PytHHOn3D documentation

David Siedel Olivier Fandeur Thomas Helfer 01/08/2020

1 Introduction

The Pythhon3D library is a PDE solver implementing the HHO method in a generic and object-oriented fashion using python, with specific applications to solid mechanics: the structure of the code follows that of a usual FEM solver, namely:

- PytHHOn3D reads and parses a mesh file as input
- it takes a set of input data to set the PDE problem to solve (boundary conditions, volumetric forces, HHO elements characteristics)
- $\bullet\,$ PytHH0n3D builds the framework specific to HHO methods, and the HHO element for each element in the mesh
- it solves a linear system given a set of fourth order tensors (one for each element) as tangent matrices
- it returns the unknown values at vertices and quadrature points

2 Model problem: small strain linear elasticity

Problem setting Let Ω a domain in \mathbb{R}^d with Lipschitz boundaries $\partial\Omega$. Let the following regular functional spaces :

Definition: Sobolev space

$$H^1(\Omega;\mathbb{R}^d) = \left\{\underline{\boldsymbol{u}} \in L^2(\Omega;\mathbb{R}^d), \nabla\underline{\boldsymbol{u}} \in L^2(\Omega;\mathbb{R}^{d\times d})\right\}$$

Definition: Divergence Sobolev space

$$H^1_{div}(\Omega;\mathbb{R}^{d\times d}) = \left\{ \underline{\tau} \in L^2(\Omega;\mathbb{R}^{d\times d}), \nabla \cdot \underline{\tau} \in L^2(\Omega;\mathbb{R}^d) \right\}$$

Problem Let the small strain linear elastic problem for u such that:

$$\begin{cases} \nabla \cdot \underline{\sigma} = -\underline{f} & \text{in } \Omega \\ \underline{\sigma} = \underline{t} \cdot \underline{n} & \text{on } \partial \Omega_{N} \\ \underline{u} = \underline{u}_{D} & \text{on } \partial \Omega_{D} \\ \underline{\sigma} = 2\mu \underline{\varepsilon} + \lambda \operatorname{Tr}(\underline{\varepsilon}) = \underline{\mathbb{C}} : \underline{\varepsilon} & \text{in } \Omega \\ \underline{\varepsilon} = \frac{1}{2} (\nabla \underline{u} + \nabla^{T} \underline{u}) = \nabla^{s} \underline{u} & \text{in } \Omega \end{cases}$$

$$(1)$$

Where

- $\sigma \in H^1_{div}(\Omega; \mathbb{R}^{d \times d})$ is the symmetric Cauchy stress tensor
- $\underline{\boldsymbol{u}} \in H^1(\Omega; \mathbb{R}^d)$ is the displacement field
- $\underline{\boldsymbol{t}} \in H^1(\Omega; \mathbb{R}^d)$ are the applied forces
- $\underline{\boldsymbol{u}}_D \in H^1(\Omega; \mathbb{R}^d)$ is the imposed displacement
- λ and μ are Lamé coefficients of the material

Weak form The weak form of (1) reads as the Principle of Virtual Works:

Theorem: Principle of Virtual Works (PVW)

find $\boldsymbol{u} \in H^1(\Omega; \mathbb{R}^d)$ such that :

$$\begin{cases} \int_{\Omega} \boldsymbol{\sigma} : \nabla^{s} \underline{\boldsymbol{v}} = \int_{\Omega} \underline{\boldsymbol{f}} \cdot \underline{\boldsymbol{v}} + \int_{\partial \Omega_{N}} \underline{\boldsymbol{t}} \cdot \underline{\boldsymbol{v}} & \forall \underline{\boldsymbol{v}} \in H^{1}(\Omega; \mathbb{R}^{d}) \\ \underline{\boldsymbol{u}} = \underline{\boldsymbol{u}}_{D} & \text{on } \partial \Omega_{D} \end{cases}$$
(2)

Discretization: the Finite Element Method (FEM) The FEM method consists in descretizing (2) in both physical and functional spaces through Lagrange polynbomials bases: defining the usual shape functions, that are polynomials of a given order k_c , such that they are of value 1 at a given node and 0 elsewhere in Ω , one can put (2) in a matricial form, which enables to find an approximation of $\underline{\boldsymbol{u}}$ in (2) numerically.

Polynomial approximation Let $\mathbb{P}^{k_c}(\Omega; \mathbb{R}^d) \subset H^1(\Omega; \mathbb{R}^d)$ a polynomial space of order k_c . Let N_k^d the dimension of $\mathbb{P}^{k_c}(\Omega; \mathbb{R}^d)$, and $(\underline{\phi}_m)_{m \leq N_{k_c}^d}$ a basis of $\mathbb{P}^{k_c}(\Omega; \mathbb{R}^d)$. N_k^d is then the number of vectors $\underline{\phi}_m$ composing

References