

# Introduction of our theoretical framework

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## INTRODUCTION

We investigate the structural changes of the  $^{18-28}\text{O}$  isotopes, both even and odd isotopes, by examining the energies of their low-lying states. This is done by developing our own shell model program and comparing the energies it generates to the NushellX@MSU program [1]. We use a frozen  $^{16}\text{O}$  core, having 8 protons and 8 neutrons at the  $sp$ -shell ( $0s_{1/2}, 0p_{3/2}, 0p_{1/2}$ ). More neutrons are then excited in the  $sd$ -shell ( $0d_{5/2}, 1s_{1/2}, 0d_{3/2}$ ), which serves as our model space. We use all possible configurations in these orbits and work in a harmonic oscillator basis with spin-orbit splitting. The model space Hamiltonian reads

$$\hat{H} = \hat{H}_0 + \hat{H}_I, \quad (1)$$

where  $\hat{H}_0$  is the unperturbed one-body Hamiltonian

$$\hat{H}_0 = \sum_{p,q} \langle p|h_0|q \rangle \hat{a}_p^\dagger \hat{a}_q, \quad (2)$$

and  $\hat{H}_I$  is the pair breaking interaction

$$\hat{H}_I = \sum_{p,q,r,s} \langle p,q|V|r,s \rangle \hat{P}_{p,q}^+ \hat{P}_{r,s}^-, \quad (3)$$

with

$$\hat{P}_{p,q}^+ = \sum_{p,q} \hat{a}_p^\dagger \hat{a}_q^\dagger, \quad \hat{P}_{r,s}^- = \sum_{r,s} \hat{a}_r \hat{a}_s, \quad (4)$$

where  $p, q, r, s$  represent all possible single-particle states (SPS) quantum numbers. The Hamiltonian is rotationally invariant

For the two-body matrix elements (TBME) (**define them as in [2]**) in Eq. (3) we use the USDB interaction TBME [2] (Tables I and II in [2], given in  $J$ -scheme, for  $T = 1, 0$  respectively), and the single-particle energies (SPE) ( $1s_{1/2}, 0d_{3/2}, 0d_{5/2}$ ) are  $(-3.2079, 2.1117, -3.9257)$ . We work in  $M$ -scheme, where the SPS are ordered as given in Table I.

### The transformation from $J$ -scheme reads?

Using the SPS we construct the appropriate Slater determinants,  $|\psi\rangle$ , according to the number of particles (even or odd) which we place in the  $sd$ -shell.

TABLE I. Single particle energies of the  $sd$ -shell in the  $M$ -scheme basis with their corresponding quantum numbers:  $N$ , the principle quantum number;  $\ell$ , the orbital angular momentum;  $J$ , the total angular momentum;  $M_j$ , the total angular momentum projection to the  $z$  axis..

index	$N$	$\ell$	$J$	$M_j$	SPE
1	1	0	1	$-1/2$	-3.20790
2	1	0	1	$+1/2$	-3.20790
3	0	2	3	$-3/2$	2.11170
4	0	2	3	$-1/2$	2.11170
5	0	2	3	$+1/2$	2.11170
6	0	2	3	$+3/2$	2.11170
7	0	2	5	$-5/2$	-3.92570
8	0	2	5	$-3/2$	-3.92570
9	0	2	5	$-1/2$	-3.92570
10	0	2	5	$+1/2$	-3.92570
11	0	2	5	$+3/2$	-3.92570
12	0	2	5	$+5/2$	-3.92570

## My TOC

### Background

1. Description for the Oxygen isotopes - the physical phenomenon we investigate.
2. The  $1s0d$ -shell model space.
3. Pairing Hamiltonian and pair breaking Hamiltonian (We use second quantization. Is this only in  $m$ -scheme?).

$$\hat{H} = \hat{H}_0 + \hat{V} \quad (5)$$

$$\hat{H}_0 = \xi \sum_{p,\sigma} (p-1) \hat{a}_{p\sigma}^\dagger \hat{a}_{p\sigma}, \quad (6)$$

$$\hat{V} = \sum_{p \leq q} \langle p|V|q \rangle \hat{P}_p^+ \hat{P}_q^-, \quad (7)$$

$$\hat{P}_p^+ = \sum_{\sigma_p, \geq 0} \hat{a}_{p,\sigma_p}^\dagger \hat{a}_{p,-\sigma_p}^\dagger, \quad \hat{P}_p^- = \sum_{\sigma_p, \geq 0} \hat{a}_{p,-\sigma_p} \hat{a}_{p,\sigma_p}. \quad (8)$$

Aren't we using

$$\hat{V} = \sum_{p \leq q} \langle pq | V | rs \rangle \hat{P}_{p,q}^+ \hat{P}_{r,s}^-? \quad (9)$$

$$\hat{P}_{p,q}^+ = \sum_{\sigma_p, \sigma_q, \geq 0} \hat{a}_{p, \sigma_p}^\dagger \hat{a}_{q, -\sigma_q}^\dagger, \quad \hat{P}_{r,s}^- = \sum_{\sigma_r, \sigma_s, \geq 0} \hat{a}_{r, -\sigma_r} \hat{a}_{s, \sigma_s} ?? \quad (10)$$

### Introduction of our theoretical framework

1. The *sd*-shell serves as our model space.
2. The Hamiltonian we are using.
  - (a) Few words on the pair breaking interaction.
  - (b) The Hamiltonian is rotationally invariant.
3. SPS coupling to  $M_{tot}$ . For  $M_{tot} = 0$  ( $M_{tot} = 1/2$ ) we obtain all possible  $J$  values for even (odd) nuclei. Explain this when comparing our results to NushellX.
4. Matrix elements in m-scheme.

- (a) Translate Eq. (19) in [2] to M-scheme?
- (b) The USD Hamiltonian is defined by 63 *sd*-shell TBME and three SPE given in Table I of [1]. This is converted to the *M*-scheme and was given to us by Morten in the *sdshellint.dat* file. To this we added the refinement of Eq. (19) in [2].

5. Slater determinants.
6. We solve the many-body Schrodinger eq.
7. The different NushellX interactions:
  - (a) USD (the original, by Winldenthal)
  - (b) USDB.
  - (c) USDA.
  - (d) CCEI (Gustav's).

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- [1] B. A. Brown and W. Rae, [Nucl. Data Sheets \*\*120\*\*, 115 \(2014\)](#).
  - [2] B. A. Brown and W. A. Richter, [Phys. Rev. C \*\*74\*\*, 034315 \(2006\)](#).