

Astronomy 401/Physics 903
Lecture 11
Galaxy Spectra I

The colors of a galaxy (the integrated light) can tell us about its stellar population, but we can learn even more information from its spectrum. Galaxies contain gas and stars; unsurprisingly, the spectrum of a galaxy looks like some mixture of the spectra of stars and clouds of gas, and it varies depending on the type of the galaxy.

1 The production of spectral lines and the Bohr atom

Before the reasons for it were understood, it was observed that hydrogen gas emitted lines with wavelengths given by

$$\frac{1}{\lambda} = R_H \left(\frac{1}{m^2} - \frac{1}{n^2} \right), \quad (1)$$

where m and n are integers with $m < n$ and R_H is the experimentally determined Rydberg constant for hydrogen.

This was explained by Niels Bohr, who proposed that the angular momentum of the orbiting electron was quantized: it could only have certain values given by multiples of \hbar : $L = n\hbar/2\pi = n\hbar$.

The kinetic and electrical potential energies of the atom can be derived classically. We will use the reduced mass μ ,

$$\mu = \frac{m_e m_p}{m_e + m_p} = 0.999 m_e \approx m_e \quad (2)$$

and the total mass $M = m_p + m_e \approx m_p$. The hydrogen atom can then be modeled as a proton of mass M that is at rest and an electron of mass μ that follows a circular orbit of radius r about the proton.

To find the kinetic energy, we start with the electric force between the proton and electron given by Coulomb's law,

$$F = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2}. \quad (3)$$

This electric force produces the centripetal acceleration of the electron, v^2/r , so, from Newton's second law,

$$\frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} = -\mu \frac{v^2}{r} \quad (4)$$

where the negative sign comes from the fact that the acceleration is directed inward. With $q_1 = e$ and $q_2 = -e$, we have

$$-\frac{1}{4\pi\epsilon_0} \frac{e^2}{r^2} = -\mu \frac{v^2}{r}. \quad (5)$$

This expression can be solved for the kinetic energy, $\frac{1}{2}\mu v^2$:

$$K = \frac{1}{2}\mu v^2 = \frac{1}{8\pi\epsilon_0} \frac{e^2}{r}. \quad (6)$$

The electrical potential energy can be found through a derivation analogous to that of gravitational potential energy. The result is

$$U = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r} = -2K. \quad (7)$$

The total energy of the atom is then

$$E = K + U = K - 2K = -K = -\frac{1}{8\pi\epsilon_0} \frac{e^2}{r}. \quad (8)$$

We then introduce the quantization of angular momentum

$$L = \mu v r = n\hbar \quad (9)$$

and use it to rewrite the kinetic energy:

$$K = \frac{1}{8\pi\epsilon_0} \frac{e^2}{r} = \frac{1}{2} \mu v^2 = \frac{1}{2} \frac{(\mu v r)^2}{\mu r^2} = \frac{1}{2} \frac{(n\hbar)^2}{\mu r^2}. \quad (10)$$

We can solve this for r to see that only certain values are allowed:

$$r_n = \frac{4\pi\epsilon_0\hbar^2}{\mu e^2} n^2 = a_0 n^2, \quad (11)$$

where $a_0 = 5.29 \times 10^{-11} \text{ m} = 0.0529 \text{ nm}$ is the Bohr radius. Only certain orbits are allowed, with $r = a_0, 4a_0, 9a_0$, etc.

We then insert this expression for the radius into our expression for the total energy of the atom to find the allowed energies:

$$E_n = -\frac{\mu e^4}{32\pi^2\epsilon_0^2\hbar^2} \frac{1}{n^2} = -13.6 \text{ eV} \frac{1}{n^2}. \quad (12)$$

The integer n is the principal quantum number. When the electron is in the ground state, $n = 1$ and it takes at least 13.6 eV to ionize the atom. When it's in the first excited state, $n = 2$ and the atom's energy is higher, $E_2 = -13.6/4 \text{ eV} = -3.40 \text{ eV}$.

Electrons can move between energy levels. If an electron goes from a higher state to a lower state, it emits a single photon with energy $E = E_{\text{high}} - E_{\text{low}}$, or

$$E = \frac{hc}{\lambda} = E_{n,\text{high}} - E_{n,\text{low}}. \quad (13)$$

Using the expression for E_n , we find

$$\frac{1}{\lambda} = \frac{\mu e^4}{64\pi^3\epsilon_0^2\hbar^3 c} \left(\frac{1}{n_{\text{low}}^2} - \frac{1}{n_{\text{high}}^2} \right), \quad (14)$$

which is the observed relationship we started with, with the combination of constants equal to the Rydberg constant. Spectral lines of hydrogen appear in different series: the Balmer series were the first observed, with $n_{\text{low}} = 2$. The Lyman series has $n_{\text{low}} = 1$ and the Paschen series has $n_{\text{low}} = 3$. The most commonly observed lines are the Balmer lines in the optical: $H\alpha$ is the strongest at 6563 Å, then $H\beta$ at 4861 Å, $H\gamma$, etc.

Electrons can also move from a lower to a higher energy level by absorbing a photon with energy equal to the energy difference between two levels.