 is

$$846 \quad \Phi_{\mu_1, \sigma_1}(0) = \frac{1}{\sigma_1 \sqrt{2\pi}} \int_{-\infty}^0 e^{-\frac{(s-\mu_1)^2}{2\sigma_1^2}} ds \quad (51)$$

847 which has derivative

$$848 \quad \frac{\partial \Phi_{\mu_1, \sigma_1}(0)}{\partial \hat{x}_p} = \frac{\sigma_1 \int_{-\infty}^0 \frac{\partial g}{\partial \hat{x}_p} e^g ds - \frac{\partial \sigma_1}{\partial \hat{x}_p} \int_{-\infty}^0 e^g ds}{\sigma_1^2 \sqrt{2\pi}}, \quad (52)$$

849 where

$$850 \quad g = \frac{-(s - \mu_1)^2}{2\sigma_1^2} \quad (53)$$

851 and

$$852 \quad \frac{\partial g}{\partial \hat{x}_p} = \frac{4\sigma_1^2(s - \mu_1) \frac{\partial \mu_1}{\partial \hat{x}_p} + 4\sigma_1 \frac{\partial \sigma_1}{\partial \hat{x}_p} (s - \mu_1)^2}{4\sigma_1^4}. \quad (54)$$

853 Now, we have all the pieces to compute the derivatives of $A = A(\delta_t)$ and

854 $B = B(\delta_t, \delta_{CWI})$ as follows

$$855 \quad \frac{\partial A}{\partial \hat{x}_p} = - \frac{-\frac{\partial \Phi_{\mu_1, \sigma_1}(0)}{\partial \hat{x}_p} \sigma_1 + (1 - \Phi_{\mu_1, \sigma_1}(0)) \frac{\partial \sigma_1}{\partial \hat{x}_p}}{(1 - \Phi_{\mu_1, \sigma_1}(0))^2 \sigma_1^2 \sqrt{2\pi}} \quad (55)$$

856 and

$$857 \quad \frac{\partial B}{\partial \hat{x}_p} = e^h \frac{\partial h}{\partial \hat{x}_p}, \quad (56)$$

858 where

$$859 \quad h = \frac{-(\delta_{CWI} - \mu_1)^2}{2\sigma_1^2} \quad (57)$$

860 and

$$861 \quad \frac{\partial h}{\partial \hat{x}_p} = \frac{4\sigma_1^2(\delta_{CWI} - \mu_1) \frac{\partial \mu_1}{\partial \hat{x}_p} + 4(\delta_{CWI} - \mu_1)^2 \sigma_1 \frac{\partial \sigma_1}{\partial \hat{x}_p}}{4\sigma_1^4}. \quad (58)$$

862 Finally, we can differentiate the likelihood for an individual event pair

$$\begin{aligned} 863 \quad \frac{\partial P(\delta_{CWIN} | \tilde{\delta}_t)}{\partial \hat{x}_p} &= \frac{\partial A(\tilde{\delta}_t)}{\partial \hat{x}_p} C(\bar{\mu}_N, \bar{\sigma}_N) \\ &\times \int_0^\infty B(\tilde{\delta}_t, \tilde{\delta}_{CWI}) D(\tilde{\delta}_{CWI}, \bar{\sigma}_N, \bar{\mu}_N) d\tilde{\delta}_{CWI} \\ &+ A(\tilde{\delta}_t) C(\bar{\mu}_N, \bar{\sigma}_N) \\ &\times \int_0^\infty \frac{\partial B(\tilde{\delta}_t, \tilde{\delta}_{CWI})}{\partial \hat{x}_p} D(\tilde{\delta}_{CWI}, \bar{\sigma}_N, \bar{\mu}_N) d\tilde{\delta}_{CWI} \end{aligned} \quad (59)$$

864 and for the logarithm we have

$$865 \quad \frac{\partial \ln [P(\delta_{CWIN}|\delta_t)]}{\partial \hat{x}_p} = \frac{1}{P(\delta_{CWIN}|\delta_t)} \frac{\partial P(\delta_{CWIN}|\delta_t)}{\partial \hat{x}_p}. \quad (60)$$

866 Thus, it follows that the derivative of L with respect to \hat{x}_p is given by

$$867 \quad \frac{\partial L(E_1, E_2, \dots, E_n)}{\partial \hat{x}_p} = - \sum_{i=p+1}^N \frac{\partial \ln [P(\delta_{CWIN}|E_p, E_i)]}{\partial \hat{x}_p} + \sum_{j=1}^{p-1} \frac{\partial \ln [P(\delta_{CWIN}|E_j, E_p)]}{\partial \hat{x}_p} \quad (61)$$

868 for a uniform prior. The change of sign in the middle (i.e. to addition) accounts for
 869 the change in order of the events under the conditional. Its inclusion here assumes
 870 the correct use of $\partial \tilde{\delta}_t / \partial \hat{x}_p$ or $\partial \tilde{\delta}_t / \partial \hat{x}_q$ when evaluating the left and right hand
 871 terms of the summation. The derivatives shown in this section appear complicated
 872 but are in practice trivial to compute numerically. Confidence in their accuracy
 873 is enhanced by demonstrating that the optimization procedure converges to the
 874 correct solution for a number of synthetic problems in 2 and 3 dimensions.