

1 **The Sensitivity to Oscillation Parameters**
2 **from a Simultaneous Beam and**
3 **Atmospheric Neutrino Analysis that**
4 **combines the T2K and SK Experiments**



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10 *Doctor of Philosophy*
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Abstract

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Acknowledgements

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Introduction

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Neutrino Oscillation Physics

¹²⁹ When first proposed, neutrinos were expected to be approximately massless
¹³⁰ fermions that only interact through weak and gravitational forces. This meant
¹³¹ they were very difficult to detect as they can pass through significant amounts
¹³² of matter without interacting. Despite this, experimental neutrino physics has
¹³³ developed many different detection techniques and observed neutrinos from
¹³⁴ both natural and artificial sources. In direct tension with Standard Model physics,
¹³⁵ neutrinos have been determined to oscillate between different lepton flavours,
¹³⁶ requiring them to have mass.

¹³⁷ The observation techniques which led to the discovery of the neutrino are doc-
¹³⁸ umented in section 2.1. The theory underpinning neutrino oscillation is described
¹³⁹ in section 2.2 and includes the approximations which can be made to simplify
¹⁴⁰ the understanding of neutrino oscillation in the two-flavour approximation. Past,
¹⁴¹ current, and future neutrino experiments are detailed in section 2.3, including the
¹⁴² reactor, atmospheric, and long-baseline accelerator neutrino sources that have
¹⁴³ been used to successfully constrain oscillation parameters. Finally, the current
¹⁴⁴ state of oscillation parameter measurements are summarised in section 2.4.

2.1 Discovery of Neutrinos

At the start of the 20th century, the electrons emitted from the β -decay of the nucleus were found to have a continuous energy spectrum [1, 2]. This observation seemingly broke the energy conservation invoked within that period's nuclear models. In 1930, Pauli provided a solution to this problem in the form of a new particle, the neutrino (originally termed "neutron"). It was theorized to be an electrically neutral spin-1/2 fermion with a mass smaller than that of the electron [3]. This neutrino was emitted with the electron in β -decay to alleviate the apparent breaking of energy conservation. As a predecessor of today's weak interaction model, Fermi's theory of β -decay developed the understanding by coupling the four constituent particles: electron, proton, neutron, and neutrino, into a quantitative model [4].

Whilst Pauli was not convinced of the ability to detect neutrinos, the first observations of the particle were made in the mid-1950s when neutrinos from a reactor were observed via the inverse β -decay (IBD) process, $\bar{\nu}_e + p \rightarrow n + e^+$ [5, 6]. The detector consisted of two parts: a neutrino interaction medium and a liquid scintillator. The interaction medium was built from two water tanks, loaded with cadmium chloride to allow for increased efficiency in the detection of neutron capture. The positron emitted from IBD annihilates, $e^+ + e^- \rightarrow 2\gamma$, generating a prompt signal and the neutron is captured on the cadmium via $n + ^{108}Cd \rightarrow ^{109*}Cd \rightarrow ^{109}Cd + \gamma$, producing a delayed signal. An increase in the coincidence rate was observed when the reactor was operating which was interpreted as interactions from neutrinos generated in the reactor.

After the discovery of the ν_e , the question of how many flavours of neutrino exist was asked. In 1962, a measurement of the ν_μ was conducted at the Brookhaven National Laboratory [7]. A proton beam was directed at a beryllium target, generating pions which then decayed via $\pi^\pm \rightarrow \mu^\pm + (\nu_\mu, \bar{\nu}_\mu)$, and the subsequent interactions of the ν_μ were observed. As the subsequent interaction of the neutrino generated muons rather than electrons, it was determined that

¹⁷⁴ the ν_μ was fundamentally different from ν_e . The final observation to be made
¹⁷⁵ was that of the ν_τ from the DONUT experiment [8]. Three neutrinos seem the
¹⁷⁶ obvious solution as it mirrors the known number of charged leptons (as they form
¹⁷⁷ weak isospin doublets) but there could be evidence of more. Several neutrino
¹⁷⁸ experiments have found anomalous results [9, 10] which could be attributed
¹⁷⁹ to “sterile” neutrinos. These hypothesised particles are not affected by gauge
¹⁸⁰ interactions in the Standard Model so their presence can only be inferred through
¹⁸¹ the observation of non-standard oscillation modes. However, cosmological
¹⁸² observations indicate the number of neutrino species $N_{eff} = 2.99 \pm 0.17$ [11], as
¹⁸³ measured from the cosmic microwave background power spectrum. LEP also
¹⁸⁴ measured the number of active neutrino flavours to be $N_\nu = 2.9840 \pm 0.0082$ [12]
¹⁸⁵ from measurements of the Z-decay width, but this does not strongly constrain
¹⁸⁶ the number of sterile neutrinos.

¹⁸⁷ 2.2 Theory of Neutrino Oscillation

¹⁸⁸ A neutrino generated with lepton flavour α can change into a different lepton
¹⁸⁹ flavour β after propagating some distance. This phenomenon is called neutrino
¹⁹⁰ oscillation and requires that neutrinos must have a non-zero mass. This behaviour
¹⁹¹ has been characterised by the Pontecorvo-Maki-Nakagawa-Sakata (PMNS) [13–
¹⁹² 15] mixing matrix which describes how the flavour and mass of neutrinos are
¹⁹³ associated. This is analogous to the Cabibbo-Kobayashi-Maskawa (CKM) [16]
¹⁹⁴ matrix measured in quark physics.

¹⁹⁵ 2.2.1 Three Flavour Oscillations

¹⁹⁶ The PMNS parameterisation defines three flavour eigenstates, ν_e , ν_μ and ν_τ
¹⁹⁷ (indexed ν_α), which are eigenstates of the weak interaction and three mass
¹⁹⁸ eigenstates, ν_1 , ν_2 and ν_3 (indexed ν_i). Each mass eigenstate is the superposition

¹⁹⁹ of all three flavour states,

$$|\nu_i\rangle = \sum_{\alpha} U_{\alpha i} |\nu_{\alpha}\rangle. \quad (2.1)$$

²⁰⁰ Where U is the 3×3 PMNS matrix which is unitary and connects the mass
²⁰¹ and flavour eigenstates.

²⁰² The weak interaction, when interacting via a W^{\pm} boson, couples to flavour
²⁰³ eigenstates so neutrinos interact with leptons of the same flavour. The prop-
²⁰⁴ agation of a neutrino flavour eigenstate, in a vacuum, can be re-written as a
²⁰⁵ plane-wave solution to the time-dependent Schrödinger equation,

$$|\nu_{\alpha}(t)\rangle = \sum_i U_{\alpha i}^{*} |\nu_i\rangle e^{-i\phi_i}. \quad (2.2)$$

²⁰⁶ The ϕ_i term can be expressed in terms of the energy, E_i , and magnitude of the
²⁰⁷ three momenta, p_i , of the neutrino, $\phi_i = E_i t - p_i x$ (t and x being time and position
²⁰⁸ coordinates). The probability of observing a neutrino of flavour eigenstate β from
²⁰⁹ one which originated as flavour α can be calculated as,

$$P(\nu_{\alpha} \rightarrow \nu_{\beta}) = |\langle \nu_{\beta} | \nu_{\alpha}(t) \rangle|^2 = \sum_{i,j} U_{\alpha i}^{*} U_{\beta i} U_{\alpha j} U_{\beta j}^{*} e^{-i(\phi_j - \phi_i)} \quad (2.3)$$

²¹⁰ The term within the exponential can be represented as,

$$\phi_j - \phi_i = E_j t - E_i t - p_j x + p_i x. \quad (2.4)$$

²¹¹ For a relativistic particle, $E_i \gg m_i$, a Taylor series expansion means,

$$p_i = \sqrt{E_i^2 - m_i^2} \approx E_i - \frac{m_i^2}{2E_i}. \quad (2.5)$$

²¹² Making the approximations that neutrinos are relativistic, the mass eigenstates
²¹³ were created with the same energy and that $x = L$, where L is the distance
²¹⁴ travelled by the neutrino, Equation 2.4 then becomes

$$\phi_j - \phi_i = \frac{\Delta m_{ij}^2 L}{2E}, \quad (2.6)$$

where $\Delta m_{ij}^2 = m_i^2 - m_j^2$. This, combined with further use of unitarity relations results in Equation 2.3 becoming

$$\begin{aligned} P(\nu_\alpha \rightarrow \nu_\beta) &= \delta_{\alpha\beta} - 4 \sum_{i>j} \Re \left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^* \right) \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right) \\ &\quad + (-) 2 \sum_{i>j} \Im \left(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^* \right) \sin \left(\frac{\Delta m_{ij}^2 L}{2E} \right) \end{aligned} \quad (2.7)$$

Where $\delta_{\alpha\beta}$ is the Kronecker delta function and the negative sign on the last term is included for the oscillation probability of antineutrinos. As an important point to note, the observation of oscillation probability requires a non-zero value of Δm_{ij}^2 , which in turn requires that neutrinos have differing masses.

Typically, the PMNS matrix is parameterised into three mixing angles, a charge parity (CP) violating phase δ_{CP} , and two Majorana phases $\alpha_{1,2}$,

$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{Atmospheric, Accelerator}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13} e^{-i\delta_{CP}} & 0 & c_{13} \end{pmatrix}}_{\text{Reactor, Accelerator}} \times \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Reactor, Solar}} \underbrace{\begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Majorana}}. \quad (2.8)$$

Where $s_{ij} = \sin(\theta_{ij})$ and $c_{ij} = \cos(\theta_{ij})$. The oscillation parameters are often grouped: (1,2) as “solar”, (2,3) as “atmospheric” and (1,3) as “reactor”. Many neutrino experiments aim to measure the PMNS parameters from a wide array of origins, as is the purpose of this thesis.

The Majorana phase, $\alpha_{1,2}$, included within the fourth matrix in Equation 2.8 is only included for completeness. For an oscillation analysis experiment, any terms containing this phase disappear due to taking the expectation value of the PMNS matrix. Measurements of these phases can be performed by experiments searching for neutrino-less double β -decay [17].

232 A two-flavour approximation can be obtained when one assumes the third
233 mass eigenstate is degenerate with another. This results in the two-flavour
234 approximation being reasonable for understanding the features of the oscillation.
235 In this two-flavour case, the mixing matrix becomes,

$$U_{\text{2 Flav.}} = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & \cos(\theta) \end{pmatrix}. \quad (2.9)$$

236 This culminates in the oscillation probability,

$$\begin{aligned} P(\nu_\alpha \rightarrow \nu_\alpha) &= 1 - \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right), \\ P(\nu_\alpha \rightarrow \nu_\beta) &= \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right). \end{aligned} \quad (2.10)$$

237 Where $\alpha \neq \beta$. For a fixed neutrino energy, the oscillation probability is
238 a sinusoidal function depending upon the distance over which the neutrino
239 propagates. The frequency and amplitude of oscillation are dependent upon
240 $\Delta m^2 / 4E$ and $\sin^2 2\theta$, respectively. The oscillation probabilities presented thus far
241 assume $c = 1$, where c is the speed of light in a vacuum. In more familiar units, the
242 maximum oscillation probability for a fixed value of θ is given at $L[\text{km}] / E[\text{GeV}] \sim$
243 $1.27 / \Delta m^2$. It is this calculation that determines the best L/E value for a given
244 experiment to be designed around for measurements of a specific value of Δm^2 .

245 2.2.2 The MSW Effect

246 The theory of neutrino oscillation in a vacuum has been described in subsec-
247 tion 2.2.1. However, the beam neutrinos and atmospheric neutrinos originating
248 from below the horizon propagate through the matter in the Earth. The coherent
249 scattering of neutrinos from a material target modifies the Hamiltonian of the
250 system which results in a change in the oscillation probability. This modification
251 is termed the Mikheyev-Smirnov-Wolfenstein (MSW) effect [18–20]. This occurs
252 because charged current scattering ($\nu_e + e^- \rightarrow \nu_e + e^-$, propagated by a W boson)
253 only affects electron neutrinos whereas the neutral current scattering ($\nu_l + l^- \rightarrow$

²⁵⁴ $\nu_l + l^-$, propagated by a Z^0 boson) interacts through all neutrino flavours equally.
²⁵⁵ In the two-flavour approximation, the effective mixing parameter becomes

$$\sin^2(2\theta) \rightarrow \sin^2(2\theta_m) = \frac{\sin^2(2\theta)}{(A/\Delta m^2 - \cos(2\theta))^2 + \sin^2(2\theta)}, \quad (2.11)$$

²⁵⁶ where $A = 2\sqrt{2}G_F N_e E$, N_e is the electron density of the medium and G_F
²⁵⁷ is Fermi's constant. It is clear that there exists a value of $A = \Delta m^2 \cos(2\theta)$ for
²⁵⁸ $\Delta m^2 > 0$, which results in a divergent mixing parameter, colloquially called the
²⁵⁹ matter resonance. This resonance regenerates the electron neutrino component of
²⁶⁰ the neutrino flux [18–20]. The density at which the resonance occurs is given by

$$N_e = \frac{\Delta m^2 \cos(2\theta)}{2\sqrt{2}G_F E}. \quad (2.12)$$

²⁶¹ At densities lower than this critical value, the oscillation probability will
²⁶² be much closer to that of vacuum oscillation. For antineutrinos, $N_e \rightarrow -N_e$
²⁶³ [21]. The resonance occurring from the MSW effect depends on the sign of Δm^2 .
²⁶⁴ Therefore, any neutrino oscillation experiment which observes neutrinos and
²⁶⁵ antineutrinos which have propagated through matter can have some sensitivity
²⁶⁶ to the ordering of the neutrino mass eigenstates.

²⁶⁷ 2.3 Neutrino Oscillation Measurements

²⁶⁸ As evidence of beyond Standard Model physics, the 2015 Nobel Prize in Physics
²⁶⁹ was awarded to the Super-Kamiokande (SK) [22] and Sudbury Neutrino Ob-
²⁷⁰ servatory (SNO) [23] collaborations for the first definitive observation of solar
²⁷¹ and atmospheric neutrino oscillation [24]. Since then, the field has seen a wide
²⁷² array of oscillation measurements from a variety of neutrino sources. As seen
²⁷³ in subsection 2.2.1, the neutrino oscillation probability is dependent on the ratio
²⁷⁴ of the propagation baseline, L , to the neutrino energy, E . It is this ratio that
²⁷⁵ determines the type of neutrino oscillation a particular experiment is sensitive to.

²⁷⁶ As illustrated in Figure 2.1, there are many neutrino sources that span a
²⁷⁷ wide range of energies. The least energetic neutrinos are from reactor and

²⁷⁸ terrestrial sources at $O(1)$ MeV whereas the most energetic neutrinos originate
²⁷⁹ from atmospheric and galactic neutrinos of $> O(1)$ TeV.



Figure 2.1: The electro-weak cross-section for $\bar{\nu}_e + e^- \rightarrow \bar{\nu}_e + e^-$ scattering on free electrons from various natural and man-made neutrino sources, as a function of neutrino energy. Taken from [25]

²⁸⁰ 2.3.1 Solar Neutrinos

²⁸¹ Solar neutrinos are emitted from fusion reaction chains at the centre of the Sun.
²⁸² The solar neutrino flux, given as a function of neutrino energy for different
²⁸³ fusion and decay chains, is illustrated in Figure 2.2. Whilst proton-proton fusion
²⁸⁴ generates the largest flux of neutrinos, the neutrinos are low energy and are
²⁸⁵ difficult to reconstruct due to the IBD interaction threshold of 1.8MeV [26].
²⁸⁶ Consequently, most experiments focus on the neutrinos from the decay of 8B
²⁸⁷ (via $^8B \rightarrow ^8Be^* + e^+ + \nu_e$), which are higher energy.

²⁸⁸ The first measurements of solar neutrinos observed a significant reduction in
²⁸⁹ the event rate compared to predictions from the Standard Solar Model [28, 29]. A
²⁹⁰ proposed solution to this “solar neutrino problem” was $\nu_e \leftrightarrow \nu_\mu$ oscillations in a



Figure 2.2: The solar neutrino flux as a function of neutrino energy for various fusion reactions and decay chains as predicted by the Standard Solar Model. Taken from [27].

291 precursory version of the PMNS model [30]. The Kamiokande [31], Gallex [32]
 292 and Sage [33] experiments confirmed the ~ 0.5 factor deficit of solar neutrinos.

293 The conclusive solution to this problem was determined by the SNO col-
 294 laboration [34]. Using a deuterium water target to observe 8B neutrinos, the
 295 event rate of charged current (CC), neutral current (NC), and elastic scattering
 296 (ES) interactions (Given in Equation 2.13) was simultaneously measured. CC
 297 events can only occur for electron neutrinos, whereas the NC channel is agnostic
 298 to neutrino flavour, and the ES reaction has a small excess sensitivity for the
 299 detection of electron neutrino interactions. This meant that there were direct
 300 measurements of the ν_e and ν_x neutrino flux. It was concluded that the CC and
 301 ES interaction rates were consistent with the deficit previously observed. Most
 302 importantly, the NC reaction rate was only consistent with the others under the

303 hypothesis of flavour transformation.

$$\begin{aligned} \nu_e + d &\rightarrow p + p + e^- & (CC) \\ \nu_x + d &\rightarrow p + n + \nu_x & (NC) \\ \nu_x + e^- &\rightarrow \nu_x + e^- & (ES) \end{aligned} \quad (2.13)$$

304 Since the SNO measurement, many experiments have since measured the
 305 neutrino flux of different interaction chains within the sun [35–37]. The most
 306 recent measurement was that of CNO-cycle neutrinos which were recently
 307 observed with 5σ significance by the Borexino collaboration [35].

308 2.3.2 Accelerator Neutrinos

309 The concept of using an artificial “neutrino beam” was first realised in 1962 [38].
 310 Since then, many experiments have adopted the same fundamental concepts.
 311 Typically, a proton beam is aimed at a target producing charged mesons that
 312 decay to neutrinos. The mesons can be sign-selected by the use of magnetic
 313 focusing horns to generate a neutrino or antineutrino beam. Pions are the primary
 314 mesons that decay and depending on the orientation of the magnetic field, a
 315 muon (anti-)neutrino beam is generated via $\pi^+ \rightarrow \mu^+ + \nu_\mu$ or $\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$.
 316 The decay of muons and kaons results in an irreducible intrinsic electron neutrino
 317 background. In T2K, this background contamination is $O(< 1\%)$ [39]. There is
 318 also an approximately $\sim 5\%$ “wrong-sign” neutrino background of $\bar{\nu}_\mu$ generated
 319 via the same decays. As the beam is generated by proton interactions (rather
 320 than anti-proton interactions), the wrong-sign component in the antineutrino
 321 beam is larger when operating in neutrino mode.

322 Tuning the proton energy in the beam and using beam focusing techniques
 323 allows the neutrino energy to be set to a value that maximises the disappear-
 324 ance oscillation probability in the L/E term in Equation 2.10. This means that
 325 accelerator experiments are typically more sensitive to the mixing parameters as
 326 compared to a natural neutrino source. However, the disadvantage compared
 327 to atmospheric neutrino experiments is the cost of building a facility to provide

328 high-energy neutrinos, with a high flux, which is required for longer baselines.
 329 Consequently, there is typically less sensitivity to matter effects and the ordering
 330 of the neutrino mass eigenstates.

331 A neutrino experiment measures

$$R(\vec{x}) = \Phi(E_\nu) \times \sigma(E_\nu) \times \epsilon(\vec{x}) \times P(\nu_\alpha \rightarrow \nu_\beta), \quad (2.14)$$

332 where $R(\vec{x})$ is the event rate of neutrinos at position \vec{x} , $\Phi(E_\nu)$ is the flux of
 333 neutrinos with energy E_ν , $\sigma(E_\nu)$ is the cross-section of the neutrino interaction and
 334 $\epsilon(\vec{x})$ is the efficiency and resolution of the detector. In order to leverage the most
 335 out of an accelerator neutrino experiment, the flux and cross-section systematics
 336 need to be constrained. This is typically done via the use of a “near detector”,
 337 situated at a baseline of $O(1)$ km. This detector observes the unoscillated neutrino
 338 flux and constrains the parameters used within the flux and cross-section model.

339 The first accelerator experiments to precisely measure oscillation parameters
 340 were MINOS [40] and K2K [41]. These experiments confirmed the ν_μ disappear-
 341 ance seen in atmospheric neutrino experiments by finding consistent parameter
 342 values for $\sin^2(\theta_{23})$ and Δm_{32}^2 . The current generation of accelerator neutrino
 343 experiments, T2K and NO ν A extended this field by observing $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ and lead
 344 the sensitivity to atmospheric mixing parameters as seen in Figure 2.6 [42]. The
 345 two experiments differ in their peak neutrino energy, baseline, and detection tech-
 346 nique. The NO ν A experiment is situated at a baseline of 810km from the NuMI
 347 beamline which delivers 2GeV neutrinos. The T2K neutrino beam is peaked
 348 around 0.6GeV and propagates 295km. Additionally, the NO ν A experiment
 349 uses functionally identical detectors (near and far) whereas T2K uses a plastic
 350 scintillator technique at the near detector and a water Cherenkov far detector.
 351 The future generation experiments DUNE [43] and Hyper-Kamiokande [44]
 352 will succeed these experiments as the high-precision era of neutrino oscillation
 353 parameter measurements develops.

354 Several anomalous results have been observed in the LSND [9] and Mini-
 355 BooNE [10] detectors which were designed with purposefully short baselines.

356 Parts of the neutrino community attributed these results to oscillations induced
357 by a fourth “sterile” neutrino [45] but several searches in other experiments,
358 MicroBooNE [46] and KARMEN [47], found no hints of additional neutrino
359 species. The solution to the anomalous results is still being determined.

360 2.3.3 Atmospheric Neutrinos

361 The interactions of primary cosmic ray protons in the Earth’s upper atmosphere
362 generate showers of energetic hadrons. These are mostly pions and kaons that
363 decay to produce a natural source of neutrinos spanning energies of MeV to
364 TeV [48]. The main decay is via,

$$\begin{aligned} \pi^\pm &\rightarrow \mu^\pm + (\nu_\mu, \bar{\nu}_\mu) \\ \mu^\pm &\rightarrow e^\pm + (\nu_e, \bar{\nu}_e) + (\nu_\mu, \bar{\nu}_\mu) \end{aligned} \tag{2.15}$$

365 such that for a single pion decay, three neutrinos can be produced. The
366 atmospheric neutrino flux energy spectra as predicted by the Bartol [49], Honda
367 [50–52], and FLUKA [53] models are illustrated in Figure 2.3. The flux distribution
368 peaks at an energy of $O(10)$ GeV. The uncertainties associated with these models
369 are dominated by the hadronic production of kaon and pions as well as the
370 primary cosmic flux.

371 Unlike long-baseline experiments which have a fixed baseline, the distance
372 atmospheric neutrinos propagate is dependent upon the zenith angle at which
373 they interact. This is illustrated in Figure 2.4. Neutrinos that are generated
374 directly above the detector ($\cos(\theta) = 1.0$) have a baseline equivalent to the
375 height of the atmosphere whereas neutrinos that interact directly below the
376 detector ($\cos(\theta) = -1.0$) have to travel a length equal to the diameter of the Earth.
377 This means atmospheric neutrinos have a baseline that varies from $O(20)$ km to
378 $O(6 \times 10^3)$ km. Any neutrino generated at or below the horizon will be subject
379 to MSW matter resonance as they propagate through the Earth.

380 Figure 2.5 highlights the neutrino flux as a function of the zenith angle for
381 different slices of neutrino energy. For medium to high-energy neutrinos (and to



Figure 2.3: Left panel: The atmospheric neutrino flux for different neutrino flavours as a function of neutrino energy as predicted by the 2007 Honda model (“This work”) [50], the 2004 Honda model (“HKKM04”)[51], the Bartol model [49] and the FLUKA model [53]. Right panel: The ratio of the muon to electron neutrino flux as predicted by all the quoted models. Both figures taken from [50].



Figure 2.4: A diagram illustrating the definition of zenith angle as used in the Super Kamiokande experiment [54].

382 a lesser degree for low-energy neutrinos), the flux is approximately symmetric
 383 around $\cos(\theta) = 0$. To the accuracy of this approximation, the systematic
 384 uncertainties associated with atmospheric flux for comparing upward-going
 385 and down-going neutrino cancels. This allows the down-going events, which are

- 386 mostly insensitive to oscillation probabilities, to act as an unoscillated prediction
 387 (similar to a near detector in an accelerator neutrino experiment).



Figure 2.5: Prediction of ν_e , $\bar{\nu}_e$, ν_μ , $\bar{\nu}_\mu$ fluxes as a function of zenith angle as calculated by the HKKM model [52]. The left, middle and right panels represent three values of neutrino energy, 0.32GeV, 1.0GeV and 3.2GeV respectively. Predictions for other models including Bartol [49], Honda [50] and FLUKA [53] are given in [54].

388 Precursory hints of atmospheric neutrinos were observed in the mid-1960s
 389 searching for $\nu_\mu + X \rightarrow X^* + \mu^\pm$ [55]. This was succeeded by the IMB-3 [56]
 390 and Kamiokande [57] experiments which measured the double ratio of muon
 391 to electron neutrinos in data to Monte Carlo, $R(\nu_\mu/\nu_e) = (\mu/e)_{Data}/(\mu/e)_{MC}$.
 392 Both experiments were found to have a consistent deficit of muon neutrinos,
 393 with $R(\nu_\mu/\nu_e) = 0.67 \pm 0.17$ and $R(\nu_\mu/\nu_e) = 0.658 \pm 0.016 \pm 0.035$, respectively.
 394 Super-Kamiokande (SK) [54] extended this analysis by fitting oscillation pa-
 395 rameters in $P(\nu_\mu \rightarrow \nu_\tau)$ which found best fit parameters $\sin^2(2\theta) > 0.92$ and
 396 $1.5 \times 10^{-3} < \Delta m^2 < 3.4 \times 10^{-3}\text{eV}^2$.

397 Since then, atmospheric neutrino experiments have been making precision
 398 measurements of the $\sin^2(\theta_{23})$ and Δm^2_{32} oscillation parameters. Atmospheric
 399 neutrino oscillation is dominated by $P(\nu_\mu \rightarrow \nu_\tau)$, where SK observed a 4.6σ
 400 discovery of ν_τ appearance [58]. Figure 2.6 illustrates the current estimates
 401 on the atmospheric mixing parameters from a wide range of atmospheric and
 402 accelerator neutrino observatories.



Figure 2.6: Constraints on the atmospheric oscillation parameters, $\sin^2(\theta_{23})$ and Δm_{32}^2 , from atmospheric and long-baseline experiments: SK [59], T2K [60], NOvA [61], IceCube [62] and MINOS [63]. Figure taken from [64].

2.3.4 Reactor Neutrinos

As illustrated in the first discovery of neutrinos (section 2.1), nuclear reactors are a very useful artificial source of electron antineutrinos. For reactors that use low-enriched uranium ^{235}U as fuel, the antineutrino flux is dominated by the β -decay fission of ^{235}U , ^{238}U , ^{239}Pu and ^{241}Pu [65] as illustrated in Figure 2.7.

Due to their low energy, reactor electron antineutrinos predominantly interact via the inverse β -decay (IBD) interaction. The typical signature contains two signals delayed by $O(200)\mu\text{s}$; firstly the prompt photons from positron annihilation, and secondly the photon emitted ($E_{tot}^\gamma = 2.2\text{MeV}$) from de-excitation after neutron capture on hydrogen. Searching for both signals improves the detector's ability to distinguish between background and signal events [67].

There are many short baseline experiments ($L \sim O(1)\text{km}$) that have measured the $\sin^2(\theta_{13})$ and Δm_{32}^2 oscillation parameters. Daya Bay [68], RENO [69] and Double Chooz [70] have all provided precise measurements, with the first discovery of a non-zero θ_{13} made by Daya Bay and RENO (and complemented by T2K [70]). The constraints on $\sin^2(\theta_{13})$ by the reactor experiments lead the field and

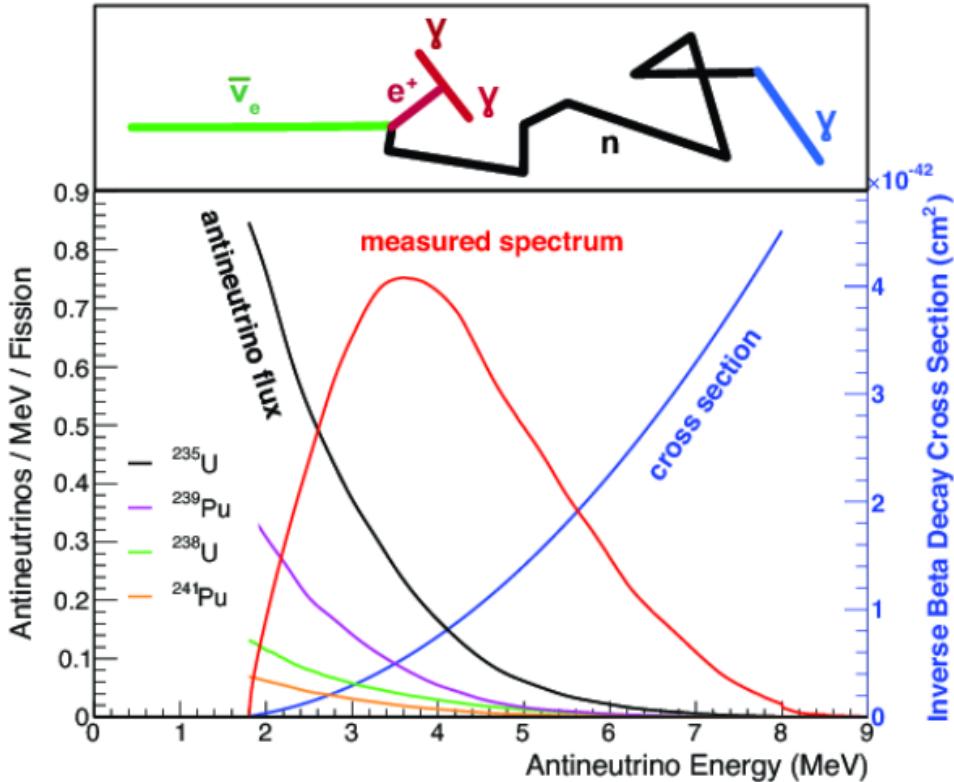


Figure 2.7: Reactor electron antineutrino fluxes for ^{235}U (Black), ^{238}U (Green), ^{239}Pu (Purple), and ^{241}Pu (Orange) isotopes. The inverse β -decay cross-section (Blue) and corresponding measurable neutrino spectrum (Red) are also given. Top panel: Schematic of Inverse β -decay interaction including the eventual capture of the emitted neutron. This capture emits a γ -ray which provides a second signal of the event. Taken from [66].

are often used as external inputs to accelerator neutrino experiments to improve their sensitivity to δ_{CP} and mass hierarchy determination. JUNO-TAO [71], a small collaboration within the larger JUNO experiment, is a next-generation reactor experiment that aims to precisely measure the isotopic antineutrino yields from the different fission chains.

Kamland [72] is the only experiment to have observed reactor neutrinos using a long baseline (flux weighted averaged baseline of $L \sim 180\text{km}$) which allows it to have sensitivity to Δm_{21}^2 . Combined with the SK solar neutrino experiment, the combined analysis puts the most stringent constraint on Δm_{21}^2 [73].

2.4 Summary Of Oscillation Parameter Measurements

Since the first evidence of neutrino oscillations, numerous measurements of the mixing parameters have been made. Many experiments use neutrinos as a tool for the discovery of new physics (diffuse supernova background, neutrinoless double beta decay and others) so the PMNS parameters are summarised in the Particle Data Group (PDG) review tables. The analysis presented in this thesis focuses on the 2020 T2K oscillation analysis presented in [74] which the 2020 PDG constraints [75] were used. These constraints are outlined in Table 2.1.

Parameter	2020 Constraint
$\sin^2(\theta_{12})$	0.307 ± 0.013
Δm_{21}^2	$(7.53 \pm 0.18) \times 10^{-5} \text{ eV}^2$
$\sin^2(\theta_{13})$	$(2.18 \pm 0.07) \times 10^{-2}$
$\sin^2(\theta_{23})$ (I.H.)	0.547 ± 0.021
$\sin^2(\theta_{23})$ (N.H.)	0.545 ± 0.021
Δm_{32}^2 (I.H.)	$(-2.546^{+0.034}_{-0.040}) \times 10^{-3} \text{ eV}^2$
Δm_{32}^2 (N.H.)	$(2.453 \pm 0.034) \times 10^{-3} \text{ eV}^2$

Table 2.1: The 2020 Particle Data Group constraints of the oscillation parameters taken from [75]. The value of Δm_{32}^2 is given for both normal hierarchy (N.H.) and inverted hierarchy (I.H.) and $\sin^2(\theta_{23})$ is broken down by whether its value is below (Q1) or above (Q2) 0.5.

The $\sin^2(\theta_{13})$ measurement stems from the electron antineutrino disappearance, $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$, and is take as the average best-fit from the combination of Daya Bay, Reno and Double Chooz. It is often used as a prior uncertainty within other neutrino oscillation experiments, typically termed the reactor constraint. The $\sin^2(\theta_{12})$ parameter is predominantly measured through electron neutrino disappearance, $P(\nu_e \rightarrow \nu_{\mu,\tau})$, in solar neutrino experiments. The long-baseline reactor neutrino experiment Kamland also has a sensitivity to this parameter and is used in a joint fit to solar data from SNO and SK, using the reactor constraint. Measurements of $\sin^2(\theta_{23})$ are made by long-baseline and atmospheric neutrino experiments. The PDG value is a joint fit of T2K, NO ν A , MINOS and IceCube DeepCore experiments. The latest T2K-only measurement, provided at Neutrino2020 and is the basis of this thesis, is given as $\sin^2(\theta_{23}) = 0.546^{+0.024}_{-0.046}$ [74].

The PDG constraint on Δm_{21}^2 is provided by the KamLAND experiment using solar and geoneutrino data. This measurement utilised a $\sin^2(\theta_{13})$ constraint from accelerator (T2K, MINOS) and reactor neutrino (Daya Bay, RENO, Double Chooz) experiments. Accelerator measurements make some of the most stringent constraints on Δm_{32}^2 although atmospheric experiments have more sensitivity to the mass hierarchy determination. The PDG performs a joint fit of accelerator and atmospheric data, in both normal and inverted hierarchies separately. The latest T2K-only result is $\Delta m_{32}^2 = 2.49^{+0.058}_{-0.082} \times 10^{-3} \text{ eV}^2$ favouring normal hierarchy [74]. The value of δ_{CP} is largely undetermined. CP-conserving values of 0 and π were rejected with $\sim 2\sigma$ intervals, as published in Nature, although more recent analyses have reduced the credible intervals to 90%. Since the 2020 PDG publication, there has been a new measurement of $\sin^2(\theta_{13}) = (2.20 \pm 0.07) \times 10^{-2}$ [76], alongside updated Δm_{32}^2 and $\sin^2(\theta_{23})$ measurements.

Throughout this thesis, several sample spectra predictions and contours are presented which require oscillation parameters to be assumed. Table 2.2 defines two sets of oscillation parameters, with “Asimov A” set being close to the preferred values from a previous T2K-only fit [77] and “Asimov B” being CP-conserving and further from maximal θ_{23} mixing.

Parameter	Asimov A	Asimov B
Δm_{12}^2	$7.53 \times 10^{-5} \text{ eV}^2$	
Δm_{32}^2	$2.509 \times 10^{-3} \text{ eV}^2$	
$\sin^2(\theta_{12})$	0.304	
$\sin^2(\theta_{13})$	0.0219	
$\sin^2(\theta_{23})$	0.528	0.45
δ_{CP}	-1.601	0.0

Table 2.2: Reference values of the neutrino oscillation parameters for two different oscillation parameter sets.

3

466

467

T2K and SK Experiment Overview

468 As the successor of the Kamiokande experiment, the Super-Kamiokande (SK)
469 collaboration has been leading atmospheric neutrino oscillation analyses for
470 over two decades. The detector has provided some of the strongest constraints
471 on proton decay and the first precise measurements of the Δm_{23}^2 and $\sin^2(\theta_{23})$
472 neutrino oscillation parameters. The history, detection technique, and operation
473 of the SK detector is described in section 3.1.

474 The Tokai-to-Kamioka (T2K) experiment was one of the first long-baseline
475 experiments to use both neutrino and antineutrino beams to precisely measure
476 the charge parity violation within the neutrino sector. The T2K experiment
477 observed the first hints of a non-zero $\sin^2(\theta_{13})$ measurement and continues to
478 lead the field with the constraints it provides on $\sin^2(\theta_{13})$, $\sin^2(\theta_{23})$, Δm_{23}^2 and
479 δ_{CP} . In section 3.2, the techniques that T2K use to generate the neutrino beam and
480 constrain systematic parameter through near detector constraints are described.

481 3.1 The Super-Kamiokande Experiment

482 The SK experiment began taking data in 1996 [78] and has had many modifi-
483 cations throughout its operation. There have been seven defined periods of
484 data taking as noted in Table 3.1. Data taking began in SK-I which ran for five

485 years. Between the SK-I and SK-II periods, approximately 55% of the PMTs were
 486 damaged during maintenance [79]. Those that survived were equally distributed
 487 throughout the detector in the SK-II era, which resulted in a reduced 19% photo-
 488 coverage. From SK-III onwards, repairs to the detector meant the full suite of
 489 PMTs was operational recovering the 40% photocoverage. Before the start of
 490 SK-IV, the data acquisition and electronic systems were upgraded. Between
 491 SK-IV and SK-V, a significant effort was placed into tank open maintenance
 492 and repair/replacement of defective PMTs, a task for which the author of this
 493 thesis was required. Consequently, the detector conditions were significantly
 494 different between the two operational periods. SK-VI marked the start of the
 495 SK-Gd era, with the detector being doped with gadolinium at a concentration
 496 of 0.01%. SK-VII, which started during the writing of this thesis, has increased
 497 the gadolinium concentration to 0.03% for continued operation [80].

498 The oscillation analysis presented within this thesis focuses on the SK-IV
 499 period of running and the data taking within it. This follows from the recent
 500 SK analysis presented in [81]. Therefore, the information presented within this
 501 section focuses on that period.

Period	Start Date	End Date	Live-time (days)
I	April 1996	July 2001	1489.19
II	October 2002	October 2005	798.59
III	July 2006	September 2008	518.08
IV	September 2008	May 2018	3244.4
V	January 2019	July 2020	461.02
VI	July 2020	May 2022	583.3
VII	May 2022	Ongoing	N/A

Table 3.1: The various SK periods and respective live-time. The SK-VI live-time is calculated until 1st April 2022. SK-VII started during the writing of this thesis.

502 3.1.1 The SK Detector

503 The basic structure of the Super-Kamiokande (SK) detector is a cylindrical tank
 504 with a diameter 39.3m and height 41.1m filled with ultrapure water [79]. A
 505 diagram of the significant components of the SK detector is given in Figure 3.1.

506 The SK detector is situated in the Kamioka mine in Gifu, Japan. The mine
 507 is underground with roughly 1km rock overburden (2.7km water equivalent
 508 overburden) [82]. At this depth, the rate of cosmic ray muons is significantly
 509 decreased to a value of $\sim 2\text{Hz}$. The top of the tank is covered with stainless
 510 steel which is designed as a working platform for maintenance, calibration, and
 511 location for high voltage and data acquisition electronics.



Figure 3.1: A schematic diagram of the Super-Kamiokande Detector. Taken from [83].

512 A smaller cylindrical structure (36.2m diameter, 33.8m height) is situated
 513 inside the tank, with an approximate 2m gap between this structure and the outer
 514 tank wall. The purpose of this structure is to support the photomultiplier tubes
 515 (PMTs). The volume inside and outside the support structure is referred to as the
 516 inner detector (ID) and outer detector (OD), respectively. In the SK-IV era, the
 517 ID and OD are instrumented by 11,129 50cm and 1,885 20cm PMTs respectively
 518 [79]. The ID contains a 32kton mass of water. Many analyses performed at SK
 519 use a “fiducial volume” defined by the volume of water inside the ID excluding
 520 some distance to the ID wall. This reduces the volume of the detector which is
 521 sensitive to neutrino events but reduces radioactive backgrounds and allows for

522 better reconstruction performance. The nominal fiducial volume is defined as the
523 area contained inside 2m from the ID wall for a total of 22.5kton water [84].

524 The two regions of the detector (ID and OD) are optically separated with
525 opaque black plastic hung from the support structure. The purpose of this is
526 to determine whether an event entered or exited the ID. This allows cosmic ray
527 muons and partially contained events to be tagged and separated from neutrino
528 events entirely contained within the ID. This black plastic is also used to cover
529 the area between the ID PMTs to reduce photon reflection from the ID walls.
530 Opposite to this, the OD is lined with a reflective material to allow photons to
531 reflect around inside the OD until collected by one of the PMTs. Furthermore,
532 each OD PMT is optically coupled with $50 \times 50\text{cm}$ plates of wavelength shifting
533 acrylic which increases the efficiency of light collection [82].

534 In the SK-IV data-taking period, the photocathode coverage of the detector, or
535 the fraction of the ID wall instrumented with PMTs, is $\sim 40\%$ [82]. The PMTs have
536 a quantum efficiency (the ratio of detected electrons to incident photons) of $\sim 21\%$
537 for photons with wavelengths of $360\text{nm} < \lambda < 390\text{nm}$ [85]. The proportion of
538 photoelectrons that produce a signal in the dynode of a PMT, termed the collection
539 efficiency, is $> 70\%$ [82]. The PMTs used within SK are most sensitive to photons
540 with wavelength $300\text{nm} \leq \lambda \leq 600\text{nm}$ [82]. One disadvantage of using PMTs as
541 the detection media is that the Earth's geomagnetic field can modify its response.
542 Therefore, a set of compensation coils is built around the inner surface of the
543 detector to mitigate this effect [86].

544 As mentioned, the SK detector is filled with ultrapure water, which in a perfect
545 world would contain no impurities. However, bacteria and organic compounds
546 can significantly degrade the water quality. This decreases the attenuation length,
547 which reduces the total number of photons that hit a PMT. To combat this, a
548 sophisticated water treatment system has been developed [82, 87]. UV lights,
549 mechanical filters, and membrane degasifiers are used to reduce the bacteria,
550 suspended particulates, and radioactive materials from the water. The flow of
551 water within the tank is also critical as it can remove stagnant bacterial growth

552 or build-up of dust on the surfaces within the tank. Gravity drifts impurities in
553 the water towards the bottom of the tank which, if left uncontrolled, can create
554 asymmetric water conditions between the top and bottom of the tank. Typically,
555 the water entering the tank is cooled below the ambient temperature of the tank
556 to control convection and inhibit bacteria growth. Furthermore, the rate of dark
557 noise hits within PMTs is sensitive to the PMT temperature [88] so controlling
558 the temperature gradients within the tank is beneficial for stable measurements.

559 SK-VI is the first phase of the SK experiment to use gadolinium dopants
560 within the ultrapure water [80]. As such, the SK water system had to be replaced
561 to avoid removing the gadolinium concentrate from the ultrapure water [89]. For
562 an inverse β -decay (IBD) interaction in a water target, the emitted neutron is
563 thermally captured on hydrogen. This process releases 2.2MeV γ ray which are
564 difficult to detect as the resulting Compton scattered electrons are very close to the
565 Cherenkov threshold, limiting detection capability. Thermal capture of neutrons
566 on gadolinium generates γ rays with higher energy (8MeV [67]) meaning they
567 are more easily detected and reconstructed. SK-VI has 0.01% Gd loading (0.02%
568 gadolinium sulphate by mass) which causes \approx 50% of neutrons emitted by IBD to
569 be captured on gadolinium[90, 91] . Whilst predominantly useful for low energy
570 analyses, Gd loading allows better $\nu/\bar{\nu}$ separation for atmospheric neutrino
571 event selections [92]. Efforts are currently in place to increase the gadolinium
572 concentrate to 0.03% for \approx 75% neutron capture efficiency on gadolinium [93].
573 The final stage of loading targets 0.1% concentrate targeting \approx 90% neutron
574 capture efficiency on gadolinium.

575 3.1.2 Calibration

576 The calibration of the SK detector is documented in [79] and summarised below.
577 The analysis presented within this thesis is dependent upon ‘high energy events’
578 (Charged particles with $O(> 100)$ MeV momenta). These are events that are
579 expected to generate a larger number of photons such that each PMT will
580 be hit with multiple photons. The reconstruction of these events depends

upon the charge deposited within each PMT and the timing response of each individual PMT. Therefore, the most relevant calibration techniques to this thesis are outlined.

Before installation, 420 PMTs were calibrated to have identical charge responses and then distributed throughout the tank in a cross-shape pattern (As illustrated by Figure 3.2). These are used as a standardised measure for the rest of the PMTs installed at similar geometric positions within SK to be calibrated against. To perform this calibration, a xenon lamp is located at the centre of the SK tank which flashes uniform light at 1Hz. This allows for geometrical effects, water quality variation, and timing effects to be measured in-situ throughout normal data-taking periods.



Figure 3.2: The location of “standard PMTs” (red) inside the SK detector. Taken from [79].

When specifically performing calibration of the detector (in out-of-data taking mode), the water in the tank was circulated to avoid top/bottom asymmetric water quality. Any non-uniformity within the tank significantly affects the PMT hit probability through scattering or absorption. This becomes a dominant effect for the very low-intensity light sources discussed later which are designed such that only one photon is incident upon a given PMT.

598 The “gain” of a PMT is defined as the ratio of the total charge of the signal
 599 produced compared to the charge of photoelectrons emitted by the photocathodes
 600 within the PMT. To calibrate the signal of each PMT, the “relative” and “absolute”
 601 gain values are measured. The relative gain is the variation of gain among each
 602 of the PMTs whereas the absolute gain is the average gain of all PMTs.

603 The relative gain is calibrated as follows. A laser is used to generate two
 604 measurements: a high-intensity flash that illuminates every PMT with a sufficient
 605 number of photons, and a low-intensity flash in which only a small number
 606 of PMTs collect light. The first measurement creates an average charge, $Q_{obs}(i)$
 607 on PMT i , whereas the second measurement ensures that each hit PMT only
 608 generates a single photoelectron. For the low-intensity measurement, the number
 609 of times each PMT records a charge larger than 1/4 photoelectrons, $N_{obs}(i)$, is
 610 counted. The values measured can be expressed as

$$\begin{aligned} Q_{obs}(i) &\propto I_H \times f(i) \times \epsilon(i) \times G(i), \\ N_{obs}(i) &\propto I_L \times f(i) \times \epsilon(i), \end{aligned} \tag{3.1}$$

611 Where I_H and I_L is the intensity of the high and low flashes, $f(i)$ is the
 612 acceptance efficiency of the i^{th} PMT, $\epsilon(i)$ is the product of the quantum and
 613 collection efficiency of the i^{th} PMT and $G(i)$ is the gain of the i^{th} PMT. The
 614 relative gain for each PMT can determined by taking the ratio of these quantities.

615 The absolute gain calibration is performed by observing fixed energy γ -rays
 616 of $E_\gamma \sim 9\text{MeV}$ emitted isotropically from neutron capture on a NiCf source
 617 situated at the centre of the detector. This generates a photon yield of about 0.004
 618 photoelectrons/PMT/event, meaning that $> 99\%$ of PMT signals are generated
 619 from single photoelectrons. A charge distribution is generated by performing
 620 this calibration over all PMTs, and the average value of this distribution is taken
 621 to be the absolute gain value.

622 As mentioned in subsection 3.1.1, the average quantum and collection effi-
 623 ciency for the SK detector PMTs is $\sim 21\%$ and $> 70\%$ respectively. However,
 624 these values do differ between each PMT and need to be calibrated accordingly.

625 Consequently, the NiCf source is also used to calibrate the “quantum \times collection”
626 efficiency (denoted “QE”) value of each PMT. The NiCf low-intensity source is
627 used as the PMT hit probability is proportional to the QE ($N_{obs}(i) \propto \epsilon(i)$ in
628 Equation 3.1). A Monte Carlo prediction which includes photon absorption,
629 scattering, and reflection is made to estimate the number of photons incident on
630 each PMT and the ratio of the number of predicted to observed hits is calculated.
631 The difference is attributed to the QE efficiency of that PMT. This technique is
632 extended to calculate the relative QE efficiency by normalizing the average of
633 all PMTs which removes the dependence on the light intensity.

634 Due to differing cable lengths and readout electronics, the timing response
635 between a photon hitting the PMT and the signal being captured by the data
636 acquisition can be different between each PMT. Due to threshold triggers (De-
637 scribed in subsection 3.1.3), the time at which a pulse reaches a threshold is
638 dependent upon the size of the pulse. This is known as the ‘time-walk’ effect
639 and also needs to be accounted for in each PMT. To calibrate the timing response,
640 a pulse of light with width 0.2ns is emitted into the detector through a diffuser.
641 Two-dimensional distributions of time and pulse height (or charge) are made
642 for each PMT and are used to calibrate the timing response. This is performed
643 in-situ during data taking with the light source pulsing at 0.03Hz.

644 The top/bottom water quality asymmetry is measured using the NiCf calibra-
645 tion data and cross-referencing these results to the “standard PMTs”. The water
646 attenuation length is continuously measured by the rate of vertically-downgoing
647 cosmic-ray muons which enter via the top of the tank.

648 Dark noise is where a PMT registers a pulse that is consistent with a single
649 photoelectron emitted from photon detection despite the PMT being in complete
650 darkness. This is predominately caused by two processes. Firstly there is
651 intrinsic dark noise which is where photoelectrons gain enough thermal energy
652 to be emitted from the photocathode, and secondly, the radioactive decay of
653 contaminants inside the structure of the PMT. Typical dark noise rate for PMTs
654 used within SK are $O(3)\text{kHz}$ [82]. This is lower than the expected number of

655 photons generated for a ‘high energy event’ (As described in subsection 3.1.4)
656 but instability in this value can cause biases in reconstruction. Dark noise is
657 related to the gain of a PMT and is calibrated using hits inside a time window
658 recorded before an event trigger [94].

659 3.1.3 Data Acquisition and Triggering

660 As the analysis presented in this thesis only uses the SK-IV period of the SK
661 experiment so this subsection focuses on the relevant points of the data acqui-
662 sition and triggering systems to that SK period. The earlier data acquisition and
663 triggering systems are documented in [95, 96].

664 Before the SK-IV period started, the existing front-end electronics were re-
665 placed with “QTC-Based Electrons with Ethernet, QBEE” systems [97]. When the
666 QBEE observes a signal above a 1/4 photoelectron threshold, the charge-to-time
667 (QTC) converter generates a rectangular pulse. The start of the rectangular pulse
668 indicates the time at which the analog photoelectron signal was received and the
669 width of the pulse indicates the total charge integrated throughout the signal.
670 This is then digitized by time-to-digital converters and sent to the “front-end”
671 PCs. The digitized signal from every QBEE is then chronologically ordered and
672 sent to the “merger” PCs. It is the merger PCs that apply the software trigger.
673 Any triggered events are passed to the “organizer” PC. This sorts the data stream
674 of multiple merger PCs into chronologically ordered events which are then saved
675 to disk. The schematic of data flow from PMTs to disk is illustrated in Figure 3.3.

676 The software trigger (described in [99]) operates by determining the number
677 of PMT hits within a 200ns sliding window, N_{200} . This window coincides with
678 the maximum time that a Cherenkov photon would take to traverse the length
679 of the SK tank [96]. For lower energy events that generate fewer photons, this
680 technique is useful for eliminating background processes like dark noise and
681 radioactive decay which would be expected to separate in time. When the value
682 of N_{200} exceeds some threshold, a software trigger is issued. There are several
683 trigger thresholds used within the SK-IV period which are detailed in Table 3.2.

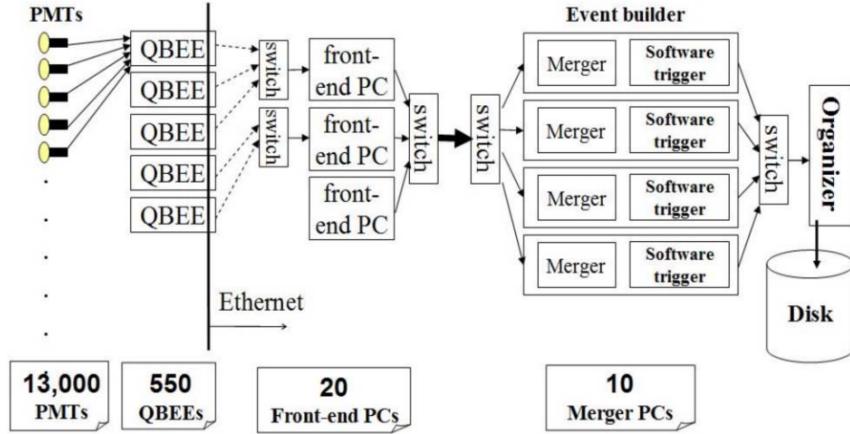


Figure 3.3: Schematic view of the data flow through the data acquisition and online system. Taken from [98].

684 If one of these thresholds is met, the PMT hits within an extended time window
 685 are also read out and saved to disk. In the special case of an event that exceeds
 686 the SHE trigger but does not exceed the OD trigger, the AFT trigger looks for
 687 delayed coincidences of 2.2MeV gamma rays emitted from neutron capture in
 688 a $535\mu\text{s}$ window after the SHE trigger. A similar but more complex “Wideband
 689 Intelligent Trigger (WIT)” has been deployed and is described in [100].

Trigger	Acronym	Condition	Extended time window (μs)
Super Low Energy	SLE	>34/31 hits	1.3
Low Energy	LE	>47 hits	40
High Energy	HE	>50 hits	40
Super High Energy	SHE	>70/58 hits	40
Outer Detector	OD	>22 hits in OD	N/A

Table 3.2: The trigger thresholds and extended time windows saved around an event which were utilised throughout the SK-IV period. The exact thresholds can change and the values listed here represent the thresholds at the start and end of the SK-IV period.

690 3.1.4 Cherenkov Radiation

691 Cherenkov light is emitted from any highly energetic charged particle traveling
 692 with relativistic velocity, β , greater than the local speed of light in a medium [101].

693 Cherenkov light is formed at the surface of a cone with characteristic pitch angle,

$$\cos(\theta) = \frac{1}{\beta n}. \quad (3.2)$$

694 where n is the refractive index of the medium. Consequently, the Cherenkov
 695 momentum threshold, P_{thres} , is dependent upon the mass, m , of the charged
 696 particle moving through the medium,

$$P_{thres} = \frac{m}{\sqrt{n^2 - 1}} \quad (3.3)$$

697 For water, where $n = 1.33$, the Cherenkov threshold momentum and energy
 698 for various particles are given in Table 3.3. In contrast, γ -rays are detected
 699 indirectly via the combination of photons generated by Compton scattering
 700 and pair production. The threshold for detection in the SK detector is typically
 701 higher than the threshold for photon production. This is due to the fact that the
 702 attenuation of photons in the water means that typically $\sim 75\%$ of Cherenkov
 703 photons reach the ID PMTs. Then the collection and quantum efficiencies
 704 described in subsection 3.1.1 result in the number of detected photons being
 705 lower than the number of photons which reach the PMTs.

Particle	Threshold Momentum (MeV)	Threshold Energy (MeV)
Electron	0.5828	0.7751
Muon	120.5	160.3
Pion	159.2	211.7
Proton	1070.0	1423.1

Table 3.3: The threshold momentum and energy for a particle to generate Cherenkov light in ultrapure water, as calculated in Equation 3.2 in ultrapure water which has refractive index $n = 1.33$.

706 The Frank-Tamm equation [102] describes the relationship between the num-
 707 ber of Cherenkov photons generated per unit length, dN/dx , the wavelength of
 708 the photons generated, λ , and the relativistic velocity of the charged particle,

$$\frac{d^2N}{dxd\lambda} = 2\pi\alpha \left(1 - \frac{1}{n^2\beta^2}\right) \frac{1}{\lambda^2}. \quad (3.4)$$

709 where α is the fine structure constant. For a 100MeV momentum electron,
 710 approximately 330 photons will be produced per centimeter in the $300\text{nm} \leq \lambda \leq$
 711 700nm region which the ID PMTs are most sensitive to [82].

712 3.2 The Tokai to Kamioka Experiment

713 The Tokai to Kamioka (T2K) experiment is a long-baseline neutrino oscillation
 714 experiment located in Japan. Proposed in the early 2000s [103, 104] to replace
 715 K2K [105], T2K was designed to observe electron neutrino appearance whilst
 716 precisely measuring the oscillation parameters associated with muon neutrino
 717 disappearance [106]. The experiment consists of a neutrino beam generated
 718 at the Japan Proton Accelerator Research Complex (J-PARC), a suite of near
 719 detectors situated 280m from the beam target, and the Super Kamiokande far
 720 detector positioned at a 295km baseline. The cross-section view of the T2K
 721 experiment is drawn in Figure 3.4.

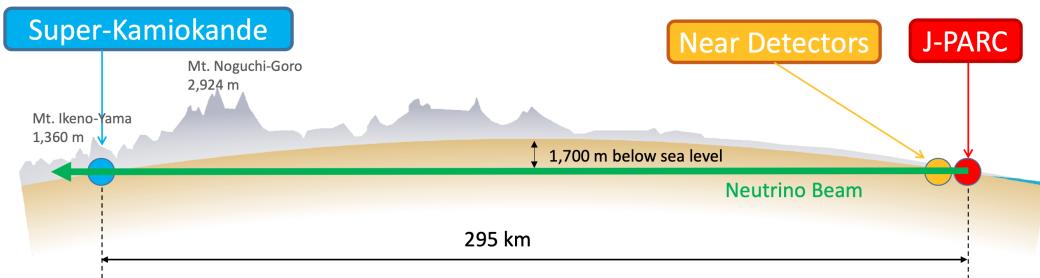


Figure 3.4: The cross-section view of the Tokai to Kamioka experiment illustrating the beam generation facility at J-PARC, the near detector situated at a baseline of 280m and the Super Kamiokande far detector situated 295km from the beam target.

722 The T2K collaboration makes world-leading measurements of the $\sin^2(\theta_{23})$
 723 , Δm_{23}^2 , and δ_{CP} oscillation parameters. Improvements in the precision and
 724 accuracy of parameter estimates are still being made by including new data
 725 samples and developing the models which describe the neutrino interactions
 726 and detector responses [107]. Electron neutrino appearance was first observed
 727 at T2K in 2014 [108] with 7.3σ significance.

The near detectors provide constraints on the beam flux and cross-section model parameters used within the oscillation analysis by observing the unoscillated neutrino beam. There are a host of detectors situated in the near detector hall (As illustrated in Figure 3.5): ND280 (subsection 3.2.2), INGRID (subsection 3.2.3), NINJA [109], WAGASCI [110], and Baby-MIND [111]. The latter three are not currently used within the oscillation analysis presented within this thesis.

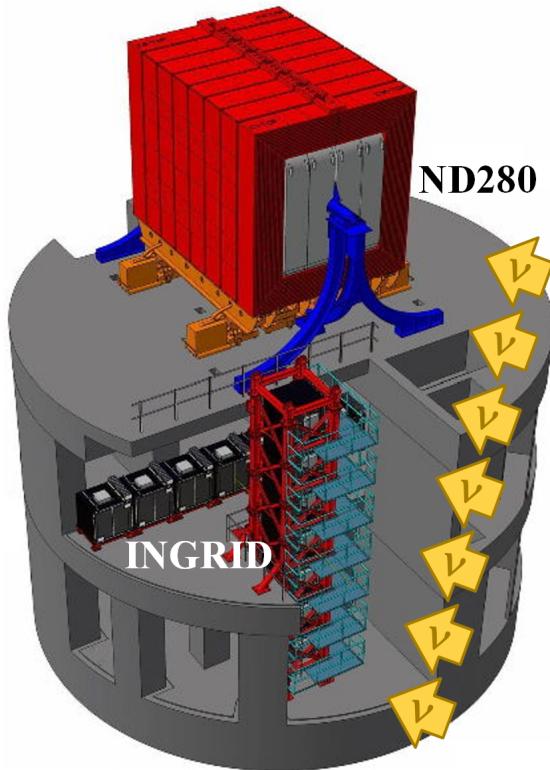


Figure 3.5: The near detector suite for the T2K experiment showing the ND280 and INGRID detectors. The distance between the detectors and the beam target is 280m.

Whilst this thesis presents the ND280 in terms of its purpose for the oscillation analysis, the detector can also make many cross-section measurements at neutrino energies of $O(1)$ GeV for the different targets within the detector [112, 113]. These measurements are of equal importance as they can lead the way in determining the model parameters used in the interaction models for the future high-precision era of neutrino physics.

There are two independent fitters, MaCh3 and BANFF, which perform the near detector fit. MaCh3 is the basis of this analysis and uses a bayesian Markov

742 Chain Monte Carlo fitting technique, whereas BANFF uses a frequentist gradient
743 descent technique. The output of each fitter is converted into a covariance
744 matrix to describe the error and correlations between all the flux and cross-
745 section parameters. This is then propagated to the far-detector oscillation analysis
746 group for use in the P-Theta and VALOR fitting framework. As MaCh3 can handle
747 both near and far detector samples, it does not use this covariance matrix and
748 instead opts for a simultaneous fit of the two detector measurements. This is an
749 analysis choice which removes the assumption of Gaussian posterior distributions
750 required when building the post-fit covariance matrix.

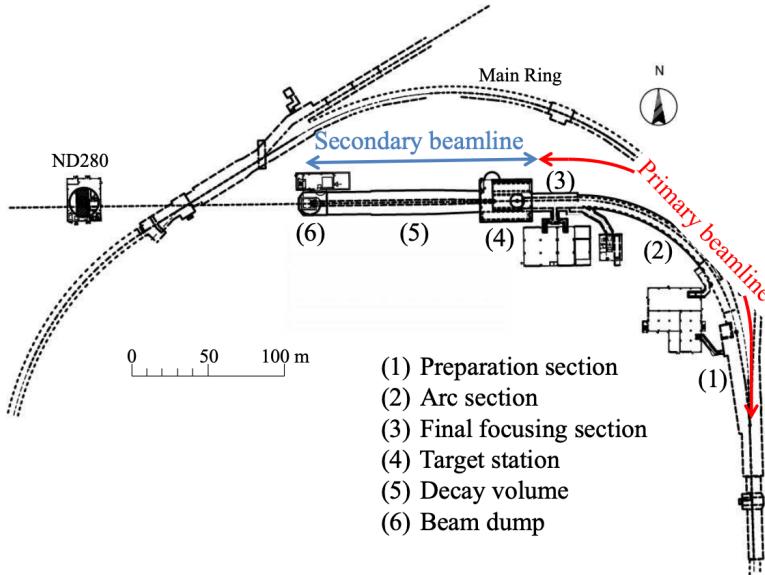
751 **DB: MaCh3 vs PTheta and Valor**

752 There are three particular tunes of the T2K flux and low energy cross section
753 model typically considered. Firstly, the “generated” tune which is the set of
754 dial values with which the Monte Carlo was generated. Secondly, the set of dial
755 values which are taken from external data measurements and used as inputs.
756 These are the “pre-fit” dial values. The reason these two sets of dial values are
757 different is that the external data measurements are continually updated but it
758 is very computationally intensive to regenerate a Monte Carlo prediction after
759 each update. The final tune is the “post-fit”, “post-ND fit” or “post-BANFF” dial
760 values. These are the values taken from the fit to the beam near detector data.

761 **3.2.1 The Neutrino Beam**

762 The neutrino beam used within the T2K experiment is described in [39, 114]
763 and summarised below. The accelerating facility at J-PARC is composed of two
764 sections; the primary and secondary beamlines. Figure 3.6 illustrates a schematic
765 of the beamline, focusing mostly on the components of the secondary beamline.
766 The primary beamline has three accelerators that progressively accelerate pro-
767 tons; a linear accelerator, a rapid-cycling synchrotron, and the main-ring (MR)
768 synchrotron. Once fully accelerated by the MR, the protons have a kinetic energy
769 of 30GeV. Eight bunches of these protons, separated by 500ns, are extracted
770 per “spill” from the MR and directed towards a graphite target (a rod of length

91.4cm and diameter 2.6cm). Spills are extracted at 0.5Hz with $\sim 3 \times 10^{14}$ protons contained per spill.



(a) Primary and secondary beamline



(b) Secondary beamline

Figure 3.6: Top panel: Bird's eye view of the most relevant part of primary and secondary beamline used within the T2K experiment. The primary beamline is the main-ring proton synchrotron, kicker magnet, and graphite target. The secondary beamline consists of the three focusing horns, decay volume, and beam dump. Figure taken from [114]. Bottom panel: The side-view of the secondary beamline including the focusing horns, beam dump and neutrino detectors. Figure taken from [115].

The secondary beamline consists of three main components: the target station, the decay volume, and the beam dump. The target station is comprised of the target, beam monitors, and three magnetic focusing horns. The proton beam interacts with the graphite target to form a secondary beam of mostly pions and kaons. The secondary beam travels through a 96m long decay volume, generating neutrinos through the following decays [39],

$$\begin{aligned}
& \pi^+ \rightarrow \mu^+ + \nu_\mu & \pi^- \rightarrow \mu^- + \bar{\nu}_\mu \\
& K^+ \rightarrow \mu^+ + \nu_\mu & K^- \rightarrow \mu^- + \bar{\nu}_\mu \\
& \quad \rightarrow \pi^0 + e^+ + \nu_e & \quad \rightarrow \pi^0 + e^- + \bar{\nu}_e \\
& \quad \rightarrow \pi^0 + \mu^+ + \nu_\mu & \quad \rightarrow \pi^0 + \mu^- + \bar{\nu}_\mu \\
& K_L^0 \rightarrow \pi^- + e^+ + \nu_e & K_L^0 \rightarrow \pi^+ + e^- + \bar{\nu}_e \\
& \quad \rightarrow \pi^- + \mu^+ + \nu_\mu & \quad \rightarrow \pi^+ + \mu^- + \bar{\nu}_\mu \\
& \mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e & \mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e
\end{aligned}$$

779 The electrically charged component of the secondary beam is focused towards
 780 the far detector by the three magnetic horns. These horns direct charged particles
 781 of a particular polarity towards SK whilst defocusing the oppositely charged
 782 particles. This allows a mostly neutrino or mostly antineutrino beam to be used
 783 within the experiment, denoted as “forward horn current (FHC)” or “reverse
 784 horn current (RHC)” respectively.

786 Figure 3.7 illustrates the different contributions to the FHC and RHC neutrino
 787 flux. The low energy flux is dominated by the decay of pions whereas kaon
 788 decay becomes the dominant source of neutrinos for $E_\nu > 3\text{GeV}$. The “wrong-
 789 sign” component, which is the $\bar{\nu}_\mu$ background in a ν_μ beam, and the intrinsic
 790 irreducible ν_e background, are predominantly due to muon decay for $E_\nu <$
 791 2GeV . As the antineutrino production cross-section is smaller than the neutrino
 792 cross-section, the wrong-sign component is more dominant in the RHC beam
 793 as compared to that in the FHC beam.

794 The beam dump, situated at the end of the decay volume, stops all charged
 795 particles other than highly energetic muons ($p_\mu > 5\text{GeV}$). The MuMon detector
 796 monitors the penetrating muons to determine the beam direction and inten-
 797 sity which is used to constrain some of the beam flux systematics within the
 798 analysis [115, 117].

799 The T2K experiment uses an off-axis beam to narrow the neutrino energy
 800 distribution. This was the first implementation of this technique in a long-
 801 baseline neutrino oscillation experiment after its original proposal [118]. Pion
 802 decay, $\pi \rightarrow \mu + \nu_\mu$, is a two-body decay. Consequently, the neutrino energy,

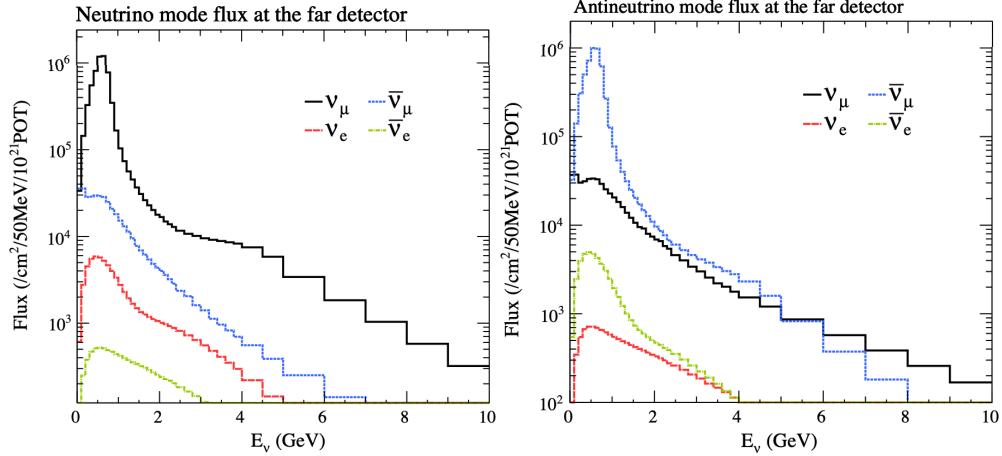


Figure 3.7: The Monte Carlo prediction of the energy spectrum for each flavour of neutrino (ν_e , $\bar{\nu}_e$, ν_μ and $\bar{\nu}_\mu$) in the neutrino dominated beam FHC mode (Left) and antineutrino dominated beam RHC mode (Right) expected at SK. Taken from [116].

803 E_ν , can be determined based on the pion energy, E_π , and the angle at which
 804 the neutrino is emitted, θ ,

$$E_\nu = \frac{m_\pi^2 - m_\mu^2}{2(E_\pi - p_\pi \cos(\theta))}, \quad (3.5)$$

805 where m_π and m_μ are the mass of the pion and muon respectively. For a fixed
 806 energy pion, the neutrino energy distribution is dependent upon the angle at
 807 which the neutrinos are observed from the initial pion beam direction. For the
 808 295km baseline at T2K, $E_\nu = 0.6$ GeV maximises the electron neutrino appearance
 809 probability, $P(\nu_\mu \rightarrow \nu_e)$, whilst minimising the muon disappearance probability,
 810 $P(\nu_\mu \rightarrow \nu_\mu)$. Figure 3.8 illustrates the neutrino energy distribution for a range of
 811 off-axis angles, as well as the oscillation probabilities most relevant to T2K.

812 3.2.2 The Near Detector at 280m

813 Whilst all the near detectors are situated in the same “pit” located at 280m from
 814 the beamline, the “ND280” detector is the off-axis detector which is situated at
 815 the same off-axis angle as the Super-Kamiokande far detector. It has two primary
 816 functions; firstly it measures the neutrino flux and secondly it counts the event
 817 rates of different types of neutrino interactions. Both of these constrain the flux



Figure 3.8: Top panel: T2K muon neutrino disappearance probability as a function of neutrino energy. Middle panel: T2K electron neutrino appearance probability as a function of neutrino energy. Bottom panel: The neutrino flux distribution for three different off-axis angles (Arbitrary units) as a function of neutrino energy.

and cross-section systematics invoked within the model for a more accurate prediction of the expected event rate at the far detector.

As illustrated in Figure 3.9, the ND280 detector consists of several sub-detectors. The most important part of the detector for this analysis is the tracker region. This is comprised of two time projection chambers (TPCs) sandwiched between three fine grain detectors (FGDs). The FGDs contain both hydrocarbon plastics and water targets for neutrino interactions and provide track reconstruction near the interaction vertex. The emitted charged particles can then propagate into the TPCs which provide particle identification and momentum reconstruction. The FGDs and TPCs are further described in subsubsection 3.2.2.1 and subsubsection 3.2.2.2 respectively. The electromagnetic calorimeter (ECAL)

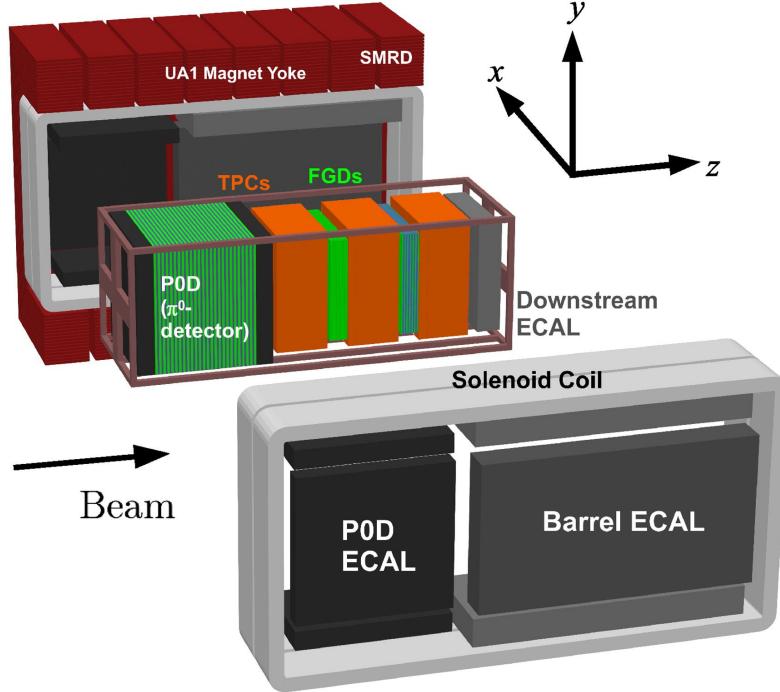


Figure 3.9: The components of the ND280 detector. The neutrino beam travels from left to right. Taken from [114].

encapsulates the tracker region alongside the π^0 detector (P0D). The ECAL measures the deposited energy from photons emitted from interactions within the FGD. The P0D constrains the cross-section of neutral current interactions which generate neutral pions, which is one of the largest backgrounds in the electron neutrino appearance oscillation channel. The P0D and ECAL detectors are detailed in subsubsection 3.2.2.3 and subsubsection 3.2.2.4 respectively. The entire detector is located within a large yoke magnet which produces a 0.2T magnetic field. This design of the magnet also includes a scintillating detector called the side muon range detector (SMRD) which is used to track high-angle muons as well as acting as a cosmic veto. The SMRD is described in subsubsection 3.2.2.5.

3.2.2.1 Fine Grained Detectors

The T2K tracker region is comprised of two fine grained detectors (FGD) and three Time Projection Chambers (TPC). A detailed description of the FGD design, construction, and assembly is found in [119] and summarised below. The FGDs are the primary target for neutrino interactions with a mass of 1.1 tonnes per FGD.

844 Alongside this, the FGDs are designed to be able to track short-range particles
845 which do not exit the FGD. Typically, short-range particles are low momentum
846 and are observed as tracks that deposit a large amount of energy per unit length.
847 This means the FGD needs good granularity to resolve these particles. The
848 FGDs have the best timing resolution ($\sim 3\text{ns}$) of any of the sub-detectors of the
849 ND280 detector. As such, the FGDs are used for time of flight measurements
850 to distinguish forward going positively charged particles from backward going
851 negatively charged particles. Finally, any tracks which pass through multiple
852 sub-detectors are required to be track matched to the FGD.

853 Both FGDs are made from square scintillator planes of side length 186cm and
854 width 2.02cm. Each plane consists of two layers of 192 scintillator bars in an
855 X or Y orientation. A wavelength shifting fiber is threaded through the centre
856 of each bar and is read out by a multi-pixel photon counter (MPPC). FGD1 is
857 the most upstream of the two FGDs and contains 15 planes of carbon plastic
858 scintillator which is a common target in external neutrino scattering data. As
859 the far detector is a pure water target, 7 of the 15 scintillator planes in FGD2
860 have been replaced with a hybrid water-scintillator target. Due to the complexity
861 of the nucleus, nuclear effects can not be extrapolated between different nuclei.
862 Therefore having the ability to take data on one target which is the same as
863 external data and another target which is the same as the far detector target is
864 beneficial for reliable model parameter estimates.

865 The integrated deposited energy is used for particle identification. The FGD
866 can distinguish protons from other charged particles by comparing the integrated
867 deposited energy from data to Monte Carlo prediction as seen in Figure 3.10.

868 3.2.2.2 Time Projection Chambers

869 The majority of particle identification and momentum measurements within
870 ND280 are provided by three Time Projection Chambers (TPCs) [120]. The
871 TPCs are located on either side of the FGDs. They are located inside of the

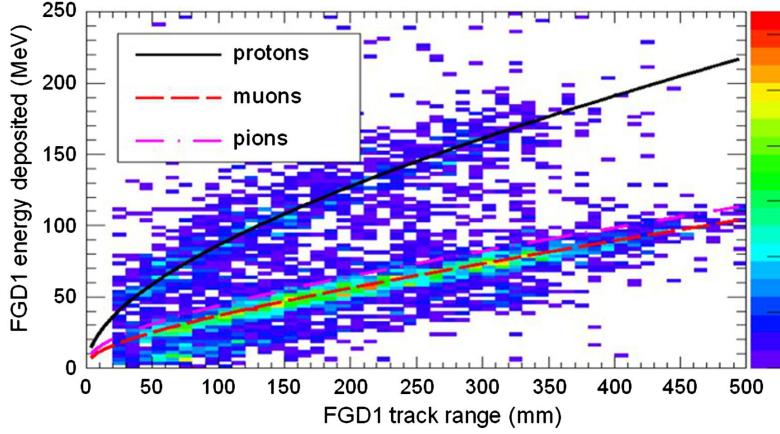


Figure 3.10: Comparison of data to Monte Carlo prediction of integrated deposited energy as a function of track length for particles that stopped in FGD1. Taken from [119].

872 magnetic field meaning the momentum of a charged particle can be determined
873 from the bending of the track.

874 Each TPC module consists of two gas-tight boxes, as shown in Figure 3.11,
875 which are made of non-magnetic material. The outer box is filled with CO₂ which
876 acts as an electrical insulator between the inner box and the ground. The inner box
877 forms the field cage which produces a uniform electric drift field of $\sim 275\text{V/cm}$
878 and is filled with an argon gas mixture. Charged particles moving through this
879 gas mixture ionize the gas and the ionised charge is drifted towards micromegas
880 detectors which measure the ionization charge. The time and position information
881 in the readout allows a three-dimensional image of the neutrino interaction.

882 The particle identification of tracks that pass through the TPCs is performed
883 using dE/dx measurements. Figure 3.12 illustrates the data to Monte Carlo
884 distributions of the energy lost by a charged particle passing through the TPC as
885 a function of the reconstructed particle momentum. The resolution is $7.8 \pm 0.2\%$
886 meaning that electrons and muons can be distinguished. This allows reliable
887 measurements of the intrinsic ν_e component of the beam.

888 3.2.2.3 π^0 Detector

889 If one of the γ -rays from a $\pi^0 \rightarrow 2\gamma$ decay is missed at the far detector, the
890 reconstruction will determine that event to be a charge current ν_e -like event.

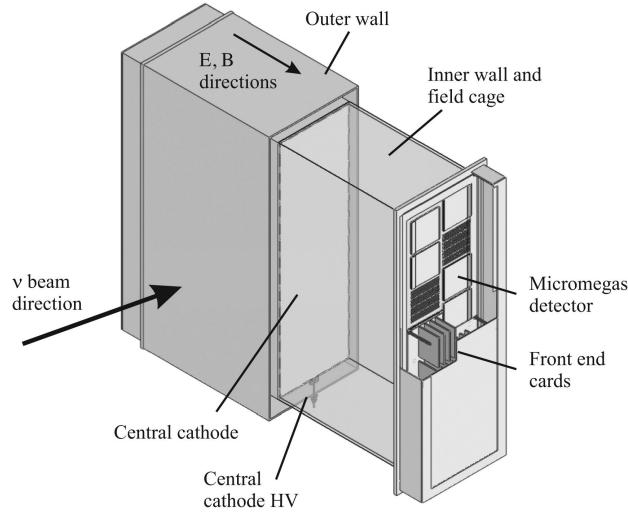


Figure 3.11: Schematic design of a Time Projection Chamber detector. Taken from [120].

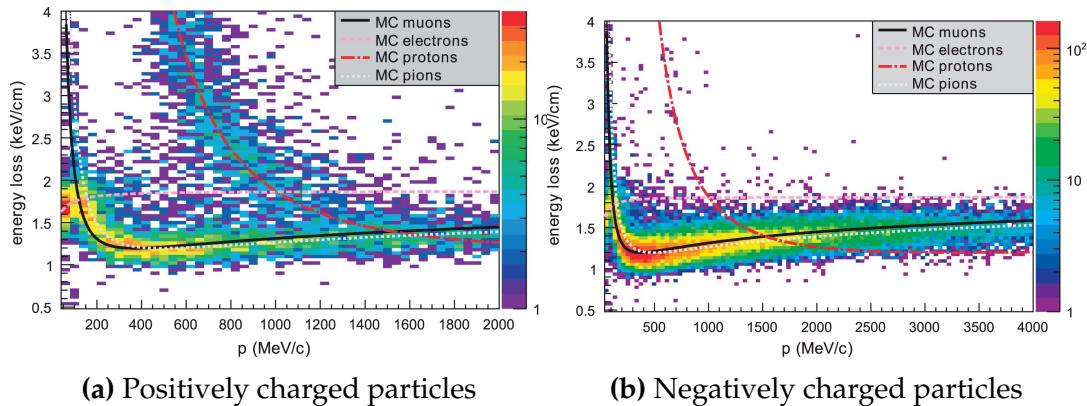


Figure 3.12: The distribution of energy loss as a function of reconstructed momentum for charged particles passing through the TPC, comparing data to Monte Carlo prediction. Taken from [120].

891 This is one of the main backgrounds hindering the electron neutrino appearance
 892 searches. The π^0 detector (P0D) measures the cross-section of the neutral current
 893 induced neutral pion production on a water target to constrain this background.

894 The P0D is a cube of approximately 2.5m length consisting of layers of scin-
 895 tillating bars, brass and lead sheets, and water bags as illustrated in Figure 3.13.

896 Two electromagnetic calorimeters are positioned at the most upstream and most
 897 downstream position in the sub-detector and the water target is situated in
 898 between them. The scintillator layers are built from two triangular bars orientated
 899 in opposite directions to form a rectangular layer. Each triangular scintillator bar
 900 is threaded with optical fiber which is read out by MPPCs. The high-Z brass and

901 lead regions produce electron showers from the photons emitted in π^0 decay.

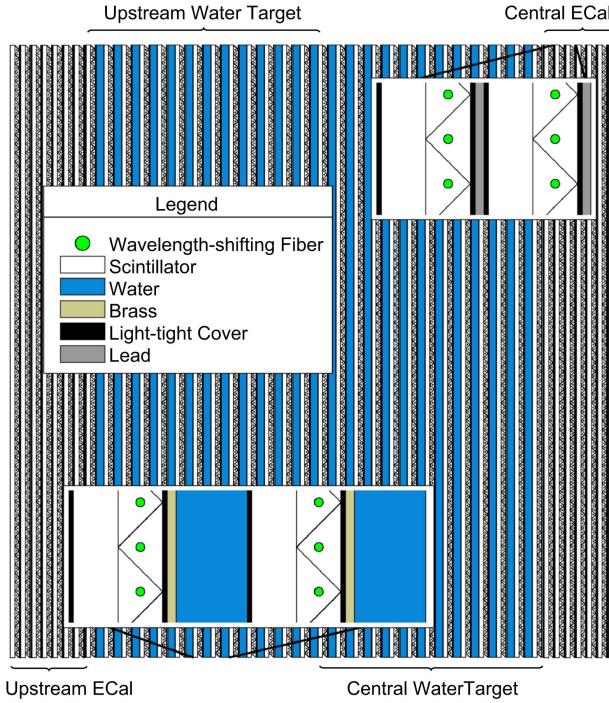


Figure 3.13: A schematic of the P0D side-view. Taken from [121].

902 The sub-detector can generate measurements of NC1 π^0 cross-sections on a
 903 water target by measuring the event rate both with and without the water target,
 904 with the cross-section on a water target being determined as the difference. The to-
 905 tal active mass is 16.1 tonnes when filled with water and 13.3 tonnes when empty.

906 3.2.2.4 Electromagnetic Calorimeter

907 The electromagnetic calorimeter [122] (ECal) encapsulates the P0D and tracking
 908 sub-detectors. Its primary purpose is to aid π^0 reconstruction from any interac-
 909 tion in the tracker. To do this, it measures the energy and direction of photon
 910 showers from $\pi^0 \rightarrow 2\gamma$ decay. It can also distinguish pion and muon tracks
 911 depending on the shape of the photon shower deposited.

912 The ECal is comprised of three sections; the P0D ECal which surrounds the
 913 P0D, the barrel ECal which encompasses the tracking region, and the downstream
 914 ECal which is situated downstream of the tracker region. The barrel and down-
 915 stream ECals are tracking calorimeters that focus on electromagnetic showers

from high-angle particles emitted from the tracking sub-detectors. Particularly in the TPC, high-angle tracks (those which travel perpendicularly to the beam-axis) can travel along a single scintillator bar resulting in very few hits. The width of the barrel and downstream ECal corresponds to ~ 11 electron radiation lengths to ensure a significant amount of the π^0 energy is contained. As the P0D has its own calorimetry which reconstructs showers, the P0D ECal determines the energy which escapes the P0D.

Each ECal is constructed of multiple layers of scintillating bars sandwiched between lead sheets. The scintillating bars are threaded with optical fiber and read out by MPPCs. Each sequential layer of the scintillator is orientated perpendicular to the previous which allows a three dimensional event reconstruction. The target mass of the P0D ECal, barrel ECal, and downstream ECal are 1.50, 4.80 and 6.62 tonnes respectively.

3.2.2.5 Side Muon Range Detector

As illustrated in Figure 3.9, the ECal, FGDs, P0D, and TPCs are enclosed within the UA1 magnet. Originally designed for the NOMAD [123] experiment and reconditioned for use in the T2K experiment [124], the UA1 magnet provides a uniform horizontal magnetic field of 0.2T with an uncertainty of 2×10^{-4} T.

Built into the UA1 magnet, the side muon range detector (SMRD)[125] monitors high-energy muons which leave the tracking region and permeate through the ECal. It additionally acts as a cosmic muon veto and trigger.

3.2.3 The Interactive Neutrino GRID

The Interactive Neutrino GRID (INGRID) detector is situated within the same “pit” as the other near detectors. It is aligned with the beam in the “on-axis” position and measures the beam direction, spread, and intensity. The detector was originally designed with 16 identical modules [114] (two modules have since been decommissioned) and a “proton” module. The design of the detector is cross-shaped with length and height 10m \times 10m as illustrated in Figure 3.14.

944 Each module is composed of iron sheets interlaced with eleven tracking
 945 scintillator planes for a total target mass of 7.1 tonnes per module. The scintillator
 946 design is an X-Y pattern of 24 bars in both orientations, where each bar contains
 947 wave-length shifting fibers which are connected to multi-pixel photon counters
 948 (MPPCs). Each module is encapsulated inside veto planes to aid the rejection
 949 of charged particles entering the module.

950 The proton module is different from the other modules in that it consists
 951 of entirely scintillator planes with no iron target. The scintillator bars are also
 952 smaller than those used in the other modules to increase the granularity of
 953 the detector and improve tracking capabilities. The module sits in the centre
 954 of the beamline and is designed to give precise measurements of quasi-elastic
 955 charged current interactions to evaluate the performance of the Monte Carlo
 956 simulation of the beamline.

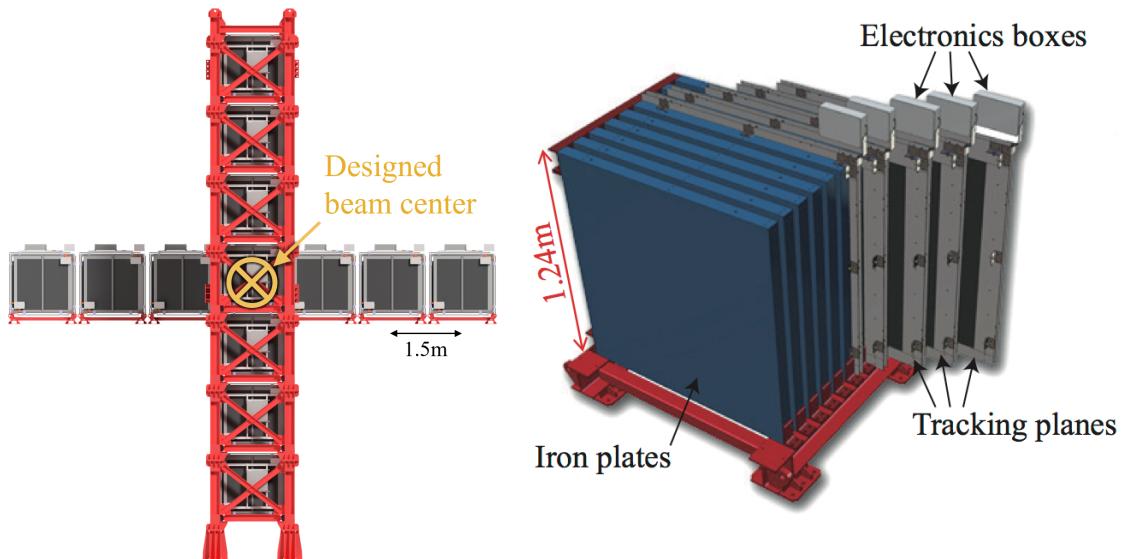


Figure 3.14: Left panel: The Interactive Neutrino GRID on-axis Detector. 14 modules are arranged in a cross-shape configuration, with the centre modules being directly aligned with the on-axis beam. Right panel: The layout of a single module of the INGRID detector. Both figures are recreated from [114].

957 The INGRID detector can measure the beam direction to an uncertainty of
 958 0.4mrad and the beam centre within a resolution of 10cm [114]. The beam
 959 direction in both the vertical and horizontal directions is discussed in [126]

⁹⁶⁰ and it is found to be in good agreement with the MUMON monitor described
⁹⁶¹ in subsection 3.2.1.

4

962

963

964

Bayesian Statistics and Markov Chain Monte Carlo Techniques

965 This thesis presents a Bayesian oscillation analysis. To extract the oscillation
966 parameters, a Markov Chain Monte Carlo (MCMC) method is used. This chapter
967 explains the theory of how parameter estimates can be determined using this
968 technique and condenses the material found in the literature [127–130].

969 The oscillation parameter determination presented within this thesis is built
970 upon a simultaneous fit to neutrino beam data in the near detector, beam data at
971 SK and atmospheric data at SK. In total, there are four oscillation parameters of in-
972 terest ($\sin^2(\theta_{23})$, $\sin^2(\theta_{13})$, Δm_{23}^2 , and δ_{CP}), two oscillation parameters to which
973 this study will not be sensitive ($\sin^2(\theta_{12})$, Δm_{12}^2) and many nuisance parameters
974 that control the systematic uncertainty models invoked within this study.

975 The MCMC technique generates a multi-dimensional probability distribution
976 across all of the model parameters used in the fit. To determine the parameter
977 estimate of a single parameter, this multi-dimensional object is integrated over all
978 other parameters. This process is called Marginalisation and is further described
979 in subsection 4.3.1. Monte Carlo techniques approximate the probability distribu-
980 tion of each parameter within the limit of generating infinite samples. As ever,
981 generating a large number of samples is time and resource-dependent. Therefore,

982 an MCMC technique is utilised within this analysis to reduce the required number
 983 of steps to sufficiently sample the parameter space. This technique is described
 984 in further detail in subsection 4.2.1.

985 **4.1 Bayesian Statistics**

986 Bayesian inference treats observable data, D , and model parameters, $\vec{\theta}$, on equal
 987 footing such that a probability model of both data and parameters is required.
 988 This is the joint probability distribution $P(D, \vec{\theta})$ and can be described by the
 989 prior distribution for model parameters $P(\vec{\theta})$ and the likelihood of the data given
 990 the model parameters $P(D|\vec{\theta})$,

$$P(D, \vec{\theta}) = P(D|\vec{\theta})P(\vec{\theta}). \quad (4.1)$$

991 The prior distribution, $P(\vec{\theta})$, describes all previous knowledge about the
 992 parameters within the model. For example, if the risk of developing health
 993 problems is known to increase with age, the prior distribution would describe the
 994 increase. For the purpose of this analysis, the prior distribution is typically
 995 the best-fit values taken from external data measurements with a Gaussian
 996 uncertainty. The prior distribution can also contain correlations between model
 997 parameters. In an analysis using Monte Carlo techniques, the likelihood of
 998 measuring some data assuming some set of model parameters is calculated
 999 by comparing the Monte Carlo prediction generated at that particular set of
 1000 model parameters to the data.

1001 It is parameter estimation that is important for this analysis and as such, we
 1002 apply Bayes' theorem [131] to calculate the probability for each parameter to have
 1003 a certain value given the observed data, $P(\vec{\theta}|D)$, which is known as the posterior
 1004 distribution (often termed the posterior). This can be expressed as

$$P(\vec{\theta}|D) = \frac{P(D|\vec{\theta})P(\vec{\theta})}{\int P(D|\vec{\theta})P(\vec{\theta})d\vec{\theta}}. \quad (4.2)$$

1005 The denominator in Equation 4.2 is the integral of the joint probability dis-
 1006 tribution over all values of all parameters used within the fit. For brevity, we
 1007 say that the posterior distribution is

$$P(\vec{\theta}|D) \propto P(D|\vec{\theta})P(\vec{\theta}). \quad (4.3)$$

1008 In subsection 4.3.1, we see that for the cases used within this analysis, it is
 1009 reasonable to know the posterior to some normalisation constant.

1010 4.1.1 Application of Prior Knowledge

1011 The posterior distribution is proportional to the prior uncertainty applied on
 1012 each parameter, as illustrated by Equation 4.3. This means that it is possible
 1013 to change the prior after the posterior distribution has been determined. The
 1014 prior uncertainty of a particular parameter can be ‘divided’ out of the posterior
 1015 distribution and the resulting distribution can be reweighted using the new
 1016 prior uncertainty that is to be applied. The methodology and implementation
 1017 of changing the prior follows that described in [132].

1018 An example implementation that is useful for this analysis is the application
 1019 of the “reactor constraint”. As discussed in section 2.4, an external constraint
 1020 on $\sin^2(\theta_{13})$ is determined from measurements taken from reactor experiments.
 1021 However, the sensitivities from just using the T2K and SK samples is equally
 1022 as important. Without this technique, two fits would have to be run, doubling
 1023 the required resources. Therefore, the key benefit for this analysis is the fact that
 1024 only a single ‘fit’ has to be performed and can be used to build the two posterior
 1025 distributions of the with and without reactor constraint applied.

1026 4.2 Monte Carlo Simulation

1027 Monte Carlo techniques are used to numerically solve a complex problem that
 1028 does not necessarily have an analytical solution. These techniques rely on

1029 building a large ensemble of samples from an unknown distribution and then
1030 using the ensemble to approximate the properties of the distribution.

1031 An example that uses Monte Carlo techniques is to calculate the area under-
1032 neath a curve. For example, take the problem of calculating the area under a
1033 straight line with gradient $M = 0.4$ and intercept $C = 1.0$. Analytically, one can
1034 calculate the area under the line is equal to 30 units for $0 \leq x \leq 10$. Using Monte
1035 Carlo techniques, one can calculate the area under this line by throwing many
1036 random values for the x and y components of each sample and then calculating
1037 whether that point falls below the line. The area can then be calculated by the
1038 ratio of points below the line to the total number of samples thrown multiplied by
1039 the total area in which samples were scattered. The study is shown in Figure 4.1
1040 highlights this technique and finds the area under the curve to be 29.9 compared
1041 to an analytical solution of 30.0. The deviation of the numerical to analytical
1042 solution can be attributed to the number of samples used in the study. The
1043 accuracy of the approximation in which the properties of the Monte Carlo samples
1044 replicate those of the desired distribution is dependent on the number of samples
1045 used. Replicating this study with a differing number of Monte Carlo samples
1046 used in each study (As shown in Figure 4.2) highlights how the Monte Carlo
1047 techniques are only accurate within the limit of a high number of samples.

1048 Whilst the above example has an analytical solution, these techniques are just
1049 as applicable to complex solutions. Clearly, any numerical solution is only as
1050 useful as its efficiency. As discussed, the accuracy of the Monte Carlo technique is
1051 dependent upon the number of samples generated to approximate the properties
1052 of the distribution. Furthermore, if the positions at which the samples are
1053 evaluated are not ‘cleverly’ picked, the efficiency of the Monte Carlo technique
1054 significantly drops. Given the example in Figure 4.1, if the region in which the
1055 samples are scattered significantly extends passed the region of interest, many
1056 calculations will be calculated but do not add to the ability of the Monte Carlo
1057 technique to achieve the correct result. For instance, any sample evaluated at
1058 a $y \geq 5$ could be removed without affecting the final result. This does bring in

- 1059 an aspect of the ‘chicken and egg’ problem in that to achieve efficient sampling,
1060 one needs to know the distribution beforehand.

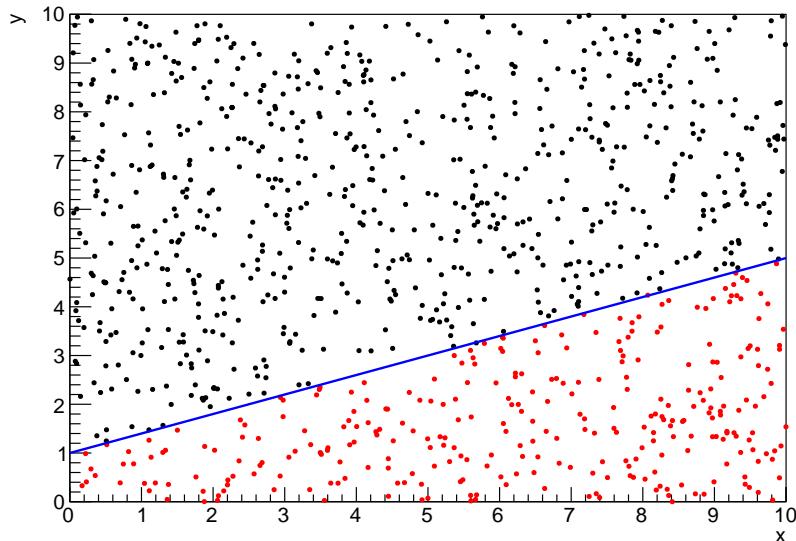


Figure 4.1: Example of using Monte Carlo techniques to find the area under the blue line. The gradient and intercept of the line are 0.4 and 1.0 respectively. The area found to be under the curve using one thousand samples is 29.9 units.

4.2.1 Markov Chain Monte Carlo

- 1061 **4.2.1 Markov Chain Monte Carlo**
- 1062 This analysis utilises a multi-dimensional probability distribution, with some
1063 dimensions being significantly more constrained than others. This could be
1064 from prior knowledge of parameter distributions from external data or un-
1065 physical regions in which parameters can not exist. Consequently, the Monte
1066 Carlo techniques used need to be as efficient as possible. For this analysis, the
1067 Markov Chain Monte Carlo (MCMC) technique is chosen. An MCMC technique
1068 is a Monte Carlo technique that uses a Markov chain to select which points at
1069 which to sample the parameter distribution. This technique performs a semi-
1070 random stochastic walk through the allowable parameter space. This builds a
1071 posterior distribution which has the property that the density of sampled points
1072 is proportional to the probability density of that parameter. This does mean that



Figure 4.2: The area under a line of gradient 0.4 and intercept 1.0 for the range $0 \leq x \leq 10$ as calculated using Monte Carlo techniques as a function of the number of samples used in each repetition. The analytical solution to the area is 30 units as given by the red line.

¹⁰⁷³ the samples produced by this technique are not statistically independent but
¹⁰⁷⁴ they will cover the space of the distribution.

¹⁰⁷⁵ A Markov chain functions by selecting the position of step \vec{x}_{i+1} based on the
¹⁰⁷⁶ position of \vec{x}_i . The space in which the Markov chain selects samples is dependent
¹⁰⁷⁷ upon the total number of parameters utilised within the fit, where a discrete point
¹⁰⁷⁸ in this space is described by the N-dimensional space \vec{x} . In a perfectly operating
¹⁰⁷⁹ Markov chain, the position of the next step depends solely on the previous step
¹⁰⁸⁰ and not on the further history of the chain (\vec{x}_0 , \vec{x}_1 , etc.). However, in solving
¹⁰⁸¹ the multi-dimensionality of the fit used within this analysis, each step becomes
¹⁰⁸² correlated with several of the steps preceding itself. This behaviour is further
¹⁰⁸³ explained in subsection 4.2.3. Providing the MCMC chain is well optimised, it will
¹⁰⁸⁴ begin to converge towards a unique stationary distribution. The period between
¹⁰⁸⁵ the chain's initial starting point and the convergence to the unique stationary
¹⁰⁸⁶ distribution is colloquially known as the burn-in period. This is discussed further
¹⁰⁸⁷ in subsection 4.2.3. Once the chain reaches the stationary distribution, all points
¹⁰⁸⁸ sampled after that point will look like samples from that distribution.

1089 Further details of the theories underpinning MCMC techniques are discussed
1090 in [128] but can be summarised by the requirement that the chain satisfies the
1091 three ‘regularity conditions’:

- 1092 • Irreducibility: From every position in the parameter space \vec{x} , there must
1093 exist a non-zero probability for every other position in the parameter space
1094 to be reached.
- 1095 • Recurrence: Once the chain arrives at the stationary distribution, every step
1096 following from that position must be samples from the same stationary
1097 distribution.
- 1098 • Aperiodicity: The chain must not repeat the same sequence of steps at any
1099 point throughout the sampling period.

1100 The output of the chain after burn-in (ie. the sampled points after the chain
1101 has reached the stationary distribution) can be used to approximate the posterior
1102 distribution and model parameters $\vec{\theta}$. To achieve the requirement that the unique
1103 stationary distribution found by the chain be the posterior distribution, one
1104 can use the Metropolis-Hastings algorithm. This guides the stochastic process
1105 depending on the likelihood of the current proposed step compared to that
1106 of the previous step. Implementation and other details of this technique are
1107 discussed in subsection 4.2.2.

1108 **4.2.2 Metropolis-Hastings Algorithm**

1109 As a requirement for MCMCs, the Markov chain implemented in this technique
1110 must have a unique stationary distribution that is equivalent to the posterior
1111 distribution. To ensure this requirement and that the regularity conditions are
1112 met, this analysis utilises the Metropolis-Hastings (MH) algorithm [133, 134].
1113 For the i^{th} step in the chain, the MH algorithm determines the position in the
1114 parameter space to which the chain moves to based on the current step, \vec{x}_i , and
1115 the proposed step, \vec{y}_{i+1} . The proposed step is randomly selected from some

1116 proposal function $f(\vec{x}_{i+1}|\vec{x}_i)$, which depends solely on the current step (ie. not
1117 the further history of the chain). The next step in the chain \vec{x}_{i+1} can be either the
1118 current step or the proposed step determined by whether the proposed step is
1119 accepted or rejected. To decide if the proposed step is selected, the acceptance
1120 probability, $\alpha(\vec{x}_i, \vec{y}_i)$, is calculated as

$$\alpha(\vec{x}_i, \vec{y}_{i+1}) = \min\left(1, \frac{P(\vec{y}_{i+1}|D)f(\vec{x}_i|\vec{y}_{i+1})}{P(\vec{x}_i|D)f(\vec{y}_{i+1}|\vec{x}_i)}\right). \quad (4.4)$$

1121 Where $P(\vec{y}_{i+1}|D)$ is the posterior distribution as introduced in section 4.1. To
1122 simplify this calculation, the proposal function is required to be symmetric such
1123 that $f(\vec{x}_i|\vec{y}_{i+1}) = f(\vec{y}_{i+1}|\vec{x}_i)$. In practice, a multi-variate Gaussian distribution is
1124 used to throw parameter proposals from. This reduces Equation 4.4 to

$$\alpha(\vec{x}_i, \vec{y}_{i+1}) = \min\left(1, \frac{P(\vec{y}_{i+1}|D)}{P(\vec{x}_i|D)}\right). \quad (4.5)$$

1125 After calculating this quantity, a random number, β , is generated uniformly
1126 between 0 and 1. If $\beta \leq \alpha(\vec{x}_i, \vec{y}_{i+1})$, the proposed step is accepted. Otherwise, the
1127 chain sets the next step equal to the current step and this procedure is repeated.
1128 This can be interpreted as if the posterior probability of the proposed step is
1129 greater than that of the current step, ($P(\vec{y}_{i+1}|D) \geq P(\vec{x}_i|D)$), the proposed
1130 step will always be accepted. If the opposite is true, ($P(\vec{y}_{i+1}|D) \leq P(\vec{x}_i|D)$),
1131 the proposed step will be accepted with probability $P(\vec{x}_i|D)/P(\vec{y}_{i+1}|D)$. This
1132 ensures that the Markov chain does not get trapped in any local minima in the
1133 potentially non-Gaussian posterior distribution. The outcome of this technique
1134 is that the density of steps taken in a discrete region is directly proportional to
1135 the probability density in that region.

1136 4.2.3 MCMC Optimisation

1137 As discussed in subsection 4.2.2, the proposal function invoked within the MH
1138 algorithm can take any form and the chain will still converge to the stationary
1139 distribution. At each set of proposed parameter values, a prediction of the same
1140 spectra has to be generated which requires significant computational resources.

1141 Therefore, the number of steps taken before the unique stationary distribution
1142 is found should be minimised as only steps after convergence add information
1143 to the oscillation analysis. Furthermore, the chain should entirely cover the
1144 allowable parameter space to ensure that all values have been considered. Tuning
1145 the distance that the proposal function jumps between steps on a parameter-by-
1146 parameter basis can both minimise the length of the burn-in period and ensure
1147 that the correlation between step \vec{x}_i and \vec{x}_j is sufficiently small.

1148 The effect of changing the width of the proposal function is highlighted in
1149 Figure 4.3. Three scenarios, each with the same underlying stationary distribution
1150 (A Gaussian of width 1.0 and mean 0.), are presented. The only difference between
1151 the three scenarios is the width of the proposal function, colloquially known as
1152 the ‘step size σ ’. Each scenario starts at an initial parameter value of 10.0 which
1153 would be considered an extreme variation. For the case where $\sigma = 0.1$, it is
1154 clear to see that the chain takes a long time to reach the expected region of the
1155 parameter. This indicates that this chain would have a large burn-in period and
1156 does not converge to the stationary distribution until step ~ 500 . Furthermore,
1157 whilst the chain does move towards the expected region, each step is significantly
1158 correlated with the previous. Considering the case where $\sigma = 5.0$, the chain
1159 approaches the expected parameter region almost instantly meaning that the
1160 burn-in period is not significant. However, there are clearly large regions of steps
1161 where the chain does not move. This is likely due to the chain proposing steps
1162 in the tails of the distribution which have a low probability of being accepted.
1163 Consequently, this chain would take a significant number of steps to fully span
1164 the allowable parameter region. For the final scenario, where $\sigma = 0.5$, you can
1165 see a relatively small burn-in period of approximately 100 steps. Once the chain
1166 reaches the stationary distribution, it moves throughout the expected region of
1167 parameter values many times, sufficiently sampling the full parameter region.
1168 This example is a single parameter varying across a continuous distribution and
1169 does not fully reflect the difficulties in the many-hundred multi-variate parameter

1170 distribution used within this analysis. However, it does give a conceptual idea of
1171 the importance of selecting the proposal function and associated step size.

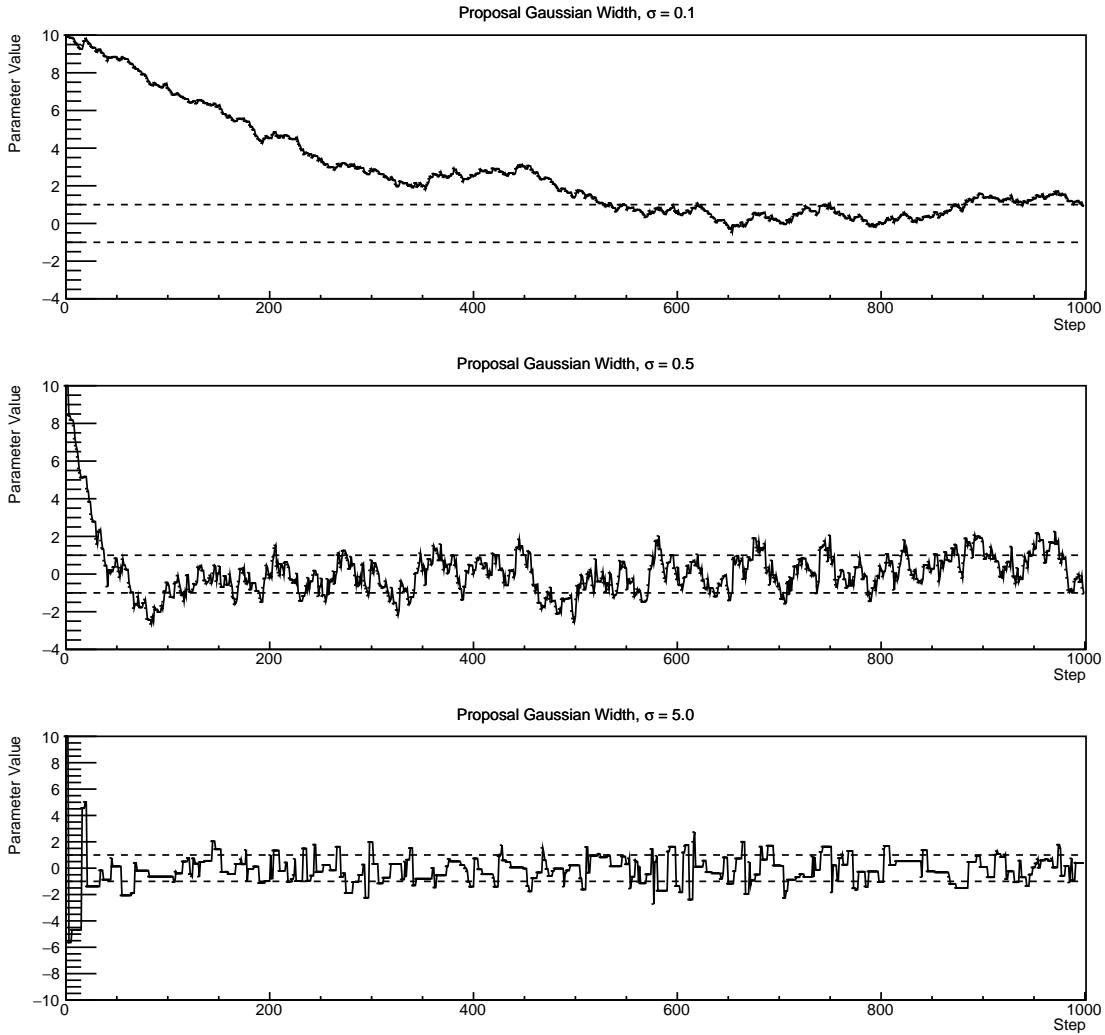


Figure 4.3: Three MCMC chains, each with a stationary distribution equal to a Gaussian centered at 0 and width 1 (As indicated by the black dotted lines). All of the chains use a Gaussian proposal function but have different widths (or ‘step size σ ’). The top panel has $\sigma = 0.1$, middle panel has $\sigma = 0.5$ and the bottom panel has $\sigma = 5.0$.

1172 As discussed, step size tuning directly correlates to the average step accep-
1173 tance rate. If the step size is too small, many steps will be accepted but the
1174 chain moves slowly. If the opposite is true, many steps will be rejected as the
1175 chain proposes steps in the tails of the distribution. Discussion in [135] suggests
1176 that the ‘ideal’ acceptance rate of a high dimension MCMC chain should be

¹¹⁷⁷ approximately $\sim 25\%$. An “ideal” step size [135] of

$$\sigma = \frac{2.4}{N_p}, \quad (4.6)$$

¹¹⁷⁸ where N_p is the number of parameters included in the MCMC fit. However,
¹¹⁷⁹ the complex correlations between systematics mean that some parameters have
¹¹⁸⁰ to be hand tuned and many efforts have been taken to select a set of parameter-
¹¹⁸¹ by-parameter step sizes to approximately reach the ideal acceptance rate.

¹¹⁸² Figure 4.3 highlights the likelihood as calculated by the fit in DB: [Link to](#)
¹¹⁸³ [AsimovA Sensitivity Section](#) as a function of the number of steps in each chain. In
¹¹⁸⁴ practice, many independent MCMC chains are run simultaneously to parallelise
¹¹⁸⁵ the task of performing the fit. This figure overlays the distribution found in each
¹¹⁸⁶ chain. As seen, the likelihood decreases from its initial value and converges
¹¹⁸⁷ towards a stationary distribution after $\sim 1 \times 10^5$ steps.

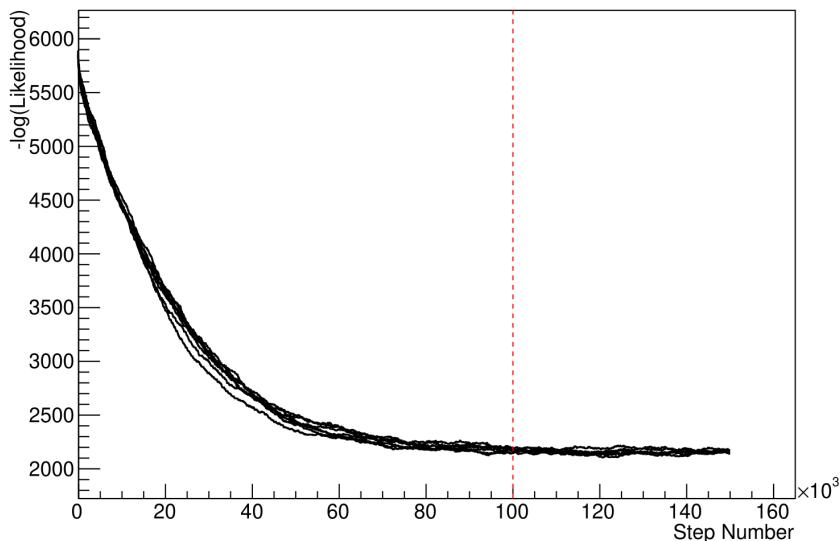


Figure 4.4: The log-likelihood from the fit detailed in DB: [Link to AsimovA Sensitivity Section](#) as a function of the number of steps accumulated in each fit. Many independent MCMC chains were run in parallel and overlaid on this plot. The red line indicates the 1×10^5 step burn-in period after which the log-likelihood becomes stable.

¹¹⁸⁸ Multiple configurations of this analysis have been performed throughout this
¹¹⁸⁹ thesis where different samples or systematics have been used. For all of these con-
¹¹⁹⁰ figurations, it was found that a burnin period of 1×10^5 was sufficient in all cases.

1191 4.3 Understanding the MCMC Results

1192 The previous sections have described how to generate the posterior probability
1193 distribution using Bayesian MCMC techniques. However, this analysis focuses
1194 on oscillation parameter determination. The posterior distribution output from
1195 the chain is a high dimension object, with as many dimensions as there are
1196 parameters included in the oscillation analysis. However, this multi-dimensional
1197 object is difficult to conceptualize so parameter estimations are often presented
1198 in one or two-dimensional projections of this probability distribution. To do this,
1199 we invoke the marginalisation technique highlighted in subsection 4.3.1.

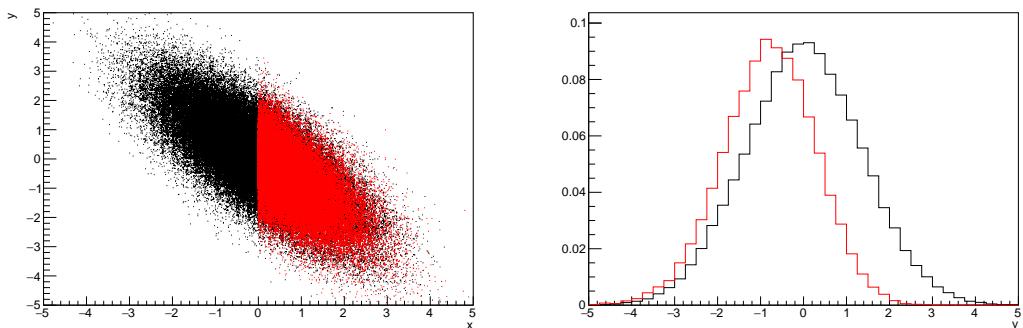
1200 4.3.1 Marginalisation

1201 The output of the MCMC chain is a highly dimensional probability distribution
1202 which is very difficult to interpret. From the standpoint of an oscillation analysis
1203 experiment, the one or two-dimensional ‘projections’ of the oscillation parameters
1204 of interest are most relevant. Despite this, the best fit values and uncertainties on
1205 the oscillation parameters of interest should correctly encapsulate the correlations
1206 to the other systematic uncertainties (colloquially called ‘nuisance’ parameters).
1207 For this joint beam and atmospheric analysis, the oscillation parameters of interest
1208 are $\sin^2(\theta_{23})$, $\sin^2(\theta_{13})$, Δm_{23}^2 , and δ_{CP} . All other parameters (Including the
1209 oscillation parameter this fit is insensitive to) are deemed nuisance parameters.
1210 To generate these projections, we rely upon integrating the posterior distribution
1211 over all nuisance parameters. This is called marginalisation. A simple example
1212 of this technique is to imagine the scenario where two coins are flipped. To
1213 determine the probability that the first coin returned a ‘head’, the exact result of
1214 the second coin flip is disregarded and simply integrated over. For the parameters
1215 of interest, $\vec{\theta}_i$, we can calculate the marginalised posterior by integrating over
1216 the nuisance parameters, $\vec{\theta}_n$. In this case, Equation 4.2 becomes

$$P(\vec{\theta}_i|D) = \frac{\int P(D|\vec{\theta}_i, \vec{\theta}_n)P(\vec{\theta}_i, \vec{\theta}_n)d\vec{\theta}_n}{\int P(D|\vec{\theta})P(\vec{\theta})d\vec{\theta}} \quad (4.7)$$

1217 Where $P(\vec{\theta}_i, \vec{\theta}_n)$ encodes the prior knowledge about the uncertainty and
 1218 correlations between the parameters of interest and the nuisance parameters.
 1219 In practice, this is simply taking the one or two-dimensional projection of the
 1220 multi-dimensional probability distribution.

1221 Whilst in principle an easy solution to a complex problem, correlations
 1222 between the interesting and nuisance parameters can bias the marginalised
 1223 results. A similar effect is found when the parameters being marginalised over
 1224 have non-Gaussian probability distributions. For example, Figure 4.5 highlights
 1225 the marginalisation bias in the probability distribution found for a parameter
 1226 when requiring a correlated parameter to have a positive parameter value. Due
 1227 to the complex nature of this oscillation parameter fit presented in this thesis,
 1228 there are correlations occurring between the oscillation parameters of interest
 1229 and the other nuisance parameters included in the fit.



1230 **Figure 4.5:** Left: The two dimensional probability distribution for two correlated
 parameters x and y . The red distribution shows the two dimensional probability
 distribution when $0 \leq x \leq 5$. Right: The marginalised probability distribution for
 the y parameter found when requiring the x to be bound between $-5 \leq x \leq 5$ and
 $0 \leq x \leq 5$ for the black and red distribution, respectively.

1230 4.3.2 Parameter Estimation and Credible Intervals

1231 The purpose of this analysis is to determine the best fit values for the oscillation
 1232 parameters that the beam and atmospheric samples are sensitive to: $\sin^2(\theta_{23})$
 1233 , $\sin^2(\theta_{13})$, Δm_{23}^2 , and δ_{CP} . Typically, the results presented take the form of
 1234 one or two-dimension marginalised probability distributions for the appearance

1235 ($\sin^2(\theta_{13})$ and δ_{CP}) and disappearance ($\sin^2(\theta_{23})$ and Δm_{23}^2) parameters. The
1236 posterior probability density taken from the output MCMC chain is binned in
1237 these parameters. The parameter best-fit point is then taken to be the value
1238 that has the highest posterior probability. This is performed in both one and
1239 two-dimensional projections.

1240 However, the single best-fit point in a given parameter is not of much use on
1241 its own. We would also like to determine the uncertainty, or credible interval,
1242 on that best-fit point. The definition of the 1σ credible interval is that we have
1243 68% belief that the parameter is within those bounds. For a more generalised
1244 definition, the credible interval is the region, R , of the posterior distribution that
1245 contains a specific fraction of the total probability, such that

$$\int_R P(\theta|D) d\theta = \alpha \quad (4.8)$$

1246 Where θ is the parameter on which we calculate the credible interval. This
1247 technique then calculates the $\alpha \times 100\%$ credible interval.

1248 In practice, this analysis uses the highest posterior density (HPD) credible in-
1249 tervals which are calculated through the following method. First, the probability
1250 distribution is area-normalised such that it has an integrated area equal to 1.0.
1251 The bins of probability are then summed from the highest to lowest until the sum
1252 exceeds the 1σ level (0.68 in this example). This process is repeated for a range of
1253 credible intervals, notably the 1σ , 2σ and 3σ along with other levels where the
1254 critical values for each level can be found in [75]. This process can be repeated
1255 for the two-dimensional probability distributions by creating two-dimensional
1256 contours of credible intervals rather than a one-dimensional result.

1257 4.3.3 Bayesian Model Comparisons

1258 Due to the matter resonance, this analysis has some sensitivity to the mass
1259 hierarchy of neutrino states (whether Δm_{23}^2 is positive or negative) and the
1260 octant of $\sin^2(\theta_{23})$. The Bayesian approach utilised within this analysis gives an
1261 intuitive method of model comparison by determining which hypothesis is most

₁₂₆₂ favourable. Taking the ratio of Equation 4.3 for the two hypotheses of normal
₁₂₆₃ hierarchy, NH , and inverted hierarchy, IH , gives

$$\frac{P(\vec{\theta}_{NH}|D)}{P(\vec{\theta}_{IH}|D)} = \frac{P(D|\vec{\theta}_{NH})}{P(D|\vec{\theta}_{IH})} \times \frac{P(\vec{\theta}_{NH})}{P(\vec{\theta}_{IH})}. \quad (4.9)$$

₁₂₆₄ The middle term defines the Bayes factor which is a data-driven interpretation
₁₂₆₅ of how strong the data prefers one hierarchy to the other. For this analysis, equal
₁₂₆₆ priors on both mass hierarchy hypotheses are chosen ($P(\vec{\theta}_{NH}) = P(\vec{\theta}_{IH}) = 0.5$).
₁₂₆₇ In practice, the MCMC chain proposes a value of $|\Delta m_{23}^2|$ and then applies a
₁₂₆₈ 50% probability that the value is sign flipped. Consequently, the Bayes factor
₁₂₆₉ can be calculated from the ratio of the probability density in either hypothesis.
₁₂₇₀ This equates to counting the number of steps taken in the normal and inverted
₁₂₇₁ hierarchies and taking the ratio. The same approach can be taken to compare the
₁₂₇₂ upper octant (UO) compared to the lower octant (LO) hypothesis of $\sin^2(\theta_{23})$.

₁₂₇₃ Whilst the value of the Bayes factor should always be shown, the Jeffreys scale
₁₂₇₄ [136] (highlighted in Table 4.1) gives an indication of the strength of preference
₁₂₇₅ for one model compared to the other. Other interpretations of the strength of
₁₂₇₆ preference of a model exist, e.g. the Kass and Raferty Scale [137].

$\log_{10}(B_{AB})$	B_{AB}	Strength of Preference
< 0.0	< 1	No preference for hypothesis A (Supports hypothesis B)
0.0 – 0.5	1.0 – 3.16	Preference for hypothesis A is weak
0.5 – 1.0	3.16 – 10.0	Preference for hypothesis A is substantial
1.0 – 1.5	10.0 – 31.6	Preference for hypothesis A is strong
1.5 – 2.0	31.6 – 100.0	Preference for hypothesis A is very strong
> 2.0	> 100.0	Decisive preference for hypothesis A

Table 4.1: Jeffreys scale for strength of preference for two models A and B as a function of the calculated Bayes factor ($B_{AB} = B(A/B)$) between the two models [136]. The original scale is given in terms of $\log_{10}(B(A/B))$ but converted to linear scale for easy comparison throughout this thesis.

₁₂₇₇ 4.3.4 Comparison of MCMC Output to Expectation

₁₂₇₈ To ensure the fit is performing well, a best-fit spectrum is produced using the
₁₂₇₉ posterior probability distribution and compared with the data, allowing easy

1280 by-eye comparisons to be made. A simple method of doing this is to perform a
1281 comparison in the fitting parameters (For instance, the reconstructed neutrino
1282 energy and lepton direction for T2K far detector beam samples) of the spectra
1283 generated by the MCMC chain to ‘data’. This ‘data’ could be true data or some
1284 variation of Monte Carlo prediction. This allows easy comparison of the MCMC
1285 probability distribution to the data. To perform this, N steps from the post burn-in
1286 MCMC chain are randomly selected (Where for all plots of this style in this thesis,
1287 $N = 3000$). From these, the Monte Carlo prediction at each step is generated by
1288 reweighting the model parameters to the values specified at that step. Due to the
1289 probability density being directly correlated with the density of steps in a certain
1290 region, parameter values close to the best fit value are most likely to be selected.

1291 In practice, for each bin of the fitting parameters has a probability distribution
1292 of event rates, with one entry per sampled MCMC step. This distribution is
1293 binned where the bin with the highest probability is selected as the mean and an
1294 error on the width of this probability distribution is calculated using the approach
1295 highlighted in subsection 4.3.2. Consequently, the best fit distribution in the fit
1296 parameter is not necessarily that which would be attained by reweighting the
1297 Monte Carlo prediction to the most probable parameter values.

1298 A similar study can be performed to illustrate the freedom of the model
1299 parameter space prior to the fit. This can be done by throwing parameter values
1300 from the prior uncertainty of each parameter. This becomes troublesome for
1301 parameters with no prior uncertainty as the range is technically infinite. Where
1302 applicable solutions to remove these have been addressed.

5

1303

1304

1305

Simulation, Reconstruction, and Event Reduction

1306 As a crucial part of the oscillation analysis, an accurate prediction of the expected
1307 neutrino spectrum at the far detector is required. This includes modeling the
1308 flux generation, neutrino interactions, and detector effects. All of the simulation
1309 packages required to do this are briefly described in section 5.1. The reconstruc-
1310 tion of neutrino events inside the far detector, including the `fitQun` algorithm,
1311 is documented in section 5.2. This also includes data quality checks of the SK-
1312 V data which the author performed for the T2K oscillation analysis presented
1313 at Neutrino 2020 [74]. Finally, section 5.3 describes the steps taken in the SK
1314 detector to trigger on events of interest whilst removing the comparatively large
1315 rate of cosmic ray muon events.

1316 5.1 Simulation

1317 In order to generate a Monte Carlo prediction of the expected event rate at the far
1318 detector, all the processes in the beam and atmospheric flux, neutrino interaction,
1319 and detector need to be modeled. Each of these parts is individually modeled
1320 and each of them is detailed below.

1321 The beamline simulation consists of three distinct parts: the initial hadron
1322 interaction modeled by FLUKA [138], the target station geometry and particle
1323 tracking performed by JNUBEAM, [139, 140] and any hadronic re-interactions
1324 simulated by GCALOR [141]. The primary hadronic interactions are $O(10)\text{GeV}$,
1325 where FLUKA matches external cross-section data better than GCALOR [142].
1326 However, FLUKA is not very adaptable so a small simulation is built to model
1327 the interactions in the target and the output is then passed to JNUBEAM and
1328 GCALOR for propagation. The hadronic interactions are tuned to data from the
1329 NA61/SHINE [143–145] and HARP [146] experiments. The tuning is done by
1330 reweighting the FLUKA and GCALOR predictions to match the external data
1331 multiplicity and cross-section measurements, based on final state particle kine-
1332 matics [142]. The culmination of this simulation package generates the predicted
1333 flux for neutrino and antineutrino beam modes which are illustrated in Figure 3.7.

1334 The atmospheric neutrino flux is simulated by the HKKM model [50, 52]. The
1335 primary cosmic ray flux is tuned to AMS [147] and BESS [148] data assuming
1336 the US-standard atmosphere '76 [149] density profile and includes geomagnetic
1337 field effects. The primary cosmic rays interact to generate pions and muons.
1338 The interaction of these secondary particles to generate neutrinos is handled by
1339 DPMJET-III [150] for energies above 32GeV and JAM [52, 151] for energies below
1340 that value **DB: Question for Giles: Why different generators for above/below**
1341 **32GeV?**. These hadronic interactions are tuned to BESS and L3 data [152, 153]
1342 using the same methodology as the tuning of the beamline simulation. The
1343 energy and cosine zenith predictions of $\nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu$ flux are given in Figure 2.3
1344 and Figure 2.5, respectively. The flux is approximately symmetrical and peaked
1345 around the horizon ($\cos(\theta_Z) = 0.0$). This is because horizontally-going pions
1346 and kaons can travel further than their vertically-going counterparts resulting
1347 in a larger probability of decaying to neutrinos. The symmetry is broken in
1348 lower-energy neutrinos due to geomagnetic effects, which modify the track of the
1349 primary cosmic rays. Updates to the HKKM model are currently ongoing [154].

Once a flux prediction has been made for all three detectors, NEUT 5.4.0 [155, 156] models the interactions of the neutrinos in the detectors. For the purposes of this analysis, quasi-elastic (QE), meson exchange (MEC), single meson production (PROD), coherent pion production (COH), and deep inelastic scattering (DIS) interactions are simulated. These interaction categories can be further broken down by whether they were propagated via a W^\pm boson in Charged Current (CC) interactions or via a Z^0 boson in Neutral Current (NC) interactions. CC interactions have a charged lepton in the final state, which can be flavour-tagged in reconstruction to determine the flavour of the neutrino. In contrast, NC interactions have a neutrino in the final state so no flavour information can be determined from the observables left in the detector after an interaction. This is the reason why neutrinos which interact through NC modes are assumed to not oscillate within this analysis. Both CC and NC interactions are modeled for all the above interaction categories, other than MEC interactions which are only modeled for CC events.

As illustrated in Figure 5.1, CCQE interactions dominate the cross-section of neutrino interactions around $E_\nu \sim 0.5\text{GeV}$. The NEUT implementation adopts the Llewellyn Smith [157] model for neutrino-nucleus interactions, where the nuclear ground state of any bound nucleons (neutrino-oxygen interactions) is approximated by a spectral-function [158] model that simulates the effects of Fermi momentum and Pauli blocking. The cross-section of QE interactions is controlled by vector and axial-vector form factors parameterised by the BBBA05 [159] model and a dipole form factor with $M_A^{QE} = 1.21\text{GeV}$ fit to external data [160], respectively. NEUT implements the Valencia [161] model to simulate MEC events, where two nucleons and two holes in the nuclear target are produced (Often called 2p2h interactions).

For neutrinos of energy $O(1)\text{GeV}$, PROD interactions become dominant. These predominantly produce charged and neutral pions although γ , kaon, and η production is also considered. To simulate these interactions, the Berger-Sehgal [162] model is implemented within NEUT. It simulates the excitation of a



Figure 5.1: The NEUT prediction of the ν_μ -H₂O cross-section overlaid on the T2K ν_μ flux. The charged current (black, solid) and neutral current (black, dashed) inclusive, charged current quasi-elastic (blue, solid), charged current 2p2h (blue, dashed), charged current single pion production (pink), and charged current multi- π and DIS (Purple) cross-sections are illustrated. Figure taken from [155].

nucleon from a neutrino interaction, production of an intermediate baryon, and the subsequent decay to a single meson or γ . Pions can also be produced through COH interactions, which occur when the incoming neutrino interacts with the entire oxygen nucleus leaving a single pion outside of the nucleus. NEUT utilises the Berger-Sehgal [163] model to simulate these COH interactions.

DIS and multi- π producing interactions become the most dominant for energies $> O(5)\text{GeV}$. PYTHIA [164] is used to simulate any interaction with invariant mass $W > 2\text{GeV}/c^2$, which produces at least one meson. For any interaction which produces at least two mesons but has $W < 2\text{GeV}/c^2$, the Bronner model is used [165]. Both of these models use parton distribution functions based on the Bodek-Yang model [166–168].

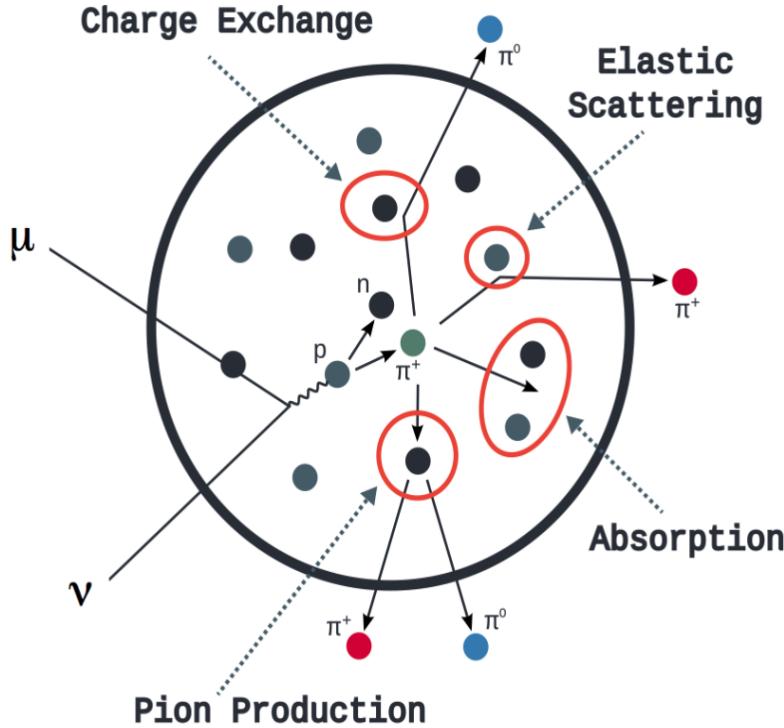


Figure 5.2: Illustration of the various processes which a pion can undergo before exiting the nucleus. Taken from [169].

1391 Any pion which is produced within the nucleus can re-interact through
 1392 final state interactions before it exits, as illustrated by the scattering, absorption,
 1393 production, and exchange interactions in Figure 5.2. These re-interactions alter
 1394 the observable particles within the detector. For instance, if the charged pion
 1395 from a CC PROD interaction is absorbed, the observables would mimic a CC QE
 1396 interaction. To simulate these effects, NEUT uses a semi-classical intranuclear
 1397 cascade model [155]. This cascade functions by stepping the pion through the
 1398 nucleus in fixed-length steps equivalent to $dx = R_N/100$, where R_N is the radius
 1399 of the nucleus. At each step, the simulation allows the pion to interact through
 1400 scattering, charged exchange, absorption, or production with an interaction-
 1401 dependent probability calculated from a fit to external data [170]. This cascade
 1402 continues until the pion is absorbed or exits the nucleus.

1403 Once the final state particle kinematics have been determined by NEUT,
 1404 they are passed into the detector simulation. The near detectors, ND280 and

1405 INGRID, are simulated using a GEANT4 package [114, 171] to simulate the detector
1406 geometry, particle tracking, and energy deposition. The response of the detectors
1407 is simulated using the elecSim package [114]. The far detector simulation is based
1408 upon the original Kamiokande experiment software which uses the GEANT3-based
1409 SKDETSIM [114, 172] package. This simulates the interactions of particles in
1410 the water as well as Cherenkov light production. The water quality and PMT
1411 calibration measurements detailed in subsection 3.1.2 are also used within this
1412 simulation to make accurate predictions of the detector response.

1413 5.2 Event Reconstruction at SK

1414 Any event which generates optical photons that occurs in SK will be observed by
1415 the PMT array, where each PMT records the time and accumulated charge. This
1416 recorded information is shown in event displays similar to those illustrated in
1417 Figure 5.3. To be useful for physics analyses, this series of PMT hit information
1418 needs to be reconstructed to determine the particle’s identity and kinematics (or
1419 track parameters): four-vertex, direction, and momenta. The reconstruction
1420 uses the fact that the charge and timing distribution of photons generated
1421 by a particular particle in an event is dependent upon its initial kinematics.
1422 Electron and muon rings are distinguished by their “fuzziness”. Muons are
1423 heavier and less affected by scattering or showering meaning they typically
1424 produce “crisp” rings. Electrons are more likely to interact via electromagnetic
1425 showering or scattering which results in larger variations of their direction from
1426 the initial direction. Consequently, electrons typically produce “fuzzier” rings
1427 compared to muons.

1428 For the purposes of this analysis, the `fiTQun` reconstruction algorithm is
1429 utilised. Its core function is to compare a prediction of the accumulated charged
1430 and timing distribution from each PMT, generated for a particular particle identity
1431 and track parameters, to that observed in the neutrino event. It determines the
1432 preferred values by minimising a likelihood function which includes information
1433 from PMTs which were hit and those that were not hit. `fiTQun` performs a

1434 simultaneous fit of particle kinematics and identity, improving both the accuracy
 1435 of the fit parameters and the rejection of neutral current π^0 events [173, 174]. The
 1436 `fitQun` algorithm is based on the key concepts of the MiniBooNE reconstruction
 1437 algorithm [175] and is described in [176] which is summarised below. The
 1438 `fitQun` algorithm improves upon the `APFit` reconstruction algorithm which has
 1439 been used for many previous SK analyses. `APFit` fits the vertex from timing
 1440 information and then fits the momentum and direction of the particle from PMT
 1441 hits within a 43 deg Cherenkov cone (which assumes an ultra-relativistic particle).
 1442 It then fits the particle identity once the track parameters have been fit.

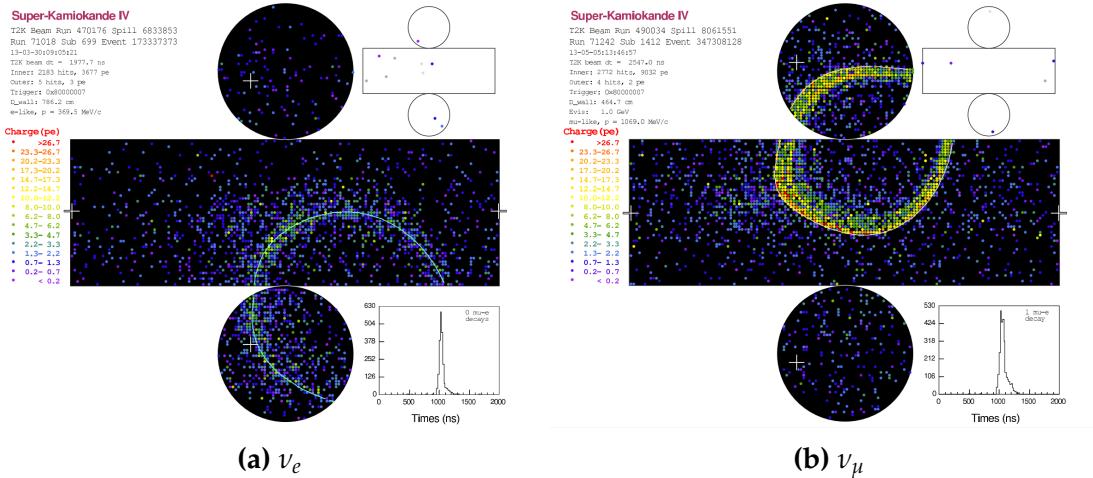


Figure 5.3: Event displays from Super Kamiokande illustrating the “crisp” ring from a muon and the typically “fuzzy” electron ring. Each pixel represents a PMT and the color scheme denotes the accumulated charge deposited on that PMT. Figures taken from [177].

1443 The `fitQun` reconstruction algorithm proceeds by:

- 1444 • **Vertex pre-fitting:** An estimate of the vertex is made using a goodness-of-fit
 1445 metric
- 1446 • **Peak finding:** The initial time of the event is determined by clustering
 1447 events by time residuals
- 1448 • **Single-ring fits:** Given the pre-fit vertex and estimated time of interaction,
 1449 a maximum likelihood technique searches for a single particle generating
 1450 light

- 1451 • **Multi-ring fits:** Seeded from the single-ring fits, hypotheses with multiple
 1452 light-producing particles are considered using the same maximum likeli-
 1453 hood technique

1454 An event in SK can consist of a primary and decay particles. For example,
 1455 a charged current muon neutrino interaction can generate two particles that
 1456 have the potential of generating Cherenkov photons: the primary muon, and
 1457 the secondary decay-electron from the muon. To ensure the particles are recon-
 1458 structed separately, each event is divided into time clusters which are called
 1459 “subevents”. Subevents after the primary subevent are considered to be decay
 1460 electrons. To find all the subevents in an event, a vertex goodness metric is
 1461 calculated for some vertex position \vec{x} and time t ,

$$G(\vec{x}, t) = \sum_i^{\text{hit PMTs}} \exp \left(-\frac{1}{2} \left(\frac{T_{Res}^i(\vec{x}, t)}{\sigma} \right)^2 \right) \quad (5.1)$$

1462 where

$$T_{Res}^i(\vec{x}, t) = t^i - t - |R_{PMT}^i - \vec{x}| / c_n \quad (5.2)$$

1463 is the residual hit time. It is the difference in time between the PMT hit time,
 1464 t^i , of the i^{th} PMT and the expected time of the PMT hit if the photon was emitted
 1465 at the start of the vertex. R_{PMT}^i is the position of the i^{th} PMT, c_n is the speed of
 1466 light in water and $\sigma = 4\text{ns}$ which is comparable to the time resolution of the
 1467 PMT. When the proposed fit values of time and vertex are close to the true values,
 1468 $T_{Res}^i(\vec{x}, t)$ tends to zero resulting in subevents appearing as spikes in the goodness
 1469 metric. The proposed fit vertex and time are grid-scanned, and the values which
 1470 maximise the goodness metric are selected as the “pre-fit vertex”. Whilst this
 1471 predicts a vertex for use in the clustering algorithm, the final vertex is fit using
 1472 the higher-precision maximum likelihood method described below.

1473 Once the pre-fit vertex has been determined, the goodness metric is scanned as
 1474 a function of t to determine the number of subevents. A peak-finding algorithm
 1475 is then used on the goodness metric, requiring the goodness metric to exceed

some threshold and drop below a reduced threshold before any subsequent additional peaks are considered. The thresholds are set such that the rate of false peak finding is minimised while still attaining good data to Monte Carlo agreement. To improve performance, the pre-fit vertex for each delayed subevent is re-calculated after PMT hits from the previous subevent are masked. This improves the decay-electron tagging performance. Once all subevents have been determined, the time window around each subevent is then defined by the earliest and latest time which satisfies $-180 < T_{Res}^i < 800\text{ns}$. The subevents and associated time windows are then used as seeds for further reconstruction.

For a given subevent, the `fitQun` algorithm constructs a likelihood based on the accumulated charge q_i and time information t_i from the i^{th} PMT,

$$L(\Gamma, \vec{\theta}) = \prod_j^{\text{unhit}} P_j(\text{unhit}|\Gamma, \vec{\theta}) \prod_i^{\text{hit}} \{1 - P_i(\text{unhit}|\Gamma, \vec{\theta})\} f_q(q_i|\Gamma, \vec{\theta}) f_t(t_i|\Gamma, \vec{\theta}), \quad (5.3)$$

where $\vec{\theta}$ defines the track parameters; vertex position, direction vector and momenta, and Γ represents the particle hypothesis. $P_i(\text{unhit}|\Gamma, \vec{\theta})$ is the probability of the i^{th} tube to not register a hit given the track parameters and particle hypothesis. The charge likelihood, $f_q(q_i|\Gamma, \vec{\theta})$, and time likelihood, $f_t(t_i|\Gamma, \vec{\theta})$, represents the probability density function of observing charge q_i and time t_i on the i^{th} PMT given the specified track parameters and particle hypothesis.

The predicted charge is calculated based on contributions from both the direct light and the scattered light. The direct light contribution is determined based on the integration of the Cherenkov photon profile along the track. PMT angular acceptance, water quality, and calibration measurements discussed in subsection 3.1.2 are included to accurately predict the charge probability density at each PMT. The scattered light is calculated in a similar way, although it includes a scattering function that depends on the vertex of the particle and the position of the PMT. The charge likelihood is calculated by comparing the prediction to the observed charge in the PMT.

1502 The time likelihood is approximated to depend on the vertex \vec{x} , direction \vec{d} ,
 1503 and time t of the track as well as the particle hypothesis. The expected time
 1504 for PMT hits is calculated by assuming unscattered photons being emitted from
 1505 the midpoint of the track, S_{mid} ,

$$t_{exp}^i = t + S_{mid}/c + |R_{PMT}^i - \vec{x} - S_{mid}\vec{d}|/c_n, \quad (5.4)$$

1506 where c is the speed of light in a vacuum. The time likelihood is then expressed
 1507 in terms of the residual difference between the PMT hit time and the expected
 1508 hit time, $t_{Res}^i = t^i - t_{exp}^i$. The particle hypothesis and momentum also affect the
 1509 Cherenkov photon distribution. These parameters modify the shape of the time
 1510 likelihood density since in reality not all photons are emitted at the midpoint of
 1511 the track. As with the charge likelihood, the contributions from both the direct
 1512 and scattered light to the time likelihood density are calculated separately, which
 1513 are both calculated from particle gun studies.

1514 The track parameters and particle identity which maximise $L(\Gamma, \vec{\theta})$ are defined
 1515 as the best-fit parameters. In practice MINUIT [178] is used to minimise the value
 1516 of $-\ln L(\Gamma, \vec{\theta})$. The `fiTQun` algorithm considers an electron-like, muon-like, and
 1517 charged pion-like hypothesis for events with a single final state particle, denoted
 1518 “single-ring events”. The particle’s identity is determined by taking the ratio of
 1519 the likelihood of each of the hypotheses. For instance, electrons and muons are
 1520 distinguished by considering the value of $\ln(L(e, \vec{\theta}_e)/L(\mu, \vec{\theta}_\mu))$ in comparison
 1521 to the reconstructed momentum of the electron hypothesis, as illustrated by
 1522 Figure 5.4. The coefficients of the discriminator between electron-like and muon-
 1523 like events are determined from Monte Carlo studies [176]. Similar distributions
 1524 exist for distinguishing electron-like events from π^0 -like events, and muon-like
 1525 events from pion-like events. The cuts are defined as,

$$\text{Electron/Muon} : \ln(L_e/L_\mu) > 0.2 \times p_e^{rec} [\text{MeV}],$$

$$\text{Electron}/\pi^0 : \ln(L_e/L_{\pi^0}) < 175 - 0.875 \times m_{\gamma\gamma} [\text{MeV}], \quad (5.5)$$

$$\text{Muon/Pion} : \ln(L_\mu/L_{\pi^\pm}) < 0.15 \times p_\mu^{rec} [\text{MeV}],$$

as taken from [179], where p_e^{rec} and p_μ^{rec} are the reconstructed momentum of the single-ring electron and muon fits, respectively. $m_{\gamma\gamma}$ represents the reconstructed invariant mass of the two photons emitted from π^0 decay. Typically, the distance between a particular entry in these two-dimensional distributions and the cut-line is termed the PID parameter and is illustrated in Figure 5.5.

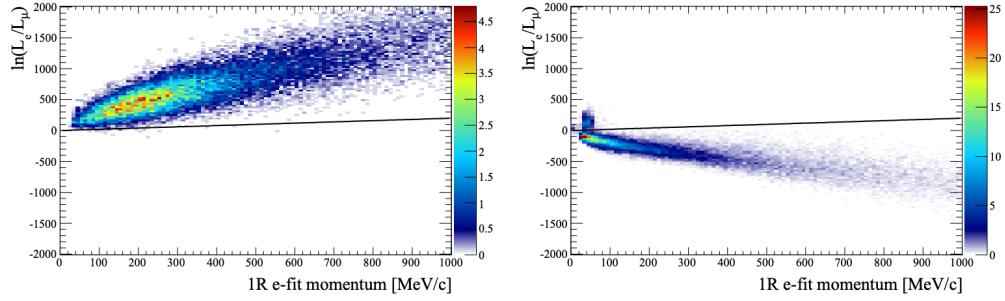


Figure 5.4: The difference of the electron-like and muon-like log-likelihood compared to the reconstructed single-ring fit momentum for atmospheric ν_e (left) and ν_μ (right) samples. The black line represents the cut used to discriminate electron-like and muon-like events, which coefficients obtained from Monte Carlo studies. Figures taken from [176].

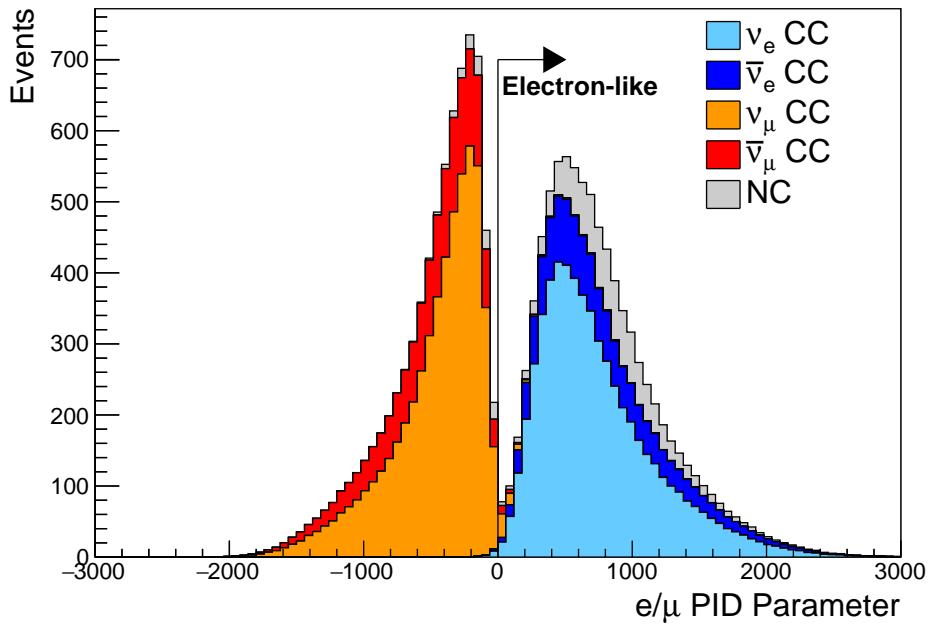


Figure 5.5: The electron/muon PID separation parameter for all sub-GeV single-ring events in SK-IV. The charged current (CC) component is broken down in four flavours of neutrino (ν_μ , $\bar{\nu}_\mu$, ν_e and $\bar{\nu}_e$). Events with positive values of the parameter are determined to be electron-like.

1531 The `fitQun` algorithm also considers a π^0 hypothesis. To do this, it performs
1532 a fit looking for two standard electron-hypothesis tracks which point to the
1533 same four-vertex. This assumes the electron tracks are generated from photon-
1534 conversion so the electron tracks actually appear offset from the proposed π^0
1535 vertex. For these fits, the conversion length, direction, and momentum of each
1536 photon are also considered as track parameters which are then fit in the same
1537 methodology as the standard single-ring hypotheses.

1538 Whilst lower energy events are predominantly single-ring events, higher
1539 energy neutrino events can generate final states with multiple particles which
1540 generate Cherenkov photons. These “multi-ring” hypotheses are also considered
1541 in the `fitQun` algorithm. When calculating the charge likelihood density, the
1542 predicted charge associated with each ring is calculated separately and then
1543 summed to calculate the total accumulated charge on each PMT. Similarly,
1544 the time likelihood for the multi-ring hypothesis is calculated assuming each
1545 ring is independent. Each track is time-ordered based on the time of flight
1546 from the center of the track to the PMT and the direct light from any ring
1547 incident on the PMT is assumed to arrive before any scattered light. To reduce
1548 computational resource usage, the multi-ring fits only consider electron-like
1549 and charged pion-like rings as the pion fit can be used as a proxy for a muon
1550 fit due to their similar mass.

1551 Multi-ring fits proceed by proposing another ring to the previous fit and
1552 then fitting the parameters in the method described above. Typically, multi-ring
1553 fits have the largest likelihood because of the additional degrees of freedom
1554 introduced. A likelihood value is calculated for the n -ring and $(n + 1)$ -ring
1555 hypotheses, where the additional ring is only included if the likelihood value
1556 is above 9.35, based on Monte Carlo studies in [180].

1557 5.2.1 Validation of Reconstruction in SK-V

1558 As an example of how the reconstruction depends on the detector conditions, the
1559 author of this thesis assessed the quality of event reconstruction for SK-V data.

1560 The detector systematics used in the T2K-only oscillation analysis are determined
1561 using data-to-Monte Carlo comparisons of the SK-IV data [181]. Due to tank-open
1562 maintenance occurring between SK-IV and SK-V, the dark rate of each PMT was
1563 observed to increase in SK-V due to light exposure for a significant time during
1564 the repairs. This increase can be seen in Figure 5.6. Run-10 of the T2K experiment
1565 was conducted in the SK-V period, so the consistency of SK-IV and SK-V data
1566 needs to be studied to determine whether the SK-IV-defined systematics can
1567 be applied to the run-10 data. This comparison study was performed using
1568 the stopping muon data set for both the SK-IV and SK-V periods. This data
1569 sample is used due to the high rate of interactions ($O(200)$ events per hour) as
1570 well as having similar energies to muons from CCQE ν_μ interactions from beam
1571 interactions. The rate of cosmic muons does depend on the solar activity cycle
1572 [182] but has been neglected in this comparison study. This is because the shape
1573 of the distributions is most important for the purposes of being compared to the
1574 detector systematics. The SK-IV and SK-V data samples consist of 2398.42 and
1575 626.719 hours of data which equates to 686k and 192k events respectively.

1576 The predicted charge calculated in the `fitQun` algorithm includes a contribu-
1577 tion from the photoelectron emission due to dark noise. Therefore, the increase
1578 in the SK-V dark rate needs to be accounted for. In practice, the average dark rate
1579 in each SK period is calculated and used as an input in the reconstruction. This is
1580 calculated by averaging the dark rate per run for each period separately, using
1581 the calibration measurements detailed in subsection 3.1.2. The average dark rate
1582 from SK-IV and SK-V were found to be 4.57kHz and 6.30kHz, respectively. The
1583 charges associated with the muon and decay electron subevents are illustrated in
1584 Figure 5.7. The photoelectron emission from dark noise is more significant for
1585 events that have lower energy. This is because this contribution becomes more
1586 comparable to the number of photoelectrons emitted from incident photons in
1587 lower-energy events. This behaviour is observed in the data, where the charge
1588 deposited by the muon subevent is mostly unaffected by the increase in dark rate,
1589 whilst the charge associated with the decay-electron is clearly affected.

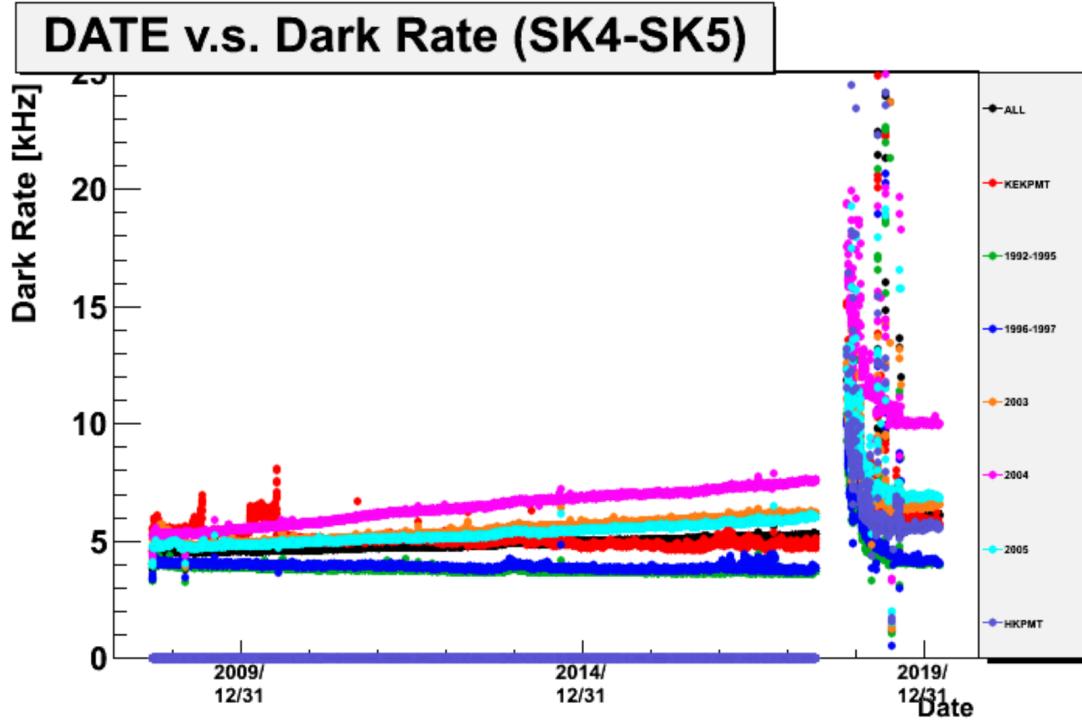


Figure 5.6: The variation of the measured dark rate as a function of date, broken down by PMT type. The SK-IV and SK-V periods span September 2008 to May 2018 and January 2019 to July 2020, respectively. The break in measurement in 2018 corresponds to the period of tank repair and refurbishment. Figure adapted from [181].

1590 The energy scale systematic is estimated from data-to-Monte Carlo differences
 1591 in the stopping muon sample in [183] and found to be 2.1%. To determine the
 1592 consistency of SK-IV and SK-V with respect to the energy scale systematic, the
 1593 muon momentum distribution is compared between the two SK periods. As
 1594 the total number of Cherenkov photons is integrated across the track length,
 1595 the reconstructed momentum divided by track length (or range) is compared
 1596 between SK-IV and SK-V as illustrated in Figure 5.8.

1597 The consistency between these distributions has been computed in two ways.
 1598 Firstly, a Gaussian is fit to the peak of each distribution separately, whose mean
 1599 is found to be $(2.272 \pm 0.003)\text{MeV/cm}$ and $(2.267 \pm 0.006)\text{MeV/cm}$ for SK-IV
 1600 and SK-V respectively. The ratio of these is equal to 1.002 ± 0.003 . The means of
 1601 the Gaussian fits are consistent with the expected stopping power of a minimum
 1602 ionising muon for a target material (water) with $Z/A \sim 0.5$ [184]. The second

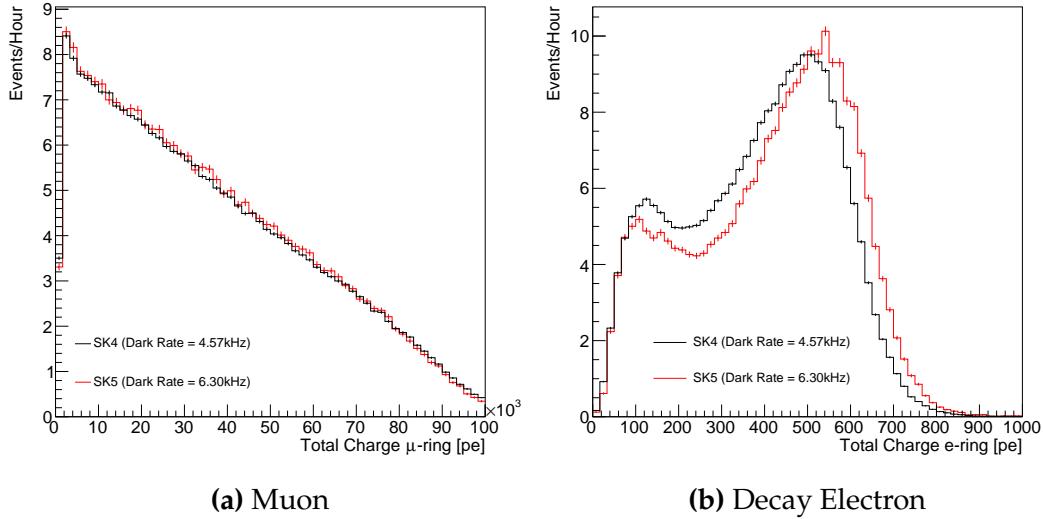


Figure 5.7: Comparison of the measured raw charge deposited per event from the stopping muon data samples between SK-IV (Blue) and SK-V (Red), split by the primary muon subevent and the associated decay electron subevent.

1603 consistency check is performed by introducing a nuisance parameter, α , which
 1604 modifies the SK-V distribution. The value of α which minimises the χ^2 value
 1605 between the SK-IV and SK-V is determined by scanning across a range of values.
 1606 This is repeated by applying the nuisance parameter as both a multiplicative
 1607 factor and an additive shift. The χ^2 distributions for different values of α is
 1608 illustrated in Figure 5.9. The values which minimise the χ^2 are found to be 0.0052
 1609 and 1.0024 for the additive and multiplicative implementations, respectively. No
 1610 evidence of shifts larger than the 2.1% uncertainty on the energy scale systematic
 1611 has been found in the reconstructed momentum distribution of SK-IV and SK-V.

1612 5.3 Event Reduction at SK

1613 Atmospheric neutrino events observed in the SK detector are categorised into
 1614 three different types of samples: fully contained (FC), partially contained (PC)
 1615 and up-going muon (Up- μ), using PMT hit signatures in the inner and outer
 1616 detector (ID and OD, respectively). To identify FC neutrino events, it is re-
 1617 quired that the neutrino interacts inside the fiducial volume of the ID and that
 1618 no significant OD activity is observed. For this analysis, an event is defined

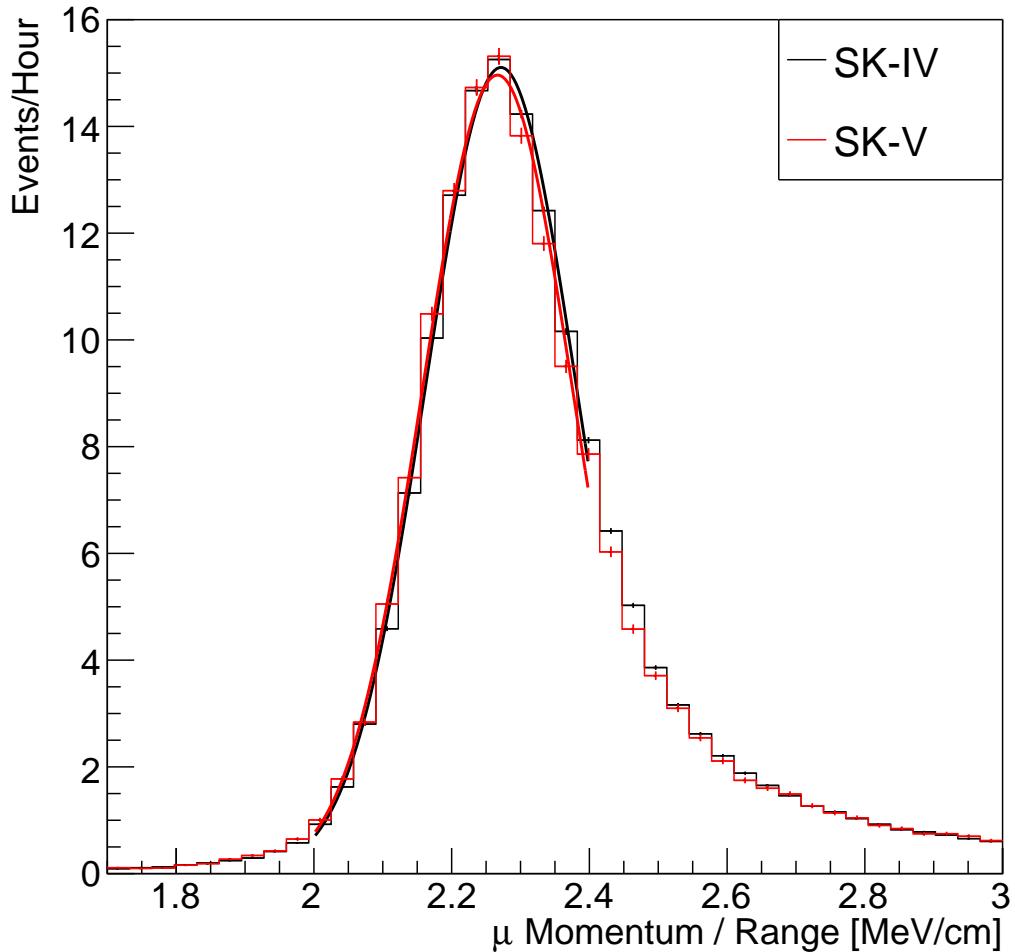


Figure 5.8: The distribution of the reconstructed momentum from the muon ring divided by the distance between the reconstructed muon and decay electron vertices as found in the stopping muon data sets of SK-IV (Black) and SK-V (Red). Only events with one tagged decay electron are considered. A Gaussian fit is considered in the range [2.0, 2.4] MeV/cm and illustrated as the solid curve.

¹⁶¹⁹ to be in the fiducial volume provided the event vertex is at least 0.5m away
¹⁶²⁰ from the ID walls. PC events have the same ID requirements but can have a
¹⁶²¹ larger signal present inside the OD. Typically, only high energy muons from
¹⁶²² ν_μ interactions can penetrate the ID wall. The Up- μ sample contains events where
¹⁶²³ muons are created from neutrino interactions in the OD water or rock below the
¹⁶²⁴ tank. They then propagate upwards through the detector. Downward-going
¹⁶²⁵ muons generated from neutrino interactions above the tank are neglected is

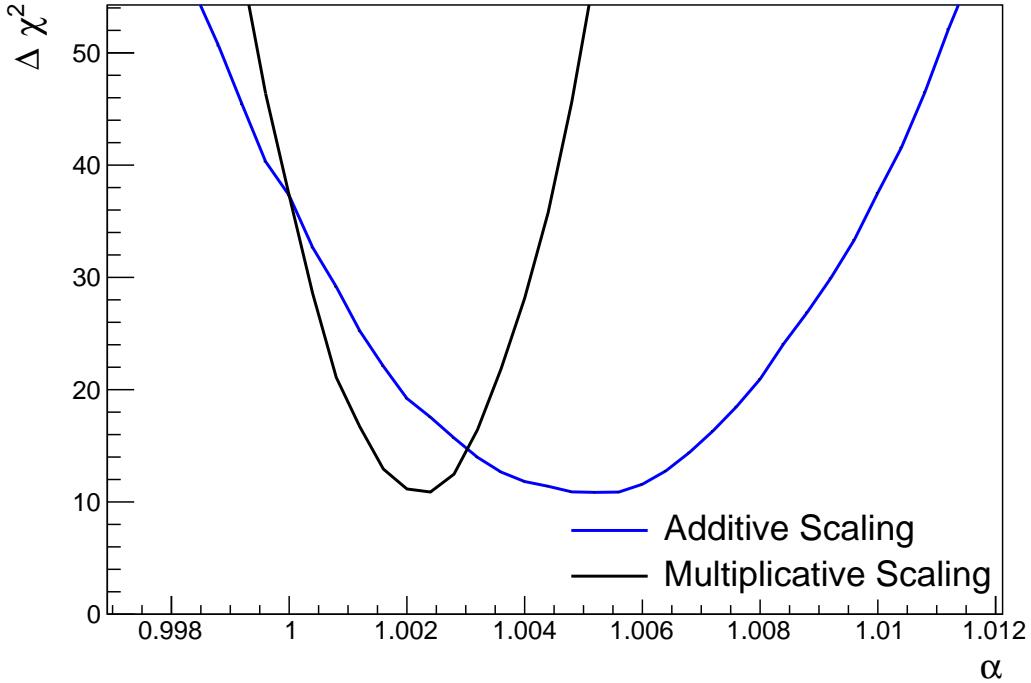


Figure 5.9: The χ^2 difference between the SK-IV and SK-V reconstructed muon momentum divided by range when the SK-V is modified by the scaling parameter α . Both additive (Blue) and multiplicative (Black) scaling factors have been considered. In practice, the additive scaling factor actually uses the value of $(\alpha - 1.0)$ but is illustrated like this so the results can be shown on the same axis range.

¹⁶²⁶ because of the difficulty in separating their signature from the cosmic muon
¹⁶²⁷ shower background. The sample categories are visually depicted in Figure 5.10.

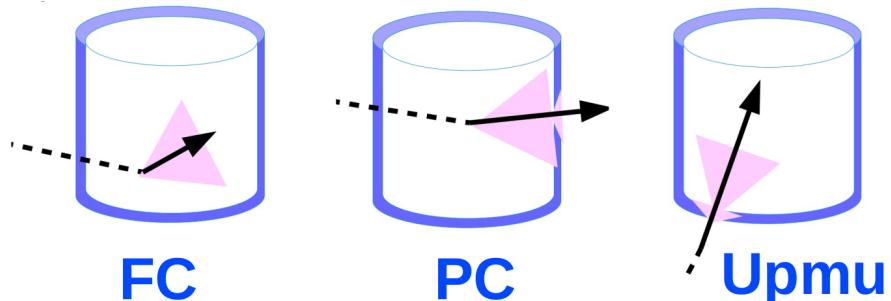


Figure 5.10: A depiction of the topology patterns for fully-contained (FC), partially-contained (PC) and up-going muon ($Up-\mu$) samples included in this analysis.

¹⁶²⁸ Based on the event characteristics, as defined by the `fitQun` event reconstruction software, the FC events are categorised by

- 1630 • **Visible Energy:** equal to the sum of the reconstructed kinetic energy of
1631 particles above the Cerenkov threshold for all rings present in the event.
1632 The purpose is to separate events into sub-GeV and multi-GeV categories.
- 1633 • **Number of observed Cerenkov rings.** The purpose is to separate single-
1634 ring and multi-ring events, where single-ring events predominantly consist
1635 of quasi-elastic interactions and multi-ring events are typically resonant
1636 pion production or deep inelastic scattering events.
- 1637 • **Particle identification parameter of the most energetic ring:** A value de-
1638 termined from the maximum likelihood value based on `fitQun`'s electron,
1639 muon, or pion hypothesis. The purpose is to separate electron-like and
1640 muon-like events.
- 1641 • **Number of decay electrons:** The purpose is to separate quasi-elastic events
1642 (which have one decay electron emitted from the muon decay) and resonant
1643 pion production events (which have two decay electrons emitted from the
1644 muon and pion).

1645 The PC and Up- μ categories are broken down into “through-going” and
1646 “stopping” samples depending on whether the muon leaves the detector. This is
1647 because the PC stopping events deposit the entire energy of the interaction into
1648 the detector, resulting in better reconstruction. The energy of events that exit the
1649 detector has to be estimated, with typically worse resolution, which introduces
1650 much larger systematic uncertainties. Through-going Up- μ samples are further
1651 broken down by whether any hadronic showering was observed in the event
1652 which typically indicates DIS interactions. The expected neutrino energy for the
1653 different categories is given in Figure 5.11. FC sub-GeV and multi-GeV events
1654 peak around 0.7GeV and 3GeV respectively, with slightly different peak energies
1655 for ν_e and $n\nu_\mu$ oscillation channels. PC and Up- μ are almost entirely comprised
1656 of ν_μ events and peak around 7GeV and 100GeV, respectively.

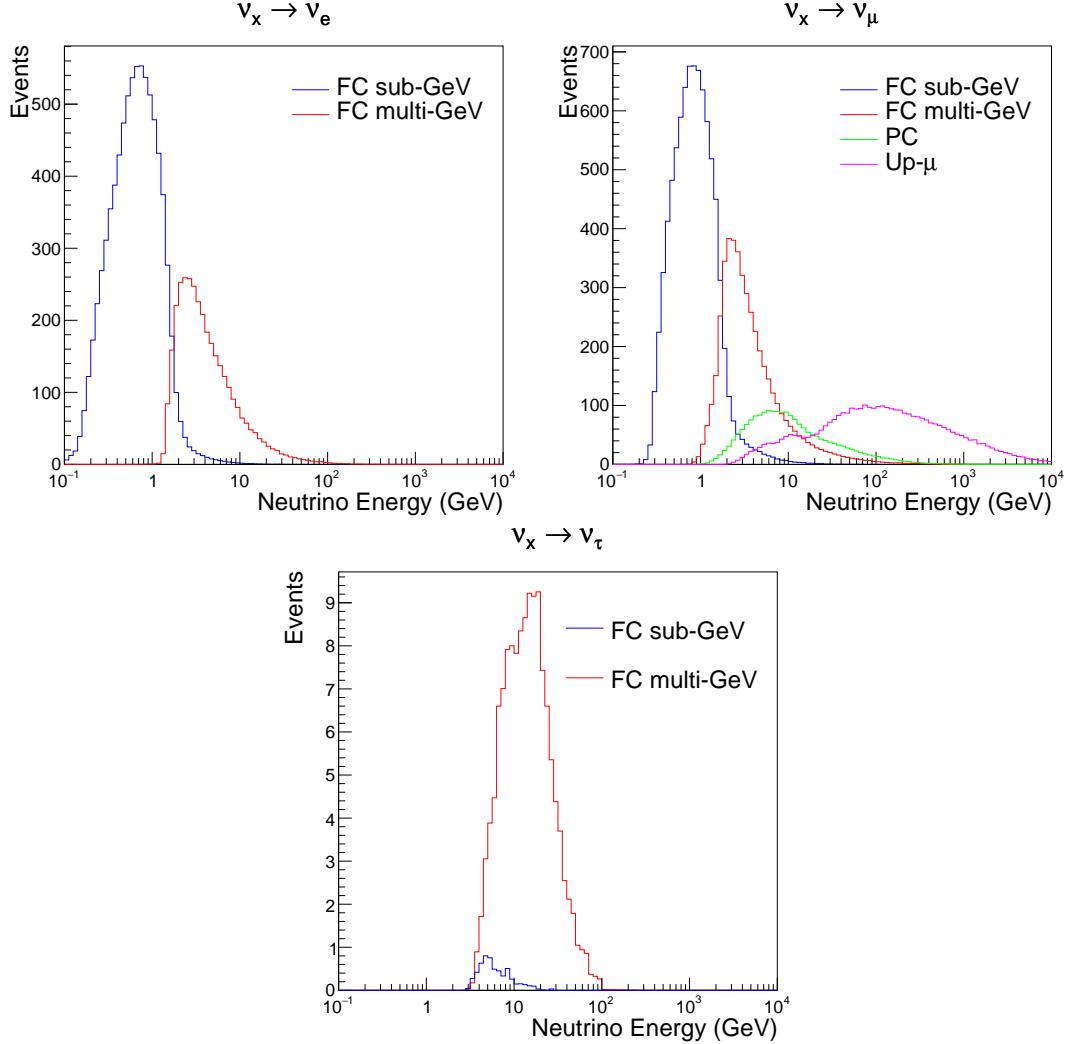


Figure 5.11: The predicted neutrino flux of the fully contained (FC) sub-GeV and multi-GeV, partially contained (PC), and upward-going muon (Up- μ) events. The prediction is broken down by the $\nu_x \rightarrow \nu_e$ prediction (top left), $\nu_x \rightarrow \nu_\mu$ prediction (top right) and $\nu_x \rightarrow \nu_\tau$ prediction (bottom). Asimov A oscillation parameters are assumed (given in Table 2.2).

In normal data-taking operations, the SK detector observes many background events alongside the beam and atmospheric neutrino signal events of physics interest for this thesis. Cosmic ray muons and flasher events, which are the spontaneous discharge of a given PMT, contribute the largest amount of background events in the energy range relevant to this thesis. Lower energy analyses like DSNB searches are also subject to radioactive backgrounds [185]. Therefore the data recorded is reduced with the aim of removing these background events. The reduction process is detailed in [54, 84] and briefly summarised below.

1665 The first two steps in the FC reconstruction remove the majority of cosmic
1666 ray muons by requiring a significant amount of ID activity compared to that
1667 measured in the OD. Events that pass this cut are typically very high momentum
1668 muons or events that leave very little activity in the OD. Consequently, a third
1669 reduction step is then applied to select cosmic-ray muons that pass the initial
1670 reduction step. A purpose-built cosmic muon fitter is used to determine the
1671 entrance (or exit) position of the muon and a cut is applied to OD activity
1672 contained within 8m of this position. Flasher events are removed in the fourth
1673 reduction step which is based on the close proximity of PMT hits surrounding
1674 the PMT producing the flash. Events that pass all these reduction steps are
1675 reconstructed with the APFit algorithm. The fifth step of the reduction uses
1676 information from the more precise fitter to repeat the previous two steps with
1677 tighter cuts. Muons below the Cherenkov threshold can not generate optical
1678 photons in the ID but the associated decay electron can due to its lower mass.
1679 These are the types of events targeted in the fifth reduction step. The final
1680 cuts require the event vertex to be within the fiducial volume (0.5m from the
1681 wall although the nominal distance is 2.0m), visible energy $E_{vis} > 30\text{MeV}$ and
1682 fewer than 16 hits within the higher energy OD cluster. The culmination of
1683 the fully contained reduction results in 8.09 events/day in the nominal fiducial
1684 volume [81]. The uncertainty in the reconstruction is calculated by comparing
1685 Monte Carlo prediction to data. The largest discrepancy is found to be 1.3%
1686 in the fourth reduction step.

1687 The PC and Up- μ events are processed through their own reduction processes
1688 detailed in [54]. Both of these samples are reconstructed with the APFit algorithm
1689 rather than `fiTQun`. This is because the efficiency of reconstructing events that
1690 leave the detector has not been sufficiently studied for reliable systematic uncer-
1691 tainties wth `fiTQun`. The PC and Up- μ samples acquire events at approximately
1692 0.66 and 1.44 events/day.

1693 Events due to beam neutrinos undergo the same reduction steps as FC events
1694 and are then subject to further cuts [186]. The GPS system which links the timing

1695 between the beam facility and SK needs to be operating correctly and there should
1696 be no activity within the detector in the previous $100\mu\text{s}$ before the trigger. The
1697 events then need to triggered between $-2\mu\text{s}$ and $10\mu\text{s}$ of the expected spill timing.

1698 The beam neutrino samples are not split by visible energy since their energy
1699 range is smaller than the atmospheric neutrino events. Following the T2K
1700 analysis in [74], only single-ring beam neutrino events are considered. Similar to
1701 atmospheric event selection, the number of decay electrons is used as a proxy for
1702 distinguishing CCQE and CCRES events. The expected neutrino energy, broken
1703 down by number of decay electrons, is given in Figure 5.12.

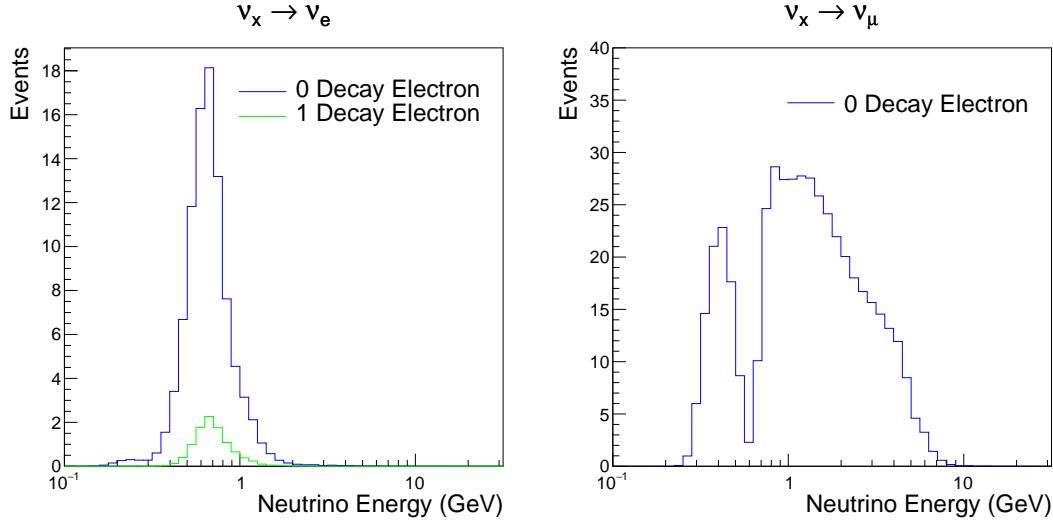


Figure 5.12: The predicted flux of beam neutrinos, as a function of neutrino energy. The predictions are broken down by the number of decay electrons associated with the particular events. Asimov A oscillation parameters are assumed (given in Table 2.2).

6

1704

1705

Sample Selections and Systematics

1706 The oscillation analysis presented within this thesis is built upon a simultaneous
1707 fit to atmospheric data at SK, neutrino beam data in the near detector, and
1708 beam data measured at SK. The definitions of these samples are documented
1709 in section 6.1, section 6.2, and section 6.3, respectively. The data collected and
1710 used within this analysis is detailed in Table 6.1. The near and far detector data
1711 corresponds to T2K runs 2-9 and runs 1-10, respectively. The accumulated POT
1712 and beam power for runs 1 – 10 are illustrated in Figure 6.1.

Data Type	Total
Near Detector FHC	1.15×10^{21} POT
Near Detector RHC	8.34×10^{20} POT
Far Detector FHC	1.97×10^{21} POT
Far Detector RHC	1.63×10^{21} POT
Atmospheric SK-IV	3244.4 days

Table 6.1: The amount of data collected in each detector used within this analysis. The data collected at the near and far detector, for both neutrino beam (FHC) and antineutrino beam (RHC), is measured as the number of protons on target (POT).

1713 The difference in POT recorded at the near and far detector is due to the
1714 difference in downtime. The SK detector is very stable with almost 100% of
1715 data recorded during beam operation. Due to various technical and operational

¹⁷¹⁶ issues, the downtime of the near detector is significantly higher due to its more
¹⁷¹⁷ complex design and operating requirements.

¹⁷¹⁸ The systematic parameters invoked within the flux, detector, and interaction
¹⁷¹⁹ models used within this analysis are documented in section 6.4. The standard
¹⁷²⁰ configuration of the joint beam and atmospheric data fit utilises far detector sys-
¹⁷²¹ tematics provided in the official inputs from the two experiments. Additionally,
¹⁷²² a correlated detector model which fits the parameters used in sample selections
¹⁷²³ to data has been developed and documented in subsection 6.4.5.

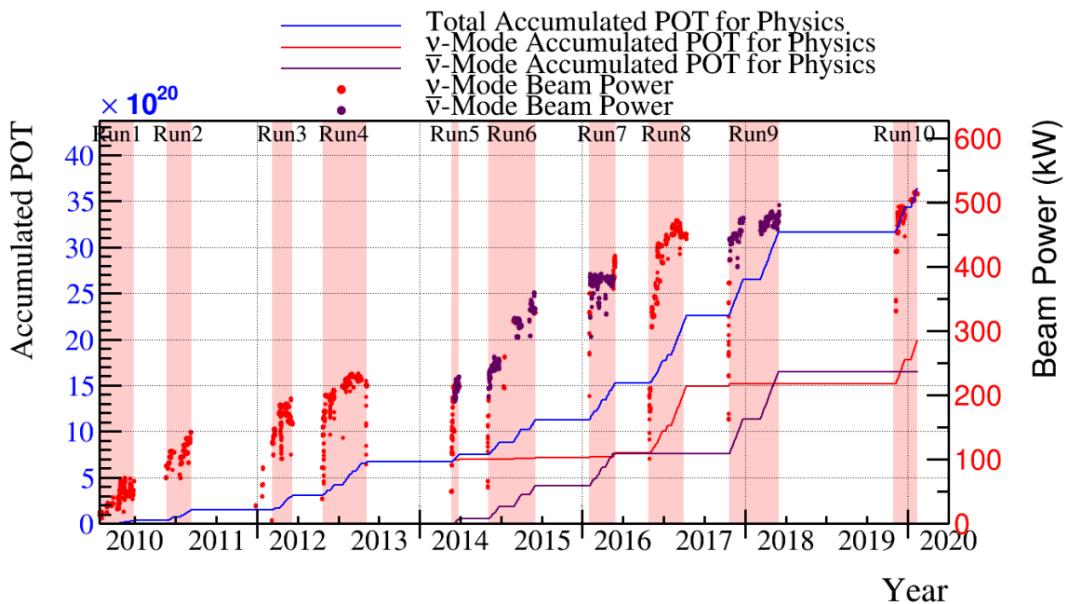


Figure 6.1: The accumulated beam data, measured as the number of protons on target (POT). The total data (blue) is given which comprises of the neutrino beam (red) and antineutrino (purple) components. The beam power for neutrino and antineutrino beams is given as the markers using the same colour scheme. The timescale runs from Run 1 which started in January 2010 until Run 10 which ended in February 2020. The ratio of accumulated data in neutrino and antineutrino beam is 54.7% : 45.3%.

1724 6.1 Atmospheric Samples

1725 The atmospheric event selection follows the official SK-IV analysis presented
1726 in [84] and is documented below. The Monte Carlo prediction used within this
1727 analysis corresponds to 500 years worth of neutrino events, which is scaled down
1728 to match the SK-IV livetime of 3244.4 days.

1729 The fully contained (FC), partially contained (PC), and upward going muon
1730 events ($\text{up-}\mu$) which pass the reduction cuts discussed in section 5.3 are further
1731 broken down into different samples based on reconstruction information. This
1732 section details the samples used within this oscillation analysis, alongside the
1733 chosen binning.

1734 FC events are first separated by the visible energy deposited within the
1735 detector. This is calculated as the sum of the reconstructed kinetic energy
1736 above the Cherenkov threshold for all rings present in the event. Events are
1737 separated by whether they were above or below $E_{\text{vis}} = 1.33\text{GeV}$. This separates
1738 “subGeV” and “multiGeV” events. Typically, lower energy events consist of
1739 charged current quasi-elastic (CCQE) interactions which are better understood
1740 and simpler to reconstruct resulting in smaller systematic uncertainties. Events
1741 are further separated by the number of rings associated with the event due to
1742 similar reasoning. As the oscillation probability is dependent upon the flavour
1743 of neutrino, electron and muon events are separated using a similar likelihood
1744 method to that discussed in section 5.2. To reduce computational resources
1745 required for the reconstruction, only electron and pion hypotheses are considered
1746 so this separation cut depends on the ratio of the electron to pion likelihoods,
1747 $\log(L_e/L_\pi)$. Finally, the number of decay electrons is used to classify events.
1748 Charged current resonant pion production (CCRES) interactions generate a final-
1749 state pion. This can decay, mostly likely through a muon, into a decay electron.
1750 Therefore any electron-like event with one decay electron or muon-like event
1751 with two decay electrons was most likely produced by a CCRES interaction.
1752 Consequently, the number of decay electrons can be used to distinguish CCQE

¹⁷⁵³ and CCRES interaction modes. Ultimately, FC subGeV events are separated
¹⁷⁵⁴ into the samples listed in Table 6.2.

Sample Name	Description
SubGeV- <i>e</i> like-0dcy	Single ring <i>e</i> -like events with zero decay electrons
SubGeV- <i>e</i> like-1dcy	Single ring <i>e</i> -like events with one or more decay electrons
SubGeV- <i>μ</i> like-0dcy	Single ring <i>μ</i> -like events with zero decay electrons
SubGeV- <i>μ</i> like-1dcy	Single ring <i>μ</i> -like events with one decay electrons
SubGeV- <i>μ</i> like-2dcy	Single ring <i>μ</i> -like events with two or more decay electrons
SubGeV- <i>π</i> 0like	Two <i>e</i> -like ring events with zero decay electrons and reconstructed π^0 mass $85 \leq m_{\pi^0} < 215$ MeV

Table 6.2: The fully contained subGeV samples, defined as events with visible energy $E_{vis} < 1.33$ GeV, used within this oscillation analysis.

¹⁷⁵⁵ In addition to the cuts discussed above, multiGeV samples also have addi-
¹⁷⁵⁶ tional cuts to separate samples which target neutrino and antineutrino events.
¹⁷⁵⁷ As discussed in section 7.1, the matter resonance only occurs for neutrinos in
¹⁷⁵⁸ normal hierarchy and antineutrinos in an inverted mass hierarchy. Therefore,
¹⁷⁵⁹ having flavour-enriched samples aids in the determination of the mass hierarchy.
¹⁷⁶⁰ For a CCRES interaction,

$$\begin{aligned}
 \bar{\nu}_e + N &\rightarrow e^+ + N' + \pi^-, \\
 \nu_e + N &\rightarrow e^- + N' + \pi^+ \\
 &\quad \downarrow \mu^+ + \nu_\mu \\
 &\quad \downarrow e^+ + \nu_e + \bar{\nu}_\mu.
 \end{aligned} \tag{6.1}$$

¹⁷⁶¹ The π^- emitted from a $\bar{\nu}_e$ interaction is more likely to be absorbed within the
¹⁷⁶² oxygen nucleus than the π^+ from ν_e interactions [187]. These pions then decay,
¹⁷⁶³ mostly through muons, to electrons. Therefore the number of tagged decay
¹⁷⁶⁴ electrons associated with an event gives an indication of whether the interaction
¹⁷⁶⁵ was due to a neutrino or antineutrino: zero for $\bar{\nu}_e$ events, and one for ν_e events.
¹⁷⁶⁶ The ability to separate neutrino from antineutrino events is illustrated in Table 6.4,
¹⁷⁶⁷ where the MultiGeV-*e*like-nue has 78% purity of CC neutrino interactions with
¹⁷⁶⁸ only 7% antineutrino background, the rest consisting of NC backgrounds.

1769 The number of decay electrons discriminator works reasonably well for single-
1770 ring events. However, this is not the case for multi-ring events. A multiGeV
1771 multiring electron-like (MME) likelihood cut was introduced in [188, 189]. This
1772 is a two-stage likelihood selection cut. Four observables are used in the first
1773 likelihood cut to distinguish $CC\nu_e$ and $CC\bar{\nu}_e$ events from background:

- 1774 • The number of decay electrons
- 1775 • The maximum distance between the vertex of the neutrino and the decay
1776 electrons
- 1777 • The energy deposited by the highest energy ring
- 1778 • The particle identification of that highest energy ring

1779 Background events consist of $CC\nu_\mu$ and NC interactions. Typically, the
1780 majority of the energy in these background events is carried by the hadronic
1781 system. Additionally, muons tend to travel further than the pions from $CC\nu_e$
1782 before decaying. Thus, the parameters used within the likelihood cut target these
1783 typical background interaction kinematics.

Sample Name	Description
MultiGeV-elike-nue	Single ring e -like events with zero decay electrons
MultiGeV-elike-nuebar	Single ring e -like events with one or more decay electrons
MultiGeV-mulike	Single ring μ -like events
MultiRing-elike-nue	Two or more ring events with leading energy e -like ring and passed both MME and $\nu/\bar{\nu}$ separation cuts
MultiRing-elike-nuebar	Two or more ring events with leading energy e -like ring and passed MME and failed $\nu/\bar{\nu}$ separation cuts
MultiRing-mulike	Two or more ring events with leading energy μ -like ring and only requires $E_{vis} > 0.6\text{GeV}$ DB: Why is this not }1.33\text{GeV
MultiRing-Other1	Two or more ring events with leading energy e -like ring and failed the MME likelihood cut

Table 6.3: The fully contained multiGeV samples used within this oscillation analysis. Both the sample name and description are given.

1784 Neutrino and antineutrino events are then separated by a second likelihood
1785 method ($\nu/\bar{\nu}$ separation) detailed in [59]. This uses the number of decay electrons,

1786 the number of reconstructed rings, and the event’s transverse momentum. The
1787 last two parameters are used because higher-energy samples tend to have more
1788 pions produced above the Cherenkov threshold which results in more rings
1789 compared to an antineutrino interaction. Furthermore, the angular distribution
1790 also tends to be more forward peaked in antineutrino interactions as compared
1791 to neutrino interactions [84]. These FC multiGeV sample definitions are de-
1792 tailed in Table 6.3.

1793 The PC and up- μ samples are split by the amount of energy deposited within
1794 the outer detector, into “stopping” and “through-going” samples. If an event
1795 leaves the detector, the energy it takes with it has to be estimated which increases
1796 the systematic uncertainty compared to events entirely contained within the
1797 inner detector. This estimation is particularly poor at high energies, thus the
1798 up- μ through-going events are not binned in reconstructed momentum. The
1799 through-going up- μ are further separated by the presence of any electromagnetic
1800 showering in the event, as the assumption of non-showering muon does not give
1801 reliable reconstruction for these types of events [54]. In total, 13 FC, 2 PC, and
1802 3 up- μ atmospheric samples are included within this analysis.

1803 The atmospheric samples are binned in direct observables: reconstructed
1804 lepton momentum and direction, as given by Table A.1. The distribution of
1805 the reconstructed lepton momentum (for samples that only have one bin in
1806 reconstructed zenith angle) and reconstructed direction for each atmospheric
1807 sample used within this analysis is illustrated in Figure 6.2. The by-mode
1808 breakdown of each of the atmospheric samples is given in Appendix A.

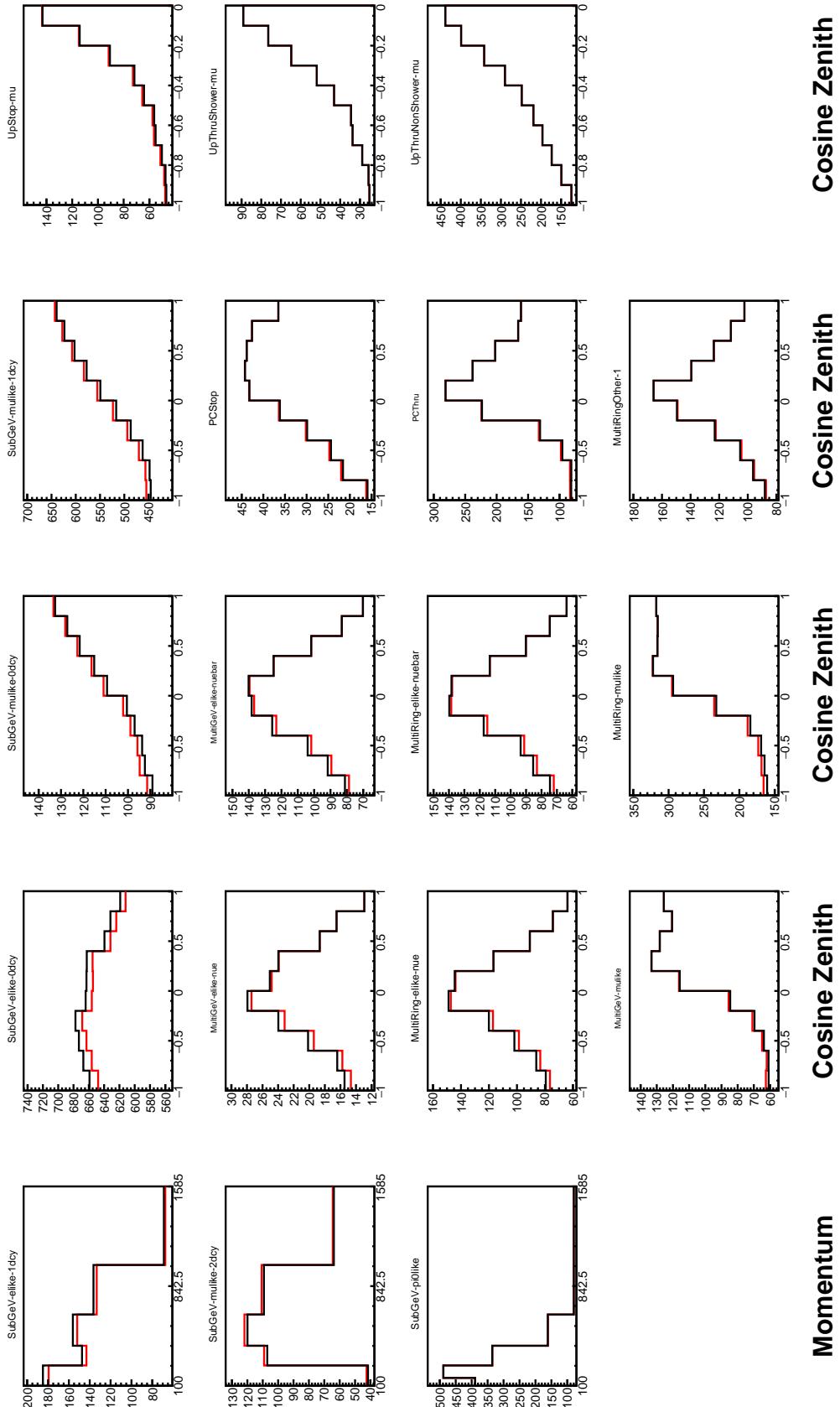


Figure 6.2: Comparison of the SK-IV atmospheric samples between predictions made with the CP-violating Asimov A (Black) and CP-conserving Asimov B (Red) oscillation parameter sets (given in Table 2.2). The subGeV samples CCRES and π^0 -like samples are given in their reconstructed lepton momentum. All other samples are presented in their reconstructed zenith angle projection.

Sample	$CC\nu_e$	$CC\bar{\nu}_e$	$CC(\nu_\mu + \bar{\nu}_\mu)$	$CC(\nu_\tau + \bar{\nu}_\tau)$	NC
SubGeV- <i>elike</i> -0dcy	72.17	23.3	0.724	0.033	3.77
SubGeV- <i>elike</i> -1dcy	86.81	1.773	7.002	0.062	4.351
SubGeV- <i>mulike</i> -0dcy	1.003	0.380	90.07	0.036	8.511
SubGeV- <i>mulike</i> -1dcy	0.023	0.	98.46	0.029	1.484
SubGeV- <i>mulike</i> -2dcy	0.012	0.	99.25	0.030	0.711
SubGeV- <i>pi0like</i>	6.923	2.368	0.928	0.011	89.77
MultiGeV- <i>elike</i> -nue	78.18	7.041	3.439	1.886	9.451
MultiGeV- <i>elike</i> -nuebar	56.68	37.81	0.174	0.614	4.718
MultiGeV- <i>mulike</i>	0.024	0.005	99.67	0.245	0.058
MultiRing- <i>elike</i> -nue	59.32	12.39	4.906	3.385	20
MultiRing- <i>elike</i> -nuebar	52.39	31.03	1.854	1.585	13.14
MultiRing- <i>mulike</i>	0.673	0.080	97.33	0.342	1.578
MultiRingOther-1	27.98	2.366	34.93	4.946	29.78
PCStop	8.216	3.118	84.45	0.	4.214
PCThrus	0.564	0.207	98.65	0.	0.576
UpStop-mu	0.829	0.370	98.51	0.	0.289
UpThruNonShower-mu	0.206	0.073	99.62	0.	0.103
UpThruShower-mu	0.128	0.054	99.69	0.	0.132

Table 6.4: The purity of each atmospheric sample used within this analysis, broken down by charged current (CC) and neutral current (NC) interactions and which neutrino flavour interacted within the detector. Asimov A oscillation parameter sets are assumed (given in Table 2.2). Electron neutrino and antineutrino events are separated to illustrate the ability of the separation likelihood cuts used within the multiGeV and multiring sample selections.

1809 6.2 Near Detector Beam Samples

1810 The near detector sample selections are documented in detail within [190] and
1811 summarised below. Samples are selected based upon the which Fine Grained
1812 Detector (FGD) that the vertex is reconstructed in as well as the operating mode
1813 of the beam: FHC or RHC. Wrong-sign neutrino samples are considered in the
1814 RHC mode in order to add additional constraints on model parameters. Samples
1815 from the wrong-sign component of the FHC beam mode are not included as they
1816 are statistically insignificant compared to those samples already listed.

1817 For additional constraints on model parameters, wrong-sign neutrino samples
1818 are also considered when the beam is operating in RHC mode.

1819 The reconstruction algorithm uses a clustering algorithm to group hits within
1820 the TPC. It then adds information from the upstream FGD to form a track
1821 which passes through both sub-detectors. In FHC(RHC), the highest momentum
1822 negative(positive) curvature track is defined as the muon candidate. Before
1823 being assigned a sample, these candidate muon events must pass CC-inclusive
1824 cuts, as defined in [191]:

- 1825 • Event Timing: The DAQ must be operational and the event must occur
1826 within the expected beam time window consistent with the beam spill
- 1827 • TPC Requirement: The muon-candidate track path must intercept one or
1828 more TPCs
- 1829 • Fiducial volume: The event must originate from within the fiducial volume
1830 defined in [192].
- 1831 • Upstream Background: Remove events that have muon tracks that originate
1832 upstream of the FGDs by requiring no high-momentum tracks within
1833 150mm upstream of the candidate vertex. Additionally, events that occur
1834 within the downstream FGD are vetoed if a secondary track starts within
1835 the upstream FGD

- 1836 • Broken track removal: All candidates where the muon candidate is broken
1837 in two are removed

- 1838 • Muon PID: Measurements of dE/dx in a TPC are used to distinguish muon-
1839 like events, from electron-like or proton-like, using a likelihood cut

1840 In addition to these cuts, RHC neutrino events also have to undergo the
1841 following cuts to aid in the separation of neutrino and antineutrino [193]:

- 1842 • TPC Requirement: The track path must intercept TPC2

- 1843 • Positive Track: The highest momentum track must have positive recon-
1844 structed charge

- 1845 • TPC1 Veto: Remove any events originating upstream of TPC1

1846 Once all CC-inclusive events have been determined, they are further split by
1847 pion multiplicity: CC0 π , CC1 π , and CCOther. These target specific interaction
1848 modes CCQE, CCRES, and other CC background interactions, respectively. Pions
1849 in the TPCs and FGDs are selected by requiring a second track to be observed,
1850 which is separate from the muon track and is in the same beam spill window and
1851 sub-detector. If the pion originated within a FGD, it must also pass through the
1852 sequential downstream TPC (TPC2 for FGD1, TPC3 for FGD2).

1853 CC0 π , CC1 π , and CCOther samples are defined with the following cuts:

1854 **DB: Understand pion cuts at ND**

- 1855 • ν_μ CC0 π **Selection**: No electrons in TPC and no charged pions or decay
1856 electrons within the TPC or FGD

- 1857 • ν_μ CC1 π **Selection**: Exactly one charged pion in either the TPC or FGD,
1858 where the number of charged pions in the FGD is equal to the number of
1859 decay electrons

- 1860 • ν_μ CCOther **Selection**: All events which are not classified into the above
1861 two selections.

Counting the three selections for each FGD in FHC and RHC running, including the wrong-sign background in RHC, 18 near detector samples are used within this analysis. These samples are binned in reconstructed lepton momentum (illustrated in Figure 6.3) and direction with respect to the beam. The binning is chosen such that each event has at least 20 Monte Carlo events in each bin [192]. This is to ensure that the bins are coarse enough to ensure the reduction of statistical errors, whilst also being fine enough to sample the high-resolution peak regions. The exact binning is detailed in [192].

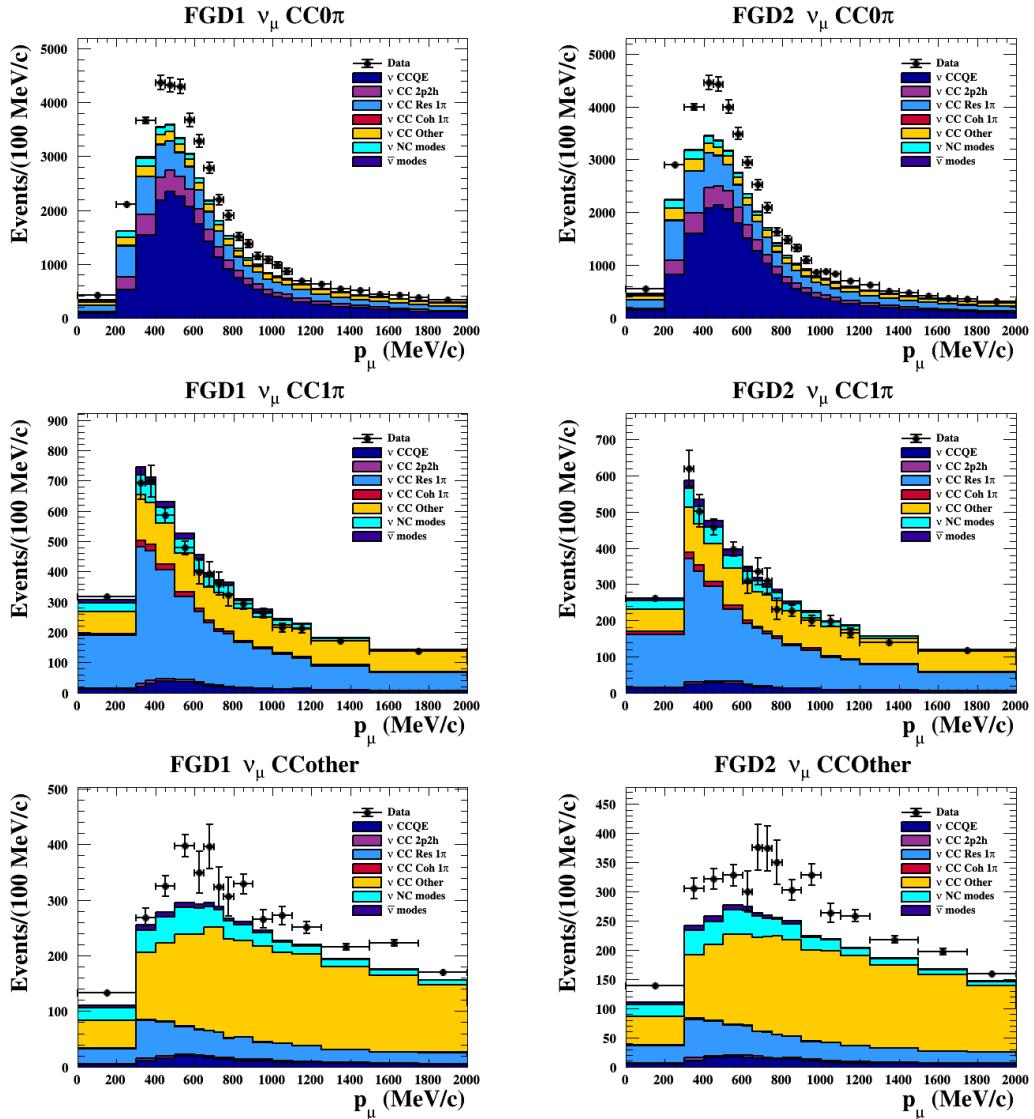


Figure 6.3: The nominal Monte Carlo predictions compared to data for the FGD1 and FGD2 samples in neutrino beam mode, broken down into the $CC\nu_\mu 0\pi$, $CC\nu_\mu 1\pi$ and $CC\nu_\mu$ Other categories. Figures taken from [190].

1870 6.3 Far Detector Beam Samples

1871 The beam neutrino events which occur at the SK detector, which pass the
 1872 reduction cuts detailed in section 5.3, are separated based on whether the beam
 1873 was operating in FHC or RHC mode. The events are then separated into three
 1874 samples: electron-like ($1Re$), muon-like ($1R\mu$), and $CC1\pi^+$ -like ($1Re1de$) which
 1875 are observed as electron-like events with an associated decay electron [181].
 1876 As discussed in section 6.1, positively charged pions emitted from neutrino
 1877 interactions are more likely to produce decay electrons than negatively charged
 1878 pions. Consequently, the $CC1\pi^+$ -like sample is only selected when the beam is
 1879 operating in FHC mode. Therefore, five beam samples measured at SK are
 1880 used in this analysis.

1881 The fiducial volume definition for beam samples is slightly different from that
 1882 used for the atmospheric samples. It uses both the distance to the closest wall
 1883 (`dWall`) and the distance to the wall along the trajectory of the particle (`toWall`).
 1884 This allows events that originate close to the wall but are facing into the tank to be
 1885 included within the analysis, which would have otherwise been removed. These
 1886 additional events are beneficial for a statistics-limited experiment. The exact
 1887 cut values for both `dWall` and `toWall` are different for each of the three types of
 1888 sample and are optimised based on T2K sensitivity to δ_{CP} [179, 194]. They are:

1889 **1Re event selection** For an event to be classified as a $1Re$ -like, the event must sat-
 1890 isfy:

- 1891 • Fully-contained and have $dWall > 80\text{cm}$ and $toWall > 170\text{cm}$
- 1892 • Total of one ring which is reconstructed as electron-like with reconstructed
 1893 momentum $P_e > 100\text{MeV}$
- 1894 • Zero decay electrons are associated with the event
- 1895 • Passes π^0 rejection cut discussed in section 5.2

1896 The zero decay electron cut removes non-CCQE interactions and the π^0
 1897 rejection cut is designed to remove neutral current π^0 background events which
 1898 can be easily reconstructed as 1Re-like events.

1899 The zero decay electron cut removes non-CCQE interactions and the π^0
 1900 rejection cut is designed to remove neutral current π^0 background events which
 1901 can be easily reconstructed as 1Re-like events.

1902 **CC1 π^+ event selection** This event selection is very similar to that of the 1Re
 1903 sample. The only differences are that the `dWall` and `toWall` criteria are changed
 1904 to $> 50\text{cm}$ and $> 270\text{cm}$, respectively, and exactly one decay electron is required
 1905 from the π^+ decay.

1906 **1R μ event selection** A 1R μ -like event is determined by the following cuts:

- 1907 • Fully-contained and have `dWall` $> 50\text{cm}$ and `toWall` $> 250\text{cm}$
- 1908 • Total of one ring which is reconstructed as muon-like with reconstructed
 momentum $P_\mu > 200\text{MeV}$
- 1910 • Fewer than two decay electrons are associated with the event
- 1911 • Passes π^+ rejection cut discussed in section 5.2

1912 All of these samples are binned in reconstructed neutrino energy. This is
 1913 possible under a particular interaction mode assumption, as the direction from
 1914 the source is known extremely well. For the 1Re-like and 1R μ -like samples,

$$E_\nu^{rec} = \frac{(M_N - V_{nuc})E_l - m_l^2/2 + M_N V_{nuc} - V_{nuc}^2/2 + (M_P^2 + M_N^2)/2}{M_N - V_{nuc} - E_l + P_l \cos(\theta_{beam})} \quad (6.2)$$

1915 Where M_N , M_P and m_l are the masses of the neutron, proton and outgoing
 1916 lepton, respectively. $V_{nuc} = 27\text{MeV}$ is the binding energy of the oxygen nucleus
 1917 [181], θ_{beam} is the angle between the beam and the direction of the outgoing
 1918 lepton, and E_l and P_l are the energy and momentum of that outgoing lepton.

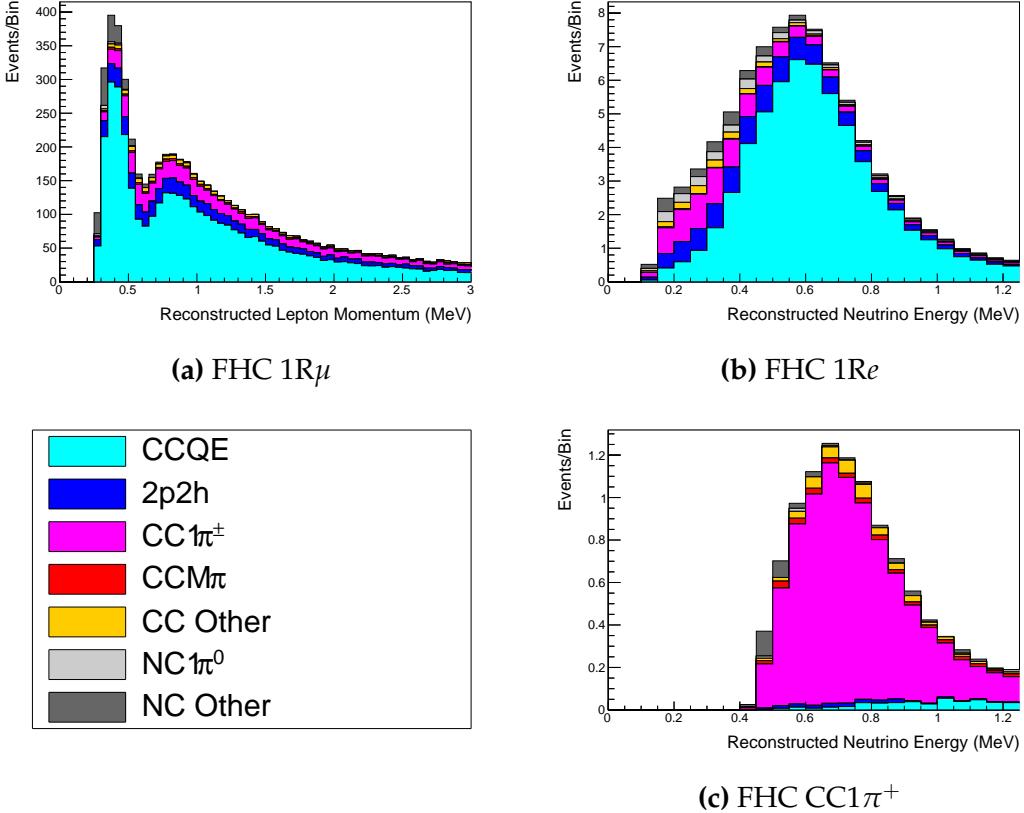


Figure 6.4: The reconstructed neutrino energy, as defined by Equation 6.2 and Equation 6.3, for the 1R μ -like, 1Re-like and CC1 π^+ -like samples. Asimov A oscillation parameter sets are assumed (given in Table 2.2). These samples are the FHC mode samples. For ease of viewing, the 1R μ sample only shows the $0 \leq E_\nu^{rec} < 3.0\text{GeV}$ but the binning extends to 30.0GeV.

1919 The reconstructed neutrino energy of the CC1 π^+ -like events also accounts
 1920 for the delta resonance produced within the interaction,

$$E_\nu^{rec} = \frac{2M_N E_l + M_{\Delta^{++}}^2 - M_N^2 - m_l^2}{2(M_N - E_l + P_l \cos(\theta_{beam}))} \quad (6.3)$$

1921 Where $M_{\Delta^{++}}$ is the mass of the delta baryon. Binding energy effects are not
 1922 considered as a two-body process with the delta baryon is assumed. This follows
 1923 the T2K oscillation analysis presented in [74], although recent developments
 1924 of the interaction model in the latest T2K oscillation analysis do include effects
 1925 from binding energy in this calculation [195].

1926 The reconstructed neutrino energy for the FHC samples is illustrated in
 1927 Figure 6.4. As expected, the 1R μ -like and 1Re-like samples are heavily dominated

by CCQE interactions, with smaller contributions from 2p2h meson exchange and resonant pion production interactions. The CC $1\pi^+$ -like sample predominantly consists of charged current resonant pion production interactions. The 1Re-like and CC $1\pi^+$ -like samples are also binned by the angle between the neutrino beam and the reconstructed lepton momentum. This is to aid in charged current and neutral current separation, as indicated in Figure 6.5. This is because the neutral current backgrounds are predominantly due to π^0 -decays, where the opening angle of the two gammas alongside the different final state kinematics produces a slightly broader angular distribution compared to the final state particles originating from charged current ν_e interactions.

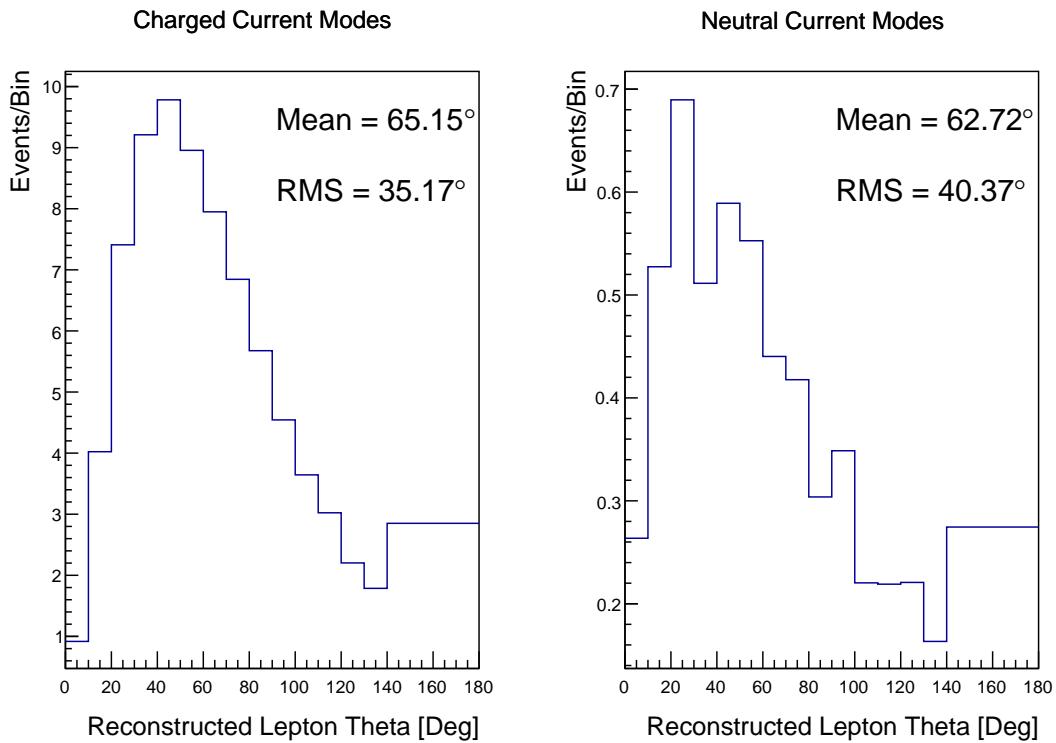


Figure 6.5: The distribution of the angle between the neutrino beam direction and the reconstructed final state lepton, for the FHC 1Re-like sample. The distribution is broken down by neutrino interaction mode into charged current (left) and neutral current (right) components. Asimov A oscillation parameter sets are assumed (given in Table 2.2). The RMS of the charged and neutral current plots are 35.17° and 40.37° , respectively.

1938 6.4 Systematic Uncertainties

1939 The systematic model parameters for this analysis are split into groups, or blocks,
1940 depending on their purpose. They consist of flux uncertainties, neutrino-matter
1941 interaction systematics, and detector efficiencies. There are also uncertainties on
1942 the oscillation parameters which this analysis is not sensitive to, namely Δm_{12}^2
1943 and $\sin^2(\theta_{12})$. These oscillation parameters uncertainties are taken from the 2020
1944 PDG measurements [75]. As described in chapter 4, each model parameter used
1945 within this analysis requires a prior uncertainty. This is provided via separate
1946 covariance matrices for each block. The covariance matrices can include prior
1947 correlations between parameters within a single block, but the separate treatment
1948 means prior correlations can not be included for parameters in different groups.
1949 Some parameters in these models have no reasonably motivated uncertainties
1950 and are assigned flat priors which do not modify the likelihood penalty. In
1951 practice, these flat prior parameters are actually assigned a Gaussian with a
1952 very large width to ensure the covariance matrix is positive definite. They are
1953 then checked at run time to determine if they contribute to the likelihood. The
1954 flux, neutrino interaction, and detector modeling simulations have already been
1955 discussed in section 5.1 and section 5.2. The uncertainties invoked within each
1956 of these models are described below.

1957 6.4.1 Beam Flux

1958 The neutrino beam flux systematics are based upon the uncertainty in the mod-
1959eling of the components of the beam simulation. This includes the model of
1960 hadron productions and reinteractions, the shape, intensity, and alignment of
1961 the beam with respect to the target, and the uniformity of the magnetic field
1962 produced by the horn, alongside other effects. The uncertainty, as a function
1963 of neutrino energy, is illustrated in Figure 6.6 which includes a depiction of
1964 the total uncertainty as well as the contribution from individual components.
1965 The uncertainty around the peak of the energy distribution ($E_\nu \sim 0.6\text{GeV}$) is

1966 dominated by uncertainties in the beam profile and alignment. Outside of this
1967 region, uncertainties on hadron production dominate the error.

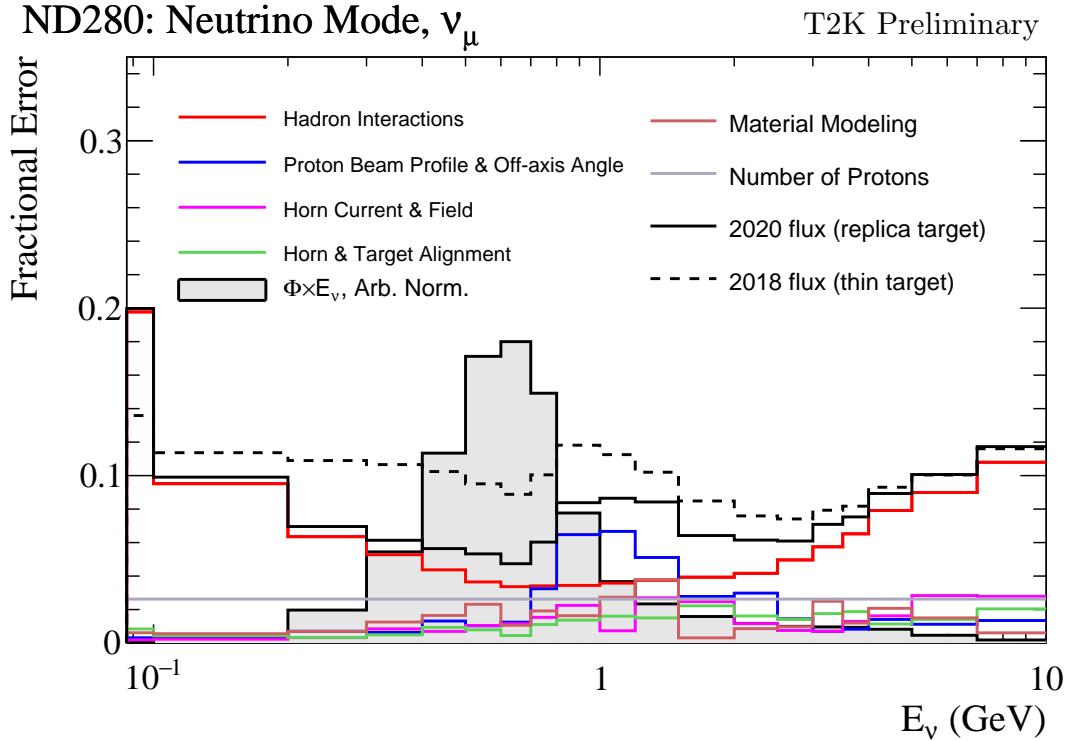


Figure 6.6: The total uncertainty evaluated on the near detector ν_μ flux prediction constrained by the replica-target data, illustrated as a function of neutrino energy. The solid(dashed) line indicates the uncertainty used within this analysis(the T2K 2018 analysis [196]). The solid histogram indicates the neutrino flux as a function of energy. Figure taken from [197].

1968 The beam flux uncertainties are described by one hundred parameters. They
1969 are split between the ND280 and SK detectors and binned by neutrino flavour:
1970 ν_μ , $\bar{\nu}_\mu$, ν_e and $\bar{\nu}_e$. The response is then broken down as a function of neutrino
1971 energy. The bin density in the neutrino energy is the same for the ν_μ in FHC
1972 and $\bar{\nu}_\mu$ in RHC beams, and narrows for neutrino energies close to the oscillation
1973 maximum of $E_\nu = 0.6\text{GeV}$. This binning is specified in Table 6.5. All of these
1974 systematic uncertainties are applied as normalisation parameters with Gaussian
1975 priors centered at 1.0 and error specified from a covariance matrix provided
1976 by the T2K beam group [197].

Neutrino Flavour	Sign	Neutrino Energy Bin Edges (GeV)
μ	Right	0., 0.4, 0.5, 0.6, 0.7, 1., 1.5, 2.5, 3.5, 5., 7., 30.
μ	Wrong	0., 0.7, 1., 1.5, 2.5, 30.
e	Right	0., 0.5, 0.7, 0.8, 1.5, 2.5, 4., 30.
e	Wrong	0., 2.5, 30.

Table 6.5: The neutrino energy binning for the different neutrino flavours. “Right” sign indicates neutrinos in the FHC beam and antineutrinos in the RHC beam. “Wrong” sign indicates antineutrinos in the FHC beam and neutrinos in the RHC beam. The binning of the detector response is identical for the FHC and RHC modes as well as at ND280 and SK.

1977 6.4.2 Atmospheric Flux

1978 The atmospheric neutrino flux is modeled by the HKKM model [50]. 16 systematic
 1979 uncertainties are applied to control the normalisation of each neutrino flavour,
 1980 energy, and direction. They are summarised below:

- 1981 • **Absolute Normalisation:** The overall normalisation of each neutrino flavour
 1982 is controlled by two independent systematic uncertainties, for $E_\nu < 1\text{GeV}$
 1983 and $E_\nu > 1\text{GeV}$, respectively. This is driven mostly by hadronic interaction
 1984 uncertainties for the production of pions and kaons [50]. The strength of
 1985 the response is dependent upon the neutrino energy. The uncertainty is
 1986 parameterized following Figure 11 in [50].
- 1987 • **Relative Normalisation:** Uncertainties on the ratio of $(\nu_\mu + \bar{\nu}_\mu) / (\nu_e + \bar{\nu}_e)$
 1988 are controlled by the difference between the HKKM model [50], FLUKA
 1989 [53] and Bartol models [49]. Three independent parameters are applied in
 1990 the energy ranges: $E_\nu < 1\text{GeV}$, $1\text{GeV} < E_\nu < 10\text{GeV}$, and $E_\nu > 10\text{GeV}$.
- 1991 • **$\nu/\bar{\nu}$ Normalisation:** The uncertainties in the π^+/π^- (and kaon equivalent)
 1992 production uncertainties in the flux of $\nu/\bar{\nu}$. The response is applied using
 1993 the same methodology as the relative normalisation parameters.
- 1994 • **Up/Down and Vertical/Horizontal Ratio:** Similar to the above two sys-
 1995 tematics, the difference between the HKKM, FLUKA, and Bartol model

1996 predictions, as a function of $\cos(\theta_Z)$, is used to control the normalisation of
 1997 events as a function of zenith angle.

- 1998 • **K/π Ratio:** Higher energy neutrinos ($E_\nu > 10\text{GeV}$) mostly originate in
 1999 kaon decay. Measurements of the ratio of K/π production [198] are used to
 2000 control the systematic uncertainty of the expected ratio of pion and kaon
 2001 production.
- 2002 • **Solar Activity:** As the 11-year solar cycle can affect the Earth's magnetic
 2003 field, the flux of primary cosmic rays varies across the same period. The
 2004 uncertainty is calculated by taking a ± 1 year variation, equating to a 10%
 2005 uncertainty for the SK-IV period.
- 2006 • **Atmospheric Density:** The height of the interaction of the primary cosmic
 2007 rays is dependent upon the atmospheric density. The HKKM assumes the
 2008 US standard 1976 [149] profile. This systematic controls the uncertainty in
 2009 that model.

2010 The total uncertainty is dominated by the absolute and relative normalisation
 2011 parameters. The effect of which is illustrated in Figure 6.7. Generally, the
 2012 uncertainty is large at low energy, reducing to $O(10\%)$ around the peak of the
 2013 flux distribution and then increasing once the neutrino energy exceeds 10GeV.

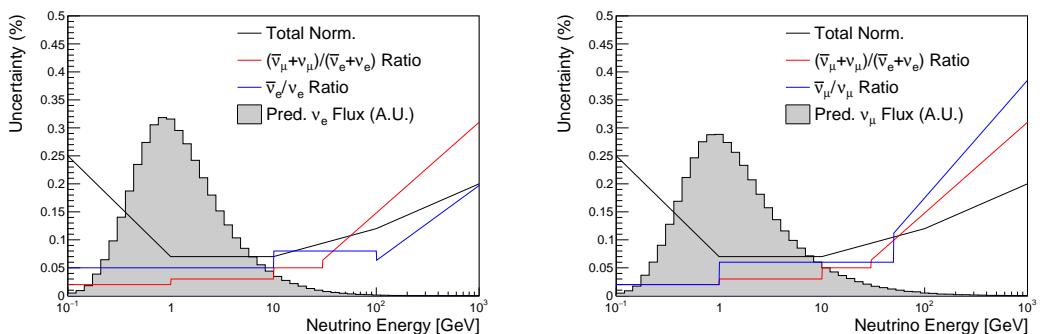


Figure 6.7: The uncertainty evaluated on the atmospheric ν_e (left) and ν_μ (right) flux predictions. The absolute normalisation and flavour ratio uncertainties are given. The solid histogram indicates the neutrino flux as a function of energy.

2014 Updates to the HKKM and Bartol models are underway [154] to use a similar
2015 tuning technique to that used in the beam flux predictions. After those updates,
2016 it may be possible to include correlations in the hadron production uncertainty
2017 systematics for beam and atmospheric flux predictions.

2018 6.4.3 Neutrino Interaction

2019 Neutrino interactions in the detectors are modeled by NEUT. The two indepen-
2020 dent oscillation analyses, T2K-only [199] and the SK-only [59], have developed
2021 separate interaction models. To maximise sensitivity out of this simultaneous
2022 beam and atmospheric analysis, a correlated interaction model has been defined
2023 in [200]. Where applicable, correlations allow the systematic uncertainties applied
2024 to the atmospheric samples to be constrained by near detector neutrino beam
2025 measurements. This can lead to stronger sensitivity to oscillation parameters
2026 as compared to an uncorrelated model.

2027 The low-energy T2K systematic model has a more sophisticated treatment
2028 of CCQE, 2p2h, and CCRES uncertainties, where extensive comparisons of
2029 this model have been performed to external data [199]. However, the model
2030 is not designed for high-energy atmospheric events, like those illustrated in
2031 Figure 5.11. Therefore the high energy systematic model from the SK-only
2032 analysis is implemented for the relevant multiGeV, PC, and up- μ samples. The
2033 T2K CCQE model is more sophisticated so it has been implemented for all
2034 samples within this analysis, where separate low-energy and high-energy dials
2035 have been implemented. The low-energy dials are constrained by the near
2036 detector measurements and are uncorrelated to their high-energy counterparts.
2037 The author of this thesis was responsible for implementing and validating the
2038 combined cross section model as documented in [200, 201].

2039 The high energy systematic model includes parameters developed from
2040 comparisons of Nieves and Rein-Seghal models which affect resonant pion
2041 producing interactions, comparisons of the GRV98 and CKMT models which
2042 control DIS interactions, and hadron multiplicity measurements which modulate

the normalisation of multi-pion producing events. The uncertainty on the ν_τ cross-section is particularly large and is controlled by a 25% normalisation uncertainty. These uncertainties are applied via normalisation or shape parameters. The former linearly scales the weight of all affected Monte-Carlo events, whereas the latter can increase or decrease a particular event's weight depending on its neutrino energy and mode of interaction. The response of the shape parameters is defined by third-order polynomial splines which return a weight for a particular neutrino energy. To reduce computational resources for the far detector fit, the response is binned by neutrino energy and sample binning: lepton momentum and cosine zenith binning for atmospheric splined responses and reconstructed neutrino energy and direction binning for beam samples. In total, 17 normalisation and 15 shape parameters are included in the high-energy model within this analysis.

Figure 6.8 indicates the predicted neutrino energy distribution for both beam and subGeV atmospheric samples. There is clearly significant overlap in neutrino energy between the subGeV atmospheric and beam samples, allowing similar kinematics in the final state particles. Figure 6.9 illustrates the fractional contribution of the different interaction modes per sample.

Comparing beam and atmospheric samples which target CCQE interactions (S.G. e-like 0de, S.G. μ -like [0,1]de, [FHC,RHC] 1R μ -like and [FHC,RHC] 1R e-like samples), there is a very similar contribution of CCQE, CC 2p2h, and CC1 π^\pm interactions. The samples which target CC1 π^\pm interactions, (S.G. e-like 0de, S.G. μ -like 2de and FHC 1R+1d.e e-like) also consist of very similar mode interactions.

As a consequence of the similarity in energy and mode contributions, correlating the systematic model between the beam and subGeV atmospheric samples ensures that this analysis attains the largest sensitivity to oscillation parameters while still ensuring neutrino interaction systematics are correctly accounted for. Due to its more sophisticated CCQE and 2p2h model, the T2K systematic model was chosen as the basis of the correlated model.

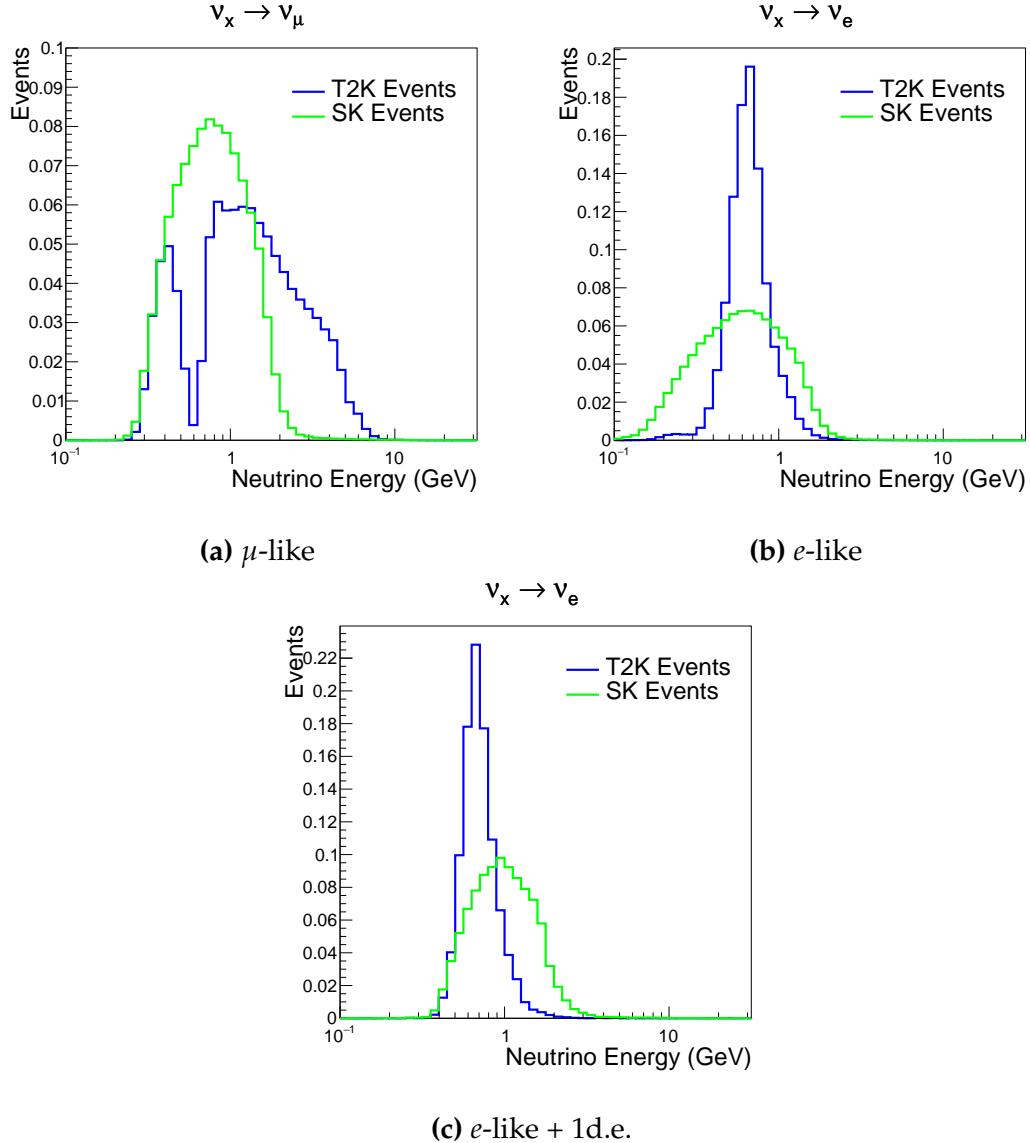


Figure 6.8: The predicted neutrino energy distribution for subGeV atmospheric and beam samples. FHC and RHC beam samples are summed together Asimov A oscillation parameters are assumed (given in Table 2.2). Beam and atmospheric samples with similar cuts are compared against one another.

The T2K systematic model [199] is applied in a similar methodology to the SK model parameters. It consists of 19 shape parameters and 24 normalisation parameters. Four additional parameters, which model the uncertainty in the binding energy, are applied in a way to shift the momentum of the lepton emitted from a nucleus. This controls the uncertainty specified on the 27MeV binding energy assumed within Equation 6.2. The majority of these parameters are

	CC QE	CC 2p2h	CC $1\pi^\pm$	CC $M\pi$	CC Other	NC $1\pi^0$	NC $1\pi^\pm$	NC $M\pi$	NC Coh.	NC Other
FHC 1R+1d.e. e-like	0.04	0.02	0.83	0.03	0.04	0.01	0.01	0.01	0.00	0.01
RHC 1R e-like	0.62	0.12	0.11	0.01	0.02	0.06	0.01	0.01	0.01	0.04
FHC 1R e-like	0.68	0.12	0.10	0.00	0.02	0.04	0.01	0.00	0.00	0.02
RHC 1R μ -like	0.62	0.13	0.17	0.02	0.03	0.00	0.02	0.00	0.00	0.00
FHC 1R μ -like	0.62	0.12	0.16	0.02	0.03	0.00	0.03	0.00	0.00	0.00
S.G. π^0 -like	0.05	0.01	0.02	0.00	0.01	0.68	0.06	0.07	0.06	0.04
S.G. μ -like 2de	0.04	0.01	0.80	0.10	0.04	0.00	0.00	0.00	0.00	0.00
S.G. μ -like 1de	0.72	0.11	0.12	0.01	0.02	0.00	0.01	0.00	0.00	0.00
S.G. μ -like 0de	0.68	0.11	0.10	0.01	0.02	0.01	0.05	0.01	0.00	0.02
S.G. e-like 1de	0.05	0.01	0.75	0.10	0.05	0.00	0.01	0.02	0.00	0.01
S.G. e-like 0de	0.73	0.11	0.10	0.01	0.02	0.02	0.00	0.00	0.00	0.00

Figure 6.9: The interaction mode contribution of each sample given as a fraction of the total event rate in that sample. Asimov A oscillation parameters are assumed (given in Table 2.2). The Charged Current (CC) modes are broken into quasi-elastic (QE), 2p2h, resonant charged pion production ($1\pi^\pm$), multi-pion production ($M\pi$), and other interaction categories. Neutral Current (NC) interaction modes are given in interaction mode categories: π^0 production, resonant charged pion production, multi-pion production, and others.

2079 assigned a Gaussian prior uncertainty. Those that have no reasonably motivated
 2080 uncertainty, or those which have not been fit to external data, are assigned a
 2081 flat prior which does not affect the penalty term.

2082 On top of the combination of the SK and T2K interaction models, several
 2083 other parameters have been specifically developed for the joint oscillation anal-
 2084 ysis. The majority of the atmospheric samples' δ_{CP} sensitivity comes from the
 2085 normalisation of subGeV electron-like events. These are modeled using a spectral
 2086 function to approximate the nuclear ground state. However, the near detector is
 2087 not able to constrain the model so an additional systematic is introduced which
 2088 models applies an alternative Continuous Random Phase Approximation (CRPA)
 2089 nuclear ground state. This dial approximates the event weights if a CRPA model

2090 had been assumed rather than a spectral function. This dial only applies to ν_e
2091 and $\bar{\nu}_e$ as the near detector does not constraint ν_e cross-section measurements.
2092 It is applied as a shape parameter.

2093 Further additions to the model have been introduced due to the inclusion of
2094 the subGeV π^0 atmospheric sample. This particularly targets charged current
2095 and neutral current π^0 producing interactions to help constrain the systematic
2096 uncertainties. Therefore, an uncertainty that affects neutral current resonant
2097 π^0 production is incorporated into this analysis. Comparisons of NEUT's NC
2098 resonant pion production predictions have been made to MiniBooNE [202] data
2099 and a consistent 16% to 21% underprediction is observed [200]. Consequently,
2100 a conservative 30% normalisation parameter is invoked.

2101 Down-going events are mostly insensitive to oscillation parameters and can
2102 act similar to the near detector within an accelerator experiment (Details will
2103 be discussed in chapter 7). This region of phase space can act as a sideband
2104 and allows the cross-section model and near detector constraint to be studied.
2105 The distribution of events in this region is calculated using the technique out-
2106 lined in subsection 4.3.4. The results are illustrated in Figure 6.10. For CCQE-
2107 targeting samples, the application of the near detector constraint is well within
2108 the statistical fluctuation of the down-going data such that no significant tension
2109 is observed between the data and the Monte Carlo prediction after the near
2110 detector constraint is applied. This is not the case for samples with target CCRES
2111 interactions. The electron-like data is consistent with the constrained prediction
2112 at high reconstructed momenta but diverges at lower momentum, whereas the
2113 muon-like sample is under-predicted throughout the range of momenta. To
2114 combat this disagreement, an additional cross-section systematic dial, specifically
2115 designed to inflate the low pion momentum systematics was developed in [200].
2116 This is a shape parameter implemented through a splined response.

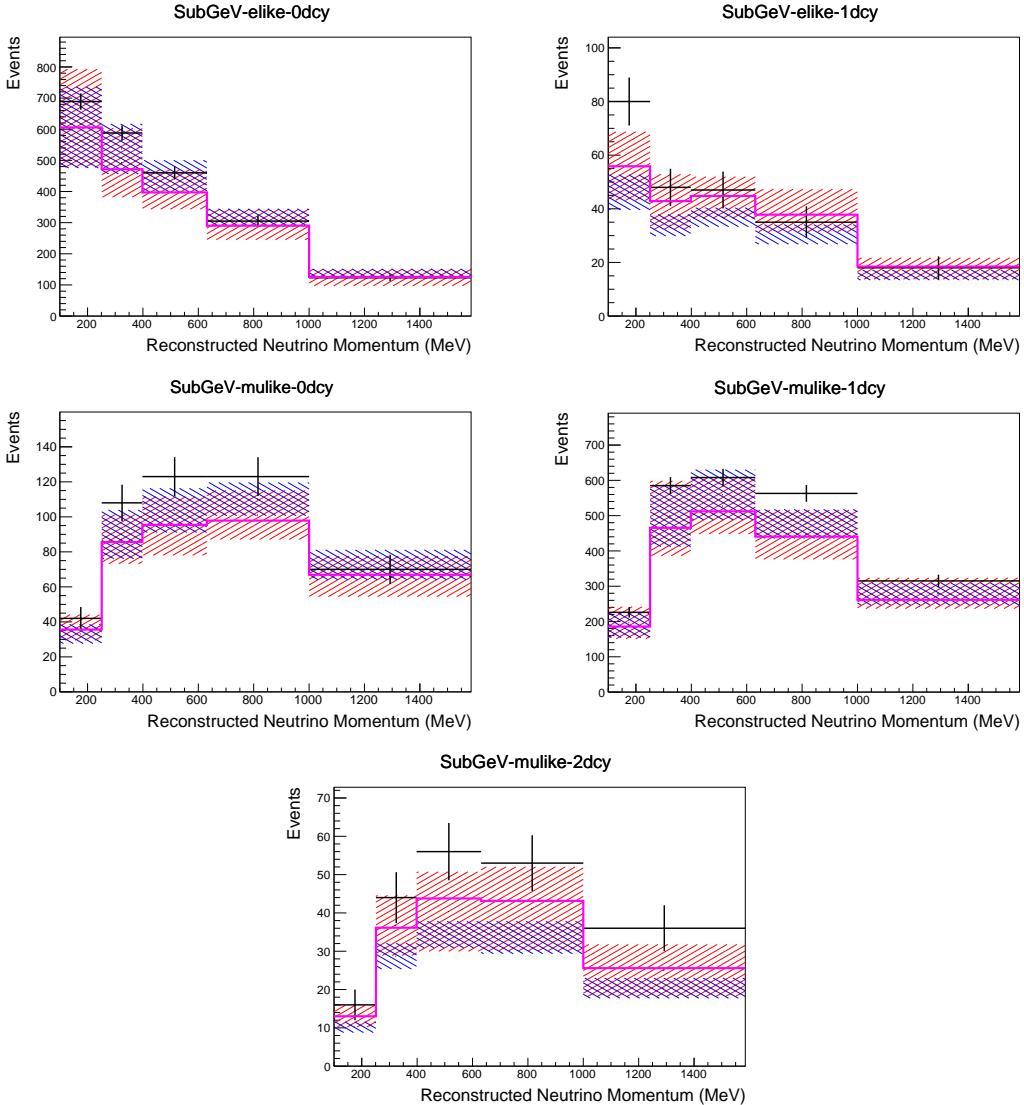


Figure 6.10: Down-going atmospheric subGeV single-ring samples comparing the mean and error of the pre-fit and post-fit Monte Carlo predictions in red and blue, respectively. The magenta histogram illustrates the Monte Carlo prediction using the generated dial values. The black points illustrate the down-going data with statistical errors given. The mean and errors of the Monte Carlo predictions are calculated by the techniques documented in subsection 4.3.4. The pre-fit spectrum is calculated by throwing the cross-section and atmospheric flux dial values from the pre-fit covariance matrix. The post-fit spectrum is calculated by sampling the cross-section dial values from an ND fit MCMC chain, whilst still throwing the atmospheric flux dials from the pre-fit covariance.

2117 6.4.4 Near Detector

2118 The systematics applied due to uncertainties arising from the response of the near
 2119 detector is documented in [129]. The response is described by 574 normalisation
 2120 parameters binned in the selected sample as well as momentum and angle,

2121 P_μ and $\cos(\theta_\mu)$, of the final-state muon. These are applied via a covariance
2122 matrix with each parameter being assigned a Gaussian prior from that covariance
2123 matrix. These normalisation parameters are built from underlying systematics,
2124 e.g. pion secondary interaction systematics, which are randomly thrown and
2125 the variation in each $P_\mu \times \cos(\theta_\mu)$ bin is determined. Two thousand throws are
2126 evaluated and a covariance matrix response is created. This allows significant
2127 correlations between FGD1 and FGD2 samples, as well as adjacent $P_\mu \times \cos(\theta_\mu)$
2128 bins. Statistical uncertainties are accounted for by including fluctuations of each
2129 event's weight from a Poisson distribution.

2130 Similar to the cross-section systematics, MaCh3 and BANFF are used to
2131 constrain the uncertainty of these systematics through independent validations.
2132 Each fitter generates a post-fit covariance matrix which is compared and passed
2133 to the far-detector oscillation analysis working group. As the analysis presented
2134 within this thesis uses the MaCh3 framework, a joint oscillation analysis fit of all
2135 three sets of samples and their respective systematics is performed.

2136 6.4.5 Far Detector

2137 Two configurations of the far detector systematic model implementation have
2138 been considered. Firstly, the far detector systematic uncertainties for beam and
2139 atmospheric samples are taken from their respective analysis inputs, denoted
2140 “official inputs” analysis, with no correlations assumed between the beam and at-
2141 mospheric samples. The beam- and atmospheric-specific inputs are documented
2142 in subsubsection 6.4.5.1 and subsubsection 6.4.5.2. Secondly, an alternative
2143 detector model has been developed which correlates the response of the SK
2144 detector systematics between the beam and atmospheric samples. Here, the
2145 distribution of parameters used for applying event cuts (e.g. electron-muon
2146 PID separation) is modified within the fit. It follows a similar methodology to
2147 the beam far detector systematics implementation but performs a joint fit of
2148 the beam and atmospheric data. This alternative implementation is detailed
2149 in subsubsection 6.4.5.3.

2150 **6.4.5.1 Beam Samples**

2151 There are 45 systematics which describe the response of the far detector to
2152 beam events [181], split into 44 normalisation parameters and one energy scale
2153 systematic. The energy scale systematic is applied as a multiplicative scaling
2154 of the reconstructed neutrino energy. It is estimated from data-to-Monte Carlo
2155 differences in the stopping muon sample in [183] and found to be 2.1%. The
2156 normalisation parameters are assigned a Gaussian error centred at one with
2157 width taken from a covariance matrix. A detailed breakdown of the generation
2158 of the covariance matrix is found in [194]. To build the covariance matrix, a fit
2159 is performed on atmospheric data which has been selected using beam sample
2160 selection cuts. These cuts use the variables, L^i , where the index i is detailed in
2161 Table 6.6. Each L^i is a smear, α , and shift, β parameter such that,

$$L_j^i \rightarrow \bar{L}_j^i = \alpha_j^i L + \beta_j^i \quad (6.4)$$

2162 Where L_j^i (\bar{L}_j^i) correspond to nominal(varied) PID cut parameters given in
2163 Table 6.6. The shift and smear parameters are nuisance parameters with no prior
2164 constraints. They are binned by final-state topology, j , where the binning is given
2165 in Table 6.7. The final-state topology binning is because the detector will respond
2166 differently to events that have one or multiple rings. For example, the detector
2167 will be able to distinguish single-ring events better than two overlapping ring
2168 events, resulting in different systematic uncertainty for one-ring events compared
2169 to two-ring events. This approach is used to allow the cut parameter distributions
2170 to be modified within the fit, allowing for better data to Monte Carlo agreement.

Cut Variable	Parameter Name
0	<code>fitQun e/mu PID</code>
1	<code>fitQun e/pi0 PID</code>
2	<code>fitQun mu/pi PID</code>
3	<code>fitQun Ring-Counting Parameter</code>

Table 6.6: List of cut variables that are included within the shift/smear fit documented in [194].

Category	Description
1e	Only one electron above Cherenkov threshold in the final state
1 μ	Only one muon above Cherenkov threshold in the final state
1e+other	One electron and one or more other charged particles above Cherenkov threshold in the final state
1 μ +other	One muon and one or more other charged particles above Cherenkov threshold in the final state
1 π^0	Only one π^0 in the final state
1 π^\pm or 1p	Only one hadron (typically charged pion or proton) in the final state
Other	Any other final state

Table 6.7: Reconstructed event topology categories on which the SK detector systematics [194] are based.

2171 The mis-modeling of π^0 events is also considered. If one of the two rings
 2172 from a π^0 event is missed, this will be reconstructed as a CC ν_e -like event. This
 2173 is one of the largest systematics hindering the electron neutrino appearance
 2174 analyses. Consequently, additional systematics have been introduced to con-
 2175 strain the mis-modeling of π^0 events in SK, binned by reconstructed neutrino
 2176 energy. To evaluate this systematic uncertainty, a set of “hybrid- π^0 ” samples is
 2177 constructed. These events are built by overlaying one electron-like ring from
 2178 the SK atmospheric neutrino samples or decay electron ring from a stopping
 2179 cosmic ray muon with one simulated photon ring. Both rings are chosen so
 2180 that momenta and opening angle follow the decay kinematics of NC π^0 events
 2181 from the T2K-MC. Hybrid- π^0 Monte Carlo samples with both rings from the
 2182 SK Monte Carlo are produced to compare with the hybrid- π^0 data samples and
 2183 the difference in the fraction of events that pass the ν_e selection criteria is used
 2184 to assign the systematic errors. In order to investigate any data to Monte Carlo
 2185 differences that may originate from either the higher energy ring or lower energy
 2186 ring, two samples are built; a sample in which the electron constitutes the higher
 2187 energy ring from the π^0 decay (called the primary sample) and another one in
 2188 which it constitutes the lower energy ring (called the secondary sample). The
 2189 standard T2K ν_e fitQun event selection criteria are used to select events.

2190 Final contributions to the covariance matrix are determined by supplemen-
 2191 tary uncertainties obtained by comparing stopping muon data to Monte Carlo

prediction, as first introduced in section 5.2. The efficiency of tagging decay electrons is estimated by the stopping muon data to Monte Carlo differences by comparing the number of one decay electron events to the number of events with one or fewer decay electrons. Similarly, the rate at which fake decay electrons are reconstructed by `f iTQun` is estimated by comparing the number of two decay electron events to the number of events with one or two reconstructed decay electrons. The two sources of systematics are added in quadrature weighted by the number of events with one true decay electron yielding a 0.2% systematic uncertainty. A fiducial volume systematic of $\pm 2.5\text{cm}$ which corresponds to a 0.5% shift in the normalisation of events is also applied. Additional normalisation uncertainties based on neutrino flavour and interaction mode are also defined in [181, 203, 204].

Two additional sources of uncertainty are included: secondary and photoneuclear interactions. These are estimated by varying the underlying parameters are building a distribution of sample event rates. These contributions are then added in quadrature to the above covariance matrix.

DB: Matrix Diagonals from SK Det Matrix

6.4.5.2 Atmospheric Samples

The detector systematics for atmospheric samples, documented in [84], are split into two sub-groups: those which are related to particle identification and ring counting systematics, and those which are related to calibration, separation, and reduction uncertainties.

The particle identification systematics consist of five parameters. The ring separation systematic enforces an anti-correlated response between the single-ring and multi-ring samples. This is implemented as a fractional increase/decrease in the overall normalisation of each sample, depending on the distance to the nearest wall from an event's vertex. The coefficients of the normalisation are estimated prior to the fit and depend on the particular atmospheric sample. Two electron-muon separation systematics are included within this model which

2221 anti-correlate the response of the electron-like and muon-like samples: one for
2222 single-ring events and another for multi-ring events.

2223 The multi-ring electron-like separation likelihood, discussed in section 6.1,
2224 encodes the ability of the detector to separate neutrino from anti-neutrino events.
2225 Two normalisation parameters vary the relative normalisation of multi-ring ν_e
2226 and $\bar{\nu}_e$ samples whilst keeping a consistent overall event rate.

2227 There are 22 systematics related to calibration measurements, including effects
2228 from backgrounds, reduction, and showering effects. They are documented in
2229 [84] and are briefly summarised in Table 6.8. They are applied via normalisation
2230 parameters, with the separation systematics requiring the conservation of event
2231 rate across all samples.

2232 6.4.5.3 Correlated Detector Model

2233 A complete uncertainty model of the SK detector would be able to determine
2234 the systematic shift on the sample spectra for a variation of the underlying
2235 parameters, e.g. PMT angular acceptance. However, this is computationally
2236 intensive, requiring Monte Carlo predictions to be made for each plausible
2237 variation. Consequently, an effective parameter model has been utilised for
2238 a correlated detector model following from the T2K-only model implementation
2239 documented in subsubsection 6.4.5.1. It correlates the detector systematics
2240 between the far-detector beam and subGeV atmospheric samples due to their
2241 similar energies and interaction types. As there are no equivalent beam samples,
2242 the multiGeV, multiring, PC, and Up- μ samples will be subject to the particle
2243 identification systematics implementation as described in subsubsection 6.4.5.2
2244 rather than using this correlated detector model. The calibration systematics also
2245 described in the aforementioned chapter still apply to all atmospheric samples.

2246 The correlated detector model utilises the same smear and shift parameters
2247 documented in subsubsection 6.4.5.1, split by final state topology. Beyond this,
2248 the shift and smear parameters are split by visible energy deposited within the
2249 detector, with binning specified in Table 6.9. This is because atmospheric events

Index	Description
0	Partially contained reduction
1	Fully contained reduction
2	Separation of fully contained and partially contained events
3	Separation of stopping and through-going partially contained events in top of detector
4	Separation of stopping and through-going partially contained events in barrel of detector
5	Separation of stopping and through-going partially contained events in bottom of detector
6	Background due to cosmic rays
7	Background due to flasher events
8	Vertex systematic moving events into and out of fiducial volume
9	Upward going muon event reduction
10	Separation of stopping and through-going in upward going muon events
11	Energy systematic in upward going muon events
12	Reconstruction of the path length of upward going muon events
13	Separation of showering and non-showering upward going muon events
14	Background of stopping upward going muon events
15	Background of non-showering through-going upward going muon events
16	Background of showering through-going upward going muon events
17	Efficiency of tagging two rings from π^0 decay
18	Efficiency of decay electron tagging
19	Background from downgoing cosmic muons
20	Asymmetry of energy deposition in tank
21	Energy scale deposition

Table 6.8: Sources of systematic errors specified within the grouped into the “calibration” systematics model.

are categorised by subGeV and multiGeV events based on visible energy, so this splitting is required when correlating the systematic model for beam and atmospheric events. Alongside the technical requirement, higher energy events will be better reconstructed due to fractionally less noise within the detector. As a result of the inclusion of visible energy binning, Equation 6.4 becomes

$$L_{jk}^i \rightarrow \bar{L}_{jk}^i = \alpha_{jk}^i L + \beta_{jk}^i, \quad (6.5)$$

where k is the visible energy bin.

The implementation of this systematic model takes the events reconstructed values of the cut parameters, modifies them by the particular shift and smear

Index	Range (MeV)
0	$30 \geq E_{vis} > 300$
1	$300 \geq E_{vis} > 700$
2	$700 \geq E_{vis} > 1330$
3	$E_{vis} \geq 1330$

Table 6.9: Visible energy binning for which the correlated SK detector systematics are based

parameter for that event, and then re-applies event selection. This causes event migration, which is a new feature incorporated into the MaCh3 framework which is only achievable due to the event-by-event reweighting scheme.

Particular care has to be taken when varying the ring counting parameter. This is because the number of rings is a finite value (one-ring, two-rings, etc.) which can not be continuously varied through this shift and smear technique. Consequently a continuous ring counting parameter, RC_i , is calculated for the i^{th} event, following the definition in [180]: the preferred likelihoods from all considered one-ring (L_{1R}) and two-ring (L_{2R}) fits are determined. The difference is computed as $\Delta_{LLH} = \log(L_{1R}) - \log(L_{2R})$. The ring counting parameter is then defined as

$$RC_i = \text{sgn}(\Delta_{LLH}) \times \sqrt{|\Delta_{LLH}|}, \quad (6.6)$$

where $\text{sgn}(x) = x/|x|$. This ring counting parameter corresponds to an intermediate likelihood value used within the `fitQun` algorithm to decide the number of rings associated with a particular event. However, fake-ring merging algorithms are applied after this likelihood value is used. Consequently, this ring counting parameter does not always exactly correspond to the number of reconstructed rings. This can be seen in Figure 6.11.

As the `fitQun` algorithm does not provide a likelihood value after the fake-ring algorithms have been applied, the ring counting parameter distribution is correlated to the final number of reconstructed rings through “maps”. These are two-dimensional distributions of the ring counting parameter and the final number of reconstructed rings. An example is illustrated in Figure 6.12. In principle,

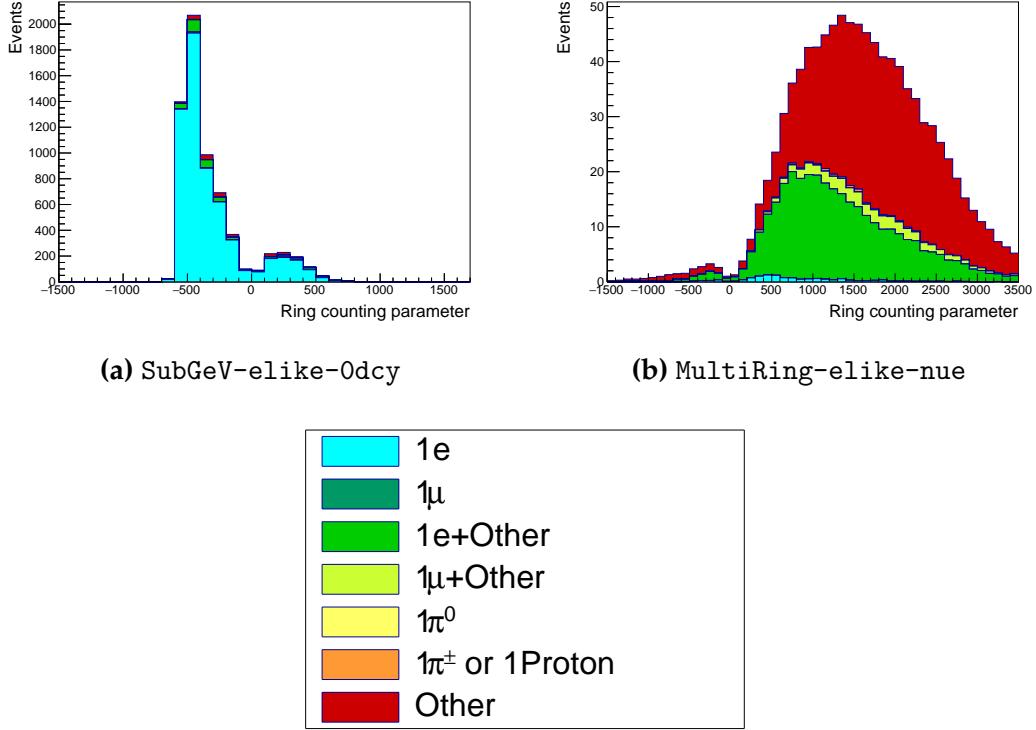


Figure 6.11: The ring counting parameter as defined in Equation 6.6 for the SubGeV-elike-0dcy and MultiRing-elike-nue samples.

the `fitQun` reconstruction algorithm should be re-run after the variation in the ring counting parameter. However, this is not computationally viable. Therefore the “maps” are used as a reweighting template.

The maps are split by final state topology and true neutrino flavour and all `fitQun`-reconstructed Monte Carlo events are used to fill them. The maps are row-normalised to represent the probability of X rings for a given RC_i value. Prior to the oscillation fit, an event’s nominal weight is calculated as $W^i(N_{Rings}^i, L_{jk}^i)$, where N_{Rings}^i is the reconstructed number of rings for the i^{th} event and $W^i(x, y)$ is the bin content in map associated with the i^{th} event, where x is the number of rings and y is the ring counting parameter. Then during the fit, the value of $R = W^i(N_{Rings}^i, \bar{L}_{jk}^i) / W^i(N_{Rings}^i, L_{jk}^i)$ is calculated as the event weight for the i^{th} event. This is the only cut variable that uses a reweighting technique rather than event migration.

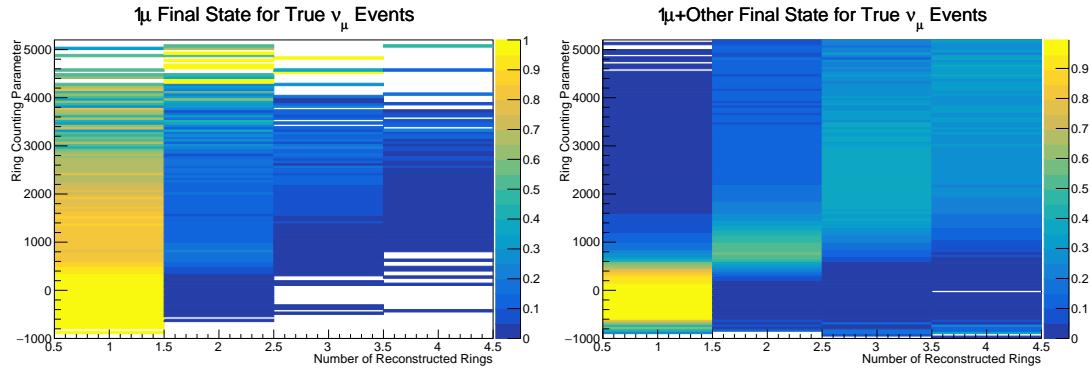


Figure 6.12: The ring counting parameter, defined in Equation 6.6, as a function of the number of reconstructed rings as found by the `fitQun` reconstruction algorithm. Left: true ν_μ events with only one muon above the Cherenkov threshold in the final state. Right: true ν_μ events with one muon and at least one other charged particle above the Cherenkov threshold in the final state.

2293 The π^0 systematics introduced in subsection 6.4.4 are applied via a covariance
 2294 matrix. This is not possible in the alternative model as no covariance matrix
 2295 is used. Thus, the implementation of the π^0 systematics has been modified.
 2296 The inputs from the hybrid π^0 sample are included via the use of “ χ^2 maps”,
 2297 which are two-dimensional histograms in α_{jk}^i and β_{jk}^i parameters over some
 2298 range. Illustrative examples of the χ^2 maps are given in Figure 6.13. Due to
 2299 their nature, the shift and smear parameters are typically very correlated. A
 2300 map is produced for each cut parameter given in Table 6.6 and for each visible
 2301 energy bin given in Table 6.9.

2302 The maps are filled through the χ^2 comparison of the hybrid π^0 Monte Carlo
 2303 and data in the particle identification parameters documented in Table 6.6. The
 2304 Monte Carlo distribution is modified by the α_{jk}^i and β_{jk}^i scaling, whilst cross-
 2305 section and flux nuisance parameters are thrown from their prior uncertainties.
 2306 The χ^2 between the scaled Monte Carlo and data is calculated and the relevant
 2307 point in the χ^2 map is filled.

2308 The implementation within this alternative detector model is to add the bin
 2309 contents of the maps, for the relevant values of the α_{jk}^i and β_{jk}^i parameters, to
 2310 the likelihood penalty. Only 1 π^0 final state topology shift and smear parameters
 2311 use this prior uncertainty.

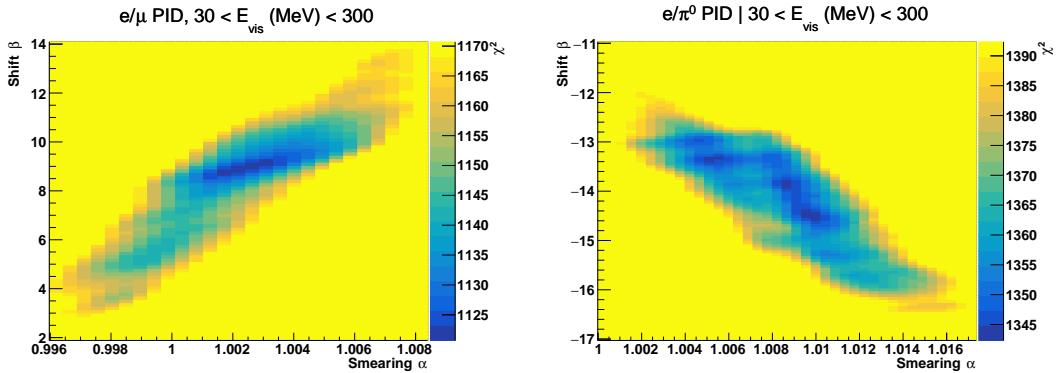


Figure 6.13: The χ^2 between the hybrid- π^0 Monte Carlo and data samples, as a function of smear (α) and shift (β) parameters, for events which have $1\pi^0$ final state topology. Left: Electron-muon separation PID parameter for events with $30 \leq E_{\text{vis}}(\text{MeV}) < 300$. Right: Electron- π^0 separation PID parameter for events with $30 \leq E_{\text{vis}}(\text{MeV}) < 300$.

Similarly, the implementation of the supplementary systematics documented in subsubsection 6.4.5.1 needs to be modified. A new framework [205] was built in tandem with the T2K-SK working group [181] so the additional parameters can be incorporated into the MaCh3 framework. These are applied as normalisation parameters, depending on the particular interaction mode, number of tagged decay electrons, and whether the primary particle generated Cherenkov light. They are assigned Gaussian uncertainties with widths described by a covariance matrix. Furthermore, the secondary interaction and photo-nuclear effects need to be accounted for in this detector model using a different implementation than that in subsubsection 6.4.5.1. This was done by including a shape parameter for each of the secondary interactions and the photo-nuclear systematic parameters.

There are a total of 224 α_{jk}^i and β_{jk}^i parameters, of which 32 have prior constraints from the hybrid π^0 samples.

One final complexity of this correlated detector model is that the two sets of samples, beam and subGeV atmospheric, use slightly different parameters to distinguish electron and muon-like events. The T2K samples use the value of $\log(L_e/L_\mu)$ whereas the atmospheric samples use the value of $\log(L_e/L_\pi)$, where L_X is the likelihood for hypothesis X. This is because the T2K fits use single-ring fiTQun fitting techniques, whereas multi-ring fits are applied to the atmospheric samples where only the electron and pion hypothesis are considered.

2332 The correlation between the two likelihood ratios is illustrated in Figure 6.14. As
 2333 discussed in section 5.2, the pion hypothesis is a very good approximation of the
 2334 muon hypothesis due to their similar mass. Consequently, using the same shift
 2335 and smear parameters correlated between the beam and subGeV atmospheric
 2336 samples is deemed a good approximation.

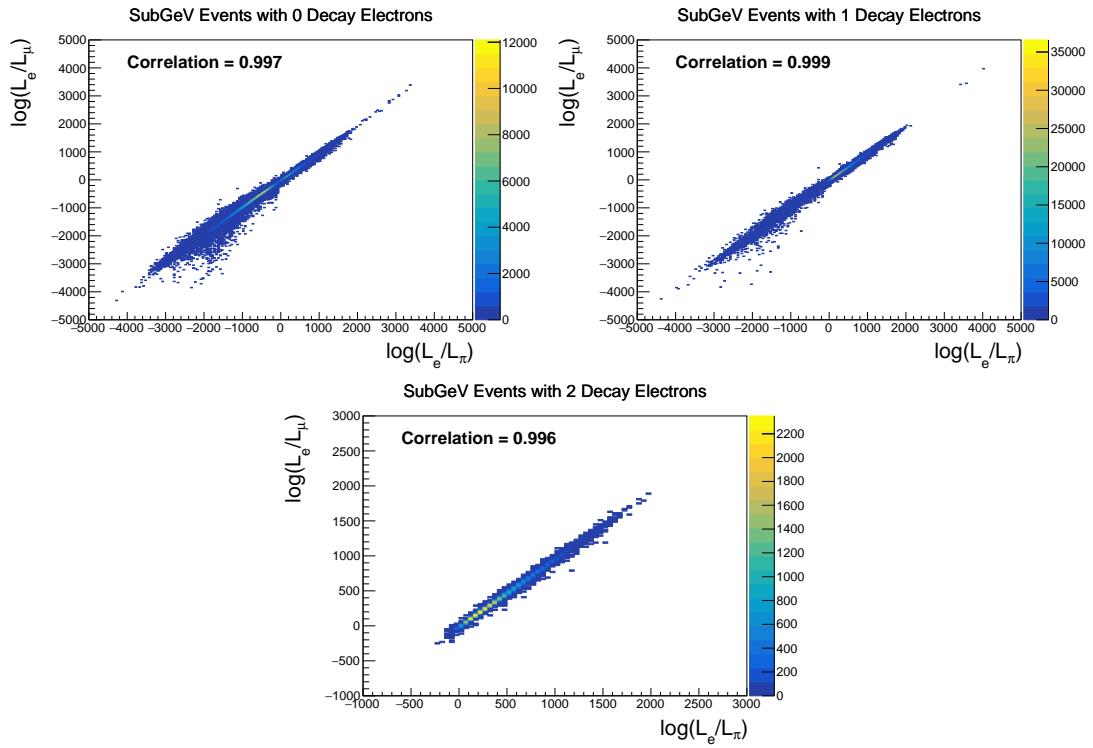


Figure 6.14: The distribution of $\log(L_e/L_\mu)$ compared to $\log(L_e/L_\pi)$ for subGeV events with zero (top left), one (top right) or two (bottom) decay electrons. The correlation in the distribution is calculated as 0.997, 0.999 and 0.996, respectively.

7

2337

2338

Oscillation Probability Calculation

2339 It is important to understand how and where the sensitivity to the oscillation
2340 parameters comes from for both atmospheric and beam samples. An overview
2341 of how these samples observe changes in δ_{CP} , Δm^2_{23} , and $\sin^2(\theta_{23})$ affect
2342 these samples is given in section 7.1. This section also explains the additional
2343 complexities involved when performing an atmospheric neutrino analysis as
2344 compared to a beam-only analysis.

2345 Without additional techniques, atmospheric sub-GeV upward-going neutrinos ($E_\nu < 1.33\text{GeV}, \cos(\theta_Z) < 0.$) can artificially inflate the sensitivity to δ_{CP}
2346 due to the quickly varying oscillation probability in this region. Therefore, a
2347 “sub-sampling” approach has been developed to reduce these biases ensuring
2348 accurate and reliable sensitivity measurements. This technique ensures that small-
2349 scale unresolvable features of the oscillation probability have been averaged over
2350 whilst the large-scale features in the oscillation probability are unaffected. The
2351 documentation and validation of this technique are found in section 7.2. The
2352 oscillation probability calculation is computationally intensive due to the large
2353 number of matrix multiplications needed. Consequently, the CUDAProb3 imple-
2354 mentation choice made within the fitting framework, as detailed in section 7.3,
2355 ensures that the analysis can be done in a timely manner.

Whilst the beam neutrinos are assumed to propagate through a constant density slab of material, the density variations through the Earth result in more complex oscillation patterns for atmospheric neutrinos. Furthermore, the uncertainty in the electron density can modify the oscillation probability for the denser core layers of the Earth. The model of the Earth used within this analysis is detailed in section 7.4. This includes information about the official SK-only methodology as well as improvements that have been made to remove some of the approximations used in that analysis. Another complexity of atmospheric neutrinos oscillation studies is that the height of production in the atmosphere is not known on an event-by-event basis. An analytical averaging technique that approximates the uncertainty of the oscillation probability has been followed, with the author of this thesis being responsible for the implementation and validation.

This implementation of an external technique is illustrated in section 7.5.

7.1 Overview

DB: Should this be moved into an earlier chapter? The selections chapter references the matter resonance which has not yet been explained at that point

The analysis presented within this thesis focuses on the determination of oscillation parameters from atmospheric and beam neutrinos. Whilst subject to the same oscillation formalism, the way in which the two samples have sensitivity to the different oscillation parameters differs significantly.

Atmospheric neutrinos have a varying baseline, or “path length”, L , such that the distance each neutrino travels before interacting is dependent upon the zenith angle, θ_Z . As primary cosmic rays can interact anywhere between the Earth’s surface and $\sim 50\text{km}$ above that, the height, h , in the atmosphere at which the neutrino was generated also affects the path length,

$$L = \sqrt{(R_E + h)^2 - R_E^2 (1 - \cos^2(\theta_Z))} - R_E \cos(\theta_Z). \quad (7.1)$$

2382 Where $R_E = 6,371\text{km}$ is the Earth's radius. Consequently, the oscillation
 2383 probability is dependent upon two parameters, $\cos(\theta_Z)$ and E_ν .

2384 The oscillation probability used within this analysis is based on [21]. The
 2385 neutrino wavefunction in the vacuum Hamiltonian evolves in each layer of
 2386 constant matter density via

$$i \frac{d\psi_j(t)}{dt} = \frac{m_j^2}{2E_\nu} \psi_j(t) - \sum_k \sqrt{2} G_F N_e U_{ej} U_{ke}^\dagger \psi_k(t), \quad (7.2)$$

2387 where m_j^2 is the square of the j^{th} vacuum eigenstate mass, E_ν is the neutrino
 2388 energy, G_F is Fermi's constant, N_e is the electron number density and U is the
 2389 PMNS matrix. The transformation $N_e \rightarrow -N_e$ and $\delta_{CP} \rightarrow -\delta_{CP}$ is applied
 2390 for antineutrino propagation. Thus, a model of the Earth's density is required
 2391 for neutrino propagation. Following the official SK-only methodology [206],
 2392 this analysis uses the Preliminary Reference Earth Model (PREM) [207] which
 2393 provides piecewise cubic polynomials as a function of the Earth's radius. This
 2394 density profile is illustrated in Figure 7.1. As the propagator requires layers of
 2395 constant density, the SK methodology approximates the PREM model by using
 2396 four layers of constant density [206], detailed in Table 7.1.

Layer	Outer Radius [km]	Density [g/cm ³]	Chemical composition (Z/A)
Inner Core	1220	13	0.468 ± 0.029
Outer Core	3480	11.3	0.468 ± 0.029
Lower Mantle	5701	5.0	0.496
Transition Zone	6371	3.3	0.496

Table 7.1: Description of the four layers of the Earth invoked within the constant density approximation of the PREM model [207].

2397 The atmospheric neutrino oscillation probabilities can be presented as two di-
 2398 mensional “oscillograms” as illustrated in Figure 7.2. The distinct discontinuities,
 2399 as a function of $\cos(\theta_Z)$, are due to the discontinuous density in the PREM model.

2400 Atmospheric neutrinos have sensitivity to δ_{CP} through the overall event
 2401 rate. Figure 7.3 illustrates the difference in oscillation probability between CP-
 2402 conserving ($\delta_{CP} = 0$) and a CP-violating ($\delta_{CP} = -1.601$) value taken from

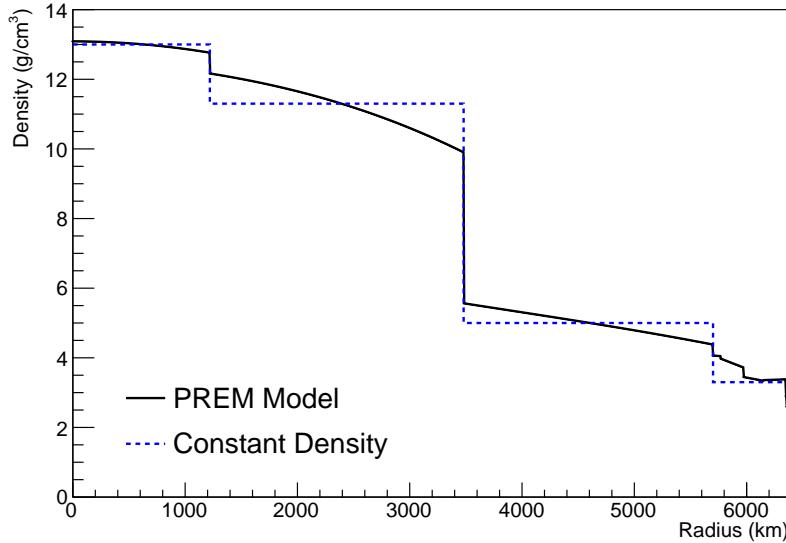


Figure 7.1: The density of the Earth given as a function of the radius, as given by the PREM model (Black), and the constant density four-layer approximation (Blue), as used in the official SK-only analysis.

2403 Asimov A oscillation parameter set (Table 2.2). The result is a complicated
 2404 oscillation pattern in the appearance probability for sub-GeV upgoing neutrinos.
 2405 The detector does not have sufficient resolution to resolve these individual
 2406 patterns so the sensitivity to δ_{CP} for atmospheric neutrinos comes via the overall
 2407 normalisation of these events.

2408 The presence of matter means that the effect δ_{CP} has on the oscillation
 2409 probability is not equal between neutrinos and antineutrinos. Furthermore,
 2410 the interaction cross-section for neutrinos is larger than antineutrinos so the
 2411 two effects have to be disentangled. These effects are further convoluted by
 2412 detector efficiencies as SK cannot distinguish neutrinos and antineutrinos well.
 2413 Furthermore, the sample selections discussed in section 6.1 have difference
 2414 efficiencies for neutrino-antineutrino selections. All of these effects lead to
 2415 a difference in the number of neutrinos detected compared to antineutrinos.
 2416 This changes how the δ_{CP} normalisation term is observed, resulting in a very
 2417 complex sensitivity to δ_{CP} .

2418 The vacuum and matter oscillation probabilities for $P(\nu_e \rightarrow \nu_e)$ and $P(\bar{\nu}_e \rightarrow$

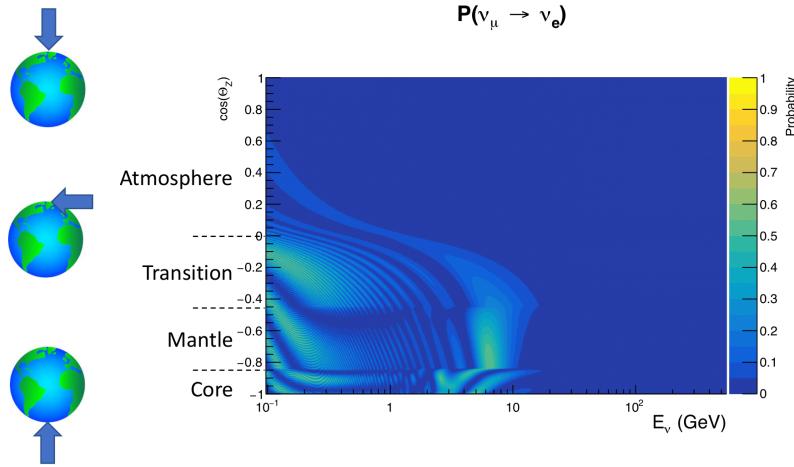


Figure 7.2: An “oscillogram” that depicts the $P(\nu_\mu \rightarrow \nu_e)$ oscillation probability as a function of neutrino energy and cosine of the zenith angle. The zenith angle is defined such that $\cos(\theta_Z) = 1.0$ represents neutrinos that travel from directly above the detector. The four-layer constant density PREM model approximation is used and Asimov A oscillation parameters are assumed (Table 2.2).

2419 $\bar{\nu}_e$) are presented in Figure 7.4, where the PREM model has been assumed.
 2420 The oscillation probability for both neutrinos and antineutrinos is affected in
 2421 the presence of matter. However, the resonance effects around $O(5)\text{GeV}$ only
 2422 occur for neutrinos in normal mass hierarchy and antineutrinos in inverse mass
 2423 hierarchy. The exact position and amplitude of the resonance depend on $\sin^2(\theta_{23})$,
 2424 further increasing the atmospheric neutrinos’ sensitivity to the parameter.

2425 As the T2K beam flux is centered at the first oscillation maximum ($E_\nu =$
 2426 0.6GeV), the sensitivity to δ_{CP} is predominantly observed as a change in the
 2427 event-rate of e-like samples in $\nu/\bar{\nu}$ modes. Figure 7.5 illustrates the $P(\nu_\mu \rightarrow \nu_e)$
 2428 oscillation probability for a range of δ_{CP} values. A circular modulation of the
 2429 first oscillation peak (in both magnitude and position) is observed when varying
 2430 throughout the allowable values of δ_{CP} . The CP-conserving values of $\delta_{CP} = 0, \pi$
 2431 have a lower(higher) oscillation maximum than the CP-violating values of $\delta_{CP} =$
 2432 $-\pi/2(\delta_{CP} = \pi/2)$. A sub-dominant shift in the energy of the oscillation peak is
 2433 also present, which aids in separating the two CP-conserving values of δ_{CP} .

2434 T2K’s sensitivity to $\sin^2(\theta_{23})$ and Δm_{23}^2 is observed as a shape-based variation

$$\mathbf{P}(\nu_\mu \rightarrow \nu_e; \delta_{CP} = -1.601) - \mathbf{P}(\nu_\mu \rightarrow \nu_e; \delta_{CP} = 0.)$$

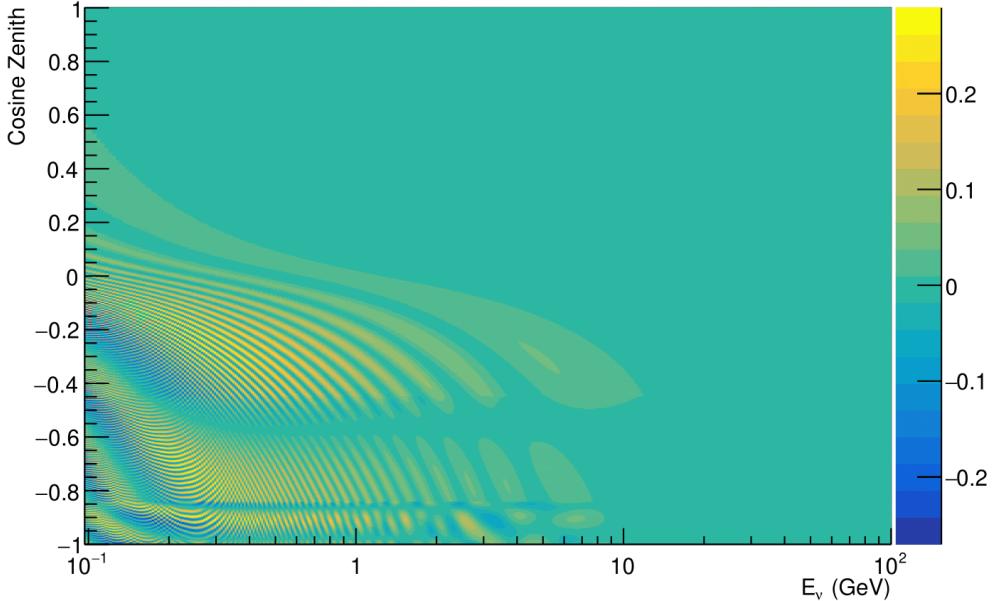


Figure 7.3: The effect of δ_{CP} for atmospheric neutrinos given in terms of the neutrino energy and zenith angle. This oscillogram compares the $P(\nu_\mu \rightarrow \nu_e)$ oscillation probability for a CP conserving ($\delta_{CP} = 0.0$) and a CP violating ($\delta_{CP} = -1.601$) value taken from the Asimov A parameter set. The other oscillation parameters assume the Asimov A oscillation parameter set given in Table 2.2.

of the muon-like samples, as illustrated in Figure 7.5. The value of Δm_{32}^2 laterally shifts the position of the oscillation dip (around $E_\nu \sim 0.6\text{GeV}$) in the $P(\nu_\mu \rightarrow \nu_\mu)$ oscillation probability. A variation of $\sin^2(\theta_{23})$ is predominantly observed as a vertical shift of the oscillation dip with second-order horizontal shifts being due to matter effects. The beam neutrinos have limited sensitivity to matter effects due to the relatively shorter baseline as well as the Earth's mantle being a relatively low-density material (as compared to the Earth's core). For some values of δ_{CP} , the degeneracy in the number of e-like events allows the mass hierarchy to be broken. This leads to a δ_{CP} -dependent mass hierarchy sensitivity which can be seen in Figure 7.6.

Whilst all oscillation channels should be included for completeness, the computational resources required to run a fit are limited and any reasonable approximations which reduce the number of oscillation probability calculations

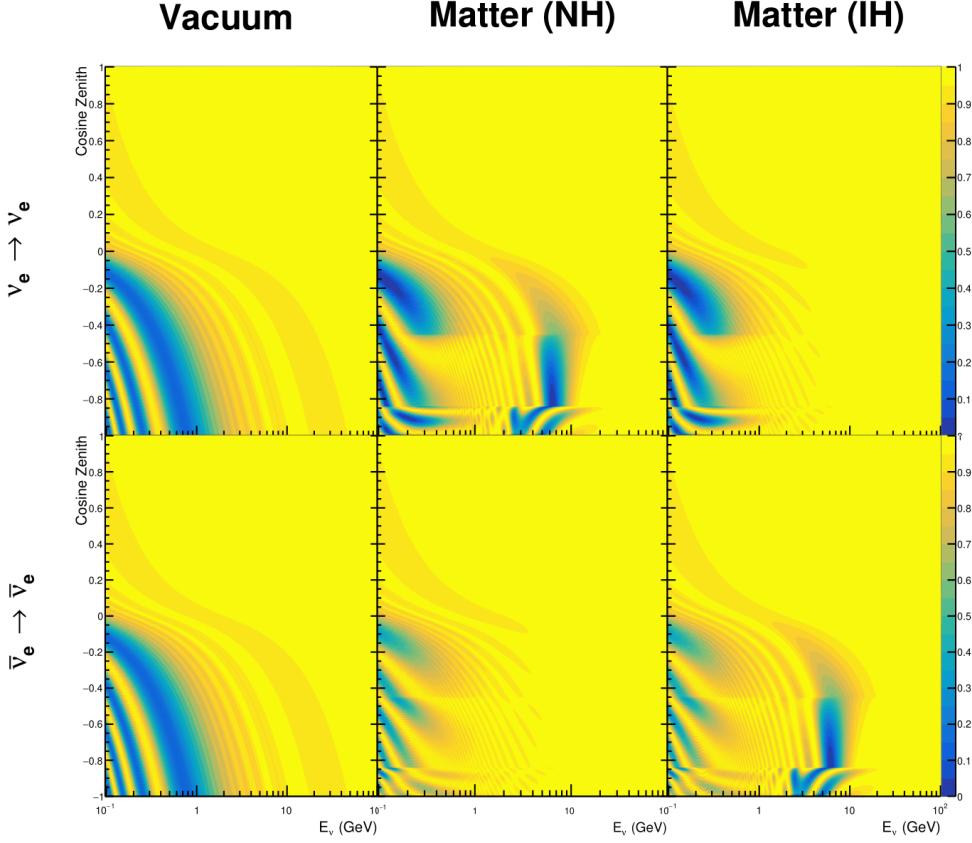


Figure 7.4: An illustration of the matter-induced effects on the oscillation probability, given as a function of neutrino energy and zenith angle. The top row of panels gives the $P(\nu_e \rightarrow \nu_e)$ oscillation probability and the bottom row illustrates the $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$ oscillation probability. The left column highlights the oscillation probability in a vacuum, whereas the middle and right column represents the oscillation probabilities when the four-layer fixed density PREM model is assumed. All oscillation probabilities assume the “Asimov A” set given in Table 2.2, but importantly, the right column sets an inverted mass hierarchy. The “matter resonance” effects at $E_\nu \sim 5\text{GeV}$ can be seen in the $P(\nu_e \rightarrow \nu_e)$ for normal mass hierarchy and $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$ for inverted hierarchy.

that need to be made should be applied. The $\nu_e \rightarrow \nu_{e,\mu,\tau}$ (and antineutrino equivalent) oscillations can be ignored for beam neutrinos as the $\nu_e/\bar{\nu}_e$ fluxes are approximately two orders of magnitude smaller than the corresponding $\nu_\mu/\bar{\nu}_\mu$ flux. Furthermore, as the peak neutrino energy of the beam is well below the threshold for charged current tau production ($E_\nu = 3.5\text{GeV}$ [58], only a small proportion of the neutrinos produced in the beam have the required energy. For the few neutrinos that have sufficient energy, the oscillation probability is very

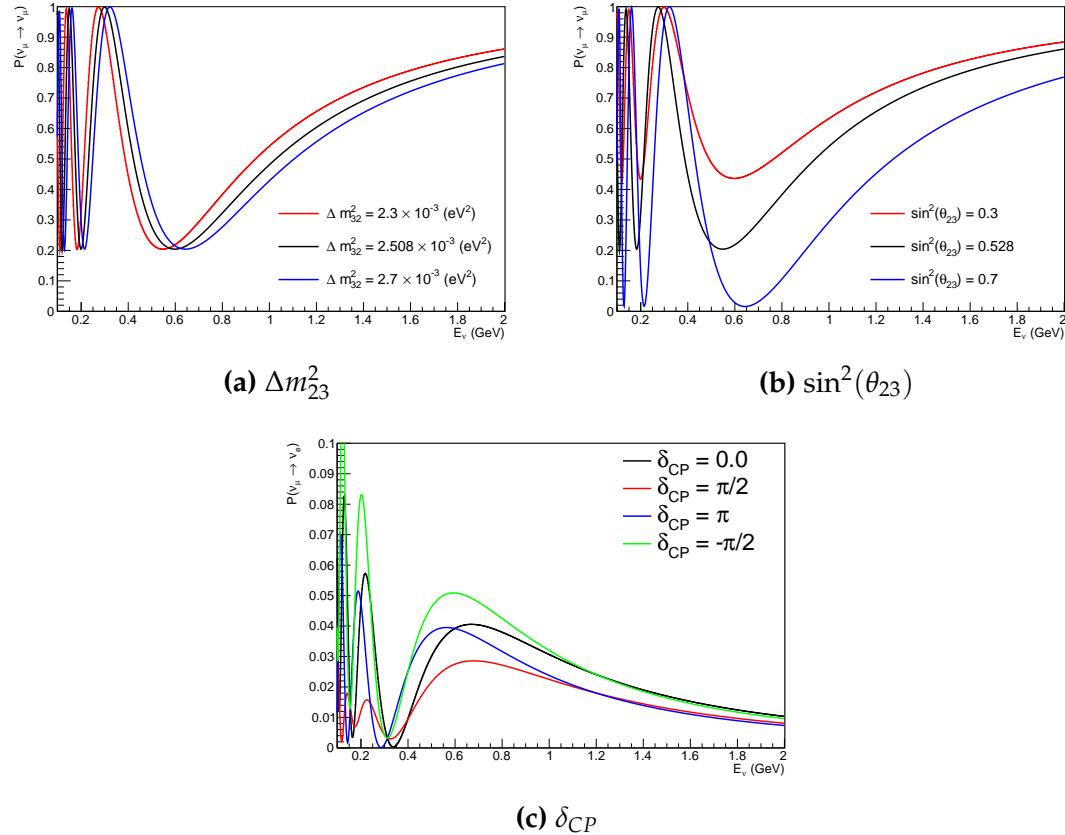


Figure 7.5: The oscillation probability for beam neutrino events given as a function of neutrino energy. All oscillation parameters assume the “Asimov A” set given in Table 2.2 unless otherwise stated. Each panel represents a change in one of the oscillation parameters whilst keeping the remaining parameters fixed.

2455 small due to the short baseline. Whilst these approximations have been made for
 2456 the beam neutrinos, the atmospheric flux of ν_e is of the same order of magnitude
 2457 as the ν_μ flux and the energy distribution of atmospheric neutrinos extends well
 2458 above the tau production threshold.

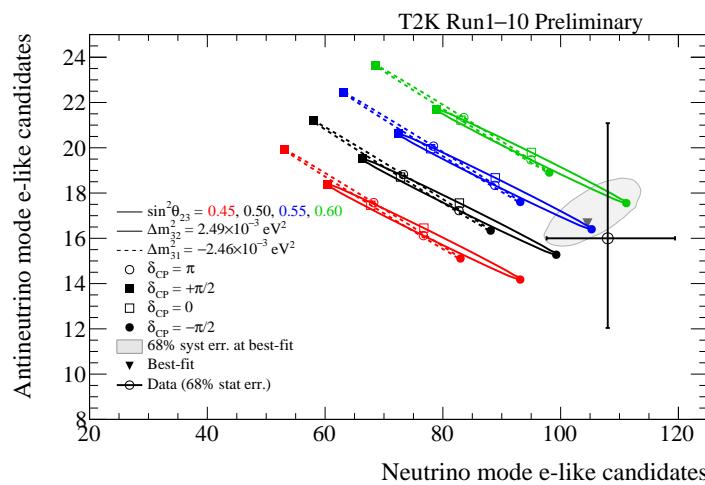


Figure 7.6: The number of electron-like events in the FHC and RHC operating mode of the beam, as a function of the oscillation probabilities. Both normal hierarchy (Solid) and inverse hierarchy (Dashed) values of Δm_{23}^2 are given.

2459 7.2 Treatment of Fast Oscillations

2460 As shown in Figure 7.7, atmospheric neutrino oscillations have a significantly
 2461 more complex structure for upgoing neutrinos with energy below 1GeV. This is
 2462 because the L/E dependence of the oscillation probability in this region induces
 2463 rapid variations for small changes in L or E . As discussed in section 7.1, this
 2464 is also the region in which atmospheric neutrinos have sensitivity to δ_{CP} . In
 2465 practice, the direction of the neutrino is inferred from the direction of the final
 2466 state particles traveling in the detector, which can be poor for low-energy neutrino
 2467 interactions. This creates a distinct difference from the beam neutrinos where
 2468 the position of the source is very precisely known.

2469 As a consequence of the unresolvable structure, an event rate consistent with
 2470 the averaged oscillation probability is observed in the subGeV upgoing region.
 2471 This creates a computational problem: A significantly large amount of Monte
 2472 Carlo statistics would be required to accurately predict the number of events
 2473 if Monte Carlo averaging was the only technique used. This section describes
 2474 the ‘sub-sampling’ approach developed for this analysis and compares it to the
 2475 methodology used within the SK-only analysis.

2476 The official SK-only analysis uses the osc3++ oscillation parameter fitter
 2477 [206]. To perform the fast oscillation averaging, it uses a ‘nearest-neighbour’
 2478 technique. For a given Monte Carlo neutrino event, the nearest twenty Monte
 2479 Carlo neighbours in reconstructed lepton momentum and zenith angle are
 2480 found and a distribution of their neutrino energies is built. The RMS, σ , of
 2481 this distribution is then used to compute an average oscillation probability for
 2482 the given neutrino Monte Carlo event.

2483 For the i^{th} event, the oscillation weight is calculated as

$$W_i = \frac{1}{5}P(E_i, \bar{L}_i) + \frac{1}{5}\sum_{\beta=-1, -0.5, 0.5, 1} P(E_i + \beta\sigma_i, L_\beta), \quad (7.3)$$

2484 where $P(E, L)$ is the oscillation probability calculation for neutrino energy
 2485 E and path length L and the two path lengths, \bar{L}_i and L_β are discussed below.

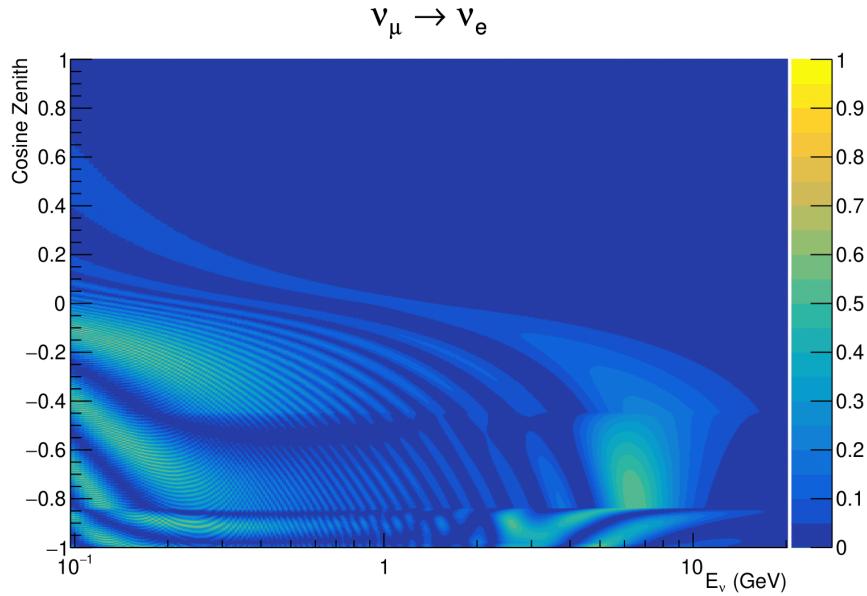


Figure 7.7: The oscillation probability $P(\nu_\mu \rightarrow \nu_e)$, given as a function of neutrino energy and zenith angle, which highlights an example of the “fast” oscillations in the sub-GeV upgoing region.

2486 All of the oscillation probability calculations are performed with a fixed zenith
2487 angle such that the same density profile is used.

2488 The uncertainty in the production height is controlled by using an “average”
2489 production height, \bar{L}_i , which represents the average path length computed using
2490 twenty production heights taken from the Honda flux model’s prediction [52].
2491 For a given event, the production heights are sampled in steps of 5% of their
2492 cumulative distribution function. L_β values are similarly calculated but instead
2493 use different combinations of four production heights,

$$\begin{aligned}
 L_{-1.0} &= \frac{1}{4}L(45, 50, 55, 60), \\
 L_{-0.5} &= \frac{1}{4}L(35, 40, 65, 70), \\
 L_{+0.5} &= \frac{1}{4}L(25, 30, 75, 68), \\
 L_{+1.0} &= \frac{1}{4}L(15, 20, 85, 89).
 \end{aligned} \tag{7.4}$$

2494 This averaging technique works because of the inference between the zenith
2495 angle and the reconstructed direction of final state particles in the detector. For

2496 low-energy neutrinos, where the resolution of the true neutrino direction is poor,
2497 σ_i will be large, resulting in significant averaging effects. Contrary to this, the
2498 inferred direction of high-energy neutrinos will be much closer to the true value,
2499 meaning that σ_i will be smaller, culminating in small averaging effects.

2500 In practice, these calculations are performed prior to the fit as only oscillation
2501 parameters at fixed points are considered. The MCMC technique used in this
2502 thesis requires oscillation probabilities to be evaluated at arbitrary parameter
2503 values, not known *a priori*. Calculating the five oscillation probabilities per
2504 event required by the SK technique is computationally infeasible, so a differencct
2505 averaging technique is used. However, the concept of the averaging technique
2506 can be taken from it.

2507 To perform a similar averaging as the SK analysis, a sub-sampling approach
2508 using binned oscillograms has been devised. A coarsly binned oscillogram is
2509 defined in $\cos(\theta_Z)$ and E_ν . For a given set of oscillation parameters, a single
2510 oscillation probability will be assigned to each coarse bin. This value will then
2511 apply to all Monte Carlo events which fall into that bin. To assign these oscillation
2512 probabilities, the probability is calculated at $N \times N$ points on a grid within a
2513 particular bin. This ensemble of oscillation probabilities is averaged to define
2514 the coarse bin's oscillation probability, assuming a flat prior in E_ν and $\cos(\theta_Z)$.
2515 Figure 7.8 illustrates the $N = 2$ example where the assigned value to a coarse
2516 bin is the average of the four fine bins which fall in that coarse bin. Whilst the
2517 coarse bin edges do not have to be linear on either axis, the sub-division of the
2518 fine bins is linear over the range of a coarse bin.

2519 The coarse binning is defined with 67×52 bins in true neutrino energy \times
2520 cosine zenith. It is picked to be identical to that provided in [208]. In general, the
2521 binning is logarithmically spaced in neutrino energy but has some hand-picked
2522 bin edges around the matter resonance to smoothly increased the bin density.
2523 This is to avoid smearing this region which can be well sampled by the Monte
2524 Carlo. The cosine zenith binning is approximately linearly spaced across the

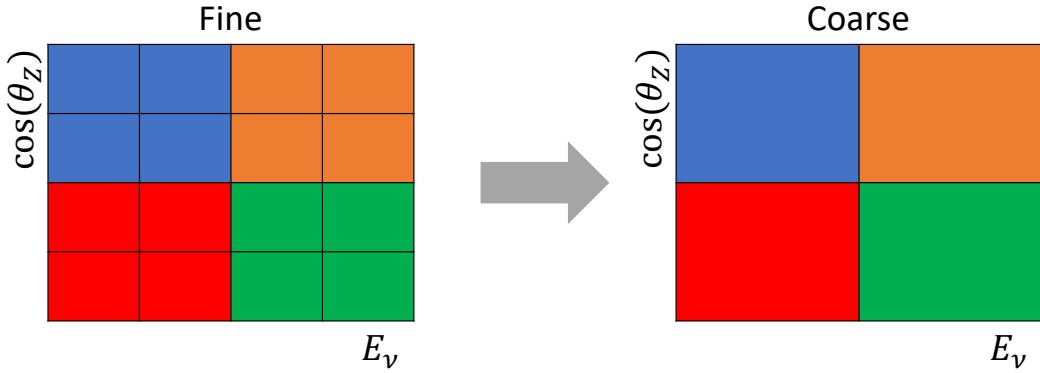


Figure 7.8: Illustration of the averaging procedure for $N = 2$. The oscillation probabilities calculated on the finer left binning are averaged to obtain the oscillation probabilities in the coarser right binning. These averaged oscillation probabilities with the coarser binning are then applied to each event during the fit.

allowable range but the values of layer transitions are hit precisely: -0.8376 (core-mantle) and -0.4464 (mantle/transition zone). Bins are spread further apart for downgoing events as this is a region unaffected by the fast oscillation wavelengths and reduces the total number of calculations required to perform the calculation.

The choice of N is justified based on two studies. Firstly, the variation of event rates of each sample is studied as a function of N . For a given set of oscillation parameters thrown from the PDG prior constraints (detailed in Table 2.1), the oscillation probabilities are calculated using a given value of N . Each sample is re-weighted and the event rate is stored. The value of N is scanned from 1, which corresponds to no averaging, to 19, which corresponds to the largest computationally viable subdivision binning. The event rate of each sample at large N is expected to converge to a stationary value due to the fine binning fully sampling the small-scale structure. Figure 7.9 illustrates this behaviour for the SubGeV_elike_0dcy sample for 9 different throws of the oscillation parameters.

Denoting the event rate for one sample for a given throw t at each N by λ_t^N , the average over all considered N values ($\bar{\lambda}_t = \frac{1}{24} \sum_{N=1}^{24} \lambda_t^N$) is computed. The variance in the event rate at each N is then calculated as

$$\text{Var}[\lambda^N] = \frac{1}{N_{\text{throws}}} \sum_{t=1}^{N_{\text{throws}}} (\lambda_t^N - \bar{\lambda}_t)^2 - \left[\frac{1}{N_{\text{throws}}} \sum_{t=1}^{N_{\text{throws}}} (\lambda_t^N - \bar{\lambda}_t) \right]^2. \quad (7.5)$$

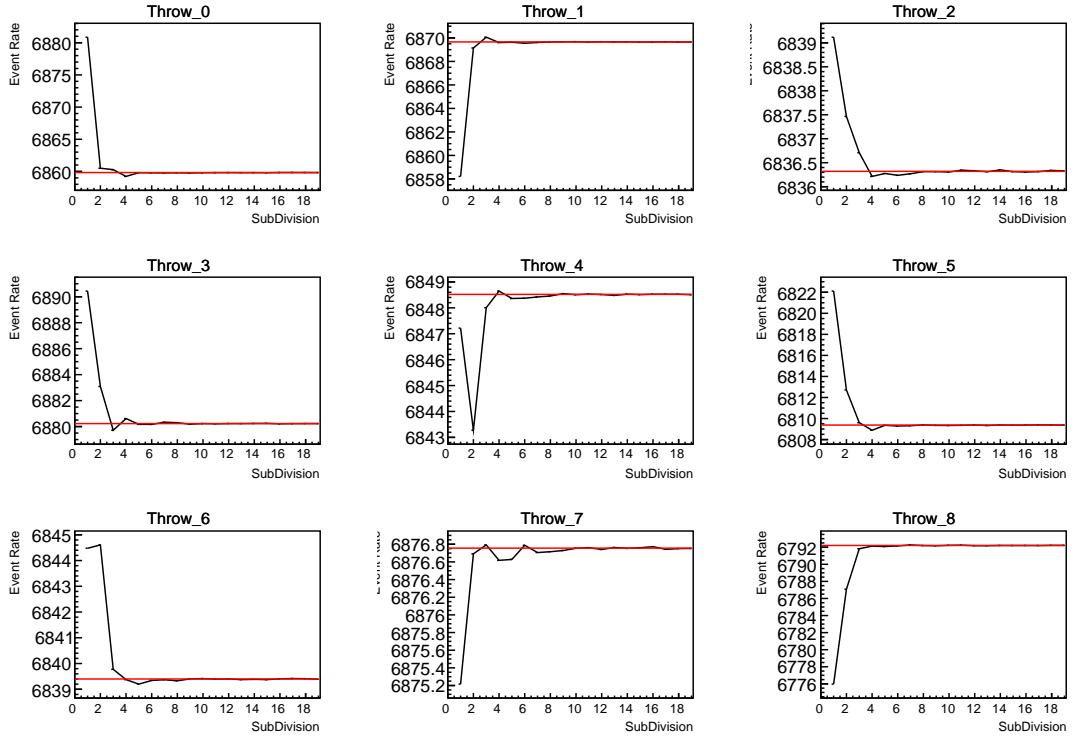


Figure 7.9: Event rate of the SubGeV_elike_0dcy sample as a function of the number of sub-divisions per coarse bin. Each subplot represents the event rate of the sample at a different oscillation parameter set thrown from the PDG priors detailed in Table 2.1. The red line in each subplot represents the mean of the event rate over the different values of sub-divisions for that particular oscillation parameter throw.

2542 In practice the following procedure is undertaken. For a particular throw,
 2543 the difference between the event rate at a particular choice of N and the mean
 2544 of the distribution is calculated. This is illustrated in Figure 7.10. This value
 2545 is then calculated for all the 2000 throws, generating a distribution of $\lambda_t^N - \bar{\lambda}_t$.
 2546 This is repeated for each of the values of N considered within this study. The
 2547 distributions of this value, for $N = \{1, 5\}$, are given in Figure 7.11. As expected,
 2548 the distribution gets narrower and tends towards zero for the higher values of N .

2549 The aim of the study is to find the lowest value of N such that this variance
 2550 is below 0.001. This utilises the width of the distributions given in Figure 7.11.
 2551 This is the typical threshold used by T2K fitters to validate systematic imple-
 2552 mentation so has been set as the same criteria. The results of this study for
 2553 each atmospheric sample used within this thesis are illustrated in Figure 7.12 for
 2554 2000 throws of the oscillation parameters. As can be seen, the variance is below

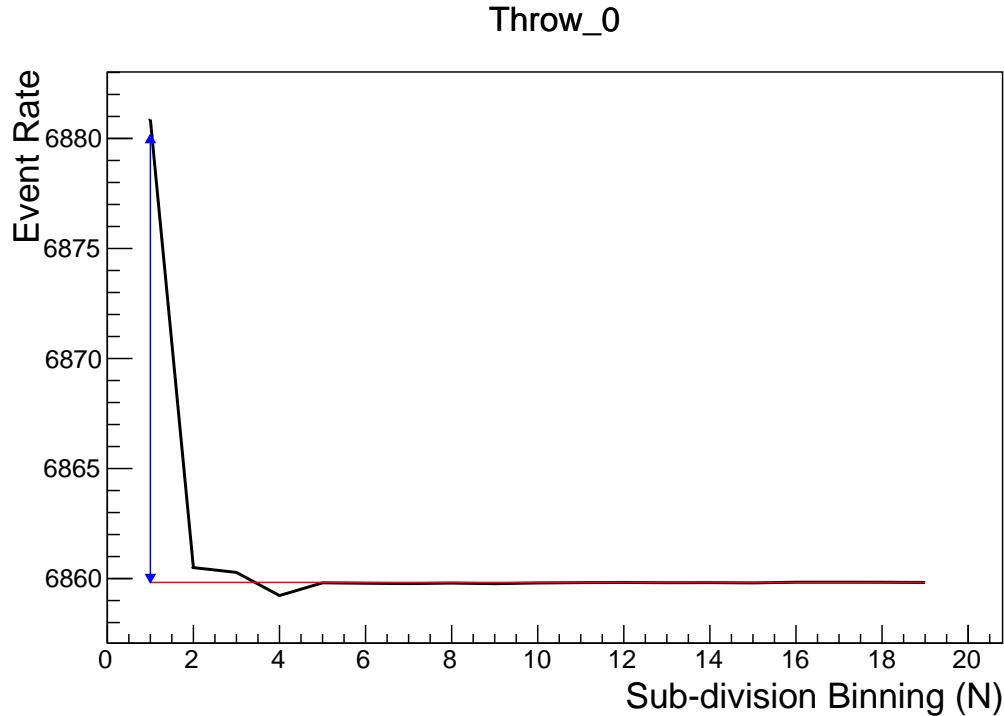


Figure 7.10: Event rate of the SubGeV_elike_0dcy sample, for a particular oscillation parameter throw, as a function of the number of sub-divisions, N , per coarse bin. The difference between the mean event rate (red), $\bar{\lambda}$, and the event rate at $N = 1$, $\lambda^{N=1}$ is defined as $\lambda^N - \bar{\lambda}$ and illustrated by the blue arrow.

2555 the threshold at $N = 10$, and is driven primarily by the SubGeV_mulike_1dcy
 2556 and SubGeV_elike_0dcy samples.

2557 The second study to determine the value of N is as follows. The likelihood
 2558 for each sample is computed against an Asimov data set created with Asimov
 2559 A oscillation parameters (Table 2.2). Following Equation 7.5, the variance of
 2560 the log-likelihood over all considered N is computed. The results are shown
 2561 in Figure 7.13.

2562 A choice of $N = 10$ sub-divisions per coarse bin has a variance in both
 2563 event rate and log-likelihood residuals less than the required threshold of 0.001.
 2564 The largest value of the likelihood variance is of order 10^{-7} , corresponding to
 2565 an error on the log-likelihood of about 3×10^{-4} which is small enough to be
 2566 negligible for the oscillation analysis.

2567 Figure 7.14 illustrates the effect of the smearing using $N = 10$. The fast oscil-

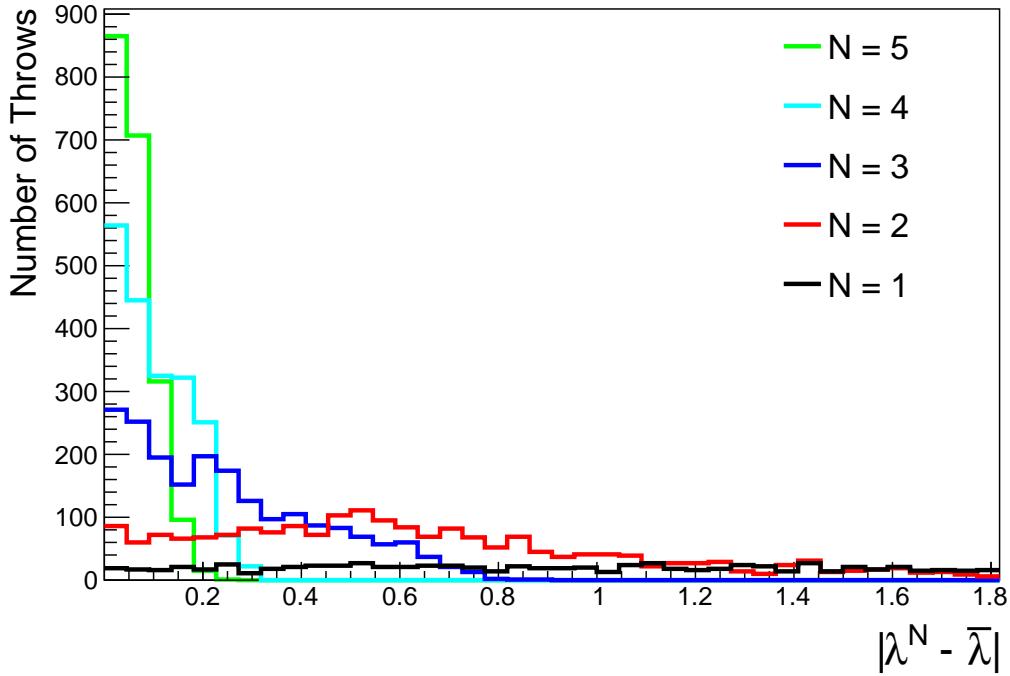


Figure 7.11: The distribution of $\lambda^N - \bar{\lambda}$ for various values of N . As expected, the distribution gets narrower for larger values of N .

2568 lations in the sub-GeV upgoing region have been replaced with a normalisation
 2569 effect whilst the large matter resonance structure remains.

2570 7.3 Calculation Engine

2571 As previously discussed in section 7.2, the calculation of oscillation probabilities
 2572 is performed at run-time. Consequently, the time per calculation is crucial for fit
 2573 performance. The initial fitting framework used for this analysis was developed
 2574 with ProbGPU [209]. This is a GPU-only implementation of the prob3 engine
 2575 [210]. It is primarily designed for neutrino propagation in a beam experiment
 2576 (single layer of constant density) with the atmospheric propagation code not
 2577 being used prior to the analysis in this thesis.

2578 Another engine, CUDAProb3 [211], has been interfaced with the fitting frame-
 2579 work used in this analysis. It has been specifically optimised for atmospheric
 2580 neutrino oscillation calculation so does not contain the code to replace the

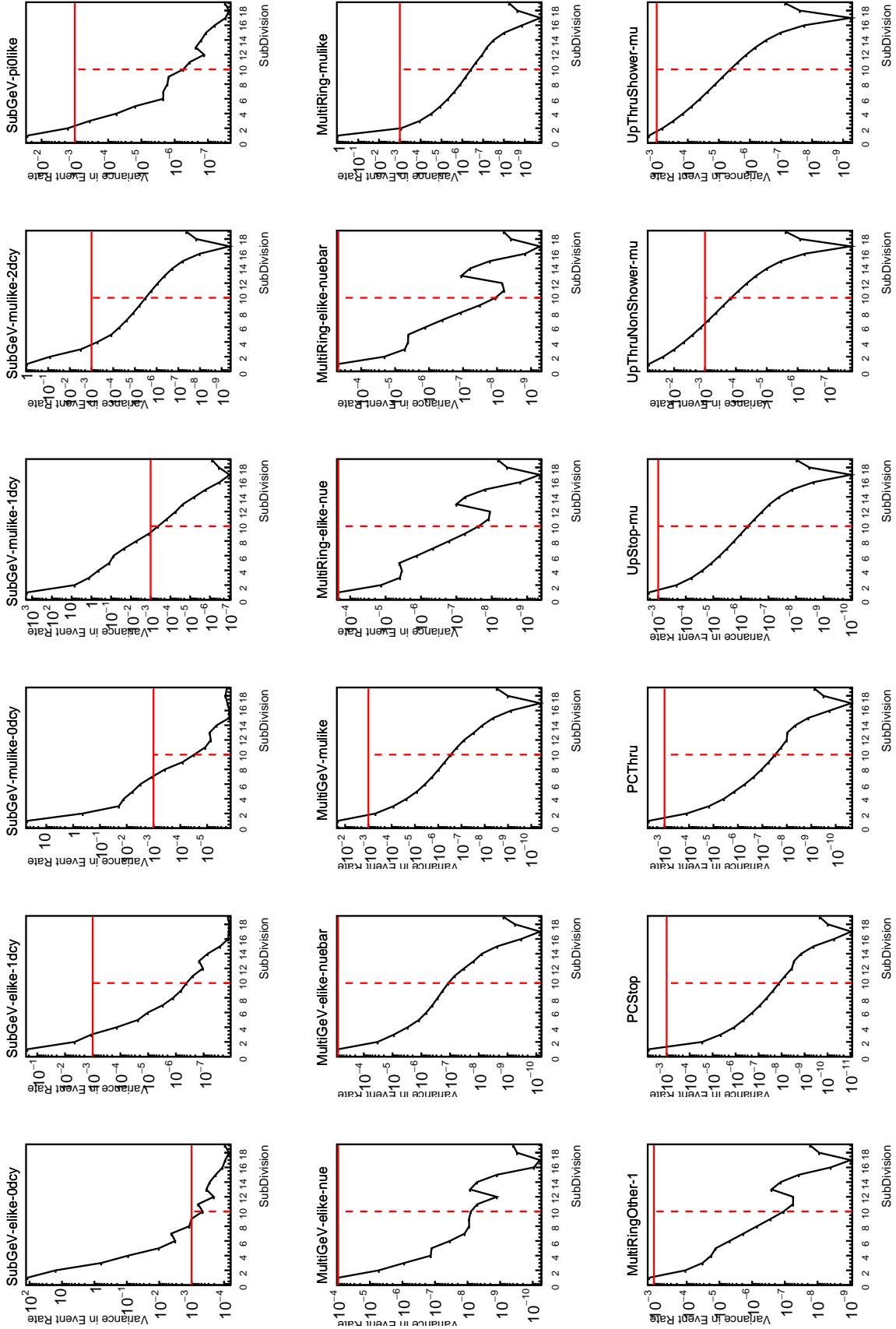


Figure 7.12: Variance of event rate for each atmospheric sample as a function of the number of sub-divisions per coarse bin. The solid red line indicates the 0.1% threshold and the dashed red line indicates the variance at a sub-division $N = 10$.

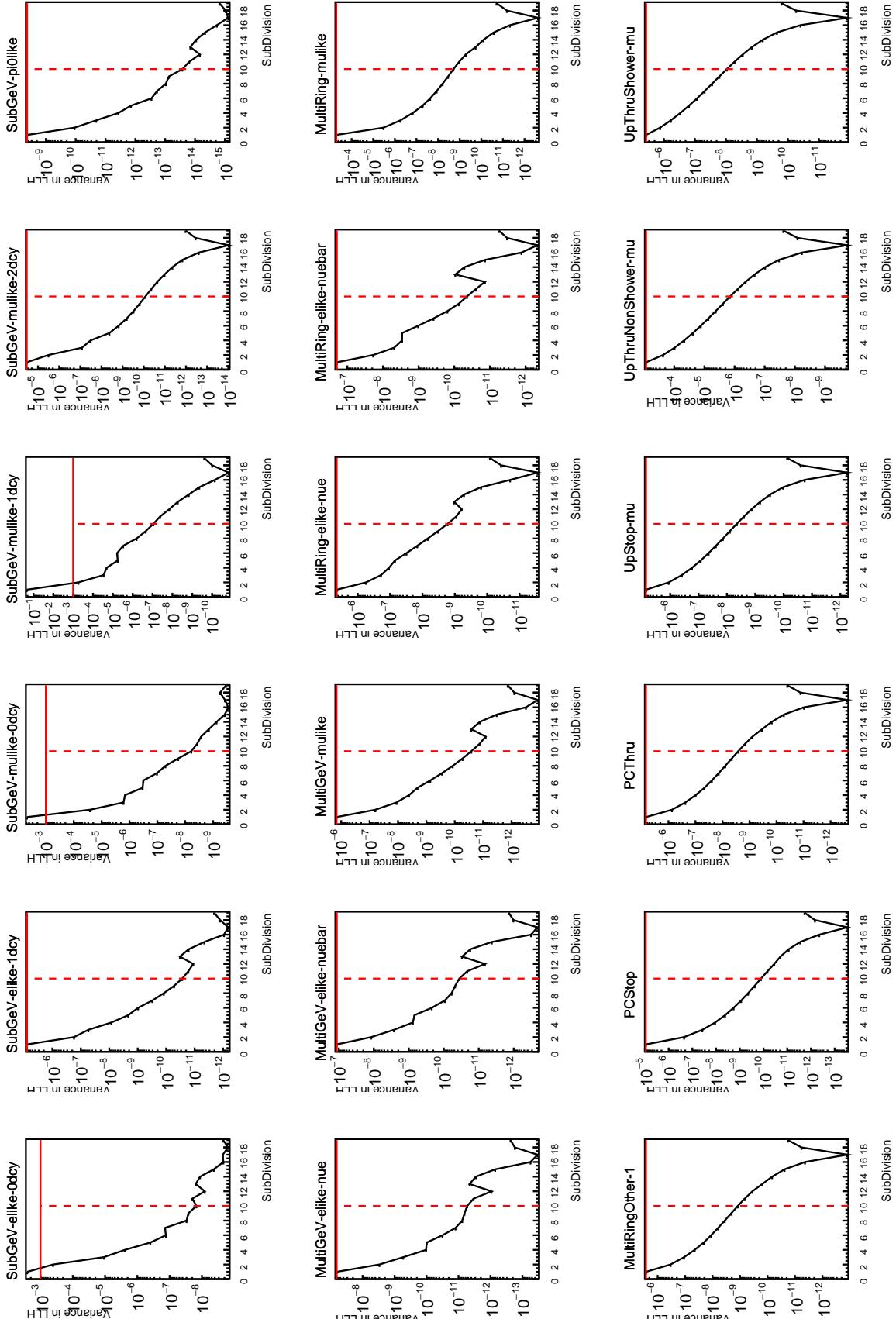


Figure 7.13: Variance of sample likelihood, when compared to ‘Asimov data’ set at Asimov A, for each atmospheric sample as a function of the number of sub-divisions per coarse bin. The solid red line indicates the 0.1% threshold and the dashed red line is a graphical indication of the variance at a sub-division $N = 10$.

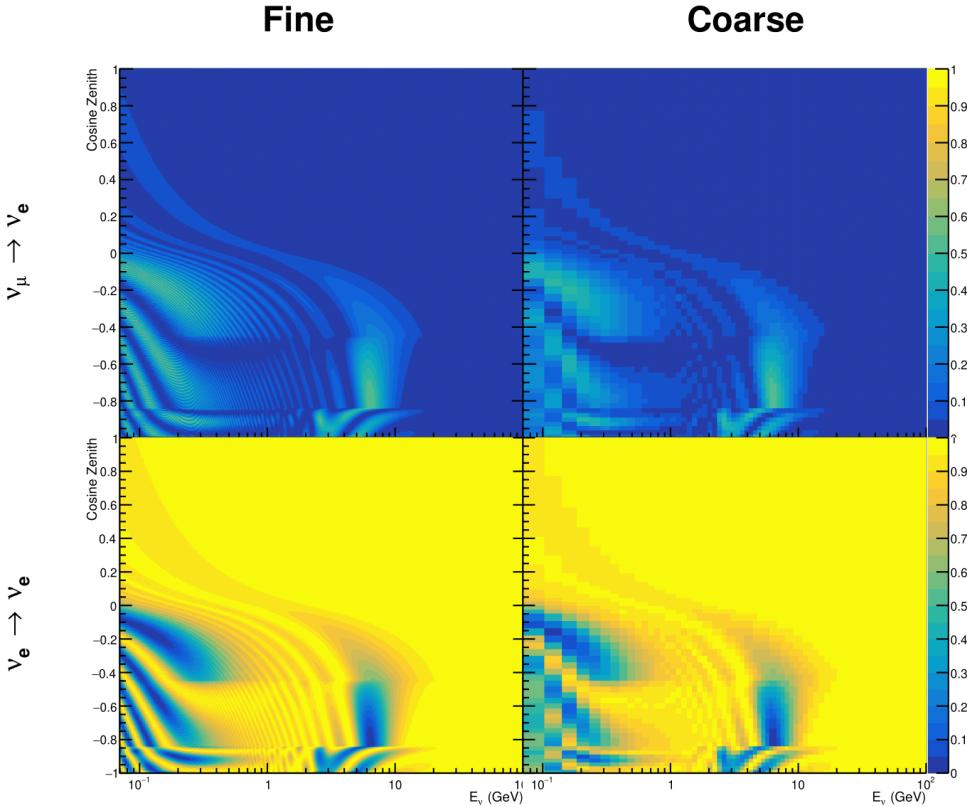


Figure 7.14: The oscillation probability, $P(\nu_\mu \rightarrow \nu_e)$ (top row) and $P(\nu_e \rightarrow \nu_e)$ (bottom row), given as a function of neutrino energy and zenith angle. The left column gives the “fine” binning used to calculate the oscillation probabilities and the right column illustrates the “coarse” binning used to reweight the Monte Carlo events. The fine binning choice is given with $N = 10$, which was determined to be below the threshold from Figure 7.12 and Figure 7.13.

beam oscillation calculation. The engine utilises object-orientated techniques as compared to the functional implementation of ProbGPU. This allows the energy and cosine zenith arrays to be kept on GPU memory, rather than having to load these arrays onto GPU memory for each calculation. Reducing the memory transfer between CPU and GPU significantly reduces the time required for calculation. This can be seen in Figure 7.15, where the GPU implementation of CUDAProb3 is approximately three times faster than the ProbGPU engine.

Another significant advantage of CUDAProb3 is that it contains a CPU multithreaded implementation which is not possible with the ProbGPU or prob3

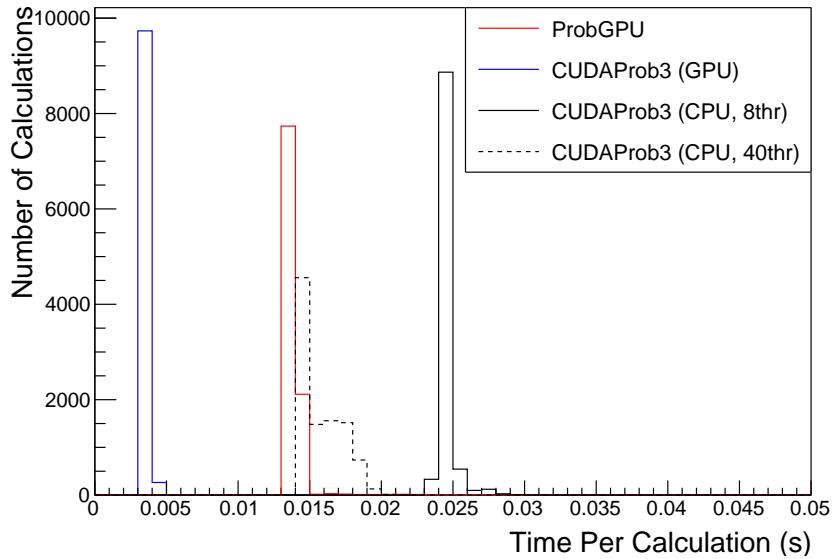


Figure 7.15: The calculation time taken to both calculate the oscillation probabilities and fill the “coarse” oscillograms, following the technique given in section 7.2, for the CUDAProb3 and ProbGPU (Red) calculation engines. CUDAProb3 has both a GPU (Blue) and CPU (Black) implementation, where the CPU implementation is multithreaded. Therefore, 8-threads (solid) and 40-threads (dashed) configurations have been tested. Prob3, which is a CPU single-thread implementation has a mean step time of 1.142s.

2590 engines. This eliminates the requirement for GPU resources when submitting
 2591 jobs to batch systems. As illustrated in Figure 7.15, the calculation speed depends
 2592 on the number of available threads. Using 8 threads (which is typical of the
 2593 batch systems being used) is approximately twice as slow as the ProbGPU engine
 2594 implementation, but would allow the fitting framework to be run on many more
 2595 resources. This fact is utilised for any SK-only fits but GPU resources are required
 2596 for any fits which include beam samples due to the ProbGPU requirement. Based
 2597 on the benefits shown by the implementation in this section, efforts are being
 2598 placed into including linear propagation for beam neutrino propagation into
 2599 the engine [212].

2600 7.4 Matter Density Profile

2601 For an experiment observing neutrinos propagating through the Earth, a model
 2602 of the Earth’s density profile is required. The model used within this analysis is

based on the Preliminary Reference Earth Model (PREM) [207], as illustrated in Figure 7.1. Table 7.1 documents the density and radii of the layers used within the constant density approximaton used by the SK-only analysis [206]. The density measurements provided in the PREM model are provided in terms of mass density, whereas neutrino oscillations are sensitive to the electron number density. This value can be computed as the product of the chemical composition, or the Z/A value, and the mass density of each layer. Currently, the only way to measure the chemical composition value for layers close to the Earth's core is through neutrino oscillations. The chemical composition of the upper layers of the Earth's Mantle and the Transition zone is well known due to it being predominantly pyrolite which has a chemical composition value of 0.496 [213]. The chemical composition dial for the core layers is set to a value of 0.468, as calculated in [214]. As this value is lesss well known, it is assigned a Gaussian error with a standard deviation equivalent to the difference in chemical composition in core and mantle layers. Figure 7.16 illustrates the effect of moving from the $Z/A = 0.5$ method which is used in the official SK-only analysis to these more precise values.

The beam oscillation probability in this thesis uses a baseline of 295km, density $2.6\text{g}/\text{cm}^3$, and chemical composition 0.5 as is done by the official T2K-only analysis [215].

For a neutrino with given E_ν , $\cos(\theta_Z)$, the oscillation probability calculation engine must be passed a list of the matter regions that the neutrino traversed, with the path length and fixed density in each region. However, a neutrino passing through the earth experiences a range of radii, and thus a range of densities, in each region. In the SK-only analysis, the earth density model used is piecewise-constant, thereby ignoring this effect. For this thesis, the density values for the calculation engine are found by averaging the earth density along the neutrino's path,

$$\langle \rho \rangle_i = \frac{1}{t_{i+1} - t_i} \int_{t_i}^{t_{i+1}} \rho(t) dt \quad (7.6)$$

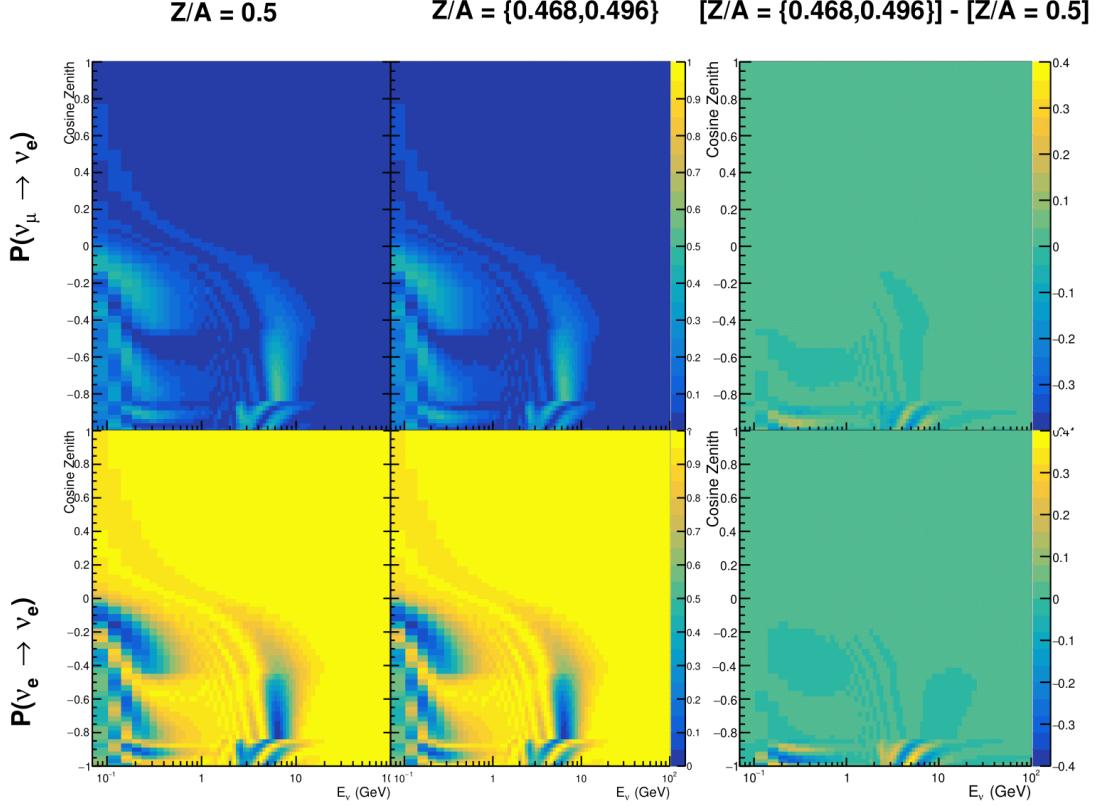


Figure 7.16: The oscillation probability, $P(\nu_\mu \rightarrow \nu_e)$ (top row) and $P(\nu_e \rightarrow \nu_e)$ (bottom row), given as a function of neutrino energy and zenith angle. The left column gives probabilities where the constant $Z/A = 0.5$ approximation which is used in the official SK-only analysis. The middle column gives the probabilities where $Z/A = [0.468, 0.498]$ values are used, as given in Table 7.1. The right column illustrates the difference in oscillation probability between the two different techniques.

where t_i are the intersection points between each layer and t is the path length of the trajectory across the layer. This leads to an improved approximation. For this averaging, the simplification of the PREM model developed in [216] is used. The layers of the prem model are combined into four to reduce calculation time, with a quadratic fit to each section. This fit was not performed by the author of the thesis and is documented in [208]. The coefficients of the quadratic fit to each layer are given in Table 7.2 with the final distribution illustrated in Figure 7.17. The quadratic approximation is clearly much closer to the PREM model as compared to the constant density approximation.

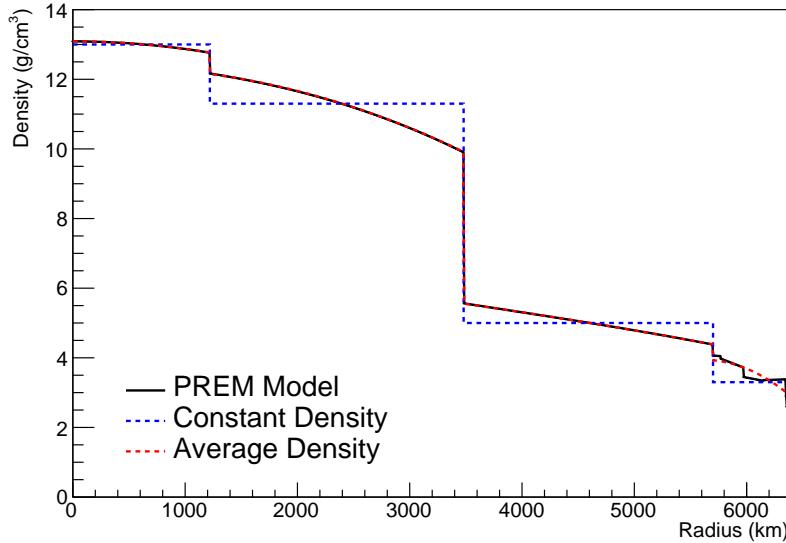


Figure 7.17: The density of the Earth given as a function of the radius, as given by the PREM model (Black), the constant density four-layer approximation (Blue), as used in the official SK-only analysis, and the quadratic approximation of the PREM model (Red).

Layer	Outer Radius [km]	Density [g/cm ³]
Inner Core	1220	$13.09 - 8.84x^2$
Outer Core	3480	$12.31 + 1.09x - 10.02x^2$
Lower Mantle	5701	$6.78 - 1.56x - 1.25x^2$
Transition Zone	6371	$-50.42 + 123.33x - 69.95x^2$

Table 7.2: The quadratic polynomial fits to the PREM model for four assumed layers of the PREM model. The fit to calculate the coefficients is given in [208], where $x = R / R_{\text{Earth}}$.

The effect of using the quadratic density per $\cos(\theta_Z)$ model is highlighted in Figure 7.18. The slight discontinuity in the oscillation probability around $\cos(\theta_Z) \sim -0.45$ in the fixed density model, which is due to the transition to mantle layer boundary, has been reduced. This is expected as the difference in density across this boundary is significantly smaller in the quadratic density model as compared to the constant density model. Whilst the difference in density across the other layer transitions is reduced, there is still a significant difference. This means the discontinuities in the oscillation probabilities remain but are significantly reduced. However, as the quadratic density approximation matches the PREM model well in this region, these discontinuities are due to the Earth model rather than an artifact of the oscillation calculation.

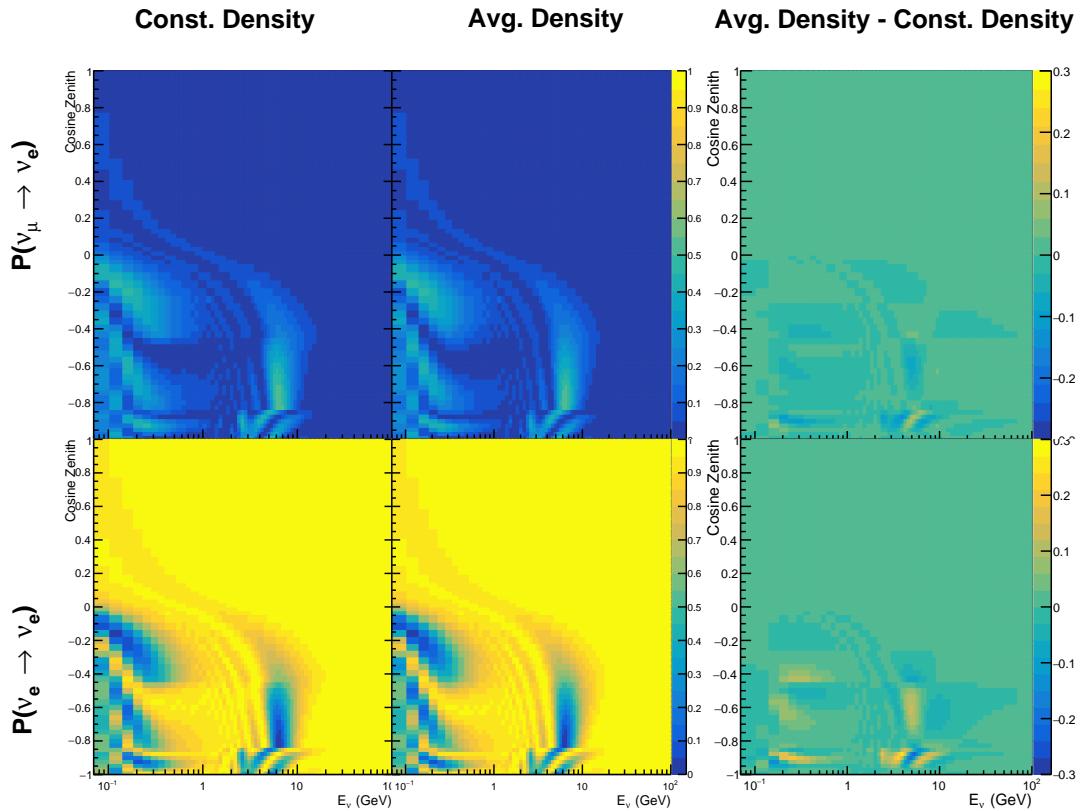


Figure 7.18: The oscillation probability, $P(\nu_\mu \rightarrow \nu_e)$ (top row) and $P(\nu_e \rightarrow \nu_e)$ (bottom row), given as a function of neutrino energy and zenith angle. The left column gives probabilities where the four-layer constant density approximation is used. The middle column gives the probabilities where the density is integrated over the trajectory, using the quadratic PREM approximation, for each $\cos(\theta_Z)$ is used. The right column illustrates the difference in oscillation probability between the two different techniques.

2651 7.5 Production Height Averaging

2652 As discussed in section 7.1, the height at which the cosmic ray flux interacts
 2653 in the atmosphere is not known on an event-by-event basis. The production
 2654 height can vary from the Earth’s surface to $\sim 50\text{km}$ above that. The SK-only
 2655 analysis methodology (described in section 7.2) for including the uncertainty
 2656 on the production height is to include variations from the Honda model when
 2657 pre-calculating the oscillation probabilities prior to the fit. This technique is not
 2658 possible for this analysis which uses continuous oscillation parameters that can
 2659 not be known prior to the fit. Consequently, an analytical averaging technique
 2660 was developed in [208]. The author of this thesis was not responsible for the
 2661 derivation of the technique but has performed the implementation and validation
 2662 of the technique for this analysis alone.

2663 Using the 20 production heights per Monte Carlo neutrino event, provided
 2664 as 5% percentiles from the Honda flux model, a production height distribution
 2665 $p_j(h|E_\nu, \cos \theta_Z)$ is built for each neutrino flavour $j = \nu_e, \bar{\nu}_e, \nu_\mu, \bar{\nu}_\mu$. In practice, a
 2666 histogram is filled with 20 evenly spaced bins in production height h between
 2667 0 and 50km. The neutrino energy and cosine zenith binning of the histogram
 2668 is the same as that provided in section 7.2. The average production height,
 2669 $\bar{h} = \int dh \frac{1}{4} \sum_j p_j(h|E_\nu, \cos(\theta_Z))$, is calculated. The production height binning
 2670 of this histogram is then translated into $\delta t(h) = t(\bar{h}) - t(h)$, where $t(h)$ is the
 2671 distance travelled along the trajectory.

2672 For the i^{th} traversed layer, the transition amplitude, $D_i(t_{i+1}, t_i)$, is computed.
 2673 The time-ordered product of these is then used as the overall transition amplitude
 2674 via

$$A(t_{n+1}, t_0) = D_n(t_{n+1}, t_n) \dots D_1(t_2, t_1) D_0(t_1, t_0), \quad (7.7)$$

2675 where,

$$\begin{aligned} D_n(t_{n+1}, t_n) &= \exp[-iH_n(t_{n+1} - t_n)] \\ &= \sum_{k=1}^3 C_k \exp[ia_k(t_{n+1} - t_n)] \end{aligned} \quad (7.8)$$

2676 is expressed as a diagonalised time-dependent solution to the Schrodinger
 2677 equation. The 0th layer is the propagation through the atmosphere and is the
 2678 only term that depends on the production height. Using the substitution $t_0 =$
 2679 $t(\bar{h}) - \delta t(h)$, it can be shown that

$$D_0(t_1, t_0) = D_0(t_1, \bar{h})D_0(\delta t). \quad (7.9)$$

2680 Thus Equation 7.7 becomes

$$\begin{aligned} A(t_{n+1}, t_0) &= D_n(t_{n+1}, t_n) \dots D_1(t_2, t_1) D_0(t_1, \bar{h}) D(\delta t) \\ &= A(t_{n+1}, \bar{h}) \sum_{k=1}^3 C_k \exp[ia_k \delta t], \\ &= \sum_{k=1}^3 B_k \exp[ia_k \delta t]. \end{aligned} \quad (7.10)$$

2681 The oscillation probability averaged over production height is then calculated
 2682 as

$$\begin{aligned} \bar{P}(\nu_j \rightarrow \nu_i) &= \int d(\delta t) p_j(\delta t | E_\nu, \cos \theta_Z) P(\nu_j \rightarrow \nu_i) \\ &= \int d(\delta t) p_j(\delta t | E_\nu, \cos \theta_Z) A(t_{n+1}, t_0) A^*(t_{n+1}, t_0) \\ &= \sum_{km} (B_k)_{ij} (B_m)^*_{ij} \int d(\delta t) p_j(\delta t | E_\nu, \cos \theta_Z) \exp[i(a_k - a_m)\delta t] \end{aligned} \quad (7.11)$$

2683 In practice, implementation in CUDAProb3 [211] is relatively straightforward
 2684 as the majority of these terms are already calculated in the standard oscillation
 2685 calculation. Figure 7.19 illustrates the results of the production height averaging.
 2686 As expected, the main effect is observed in the low-energy downward-going
 2687 and horizontal-going events. Upward-going events have to travel the radius
 2688 of the Earth, $R_E = 6371\text{km}$, where the production height uncertainty is a small
 2689 fraction of the total path length.

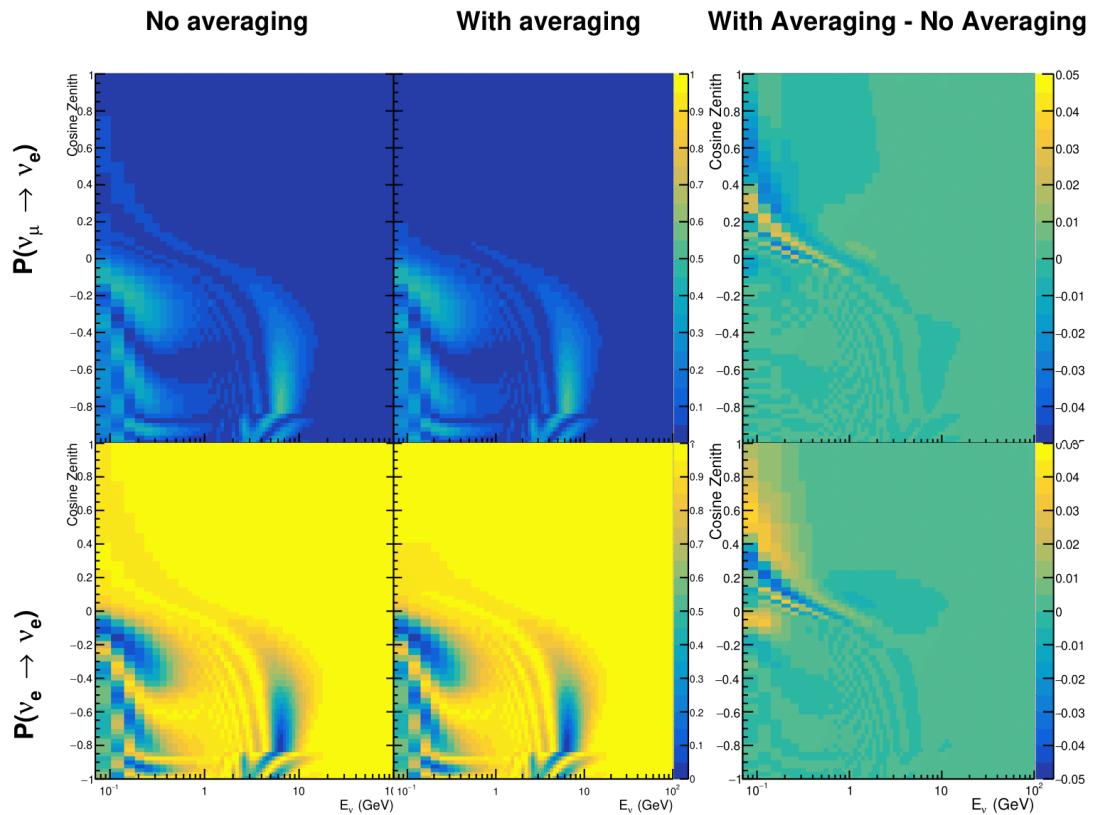


Figure 7.19: The oscillation probability, $P(\nu_\mu \rightarrow \nu_e)$ (top row) and $P(\nu_e \rightarrow \nu_e)$ (bottom row), given as a function of neutrino energy and zenith angle. The left column gives probabilities where a fixed production height of 25km is used. The middle column gives the probabilities where the production height is analytically averaged. The right column illustrates the difference in oscillation probability between the two different techniques.

8

2690

2691

Oscillation Analysis

2692 Using the samples and systematics defined in chapter 6, this chapter documents
2693 a simultaneous beam and atmospheric oscillation analysis from the T2K and SK
2694 experiments. The MaCh3 Bayesian MCMC framework introduced in chapter 4
2695 is used for all studies performed within this thesis.

2696 The MaCh3 framework used throughout this thesis has been validated through
2697 many tests. The code which handles the beam far detector samples was developed
2698 by the author and validated by comparison to the 2020 T2K analysis [74]. The
2699 sample event rates and likelihood evaluations of beam samples generated by
2700 the framework used within this thesis were compared to those from the T2K
2701 analysis by the author of this thesis. Variations of the sample predictions were
2702 compared at $\pm 1\sigma$ and $\pm 3\sigma$ and good agreement was found in all cases. A similar
2703 study, led by Dr. C. Wret was used to validate the near detector portion of
2704 the code [201]. The implementation of the atmospheric samples within MaCh3
2705 was completed and cross-checked by the author of this thesis against the P-
2706 Theta framework (Introduced in section 3.2). Both fitters are provided with
2707 the same inputs and can therefore cross-validate each other. These validations
2708 compared the event rate and likelihood calculation. Documentation of all the
2709 above validations can be found in [201].

2710 8.1 Monte Carlo Prediction

2711 Using the three sets of dial values (generated, pre-BANFF, and post-BANFF tunes)
 2712 defined in subsection 6.4.3, the predicted event rates for each sample are defined
 2713 in Table 8.1. The oscillated (AsimovA defined in Table 2.2) and un-oscillated
 2714 event rates are calculated for each tune.

Sample	Total Predicted Events					
	Generated		Pre-fit		Post-fit	
	Osc	UnOsc	Osc	UnOsc	Osc	UnOsc
SubGeV- <i>elike</i> -0dcy	7121.0	7102.6	6556.8	6540.0	7035.2	7015.7
SubGeV- <i>elike</i> -1dcy	704.8	725.5	693.8	712.8	565.7	586.0
SubGeV- <i>mulike</i> -0dcy	1176.5	1737.2	1078.6	1588.1	1182.7	1757.1
SubGeV- <i>mulike</i> -1dcy	5850.7	8978.1	5351.7	8205.1	5867.0	9009.9
SubGeV- <i>mulike</i> -2dcy	446.9	655.2	441.6	647.7	345.9	505.6
SubGeV- <i>pi0like</i>	1438.8	1445.4	1454.9	1461.1	1131.1	1136.2
MultiGeV- <i>elike</i> -nue	201.4	195.6	201.1	195.3	202.6	196.7
MultiGeV- <i>elike</i> -nuebar	1141.5	1118.3	1060.7	1039.5	1118.5	1095.7
MultiGeV- <i>mulike</i>	1036.7	1435.8	963.1	1334.1	1015.2	1405.9
MultiRing- <i>elike</i> -nue	1025.1	982.2	1026.8	984.3	1029.8	986.4
MultiRing- <i>elike</i> -nuebar	1014.8	984.5	991.0	962.0	1008.9	978.5
MultiRing- <i>mulike</i>	2510.0	3474.4	2475.6	3425.8	2514.6	3480.4
MultiRingOther-1	1204.5	1279.1	1205.8	1280.3	1207.4	1281.0
PCStop	349.2	459.2	338.4	444.7	346.8	456.1
PCThru	1692.8	2192.5	1661.5	2149.8	1689.2	2187.8
UpStop-mu	751.2	1295.0	739.7	1271.6	750.4	1293.0
UpThruNonShower-mu	2584.4	3031.6	2577.9	3019.4	2586.8	3034.0
UpThruShower-mu	473.0	488.6	473.2	488.7	473.8	489.4
FHC1Rmu	328.0	1409.2	301.1	1274.7	345.1	1568.0
RHC1Rmu	133.0	432.3	122.7	396.2	135.0	443.9
FHC1Re	84.6	19.2	77.4	18.2	93.7	19.7
RHC1Re	15.7	6.4	14.6	6.1	15.9	6.3
FHC1Re1de	10.5	3.2	10.3	3.1	8.8	2.9

Table 8.1: The Monte Carlo predicted event rate of each far detector sample used within this analysis. Three model parameter tunes are considered, as defined in subsection 6.4.3. Un-oscillated and oscillated predictions are given, where the oscillated predictions assume Asimov A oscillation parameters provided in Table 2.2.

2715 Generally, the samples which target CCQE interaction modes observe a
 2716 decrease in prediction when using the pre-fit dial values. This is in accordance
 2717 with the Monte Carlo being produced assumed $M_A^{QE} = 1.21\text{GeV}$ [160] whilst

2718 the pre-fit dial value is set to $M_A^{QE} = 1.03\text{GeV}$ as suggested by [199]. Further-
 2719 more, the predicted event rates of samples that target CCRES interaction modes
 2720 are significantly reduced when considering the post-BANFF fit. This follows
 2721 the observations in subsection 6.4.3. The strength of the accelerator neutrino
 2722 experiment can be seen in the remarkable difference between the oscillated and
 2723 unoscillated predictions in the FHC1Rmu and RHC1Rmu samples. There is a very
 2724 clear decrease in the expected event rate between the oscillated and un-oscillated
 2725 predictions which is not as obvious as in the atmospheric samples. This is due to
 2726 the fact that the beam energy is tuned to the maximum disappearance probability
 2727 which is not the case for the naturally generated atmospheric neutrinos.

2728 8.2 Likelihood Calculation

2729 This analysis performs a joint oscillation parameter fit of the ND280 beam
 2730 samples, the T2K far detector beam samples, and the SK atmospheric samples.

2731 Once the Monte Carlo predictions of each beam and atmospheric sample
 2732 have been built (see chapter 6) a likelihood needs to be constructed. This is done
 2733 by comparing the binned Monte Carlo prediction to binned data. The Monte
 2734 Carlo prediction is calculated at a particular point, $\vec{\theta}$, in the model parameter
 2735 space such that $N_i^{MC} = N_i^{MC}(\vec{\theta})$, where N_i represents the bin content of the
 2736 i^{th} bin. The data and Monte Carlo spectra are represented by N_i^D and N_i^{MC} ,
 2737 respectively. The bin contents for the beam near detector, beam far detector and
 2738 atmospheric samples are denoted with *ND*, *FD*, and *Atm*, respectively. The
 2739 binning index, i , runs over all the bins within a sample. Taking the FHC1Rmu far
 2740 detector sample as an example, the binning index runs over all the reconstructed
 2741 neutrino energy bins. The likelihood calculation between the data and the Monte
 2742 Carlo prediction for a particular bin follows a Poisson distribution, where the
 2743 data is treated as a fluctuation of the simulation.

2744 The data can consist of either real data or an ‘Asimov’ Monte Carlo prediction,
 2745 which is typically used for sensitivity studies and denoted ‘Asimov data’. The
 2746 process for building Asimov data is as follows. The Monte Carlo prediction is

2747 reweighted using a particular set of oscillation parameters (potentially those
 2748 listed in Table 2.2) and systematic parameter tune. The resulting spectra for each
 2749 sample is then defined to be the Asimov data for that sample. Whilst this results
 2750 in unphysical non-integer data predictions, it eliminates statistical fluctuations
 2751 from the data. Therefore, the results of a fit to Asimov data should not include any
 2752 biases from statistical fluctuations. Furthermore, these results should produce
 2753 posterior probability distributions consistent with the parameters which were
 2754 used to make the data prediction. That is to say, the fit results should return the
 2755 known parameters. Any biases seen would be attributed to correlations between
 2756 each oscillation parameter and correlations between oscillation and systematic
 2757 parameters. Consequently, Asimov fit results present the maximum precision
 2758 at which the oscillation parameters could be measured to.

2759 Following the T2K analysis presented in [74], the likelihood contribution for
 2760 the near detector samples also includes a Monte Carlo statistical uncertainty term,
 2761 derived from the Barlow and Beeston statistical treatment [217, 218]. In addition
 2762 to treating the data as a Poisson fluctuation of the Monte Carlo prediction, it
 2763 includes a contribution to the likelihood that which treats the generated Monte
 2764 Carlo prediction as a statistical fluctuation of the actual true simulation assuming
 2765 an infinite amount of statistics had been created. The technical implementation
 2766 of this additional likelihood term is documented in [190] and briefly summarised
 2767 as follows. The term is defined as,

$$\frac{(\beta_i - 1)^2}{2\sigma_{\beta_i}^2}, \quad (8.1)$$

2768 where β_i represents a scaling parameter for the i^{th} bin that relates the bin
 2769 content for the amount of Monte Carlo actually generated N_i^{MC} to the bin content
 2770 if an infinite amount of Monte Carlo statistics had been generated $N_{i,true}^{MC}$, such
 2771 that $N_{i,true}^{MC} = \beta_i \times N_i^{MC}$. In the case where a sufficient amount of Monte Carlo
 2772 statistics had been generated, $\beta_i = 1$. An analytical solution for β_i is given in
 2773 [190]. Additionally, $\sigma_{\beta_i} = \sqrt{\sum_i w_i^2 / N_i^{MC}}$ where $\sqrt{\sum_i w_i^2}$ represents the sum of
 2774 the square of the weights of the Monte Carlo events which fall into bin i .

An additional contribution to the likelihood comes from the variation of
 the systematic model parameters. For those parameters with well-motivated
 uncertainty estimates, a covariance matrix, V , describes the prior knowledge of
 each parameter as well as any correlations between the parameters. Due to a
 technical implementation, a single covariance matrix describes each “block” of
 model parameters, e.g. beam flux systematics. The covariance matrix associated
 with the k^{th} block is denoted V^k . This substitution results in $\vec{\theta} = \sum_k^{N_b} \vec{\theta}^k$ and
 $V = \sum_k^{N_b} V^k$ where N_b denotes the number of blocks. A single covariance
matrix is provided for: the oscillation parameters, the beam flux parameters,
the atmospheric flux parameters, the neutrino interaction systematics, the near
detector parameters, the beam far detector systematics, and the atmospheric
far detector systematics. All of these are detailed in section 6.4. The number
of parameters in the k^{th} block is defined as $n(k)$.

The equation for the likelihood \mathcal{L} includes all the terms discussed above.
 It is defined as,

$$\begin{aligned}
 -\ln(\mathcal{L}) = & \\
 & \sum_i^{\text{NDbins}} N_i^{\text{ND},MC}(\vec{\theta}) - N_i^{\text{ND},D} + N_i^{\text{ND},D} \times \ln \left[N_i^{\text{ND},D} / N_i^{\text{ND},MC}(\vec{\theta}) \right] + \frac{(\beta_i - 1)^2}{2\sigma_{\beta_i}^2} \\
 & + \sum_i^{\text{FDbins}} N_i^{\text{FD},MC}(\vec{\theta}) - N_i^{\text{FD},D} + N_i^{\text{FD},D} \times \ln \left[N_i^{\text{FD},D} / N_i^{\text{FD},MC}(\vec{\theta}) \right] \\
 & + \sum_i^{\text{Atmbins}} N_i^{\text{Atm},MC}(\vec{\theta}) - N_i^{\text{Atm},D} + N_i^{\text{Atm},D} \times \ln \left[N_i^{\text{Atm},D} / N_i^{\text{Atm},MC}(\vec{\theta}) \right] \\
 & + \frac{1}{2} \sum_k^{N_b} \sum_i^{n(k)} \sum_j^{n(k)} (\vec{\theta}^k)_i (V^k)_{ij}^{-1} (\vec{\theta}^k)_j.
 \end{aligned} \tag{8.2}$$

The negative log-likelihood value is determined at each step of the MCMC
 to build the posterior distribution defined in chapter 4. This value is minimised
 when the Monte Carlo prediction tends towards the data spectrum.

2793 8.2.1 Likelihood Scans

2794 Using the definition of the likelihood presented in section 8.2, the response of each
2795 sample to a variation of a particular parameter can be studied. Figure 8.1 presents
2796 the variation of all the samples (beam and atmospheric) at the far detector to the
2797 oscillation parameters of interest: δ_{CP} , $\sin^2(\theta_{13})$, $\sin^2(\theta_{23})$, and Δm_{32}^2 . These plots
2798 are colloquially called ‘likelihood scan’ (or ‘log-likelihood scans’). The process
2799 of making these plots is as follows. An Asimov data set (following technique
2800 detailed in section 8.2) is built using the AsimovA oscillation parameters and
2801 pre-fit systematic tune. The Monte Carlo is then reweighted using the value of
2802 the oscillation parameter at each point on the x-axis of the scan. The likelihood
2803 is then calculated between the Asimov data and Monte Carlo prediction and
2804 plotted. This process identifies which samples drive the determination of the
2805 oscillation parameters in the joint fit.

2806 Due to the caveat of fixed systematic parameters and the correlations between
2807 oscillation parameters being ignored when creating these likelihood scans, the
2808 value of $\chi^2 = 1$ (or $-2 \times \ln(\text{Likelihood}) = 1$) does not equate to the typical
2809 1σ sensitivity. However, it does give an indication of which samples respond
2810 the strongest to a variation in a particular oscillation parameter. The point at
2811 which the likelihood tends to zero illustrates the value of the parameter used
2812 to build the Asimov data prediction.

2813 The sensitivity to $\sin^2(\theta_{23})$ is mostly dominated by the beam muon-like
2814 samples. The response of an individual atmospheric sample is small but non-
2815 negligible such that the summed response over all atmospheric samples becomes
2816 comparable to that of the muon-like beam samples. Consequently, the sensitivity
2817 of the joint fit to $\sin^2(\theta_{23})$ would be expected to be greater than the beam-only
2818 analysis. The only sample which respond to the $\sin^2(\theta_{13})$ oscillation parameter is
2819 the electron-like beam samples. Consequently, no increase in sensitivity beyond
2820 that of the T2K-only analysis would be expected from the joint fit. Regardless, the
2821 sensitivity of the beam sample is significantly weaker than the external reactor
2822 constraint so prior knowledge will dominate any measurement that is included

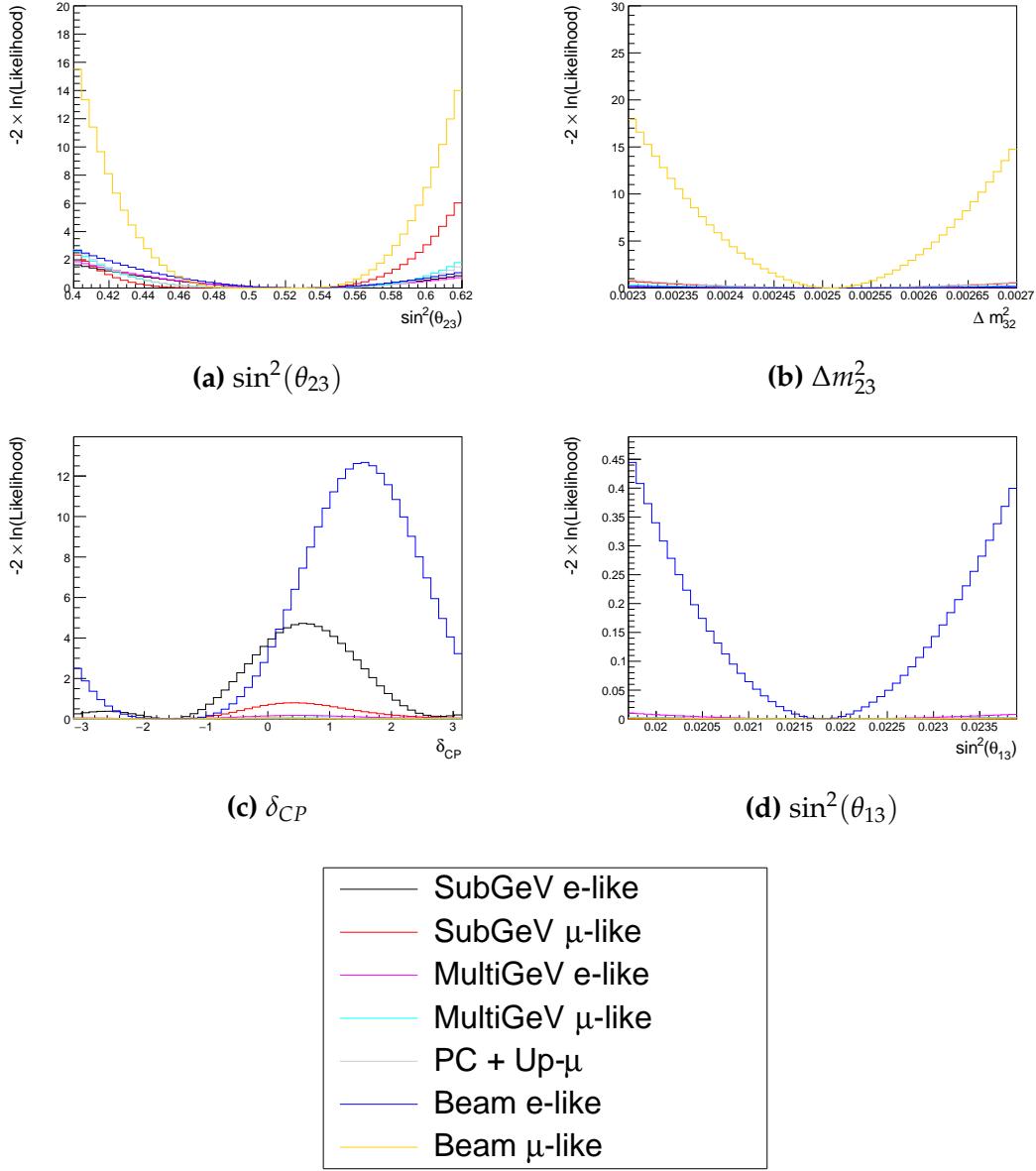


Figure 8.1: The response of the likelihood, as defined in section 8.2, illustrating the response of the samples to a variation of an oscillation parameter.

within this thesis. The Δm_{21}^2 and $\sin^2(\theta_{12})$ parameters are not considered as there is simply no sensitivity in any sample considered within this analysis. The response to Δm_{32}^2 is completely dominated by the beam muon-like samples. This is because the beam neutrino energy can be specifically tuned to match the maximal disappearance probability. Despite this, improvements to the $|\Delta m_{32}^2|$ sensitivity may be expected due to additional mass hierarchy determination

2829 added by the atmospheric samples.

2830 Two-dimensional scans of the appearance ($\sin^2(\theta_{13})-\delta_{CP}$) and disappearance
2831 ($\sin^2(\theta_{23})-\Delta m_{32}^2$) parameters are illustrated in Figure 8.2 and Figure 8.3, respec-
2832 tively. The caveat of fixed systematic parameters and correlations between other
2833 oscillation parameters being neglected still apply.

2834 The appearance log-likelihood scans show the distinct difference in how
2835 the beam and atmospheric samples respond. The beam samples have an ap-
2836 proximately constant width of the 2σ and 3σ contours, throughout all ranges
2837 of δ_{CP} . The response of the atmospheric samples to $\sin^2(\theta_{13})$ is very strongly
2838 correlated to the value of δ_{CP} being evaluated, with the strongest constraints
2839 around $\delta_{CP} \sim 1$. Consequently, this difference allows some of the degeneracy in
2840 a beam-only fit to be broken. Comparing the beam-only and joint fit likelihood
2841 scans, the 2σ continuous contour in δ_{CP} for beam samples is broken when the
2842 atmospheric samples are added. This may result in a stronger sensitivity to
2843 δ_{CP} . Similarly, the width of the 3σ contours also becomes dependent upon the
2844 value of δ_{CP} . Whilst these are encouraging results for the joint fit, these are
2845 not sensitivity measurements as the systematic parameters are fixed and the
2846 correlation between oscillation parameters is neglected. An interesting point to
2847 note is that the atmospheric samples have little sensitivity to $\sin^2(\theta_{13})$ on their
2848 own, as evidenced in Figure 8.1, but can improve sensitivity to the parameter
2849 when combined within the simultaneous fit.

2850 The response of the atmospheric samples in Figure 8.2 shows an interesting
2851 behaviour when considering the application of the reactor constraint. At higher
2852 values of $\sin^2(\theta_{13})$, two lobes appear around $\delta_{CP} \sim -\pi/2$ and $\delta_{CP} \sim 2.4$. If this
2853 distribution was projected onto the δ_{CP} axis, these lobes would mean the posterior
2854 distribution would have a significant dip between these values. However, the
2855 region of $\sin^2(\theta_{13})$ near the reactor constraint ($\sin^2(\theta_{13}) = (2.18 \pm 0.08) \times 10^{-2}$)
2856 is flatter across the range of δ_{CP} . Therefore, if we were to project only this
2857 region onto the δ_{CP} axis, the dip between the peaks would not be as significant.
2858 If this behaviour was to be seen in the results of a fit, these marginalisation

2859 effects would actually conspire to reduce the sensitivity to δ_{CP} if the reactor
2860 constraint was to be applied.

2861 The disappearance log-likelihood scans in $\sin^2(\theta_{23})$ - Δm_{32}^2 space (Figure 8.3)
2862 show the expected behaviour when considering the one-dimensional scans
2863 already discussed. The uncertainty on the width of $|\Delta m_{32}^2|$ is mostly driven
2864 by the beam-only sensitivities. However, the width of this contour in the
2865 inverted mass region ($\Delta m_{32}^2 < 0$) is significantly reduced due to the ability
2866 of the atmospheric samples to select the correct (normal) mass hierarchy. The
2867 width of the uncertainty in $\sin^2(\theta_{23})$ is also reduced compared to the beam-only
2868 sensitivities, with a further decrease in the inverted hierarchy region due to the
2869 better mass hierarchy determination.

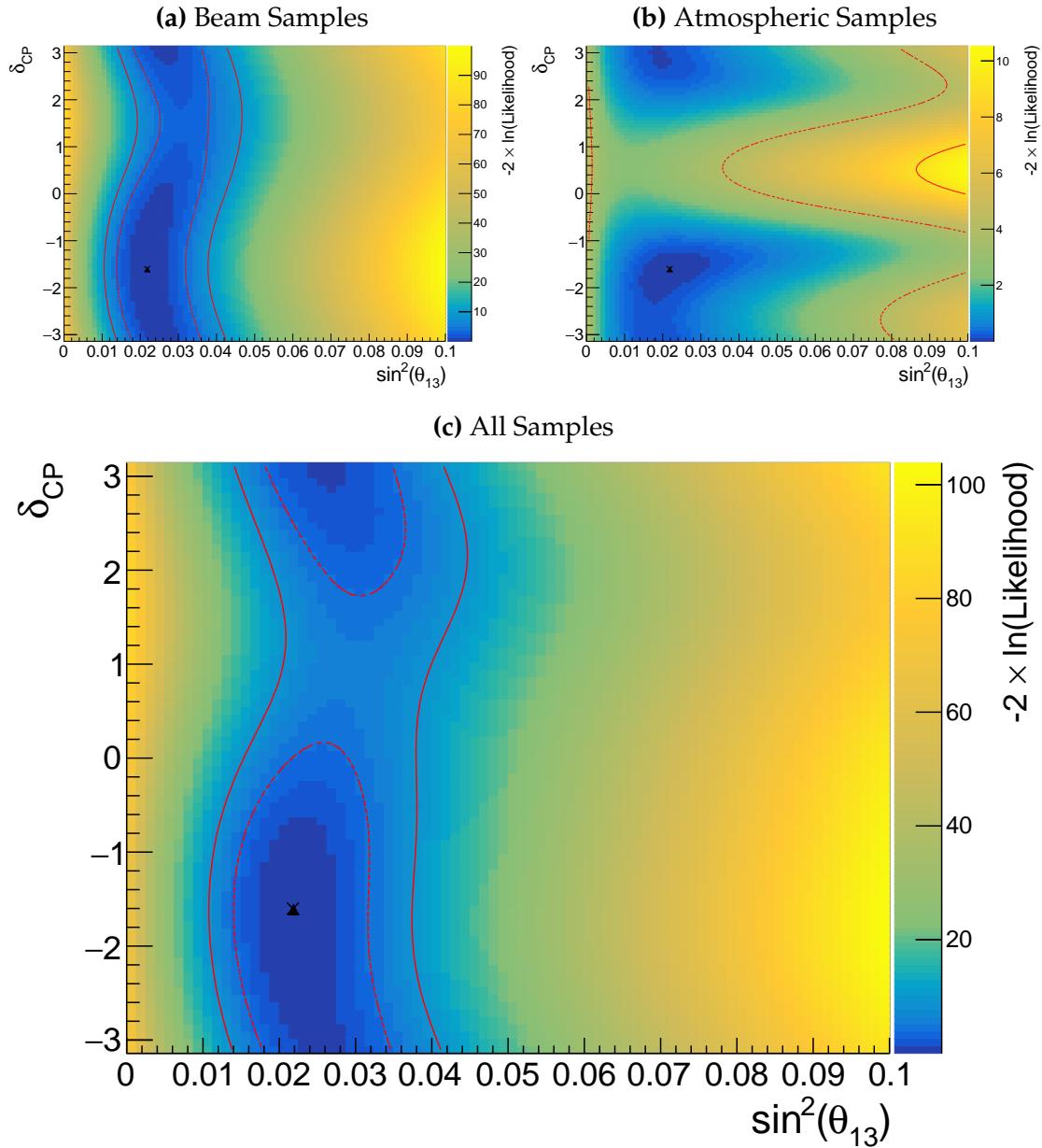


Figure 8.2: Two-dimensional log-likelihood scan of the appearance ($\sin^2(\theta_{13})$ - δ_{CP}) parameters showing the response of the beam samples (top left), atmospheric samples (top right) and the summed response (bottom). The Asimov A oscillation parameters, defined in Table 2.2, are known to be the true point (Black Cross). The position of the smallest log-likelihood is highlighted with the triangle. Prior uncertainty terms of the oscillation parameters are neglected. The two(three) sigma contour levels are illustrated with the dashed(solid) red line.

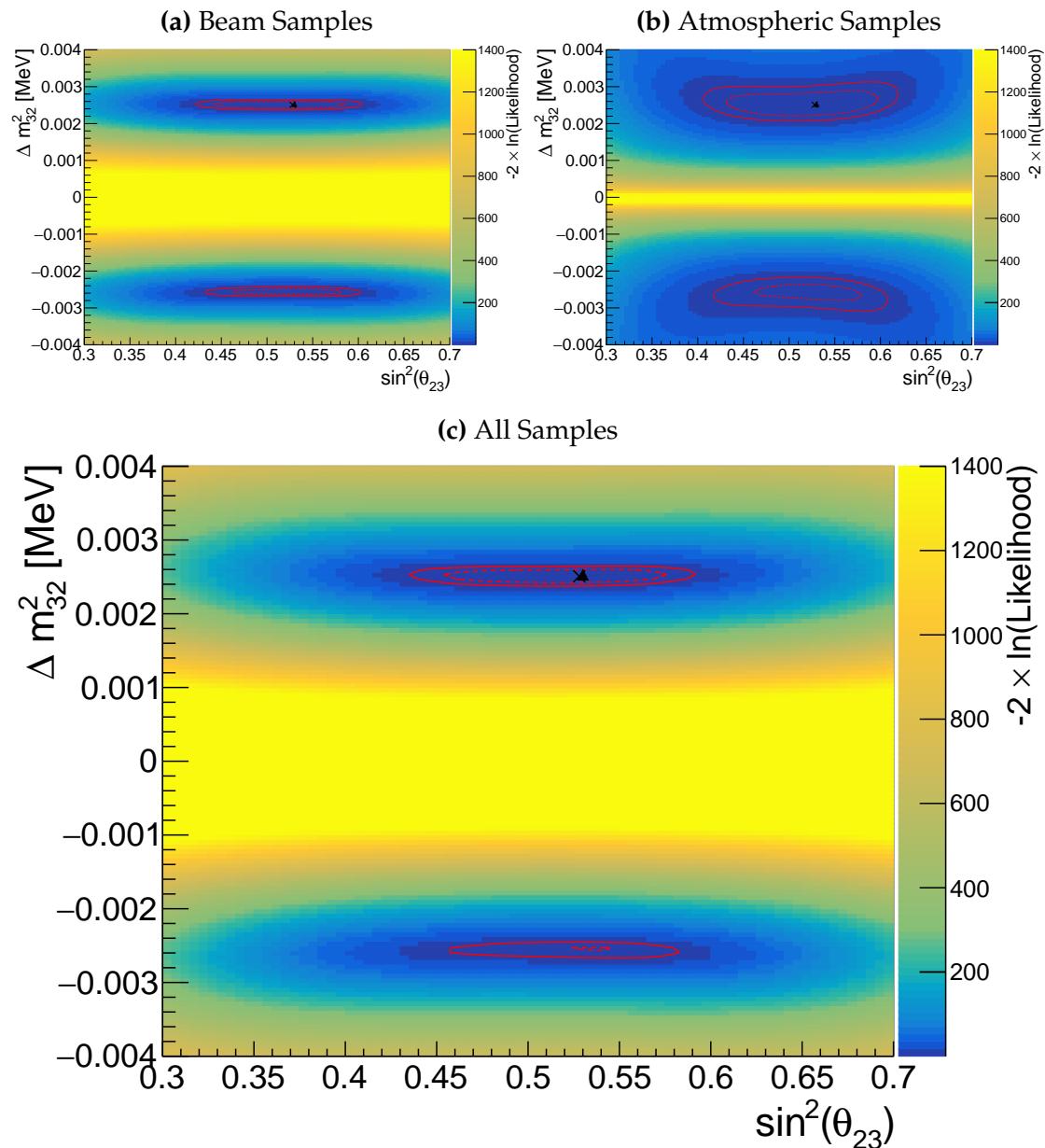


Figure 8.3: Two-dimensional log-likelihood scan of the disappearance ($\sin^2(\theta_{23})$)- Δm_{32}^2) parameters showing the response of the beam samples (top left), atmospheric samples (top right) and the summed response (bottom). The Asimov A oscillation parameters, defined in Table 2.2, are known to be the true point (Black Cross). The position of the smallest log-likelihood is highlighted with the triangle. Prior uncertainty terms of the oscillation parameters are neglected. The two(three) sigma contour levels are illustrated with the dashed(solid) red line.

The likelihood scans illustrated thus far only consider the sensitivity of this analysis for a fixed set of true oscillation parameters, namely Asimov A defined in Table 2.2. Whilst computational infeasible to run many fits at different parameter sets, it is possible to calculate the likelihood response to different Asimov data sets. Figure 8.4 and Figure 8.5 illustrate how the sensitivity changes for differing true values of δ_{CP} and $\sin^2(\theta_{23})$, respectively. For both of these plots, the other oscillation parameters are fixed at their Asimov A values. Consequently, the caveat of fixed systematic parameters and correlations between other oscillation parameters being neglected still applies.

To explain how these plots are made, consider Figure 8.4. This plot is built by considering multiple one-dimensional log-likelihood scans, each creating an Asimov data with the value of δ_{CP} taken from the x-axis. The likelihood to a particular Asimov data set is calculated after reweighting the Monte Carlo prediction to each value of δ_{CP} on the y-axis.

Figure 8.4 illustrates the sensitivity to δ_{CP} . Notably, the 1σ intervals contain regions in the off-diagonal for which the beam and atmospheric samples have broken and discontinuous contours. This indicates that there are regions of δ_{CP} which are degenerate. For example, for the x-axis value of $\delta_{CP} = 0$, the beam samples sensitivity would include two discontinuous regions excluded from the 1σ interval: $\delta_{CP} \sim 0$ and $\delta_{CP} \sim \pi$. The offset in δ_{CP} between the beam and atmospheric samples allows the joint fit to have increased sensitivity in these regions. Consequently, the difference between the beam-only and joint beam-atmospheric fit should be studied using multiple Asimov data sets.

Despite the increased sensitivity at 1σ , the 2σ intervals from the joint fit are more similar to the two independent sensitivities and the off-diagonal degeneracies mostly remain. This indicates that the joint fit has the strength to aid parameter determination but can not entirely break the degeneracies in δ_{CP} at higher confidence levels.

Figure 8.5 illustrates a similar analysis as above, although the value of $\sin^2(\theta_{23})$ is varied and δ_{CP} is fixed to the Asimov A parameter value. Due to the beam

parameters and baseline being tuned to specifically target this oscillation parameter, the average sensitivity of the beam samples is stronger than the atmospheric samples. However, the degeneracy around maximal mixing ($\sin^2(\theta_{23}) = 0.5$) is significantly more peaked in the beam samples compared to the atmospheric samples. This behaviour is strengthened when considering the 2σ intervals, to the point where two distinct discontinuous regions of the 2σ intervals exist around the Asimov point $\sin^2(\theta_{23}) \sim 0.41, 0.6$. Given the caveat of only considering likelihood scans, the joint analysis would mostly eliminate the discontinuous intervals in these regions. This means that the joint fit could feasibly have an increased preference for the correct octant hypothesis.

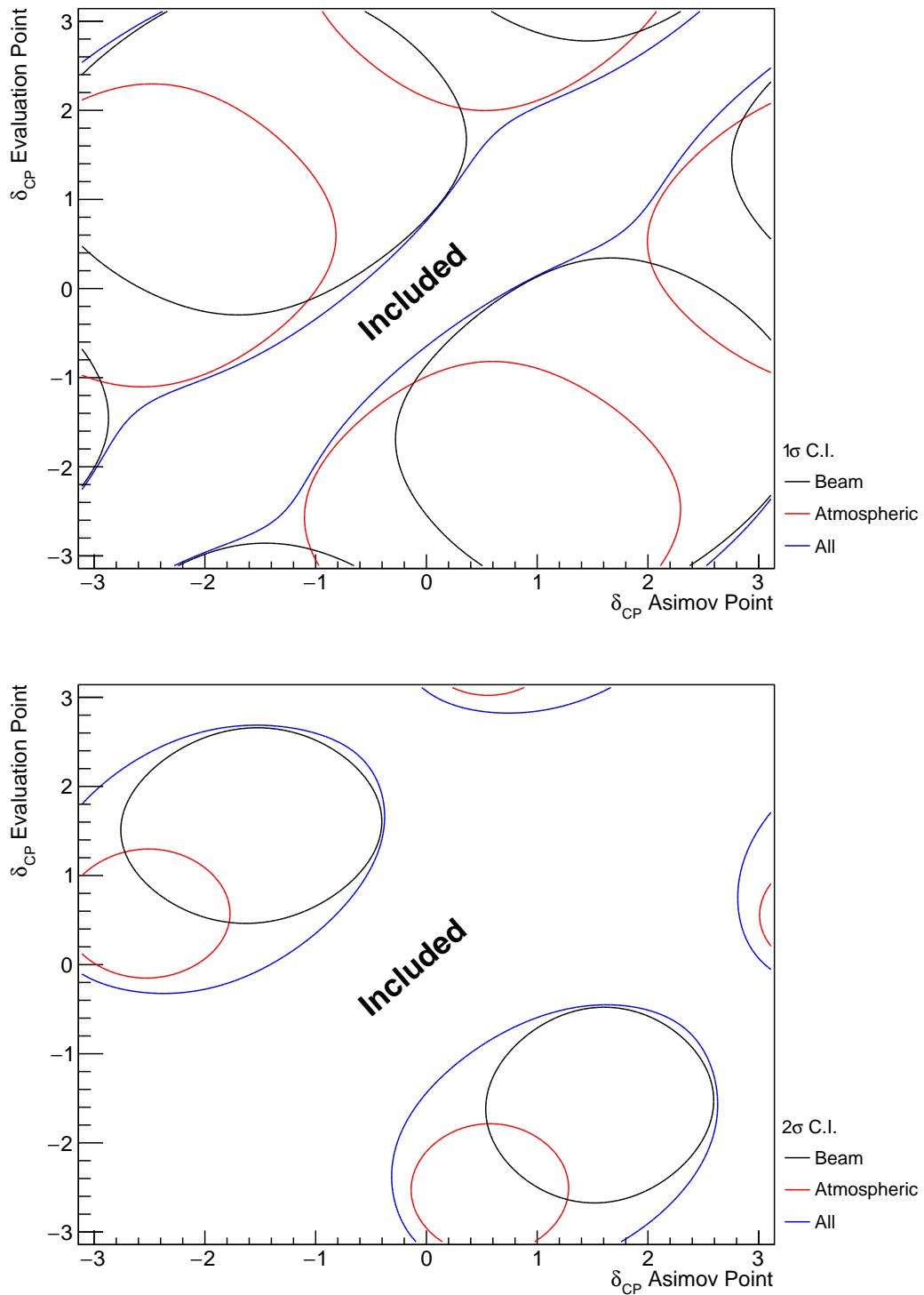


Figure 8.4: A series of one-dimensional likelihood scans over δ_{CP} , where an Asimov data set is built for each value of δ_{CP} on the x-axis and the likelihood is evaluated for each value of δ_{CP} on the y-axis. The diagonal represents the minimum log-likelihood and defines the region included within the 1σ (Top) and 2σ (Bottom) confidence intervals. The beam (black) and atmospheric (red) samples are individually plotted and the joint fit (blue) is the sum of the two.

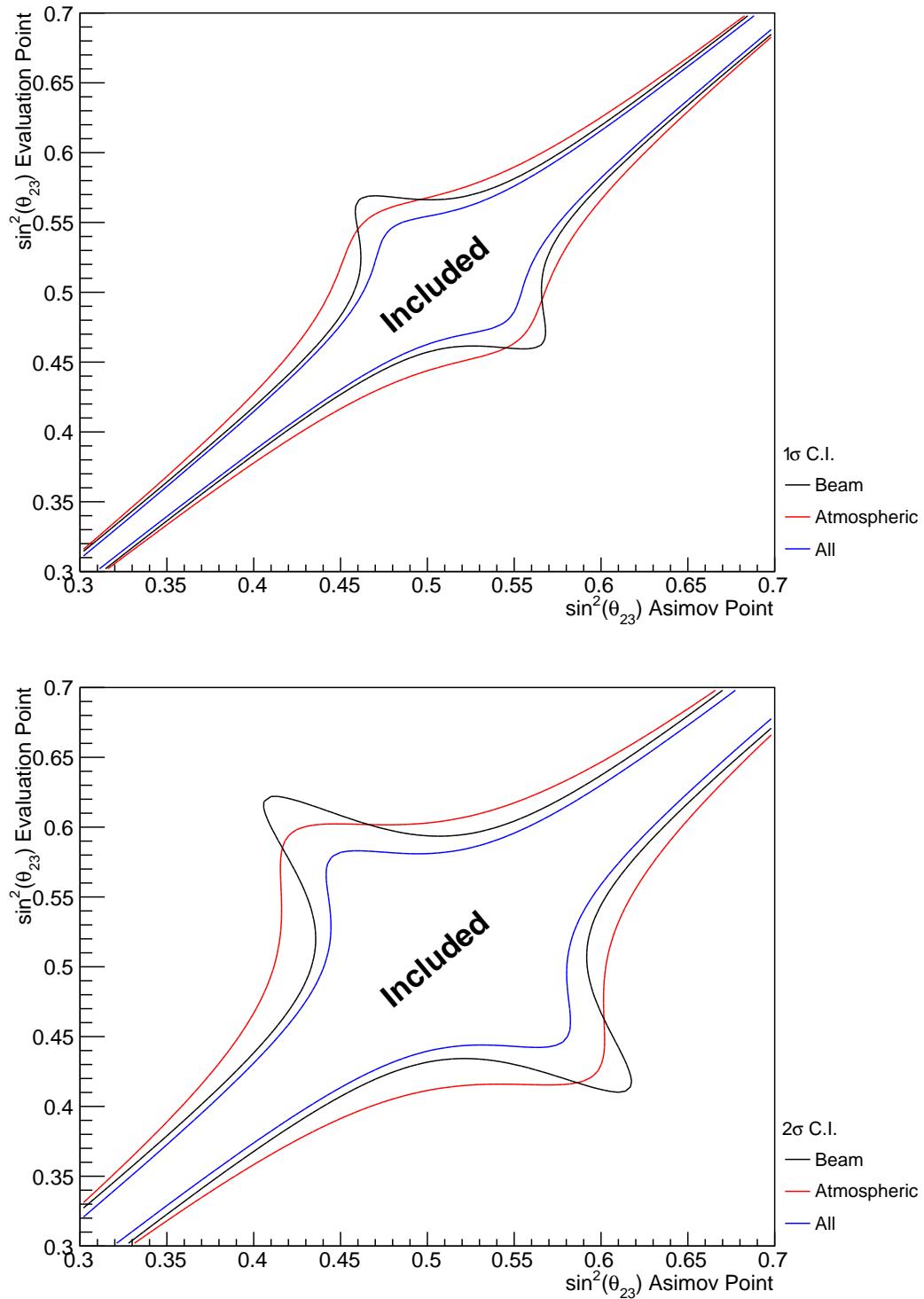


Figure 8.5: A series of one-dimensional likelihood scans over $\sin^2(\theta_{23})$, where an Asimov data set is built for each value of $\sin^2(\theta_{23})$ on the x-axis and the likelihood is evaluated for each value of $\sin^2(\theta_{23})$ on the y-axis. The diagonal represents the minimum log-likelihood and defines the region included within the 1σ (Top) and 2σ (Bottom) confidence intervals. The beam (black) and atmospheric (red) samples are individually plotted and the joint fit (blue) is the sum of the two.

Alongside oscillation parameters (Figure 8.1), the sensitivity to systematic parameters can also be studied for the joint fit. As some of these parameters are correlated between the beam and atmospheric events, the response of the atmospheric samples can modify the constraint. This means the systematics can have additional constraints than what they would from a beam-only analysis. Therefore, the response from the beam and the atmospheric samples to various systematic parameters has been compared in Figure 8.6. The Asimov data set has been created using the AsimovA oscillation parameter and the pre-fit systematic tune. For example, the systematic parameter controlling the effective axial mass coupling in CCQE interactions, M_A^{QE} , is clearly dominated by the ND constraint. An example where the response of the atmospheric sample is approximately similar to the near detector constraint is the 2p2h_Cto0 normalisation systematic. This systematic models the scaling of the 2p2h interaction cross-section on a carbon target to an oxygen target. There are also systematics which have no near detector constraint. For example, the systematic parameters which describe the normalisation of the NC1Gamma and NCOther interaction modes. The atmospheric samples are significantly more sensitive to these systematics than the beam samples due to their similar interaction contributions but relatively higher statistics (Table 8.1). As an example of how the atmospheric samples can help constrain systematic parameters used within the T2K-only analysis, these NC background events in beam electron-like samples will be considerably more constrained with the additional sensitivity of atmospheric samples. This would be expected to reduce the overall uncertainty of the beam electron-like event rates in the joint analysis compared to the beam-only studies. This could modify the sensitivity of the beam samples due to the more constrained background events.

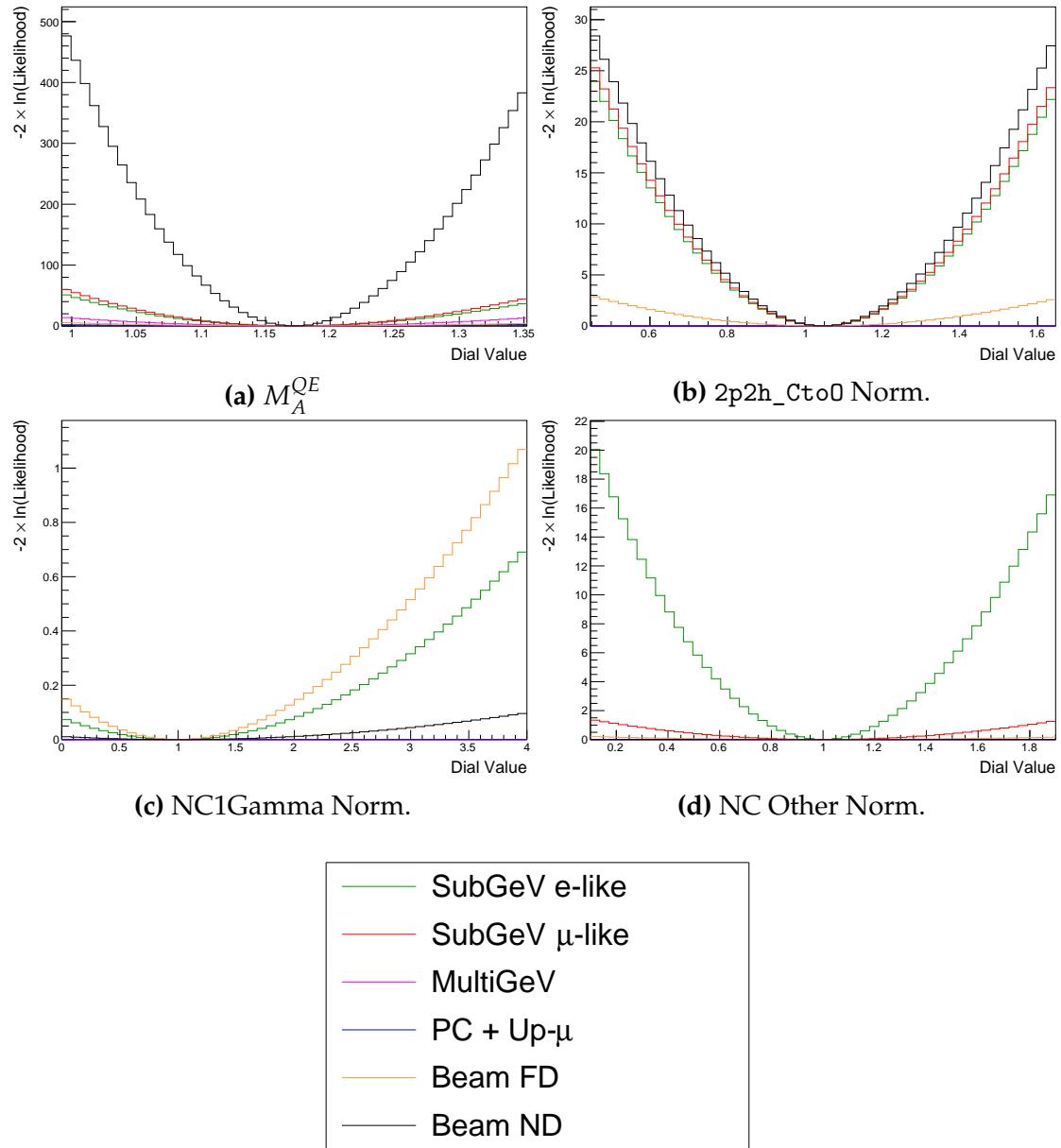


Figure 8.6: The response of the likelihood, as defined in section 8.2, illustrating the response of the samples to the various cross-section systematic parameters.

2935 8.3 Sensitivities

2936 The sensitivities of the joint T2K and SK oscillation analysis are presented in
2937 the form of Asimov fits. This technique builds an Asimov data set (follow-
2938 ing section 8.2) using the AsimovA oscillation parameters and post-BANFF
2939 systematic tune.

2940 In practice, the Asimov fits presented within this analysis are modified from
2941 the above definition. An Asimov prediction of both beam and atmospheric far
2942 detector samples is fit whilst the true data is used for near detector samples.
2943 The Asimov predictions at the far detector are built using the BANFF tuning (as
2944 discussed in section 3.2). These modifications mean that the results are equivalent
2945 to performing a far detector Asimov fit using inputs from the BANFF data fit.
2946 Consequently, this allows the results to be cross-checked with the results from
2947 the P-Theta analysis. The comparison has been performed and is documented in
2948 [219]. No significant discrepancies were found between the fitters.

2949 This section proceeds with the following studies. Firstly, the sensitivity of
2950 the atmospheric samples after the T2K cross-section has been applied to the low-
2951 energy events is detailed in subsection 8.3.1. This includes studying the choice of
2952 applying the 2020 PDG reactor constraint [75] to the atmospheric samples, which
2953 is documented in subsection 8.3.2. Additionally, the effect of applying the near-
2954 detector constraints onto the atmospheric samples is discussed in subsection 8.3.3.
2955 The main result is the sensitivity of the simultaneous beam and atmospheric fit.
2956 The sensitivities, both with and without the application of the reactor constraint,
2957 are presented in subsection 8.3.4 and subsection 8.3.5, respectively. To indicate
2958 the benefit of the joint analysis, the sensitivities are compared to the 2020 T2K
2959 sensitivities [74, 181] in subsection 8.3.6 and subsection 8.3.7. As shown in
2960 subsection 8.2.1, the response of the beam and atmospheric samples change
2961 depending upon the true set of oscillation parameters assumed. Therefore,
2962 subsection 8.3.8 documents the sensitivities at an alternative oscillation parameter

²⁹⁶³ set. It is important to note that these results have been published at the Neutrino
²⁹⁶⁴ 2022 conference on behalf of the T2K and SK collaborations [107].

2965 8.3.1 Atmospheric-Only Sensitivity Without Reactor Constraint

2966 This section presents the results of an Asimov fit using samples from the near
 2967 detector and only atmospheric samples from the far detector. The results are
 2968 presented as one-dimensional or two-dimensional histograms which have been
 2969 marginalised over all other parameters using the technique outlined in sub-
 2970 section 4.3.1. Each histogram displays the posterior probability density and
 2971 illustrates the credible intervals, calculated using the technique in subsection 4.3.2.
 2972 For these fits in this subsection, a flat penalty term is used for $\sin^2(\theta_{13})$ such the
 2973 reactor constraint is not applied. The Asimov data is generated assuming the
 2974 AsimovA oscillation parameter set defined in Table 2.2 and the post-BANFF
 2975 systematic parameter tune.

Without Reactor Constraint, Both Hierarchies

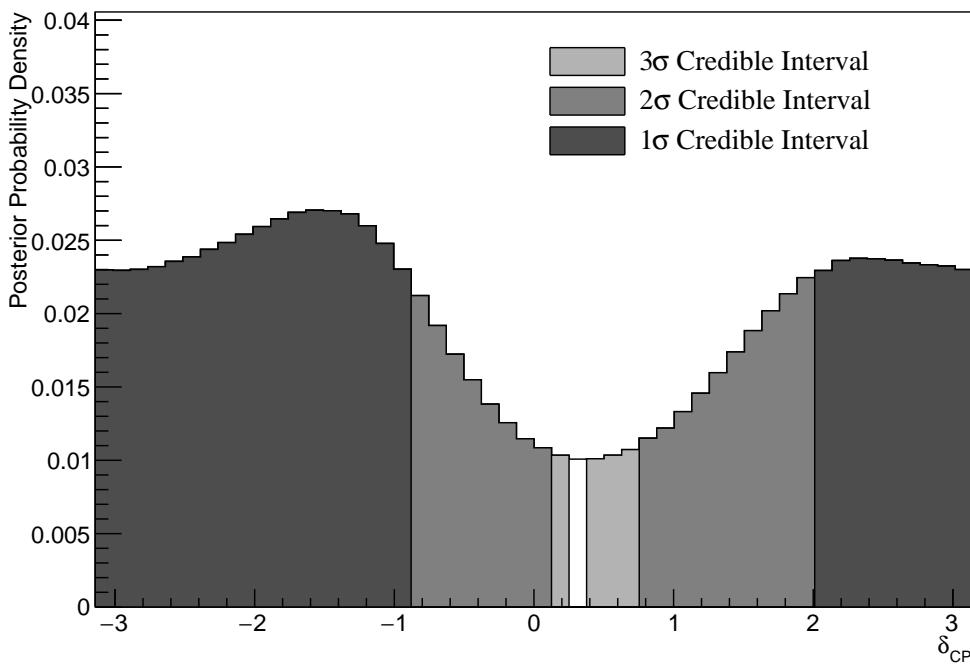


Figure 8.7: The one-dimensional posterior probability density distribution in δ_{CP} , marginalised over both hierarchies, from the SK atmospheric-only fit. The reactor constraint is not applied.

2976 Figure 8.7 illustrates the posterior probability density for δ_{CP} , marginalised
 2977 over both hierarchies. If instead, only steps in the normal hierarchy were
 2978 considered, the shape of the contours would change. The fit favours the known

2979 oscillation parameter ($\delta_{CP} = -1.601$) although the posterior probability is very
 2980 flat through the range of $-\pi < \delta_{CP} < -1$ and $2 < \delta_{CP} < \pi$. There is also a region
 2981 around $\delta_{CP} \sim 0.4$ which is disfavoured at 2σ . This indicates that the SK samples
 2982 can rule out some parts of the CP conserving parameter space reasonably well,
 2983 near $\delta_{CP} \sim 0.4$, when the true value of $\delta_{CP} \sim \pi/2$.

Without Reactor Constraint, Both Hierarchies

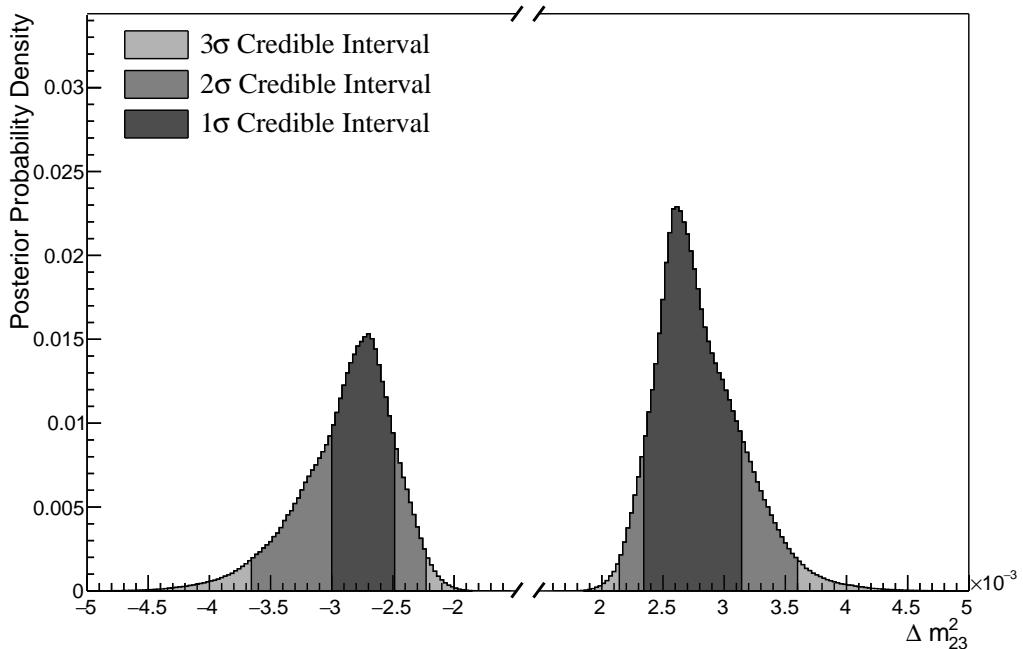


Figure 8.8: The one-dimensional posterior probability density distribution in Δm_{32}^2 , marginalised over both hierarchies, from the SK atmospheric-only fit. The reactor constraint is not applied.

2984 The posterior probability density in Δm_{32}^2 is given in Figure 8.8. This dis-
 2985 tribution includes steps in both the normal hierarchy (NH, $\Delta m_{32}^2 > 0$) and the
 2986 inverse hierarchy (IH, $\Delta m_{32}^2 < 0$). The highest posterior probability density is
 2987 found within the NH, which agrees with the known oscillation parameter value.
 2988 However, all of the credible intervals span both of the hierarchies hypotheses.
 2989 If instead, only steps in the normal hierarchy were considered, the shape of the
 2990 contours would change. The known oscillation parameter is $2.509 \times 10^{-3} \text{ eV}^2$,
 2991 which is contained within the 1σ credible interval.

	LO ($\sin^2 \theta_{23} < 0.5$)	UO ($\sin^2 \theta_{23} > 0.5$)	Sum
NH ($\Delta m_{32}^2 > 0$)	0.17	0.40	0.58
IH ($\Delta m_{32}^2 < 0$)	0.13	0.29	0.42
Sum	0.31	0.69	1.00

Table 8.2: The distribution of steps in an SK atmospheric-only fit, presented as the fraction of steps in the upper (UO) and lower (LO) octants and the normal (NH) and inverted (IH) hierarchies. The reactor constraint is not applied. The Bayes factors are calculated as $B(\text{NH}/\text{IH}) = 1.37$ and $B(\text{UO}/\text{LO}) = 2.24$.

Following the discussion in subsection 4.3.3, the Bayes factor for hierarchy preference can be calculated by determining the fraction of steps that fall into the NH and the IH regions, as an equal prior is placed on both hypotheses. A similar calculation can be performed by calculating the fraction of steps which fall in the lower octant (LO, $\sin^2 \theta_{23} < 0.5$) or upper octant (UO, $\sin^2 \theta_{23} > 0.5$). The fraction of steps, broken down by hierarchy and octant, are given in Table 8.2. The Bayes factor for preferred hierarchy model is $B(\text{NH}/\text{IH}) = 1.37$. Table 4.1 states this value of the Bayes factor indicates a weak preference for the normal hierarchy model. The Bayes factor for choice of octant is $B(\text{UO}/\text{LO}) = 2.24$. This is also classified as a weak preference for the UO. Both of these show that the fit is returning the correct choice of models (NH and UO) for the known Asimov A oscillation parameters defined in Table 2.2.

The 1σ credible intervals, broken down by hierarchy, and position in parameter space of the highest posterior probability density is given in Table 8.3. These are taken from the one-dimensional projections of the oscillation parameters, marginalised over all other parameters within the fit. For the known Asimov value of $\delta_{CP} = -1.601$, the 1σ credible interval rules out a region between $\delta_{CP} = -0.86$ and $\delta_{CP} = 1.96$, when marginalising over both hierarchies. The position of the highest posterior density is $\delta_{CP} = -1.57$ which is clearly compatible with the known oscillation parameter value.

The sensitivity of the atmospheric samples to $\sin^2(\theta_{13})$ is presented in Figure 8.9. The likelihood scans presented in Figure 8.1 suggest that the sensitivity

Parameter	Interval	HPD
δ_{CP} , (BH)	$[-\pi, -0.86], [1.96, \pi]$	-1.57
δ_{CP} , (NH)	$[-\pi, -0.86], [1.88, \pi]$	-1.57
δ_{CP} , (IH)	$[-\pi, -0.94], [1.96, \pi]$	-1.57
Δm_{32}^2 (BH) [$\times 10^{-3}\text{eV}^2$]	$[-3.00, -2.50], [2.35, 3.15]$	2.65
Δm_{32}^2 (NH) [$\times 10^{-3}\text{eV}^2$]	$[2.39, 3.04]$	2.64
Δm_{32}^2 (IH) [$\times 10^{-3}\text{eV}^2$]	$[-3.15, -2.45]$	-2.70
$\sin^2(\theta_{23})$ (BH)	$[0.476, 0.59]$	0.542
$\sin^2(\theta_{23})$ (NH)	$[0.476, 0.59]$	0.554
$\sin^2(\theta_{23})$ (IH)	$[0.476, 0.59]$	0.542

Table 8.3: The position of the highest posterior probability density (HPD) and width of the 1σ credible interval for the SK atmospheric-only fit. The reactor constraint is not applied. The values are presented by which hierarchy hypothesis is assumed: marginalised over both hierarchies (BH), normal hierarchy only (NH), and inverted hierarchy only (IH).

3014 to $\sin^2(\theta_{13})$ will be small. This behaviour is also seen in the fit results, where the
 3015 width of the 1σ credible intervals span the region of $\sin^2(\theta_{13}) = [0.008, 0.08]$. This
 3016 is more than an order of magnitude worse than the constraint from reactor
 3017 experiments [75].

3018 As previously discussed, the correlations between oscillation parameters are
 3019 also important to understand how the atmospheric samples respond. Figure 8.10
 3020 illustrates the two dimensional $\sin^2(\theta_{13}) - \delta_{CP}$ sensitivity, marginalised over all
 3021 other parameters. The displayed contours are calculated by marginalising over
 3022 both hierarchies. The shape of the 1σ credible interval shows that the constraining
 3023 power of the fit on δ_{CP} is dependent upon the value of $\sin^2(\theta_{13})$. Furthermore,
 3024 they show a strong resemblance to the likelihood scans illustrated in Figure 8.2.
 3025 Whilst the atmospheric samples do not strongly constrain the value of $\sin^2(\theta_{13})$,
 3026 the value of $\sin^2(\theta_{13})$ does impact the atmospheric sensitivity to δ_{CP} . A value of
 3027 $\sin^2(\theta_{13}) \sim 0.02$ would select a continuous contour over all values of δ_{CP} . This
 3028 shows the effect of the marginalisation effect previously described.

Without Reactor Constraint, Both Hierarchies

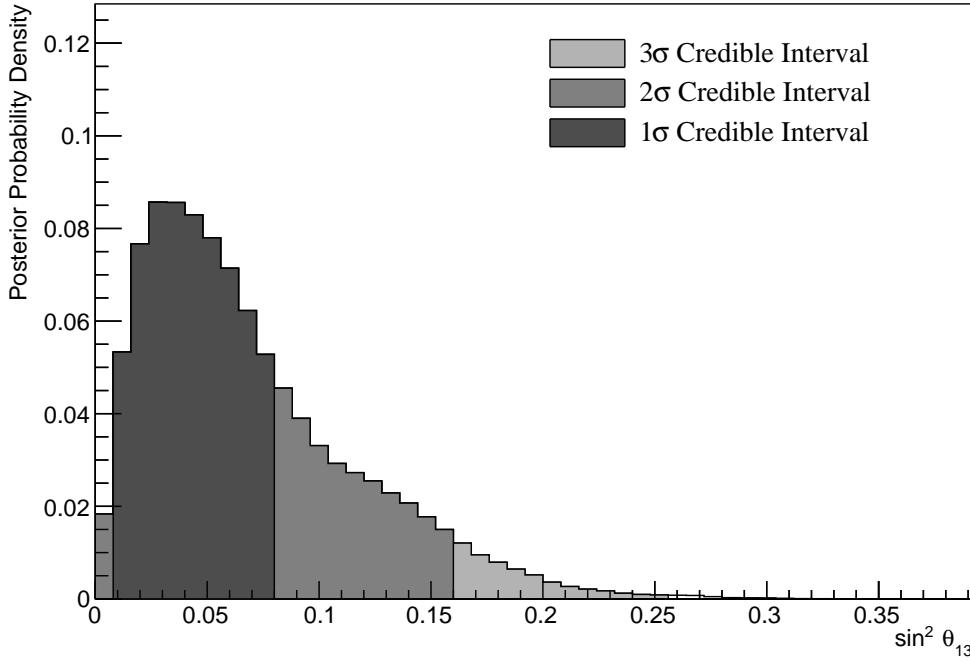


Figure 8.9: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{13})$, marginalised over both hierarchies, from the SK atmospheric-only fit. The reactor constraint is not applied.

3029 The $\sin^2(\theta_{23}) - \Delta m_{32}^2$ disappearance contours are illustrated in Figure 8.11.

3030 As expected, the area contained in the inverted hierarchy 1σ credible interval is
 3031 slightly smaller than that in the normal hierarchy. This follows from the Bayes
 3032 factor showing a weak preference for NH meaning that more of the steps will exist
 3033 in the $\Delta m_{32}^2 > 0$ region. The known oscillation parameters of $\sin^2(\theta_{23}) = 0.528$
 3034 and $\Delta m_{32}^2 = 2.509 \times 10^{-3} \text{ eV}^2$ are contained within the 1σ credible interval.

3035 Figure 8.12 illustrates the two-dimensional projections for each permutation of
 3036 oscillation parameters which this analysis is sensitive to: δ_{CP} , $\sin^2(\theta_{13})$, $\sin^2(\theta_{23})$,
 3037 and Δm_{32}^2 . The purpose of this plot is to illustrate the correlations between the
 3038 oscillation parameters. The contours are calculated whilst marginalising over
 3039 both hierarchies, however, only the NH is illustrated when plotting the Δm_{32}^2
 3040 parameter. As expected the correlations play a significant role in these sensitivity
 3041 measurements, especially the choice of the $\sin^2(\theta_{13})$ constraint. The application
 3042 of reactor constraint would be expected to alter both the width and position of the

Without Reactor Constraint, Both Hierarchies

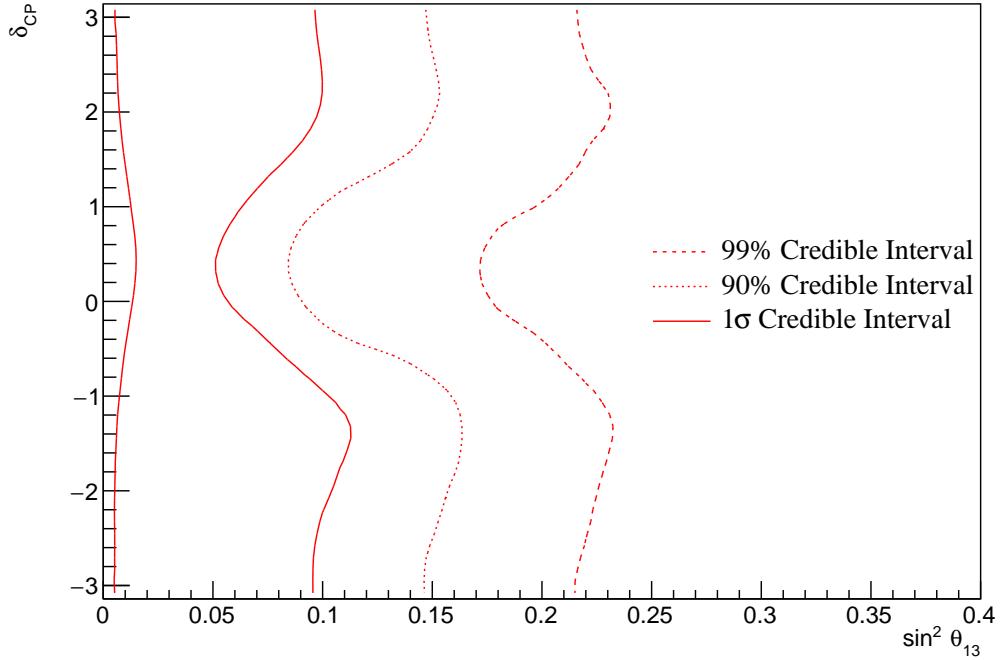


Figure 8.10: The two-dimensional posterior probability density distribution in δ_{CP} – $\sin^2(\theta_{13})$, marginalised over both hierarchies, from the SK atmospheric-only fit. The reactor constraint is not applied.

3043 Δm_{32}^2 , δ_{CP} , and $\sin^2(\theta_{23})$ constraints. The majority of the octant model preference
 3044 comes from the region of $\sin^2(\theta_{13}) \sim 0.03$ such that the application of the reactor
 3045 constraint would not be expected to significantly change the octant preference.
 3046 The reactor constraint would result in lower values of $|\Delta m_{32}^2|$. Interestingly, the
 3047 distribution of steps in the δ_{CP} - $\sin^2(\theta_{13})$ plot is slightly flatter in the region of the
 3048 reactor constraint. Both the posterior distribution from this fit and the distribution
 3049 in Figure 8.2 show a region of low negative log-likelihood extending out towards
 3050 higher values of $\sin^2(\theta_{13})$ in the $\delta_{CP} \sim -\pi/2$ and $\delta_{CP} \sim 2$ region. Consequently,
 3051 the reactor constraint could feasibly reduce the sensitivity of the atmospheric
 3052 samples to δ_{CP} , due to the previously discussed marginalisation effects.

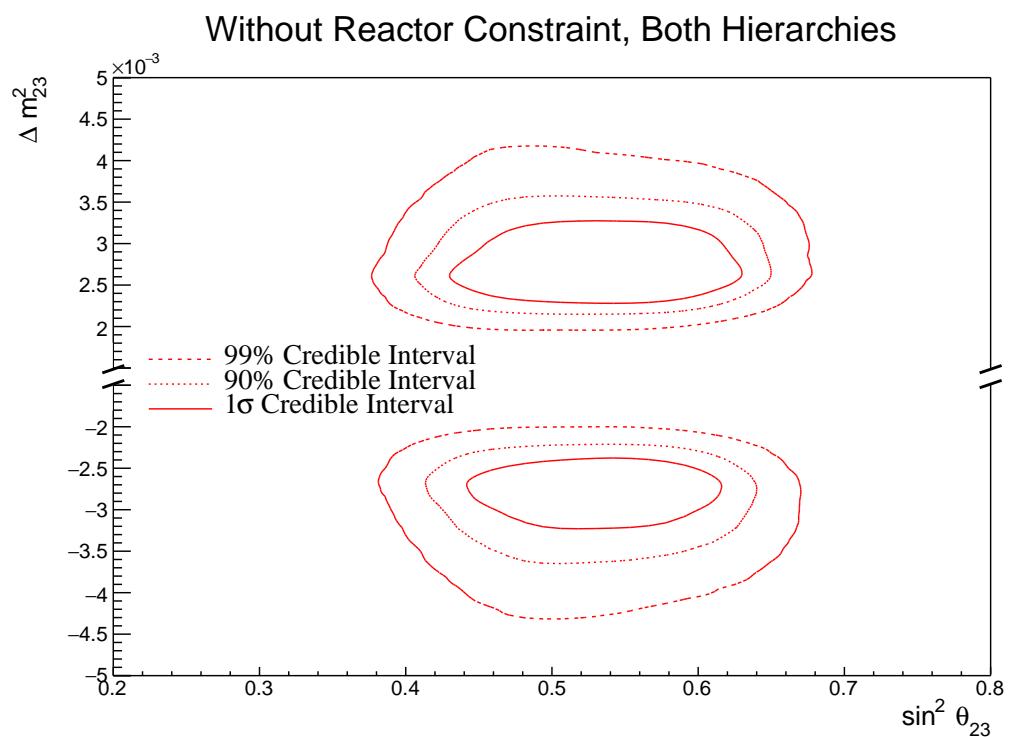


Figure 8.11: The two-dimensional posterior probability density distribution in $\Delta m_{32}^2 - \sin^2(\theta_{23})$, marginalised over both hierarchies, from the SK atmospheric-only fit. The reactor constraint is not applied.

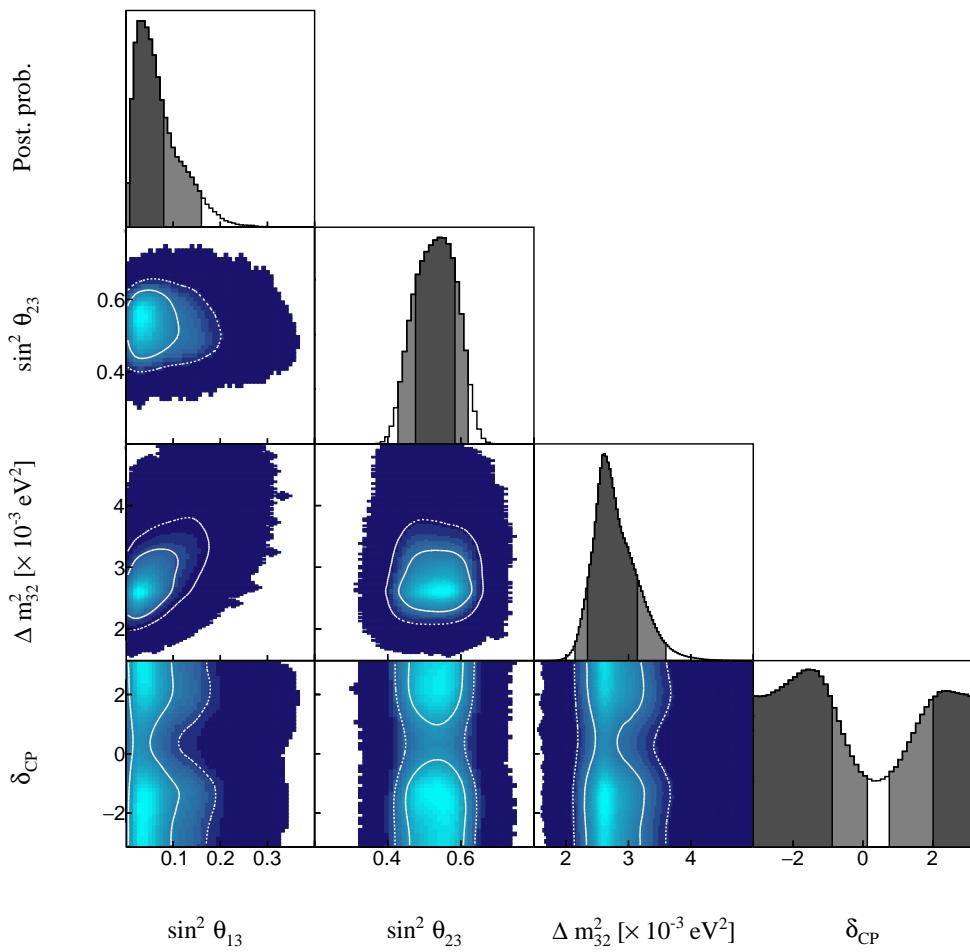


Figure 8.12: The posterior probability density distribution from the SK atmospheric-only fit. The reactor constraint is not applied. The distribution is given for each two-dimensional permutation of the oscillation parameters of interest. The one-dimensional distribution of each parameter is also given.

3053 8.3.2 Atmospheric-Only Sensitivity With Reactor Constraint

3054 The results in subsection 8.3.1 discuss the atmospheric sensitivity when the reactor
 3055 constraint is not applied. The correlations illustrated in Figure 8.12 indicate that
 3056 the marginalisation effects could contribute to differing sensitivities when the
 3057 external reactor constraint is applied. Using the technique discussed in subsec-
 3058 tion 4.1.1, the posterior distribution of the fit in subsection 8.3.1 can be reweighted
 3059 to include the reactor constraint of $\sin^2(\theta_{13}) = (2.18 \pm 0.08) \times 10^{-2}$ [75].

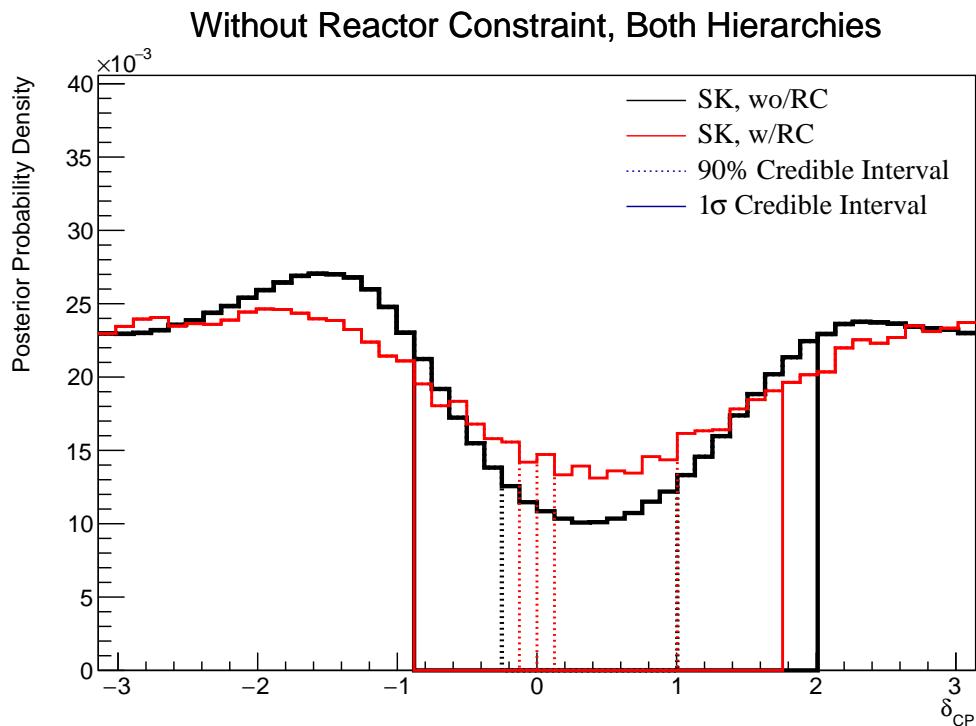


Figure 8.13: The one-dimensional posterior probability density distribution in δ_{CP} compared between the SK atmospheric-only fit (Black) and the SK atmospheric fit with the reactor constraint applied (Red). The distributions are marginalised over both hierarchies.

3060 Figure 8.13 illustrates the sensitivity to δ_{CP} of the atmospheric fit with reactor
 3061 constraint applied. The distribution is less peaked than the previous results.
 3062 This is due to the expected marginalisation effect previously discussed. The
 3063 width of the 1σ credible interval is increased when the reactor constraint is
 3064 applied, indicating less sensitivity to δ_{CP} in the region of $\sin^2(\theta_{13})$ preferred
 3065 by the reactor constraint.

Without Reactor Constraint, Both Hierarchies

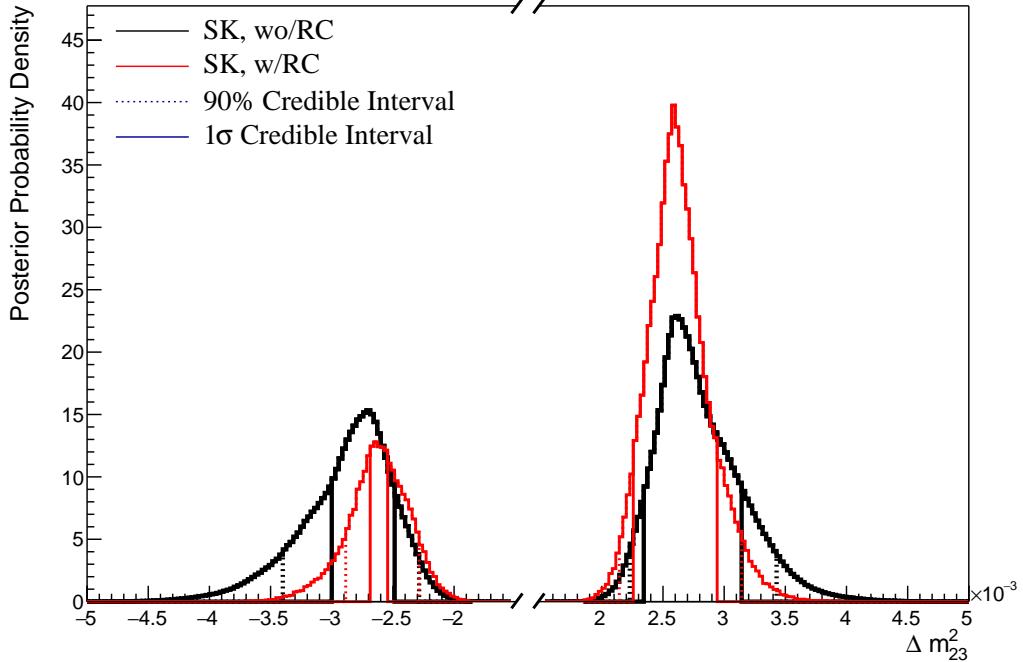


Figure 8.14: The one-dimensional posterior probability density distribution in Δm_{32}^2 compared between the SK atmospheric-only fit (Black) and the SK atmospheric fit with the reactor constraint applied (Red). The distributions are marginalised over both hierarchies.

The reactor constraint increases the sensitivity of the atmospheric samples to Δm_{32}^2 as illustrated in Figure 8.14. The 1σ credible interval in Δm_{32}^2 is determined to be $[-2.70, -2.55] \times 10^{-3}\text{eV}^2$ and $[2.25, 2.95] \times 10^{-3}\text{eV}^2$. The width of the IH credible interval is reduced by $\sim 70\%$ when the reactor constraint is applied. Due to the marginalisation effects observed in Figure 8.12, the favoured region of Δm_{32}^2 moves closer to zero for both hierarchies. A clear explanation of this behaviour is illustrated in Figure 8.15 which illustrates the posterior distribution in the $\Delta m_{32}^2 - \sin^2(\theta_{13})$ parameters, marginalised over both hierarchies. The correlation between Δm_{32}^2 and $\sin^2(\theta_{13})$ is such that lower values of $\sin^2(\theta_{13})$ tend towards lower values of $|\Delta m_{32}^2|$. This moves the posterior distribution towards the known oscillation parameter $\Delta m_{32}^2 = 2.509 \times 10^{-3}\text{eV}^2$.

Table 8.4 presents the fraction of steps in each hierarchy and octant model for the fit after the reactor constraint has been applied. The reactor constraint significantly increases the NH preference, increasing the Bayes factor from

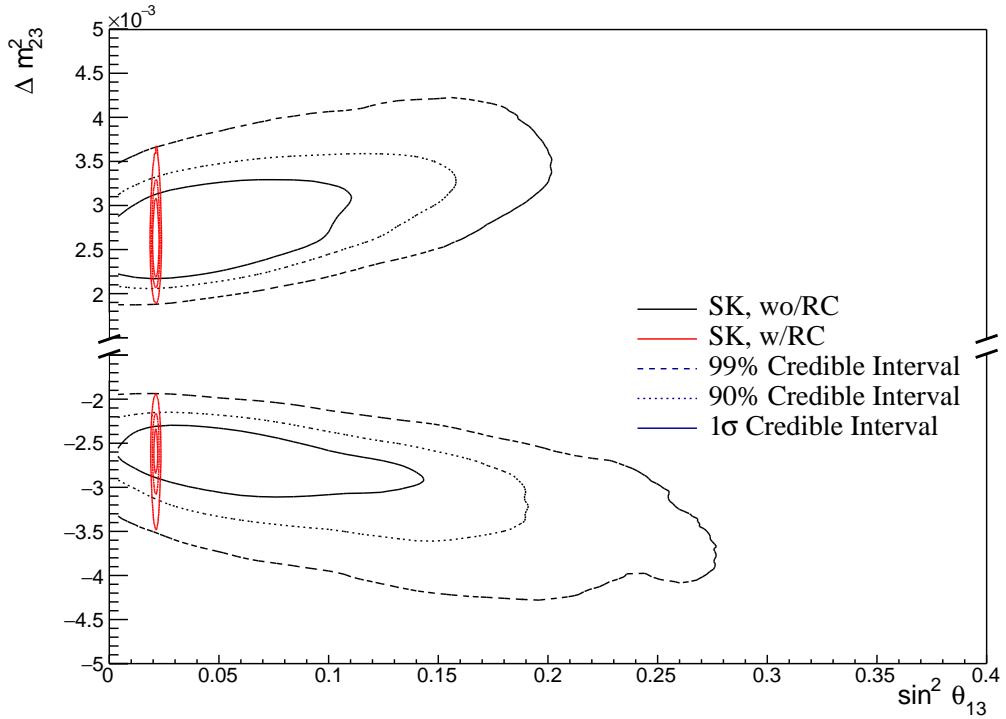


Figure 8.15: The two-dimensional posterior probability density distribution in Δm_{32}^2 – $\sin^2(\theta_{13})$ compared between the SK atmospheric-only fit (Black) and the SK atmospheric fit with the reactor constraint (Red). The distributions are marginalised over both hierarchies.

	LO ($\sin^2 \theta_{23} < 0.5$)	UO ($\sin^2 \theta_{23} > 0.5$)	Sum
NH ($\Delta m_{32}^2 > 0$)	0.21	0.53	0.74
IH ($\Delta m_{32}^2 < 0$)	0.08	0.18	0.26
Sum	0.29	0.71	1.00

Table 8.4: The distribution of steps in an SK atmospheric with reactor constraint fit, presented as the fraction of steps in the upper (UO) and lower (LO) octants and the normal (NH) and inverted (IH) hierarchies. The Bayes factors are calculated as $B(\text{NH}/\text{IH}) = 2.86$ and $B(\text{UO}/\text{LO}) = 2.39$.

3080 $B(\text{NH}/\text{IH}) = 1.37$ to $B(\text{NH}/\text{IH}) = 2.86$ when the reactor constraint is applied.

3081 This is still defined as a weak preference for NH hypothesis according to Jeffrey's

3082 scale (see Table 4.1), however, it is a stronger preference than when the constraint

3083 is not applied. The preference for the correct octant model is slightly increased

3084 by the application of the reactor constraint which is consistent with expectation.

3085 However, the conclusion that would be made does not significantly change.

3086 8.3.3 Application of Near Detector Constraints for Atmospheric 3087 Samples

3088 The choice of applying the near detector constraints to the low-energy atmo-
3089 spheric samples was introduced in subsection 6.4.3. This subsection illustrates
3090 the effect of that choice on the sensitivities of the atmospheric samples to the
3091 oscillation parameters. This Asimov data was generated assuming the ‘AsimovA’
3092 oscillation parameter set defined in Table 2.2 and the post-BANFF systematic
3093 parameter tune.

3094 The change in sensitivity on δ_{CP} is given in Figure 8.16. The reactor constraint
3095 is not applied in either of the fits within this comparison. The shape of the
3096 posterior is similar although less peaked at the Asimov point ($\delta_{CP} = -1.601$)
3097 and more symmetric between the regions of $\delta_{CP} = -1.601$ and $\delta_{CP} \sim 2.5$. The
3098 width of the 1σ credible intervals are approximately the same (identical to within
3099 a bin width) and the same conclusion holds for the higher credible intervals. The
3100 change in sensitivity to other oscillation parameters has been studied and no
3101 significant discrepancies were found. As expected, the sensitivities are statistics
3102 dominated and the exact choice of systematic model and constraint does not
3103 significantly affect the physics conclusions one would make from this analysis.

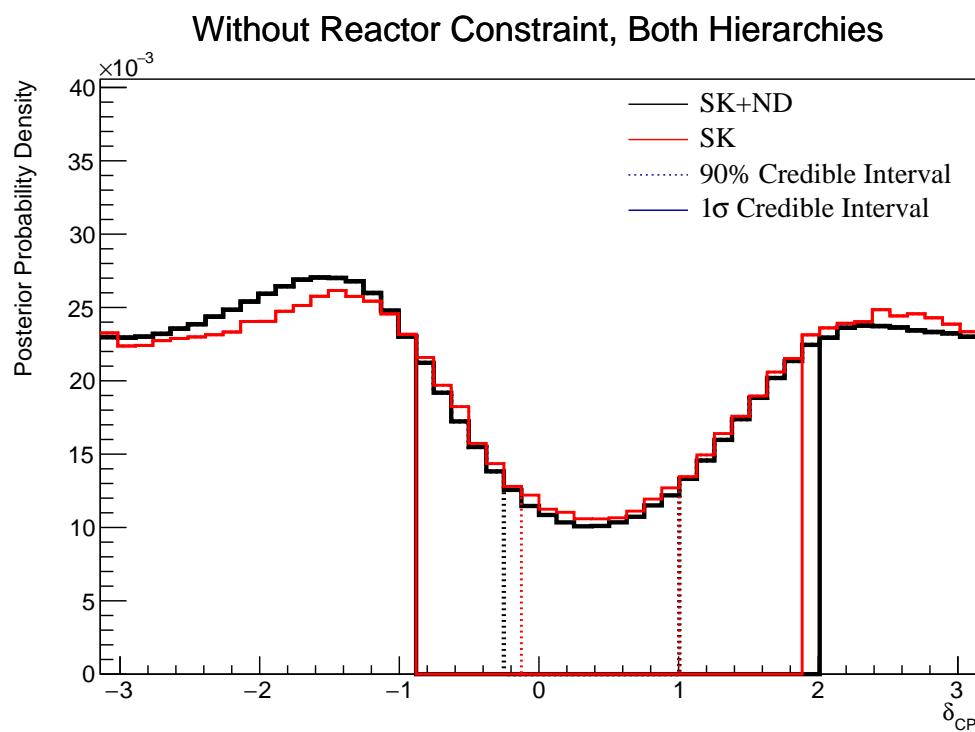


Figure 8.16: The one-dimensional posterior probability density distribution in δ_{CP} compared between the SK atmospheric-only fit where the near detector constraint is (Black) and is not (Red) applied. The distributions are marginalised over both hierarchies.

3104 8.3.4 Atmospheric and Beam Sensitivity without Reactor Con- 3105 straint

3106 This section presents the sensitivities of the simultaneous beam and atmospheric
 3107 analysis where the reactor constraint is not applied. Similar to the previous
 3108 studies, the Asimov data is built assuming the post-BANFF cross-section tune
 3109 and Asimov A oscillation parameters defined in Table 2.2. This fit uses all 18 near
 3110 detector beam samples, 5 far detector beam samples, and 18 atmospheric samples.
 3111 The sensitivity to δ_{CP} , marginalised over both hierarchies, is given in Figure 8.17.

Without Reactor Constraint, Both Hierarchies

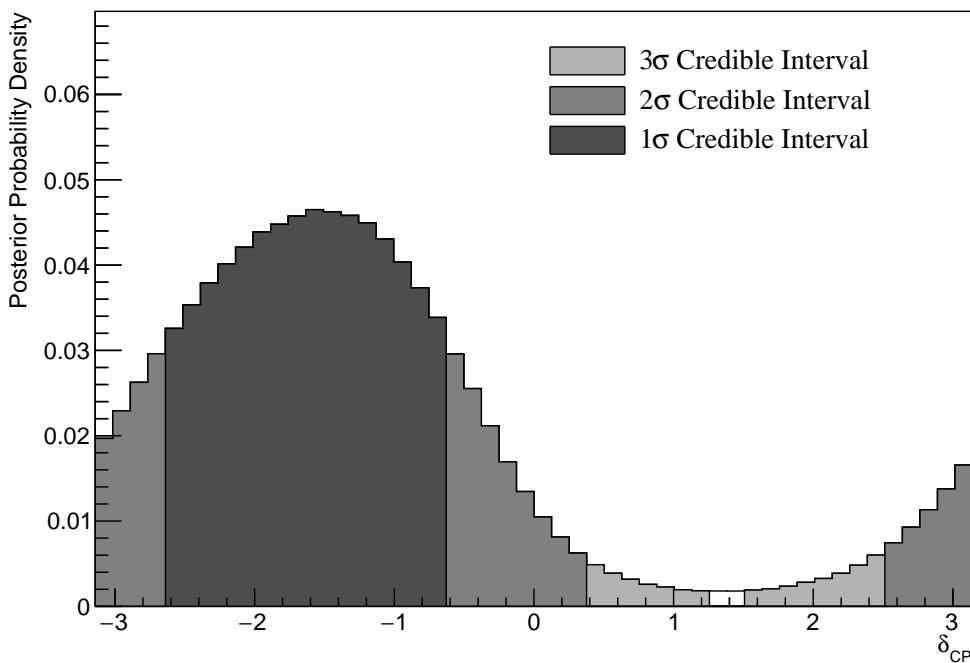


Figure 8.17: The one-dimensional posterior probability density distribution in δ_{CP} , marginalised over both hierarchies, from the joint beam and atmospheric fit. The reactor constraint is not applied.

3112 The credible intervals and highest posterior distribution for each oscillation
 3113 parameter is given in Table 8.5. The highest posterior probability density is
 3114 $\delta_{CP} = -1.58$ and is compatible with the known Asimov A value of $\delta_{CP} = -1.601$.
 3115 The CP-conserving values of $\delta_{CP} = 0, \pi, -\pi$ are disfavoured at 1 σ credible
 3116 interval. There is also a region around $\delta_{CP} = 1.4$ which is disfavoured at more
 3117 than 3 σ . Whilst these conclusions can only be made at this particular Asimov

point, it does show that if the true value of δ_{CP} was CP-violating, this joint analysis would be able to disfavour CP conserving values at over 1σ without any external constraints. The highest posterior probability density does move further away from the Asimov point when only steps in the NH region are considered. This is due to the correlations between the value of δ_{CP} and the mass hierarchy, as will be later discussed.

Parameter	Interval	HPD
δ_{CP} , (BH)	$[-2.64, -0.63]$	-1.57
δ_{CP} , (NH)	$[-2.76, -0.63]$	-1.45
δ_{CP} , (IH)	$[-2.39, -0.88]$	-1.57
Δm_{32}^2 (BH) [$\times 10^{-3}\text{eV}^2$]	[2.46, 2.58]	2.49
Δm_{32}^2 (NH) [$\times 10^{-3}\text{eV}^2$]	[2.48, 2.56]	2.51
Δm_{32}^2 (IH) [$\times 10^{-3}\text{eV}^2$]	$[-2.60, -2.52]$	-2.55
$\sin^2(\theta_{23})$ (BH)	[0.48, 0.55]	0.509
$\sin^2(\theta_{23})$ (NH)	[0.48, 0.55]	0.509
$\sin^2(\theta_{23})$ (IH)	[0.48, 0.55]	0.521

Table 8.5: The position of the highest posterior probability density (HPD) and width of the 1σ credible interval for the joint beam and atmospheric fit. The reactor constraint is not applied. The values are presented by which hierarchy hypothesis is assumed: marginalised over both hierarchies (BH), normal hierarchy only (NH), and inverted hierarchy only (IH).

The sensitivity to Δm_{32}^2 is illustrated in Figure 8.18, marginalised over both hierarchies. Notably, the 1σ credible interval is entirely contained within the normal hierarchy region, as illustrated in Table 8.5. This illustrates reasonable sensitivity to the mass hierarchy model. This is also reflected in the 1σ credible intervals being approximately the same when they are made considering both hierarchies and when considering only the NH. The known oscillation parameter is $\Delta m_{32}^2 = 2.509 \times 10^{-3}\text{eV}^2$. The normal hierarchy distribution favours this value with the highest posterior probability density of $\Delta m_{32}^2 = 2.51 \times 10^{-3}\text{eV}^2$.

The fraction of steps in each of the mass hierarchy regions and octants of $\sin^2(\theta_{23})$ is given in Table 8.6. The Bayes factors are determined to be $B(\text{NH}/\text{IH}) =$

3134 3.67 and $B(\text{UO}/\text{LO}) = 1.74$. Jeffrey's scale (presented in Table 4.1) states that
3135 this value of the hierarchy Bayes factor illustrates substantial evidence for the
3136 normal hierarchy hypothesis. This corresponds to the correct hypothesis given
3137 the known oscillation parameters. It is a stronger statement than the atmospheric-
3138 only analysis can provide. It is important to note that this is a substantial
3139 preference that requires no external constraints required. The Bayes factor for
3140 octant determination represents a weak preference for the upper octant but does
3141 select the correct octant model.

	LO ($\sin^2 \theta_{23} < 0.5$)	UO ($\sin^2 \theta_{23} > 0.5$)	Sum
NH ($\Delta m_{32}^2 > 0$)	0.29	0.50	0.79
IH ($\Delta m_{32}^2 < 0$)	0.08	0.13	0.21
Sum	0.37	0.63	1.00

Table 8.6: The distribution of steps in a joint beam and atmospheric fit, presented as the fraction of steps in the upper (UO) and lower (LO) octants and the normal (NH) and inverted (IH) hierarchies. The reactor constraint is not applied. The Bayes factors are calculated as $B(\text{NH}/\text{IH}) = 3.67$ and $B(\text{UO}/\text{LO}) = 1.74$.

3142 The sensitivity to $\sin^2(\theta_{23})$ is presented in Figure 8.19. There is a clear
3143 preference for the upper octant but the peak of the distribution is relatively
3144 flat. It peaks at $\sin^2(\theta_{23}) = 0.509$ which is in the region of the known value of
3145 $\sin^2(\theta_{23}) = 0.528$. The difference in the highest posterior distribution and the
3146 width of the credible interval is relatively unchanged when considering different
3147 hierarchy models showing no strong correlation between $\sin^2(\theta_{23})$ and $|\Delta m_{32}^2|$.
3148 The sensitivity presented as a function of the appearance parameters ($\sin^2(\theta_{13}) -$
3149 δ_{CP}) is given in Figure 8.20. As expected, the contours follow that given in
3150 Figure 8.2, where the 2σ credible intervals have a closed contour excluding the
3151 region around $\delta_{CP} \sim 1.2$. The width of the 3σ credible interval is also clearly
3152 dependent upon the value of δ_{CP} . Close to the Asimov point, $\delta_{CP} = -1.601$, the
3153 width of the 3σ credible interval approximately spans $\sin^2(\theta_{13}) = [0.013, 0.04]$.
3154 This is reduced to a region of $\sin^2(\theta_{13}) = [0.023, 0.042]$ at the most disfavoured
3155 value of δ_{CP} . This follows the behaviour shown in the likelihood scans. The 1σ

Without Reactor Constraint, Both Hierarchies

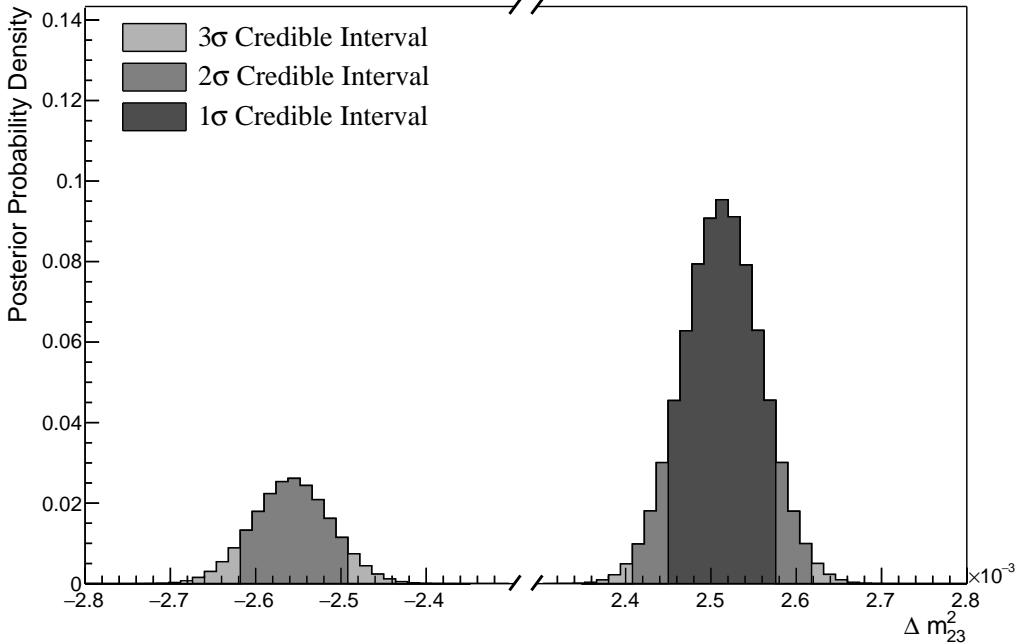


Figure 8.18: The one-dimensional posterior probability density distribution in Δm_{32}^2 , marginalised over both hierarchies, from the joint beam and atmospheric fit. The reactor constraint is not applied.

credible interval is consistent with both the known oscillation parameter and the reactor constraint ($\sin^2(\theta_{13}) = (2.18 \pm 0.08) \times 10^{-2}$). Application of the reactor constraint would be expected to decrease the width of the 1σ credible intervals of δ_{CP} due to the triangular shape of the posterior probability.

The sensitivity in terms of the ‘disappearance’ parameters marginalised over both hierarchies is given in Figure 8.21. The area contained within the IH credible intervals is significantly smaller than those in the NH region. This is reflected in the IH credible intervals being tighter in the $\sin^2(\theta_{23})$ dimension. No significant correlation is observed between the value of $\sin^2(\theta_{23})$ and $|\Delta m_{32}^2|$.

The two-dimensional posterior distribution for each permutation of the oscillation parameters of interest is given in Figure 8.22. The most notable observation is that the $\sin^2(\theta_{13})$ and $\sin^2(\theta_{23})$ are anti-correlated. If the value of $\sin^2(\theta_{13})$ was known to be closer to the known oscillation parameter value, the preferred value of $\sin^2(\theta_{23})$ would increase furthering the preference for the UO. That

Without Reactor Constraint, Both Hierarchies

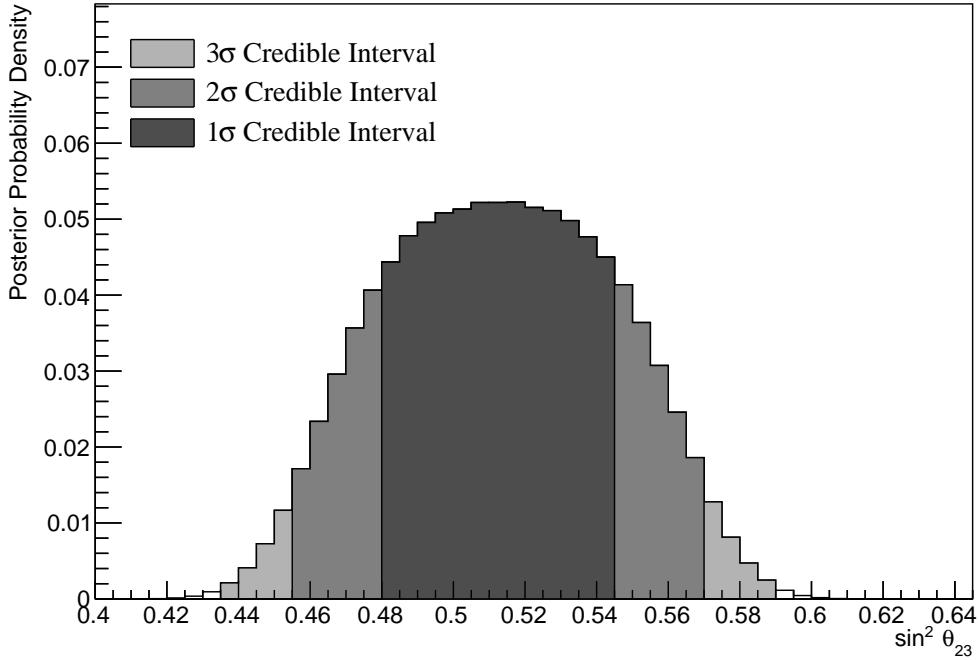


Figure 8.19: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$, marginalised over both hierarchies, from the joint beam-atmospheric fit. The reactor constraint is not applied.

would move the highest posterior probability closer in line with the Asimov value. This also means that the preference for the UO would be increased if the reactor constraint was to be applied.

Furthermore, the δ_{CP} and $|\Delta m_{32}^2|$ oscillation parameters are anti-correlated, such that higher values of $|\Delta m_{32}^2|$ prefer lower values of δ_{CP} . Whilst this is an interesting result on its own, the width of the Δm_{32}^2 contours also depend on $\sin^2(\theta_{13})$. This introduces another correlation effect that could modify the sensitivity to δ_{CP} once the reactor constraint is applied.

The correlation between $\sin^2(\theta_{13})$ and Δm_{32}^2 can be seen in Figure 8.23. A much larger fraction of the posterior distribution is contained in the NH for lower values of $\sin^2(\theta_{13})$. Consequently, the application of the reactor constraint would be expected to significantly increase the preference for NH.

Without Reactor Constraint, Both Hierarchies

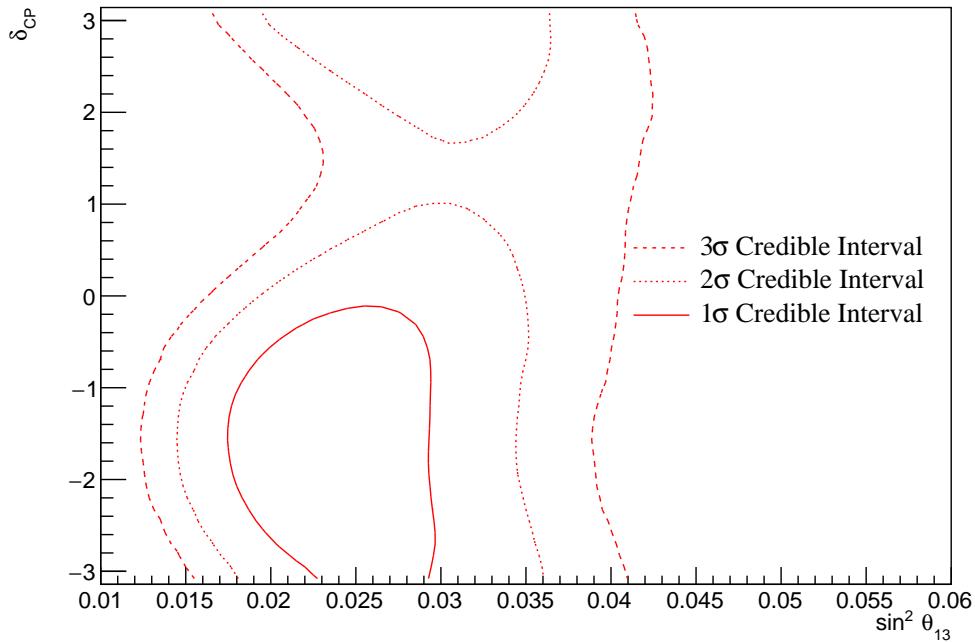


Figure 8.20: The two-dimensional posterior probability density distribution in δ_{CP} – $\sin^2(\theta_{13})$, marginalised over both hierarchies, from the joint beam and atmospheric fit. The reactor constraint is not applied.

Without Reactor Constraint, Both Hierarchies

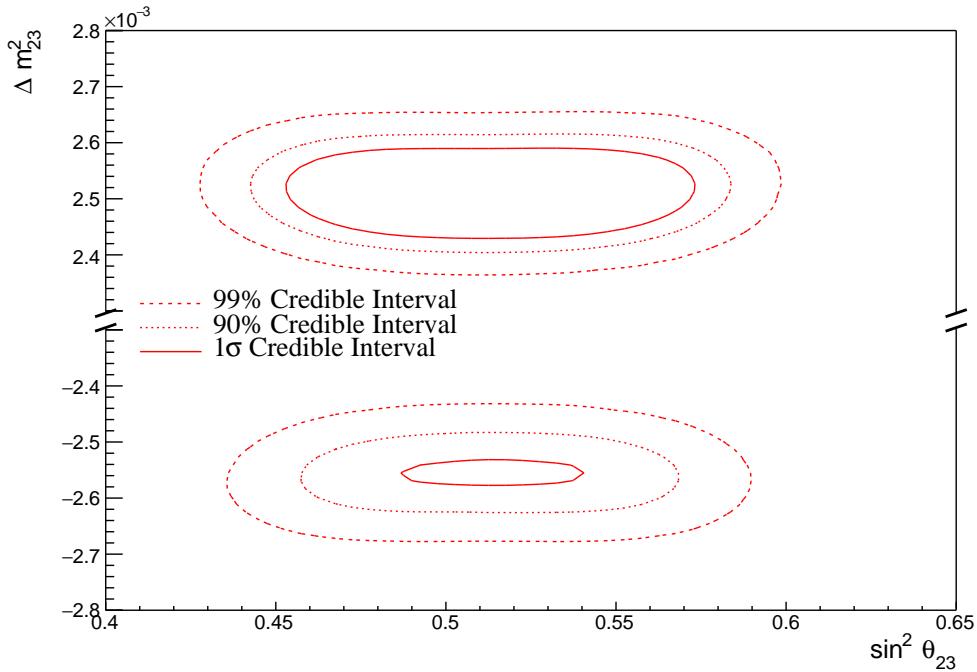


Figure 8.21: The two-dimensional posterior probability density distribution in Δm_{32}^2 – $\sin^2(\theta_{23})$, marginalised over both hierarchies, from the joint beam and atmospheric fit. The reactor constraint is not applied.

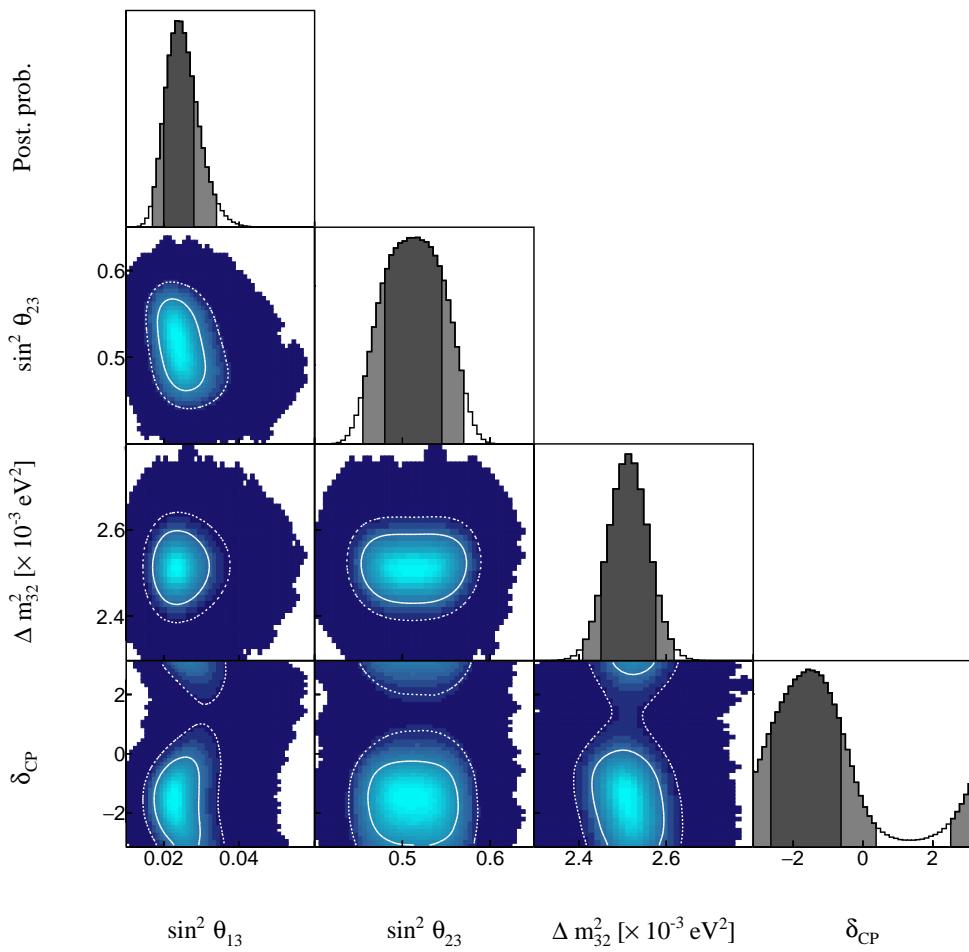


Figure 8.22: The posterior probability density distribution from the joint beam and atmospheric fit. The reactor constraint is not applied. The distribution is given for each two-dimensional permutation of the oscillation parameters of interest. The one-dimensional distribution of each parameter is also given.

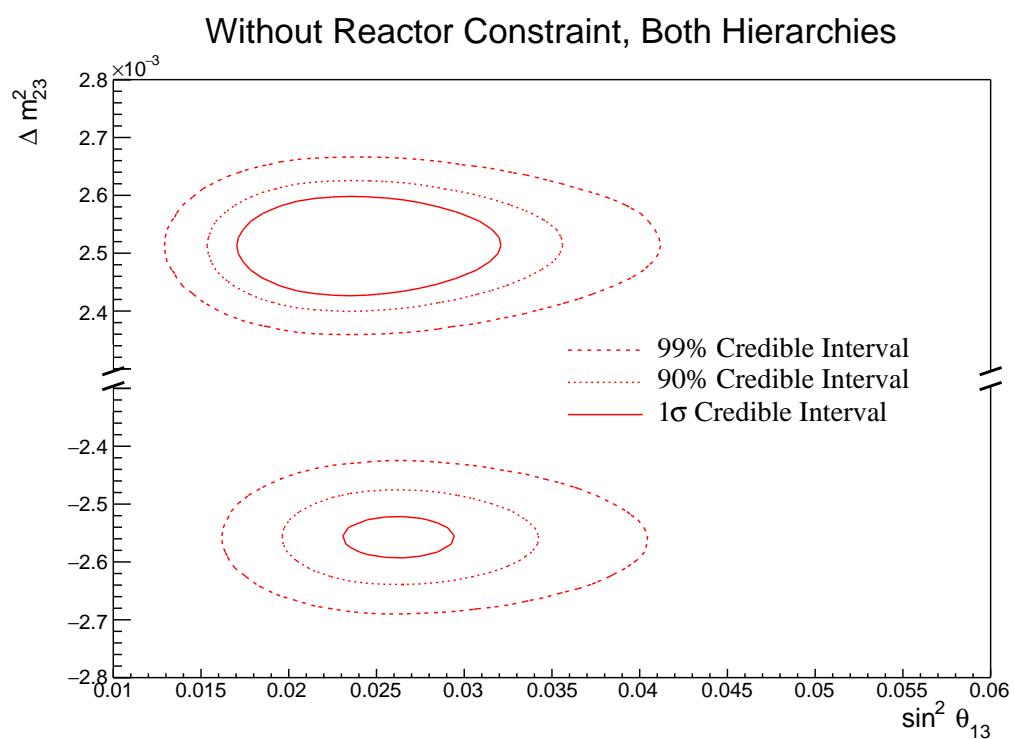


Figure 8.23: The two-dimensional posterior probability density distribution in $\Delta m_{32}^2 - \sin^2(\theta_{13})$, marginalised over both hierarchies, from the joint beam and atmospheric fit. The reactor constraint is not applied.

3182 8.3.5 Atmospheric and Beam Sensitivity with Reactor Constraint

3183 This section presents the sensitivities of the joint beam and atmospheric fit when
 3184 the reactor constraint is applied to $\sin^2(\theta_{13})$. As with the previous studies, the
 3185 Asimov data is made using the AsimovA oscillation parameter set defined in
 3186 Table 2.2 and the post-BANFF systematic parameter tune.

With Reactor Constraint, Both Hierarchies

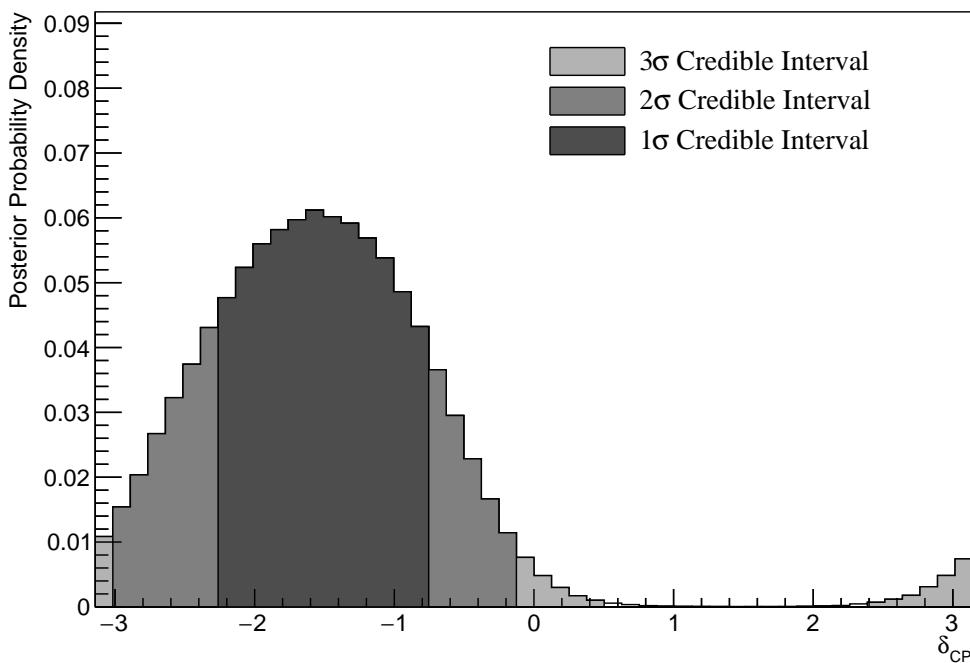


Figure 8.24: The one-dimensional posterior probability density distribution in δ_{CP} , marginalised over both hierarchies, from the joint beam and atmospheric fit where the reactor constraint is applied.

3187 Figure 8.24 illustrates the sensitivity to δ_{CP} , marginalised over both hierarchies.
 3188 The CP-conserving values of $\delta_{CP} = -\pi, 0, \pi$ are disfavoured at 2σ . Furthermore,
 3189 the 3σ credible interval excludes the region of $\delta_{CP} = [0.50, 2.39]$. Thus clearly
 3190 disfavouring the region of $\delta_{CP} = \pi/2$ at more than 3σ for this particular set
 3191 of known oscillation parameters. The width of the 1σ credible intervals and
 3192 the position of the highest posterior probability density is given in Table 8.7.
 3193 The highest posterior probability density in δ_{CP} is calculated as $\delta_{CP} = -1.57$
 3194 showing no significant biases in the determination of the known oscillation

parameters. The posterior distribution is more peaked around the known oscillation parameter value of $\delta_{CP} = -1.601$, as compared to the sensitivities when the reactor constraint is not applied (subsection 8.3.4). This follows from the correlations shown in Figure 8.20, where a lower value of $\sin^2(\theta_{13})$ results in tighter constraints on δ_{CP} .

The effect of applying the reactor constraint for δ_{CP} in the joint beam-atmospheric fit is presented in Figure 8.25. The posterior distribution from the two fits are marginalised over both hierarchies. Clearly, the reactor constraint improves the ability of the fit to select the known oscillation parameter as the shape of the distribution is much more peaked. This is also evidenced by the tightening of the 1σ and 90% credible intervals. Additionally, the disfavoured region of $1 < \delta_{CP} < 2$ is wider when the reactor constraint is applied.

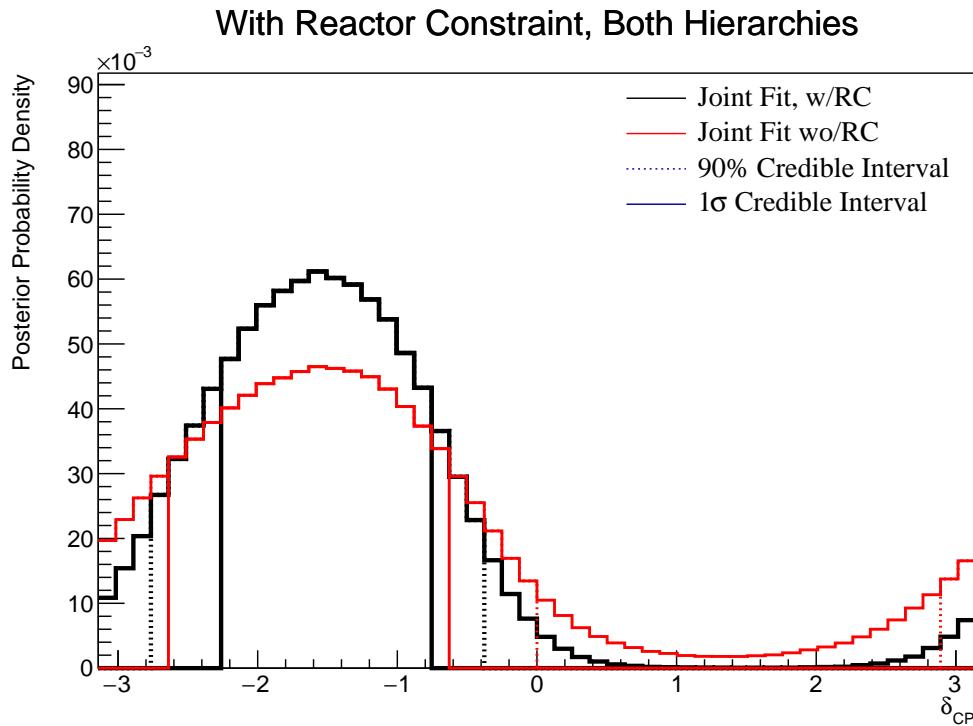


Figure 8.25: The one-dimensional posterior probability density distribution in δ_{CP} compared between the joint beam-atmospheric fit (Red) and the joint beam and atmospheric fit with the reactor constraint (Black). The distributions are marginalised over both hierarchies.

The sensitivity to $\sin^2(\theta_{23})$, marginalised over both hierarchies, is given in Figure 8.26. The highest posterior probability density is located at $\sin^2(\theta_{23}) = 0.527$

Parameter	Interval	HPD
δ_{CP} , (BH)	[-2.26, -0.75]	-1.57
δ_{CP} , (NH)	[-2.26, -0.75]	-1.57
δ_{CP} , (IH)	[-2.13, -1.00]	-1.57
Δm_{32}^2 (BH) [$\times 10^{-3}\text{eV}^2$]	[2.46, 2.52]	2.49
Δm_{32}^2 (NH) [$\times 10^{-3}\text{eV}^2$]	[2.48, 2.56]	2.51
Δm_{32}^2 (IH) [$\times 10^{-3}\text{eV}^2$]	[-2.60, -2.52]	-2.55
$\sin^2(\theta_{23})$ (BH)	[0.49, 0.55]	0.527
$\sin^2(\theta_{23})$ (NH)	[0.49, 0.55]	0.527
$\sin^2(\theta_{23})$ (IH)	[0.50, 0.56]	0.539

Table 8.7: The position of the highest posterior probability density (HPD) and width of the 1σ credible interval for the joint beam and atmospheric fit where the reactor constraint is applied. The values are presented by which hierarchy hypothesis is assumed: marginalised over both hierarchies (BH), normal hierarchy only (NH), and inverted hierarchy only (IH).

which agrees with the known value of $\sin^2(\theta_{23}) = 0.528$. The distribution clearly favours the UO with almost the entirety of the 1σ credible interval contained in the region. Figure 8.27 highlights the sensitivity of the joint fit both with and without the reactor constraint. The fit where the reactor constraint is applied selects the known value much better ($\sin^2(\theta_{23}) = 0.528$). Furthermore, the reactor constraint increases the UO preference which is evidenced by the distribution moving further away from the octant boundary. This indicates that there are marginalisation effects between the two mixing parameters. This follows from the correlation illustrated between $\sin^2(\theta_{23}) - \sin^2(\theta_{13})$ in Figure 8.22. The posterior distribution of the fit with reactor constraint is more peaked compared to the flatter distribution when the reactor constraint is not applied.

The fraction of steps contained within the two hierarchy and two octant models is given in Table 8.8. The reactor constraint significantly reduces the fraction of steps that are contained within the IH-LO region from 0.08 to 0.02, whilst significantly increasing the fraction of steps within the NH-UO region from 0.53 to 0.64. The application of the reactor constraint increases the Bayes factor

With Reactor Constraint, Both Hierarchies

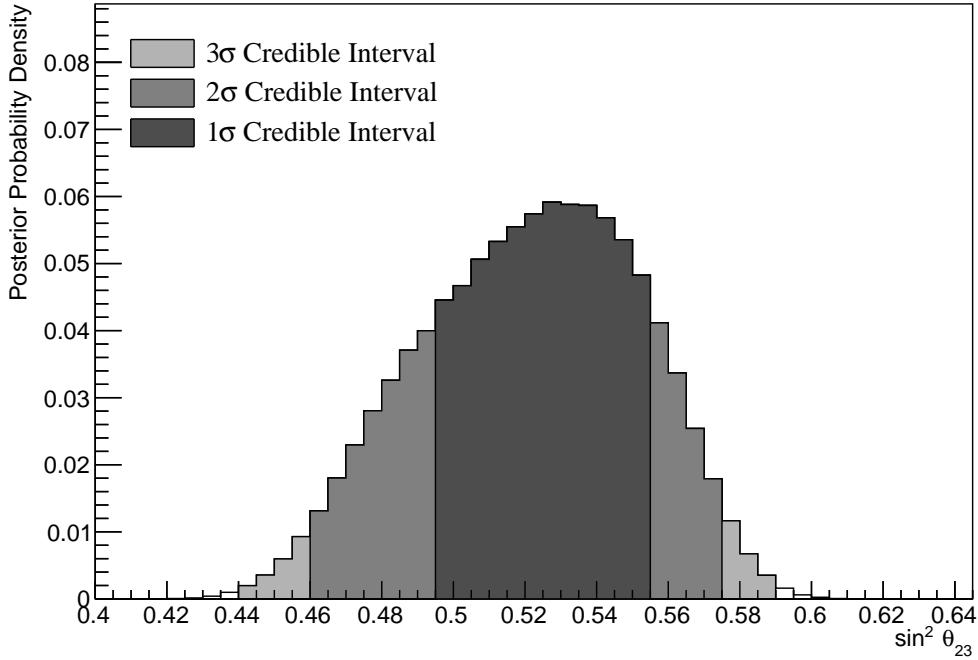


Figure 8.26: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$, marginalised over both hierarchies, from the joint beam and atmospheric fit where the reactor constraint is applied.

from $B(\text{NH}/\text{IH}) = 3.67$ to $B(\text{NH}/\text{IH}) = 7.29$. There is a very clear preference for the NH, with the Jeffreys scale stating a substantial preference for both fits (see subsection 4.3.3). The Bayes factor for UO preference is calculated as $B(\text{UO}/\text{LO}) = 2.86$. Whilst still a weak preference, this is certainly a stronger statement than the sensitivity when the reactor constraint is not applied.

	LO ($\sin^2 \theta_{23} < 0.5$)	UO ($\sin^2 \theta_{23} > 0.5$)	Sum
NH ($\Delta m_{32}^2 > 0$)	0.24	0.64	0.88
IH ($\Delta m_{32}^2 < 0$)	0.02	0.10	0.12
Sum	0.26	0.74	1.00

Table 8.8: The distribution of steps in a joint beam and atmospheric with the reactor constraint fit applied, presented as the fraction of steps in the upper (UO) and lower (LO) octants and the normal (NH) and inverted (IH) hierarchies. The Bayes factors are calculated as $B(\text{NH}/\text{IH}) = 7.29$ and $B(\text{UO}/\text{LO}) = 2.86$.

The sensitivity to Δm_{32}^2 , with the reactor constraint applied, is presented in

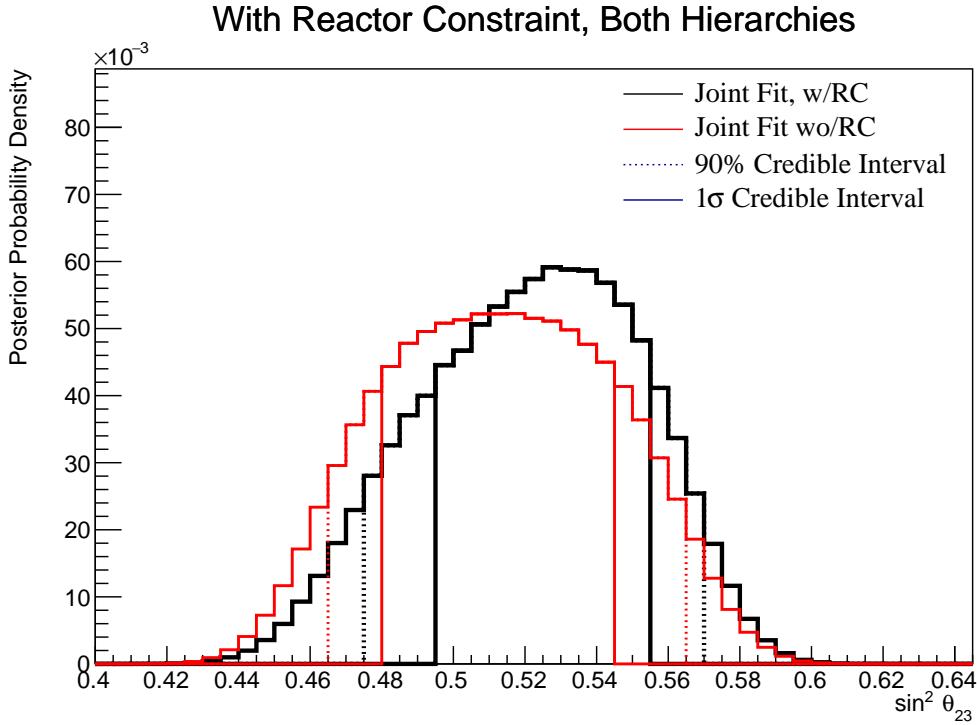


Figure 8.27: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Red) and the joint beam and atmospheric fit with the reactor constraint (Black). The distributions are marginalised over both hierarchies.

3231 Figure 8.28. The posterior distribution is marginalised over both hierarchies. As
 3232 expected, the 1σ credible interval is entirely contained within the NH region. The
 3233 position of the highest posterior probability density is given as $2.49 \times 10^{-3} \text{ eV}^2$,
 3234 illustrating no significant bias between the fit results and the known oscillation
 3235 parameters. The application of the reactor constraint does move significantly the
 3236 position of the credible intervals but does reduce their width.

3237 The sensitivity to the appearance parameters ($\sin^2(\theta_{13}) - \delta_{CP}$) is given in
 3238 Figure 8.29. The distribution is mostly uncorrelated between the two parameters
 3239 and is centered at the known oscillation parameters. The 1σ credible interval
 3240 excludes $\delta_{CP} = 0$ and $\delta_{CP} = (-)\pi$. Furthermore, the 3σ credible intervals
 3241 exclude the region of $\delta_{CP} = \pi/2$.

3242 The sensitivity to the disappearance parameters ($\sin^2(\theta_{23}) - \Delta m_{32}^2$) is illus-
 3243 trated in Figure 8.30. As expected from the one-dimensional distribution, the 1σ
 3244 credible interval is entirely contained within the NH region. Both the NH and

With Reactor Constraint, Both Hierarchies

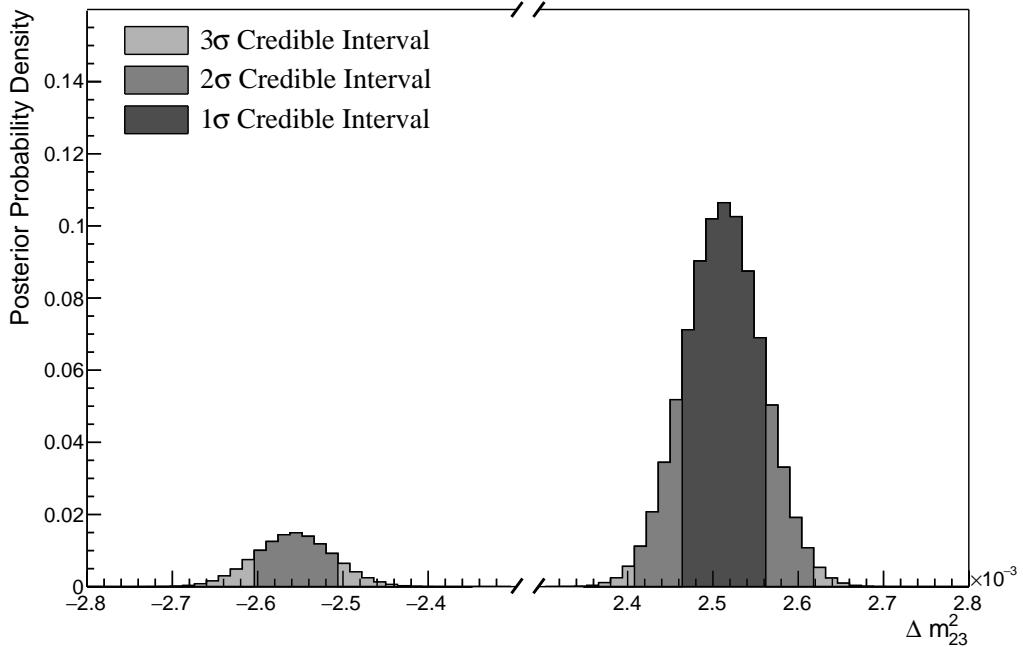


Figure 8.28: The one-dimensional posterior probability density distribution in Δm_{32}^2 , marginalised over both hierarchies, from the joint beam and atmospheric fit where the reactor constraint is applied.

3245 IH regions favour the UO, with a visually similar preference in both hierarchies.
 3246 The width of the Δm_{32}^2 1 σ credible interval does not significantly depend upon
 3247 the value or octant of $\sin^2(\theta_{23})$. This shows that there are no strong correlations
 3248 between these two parameters.

3249 Figure 8.31 illustrates the posterior distribution for each permutation of
 3250 two oscillation parameters of interest. The application of the reactor constraint
 3251 significantly reduces the correlations previously seen in Figure 8.22. There is
 3252 still a small correlation between δ_{CP} and Δm_{32}^2 . The application of the reactor
 3253 constraint has not significantly affected this correlation. The width of the 1 σ
 3254 credible interval in Δm_{32}^2 is wider for a value of $\delta_{CP} = 0$ as compared to a value
 3255 of $\delta_{CP} = \pi$. Similarly, the width of the 1 σ credible interval in δ_{CP} is smaller
 3256 for lower values of $\sin^2(\theta_{23})$.

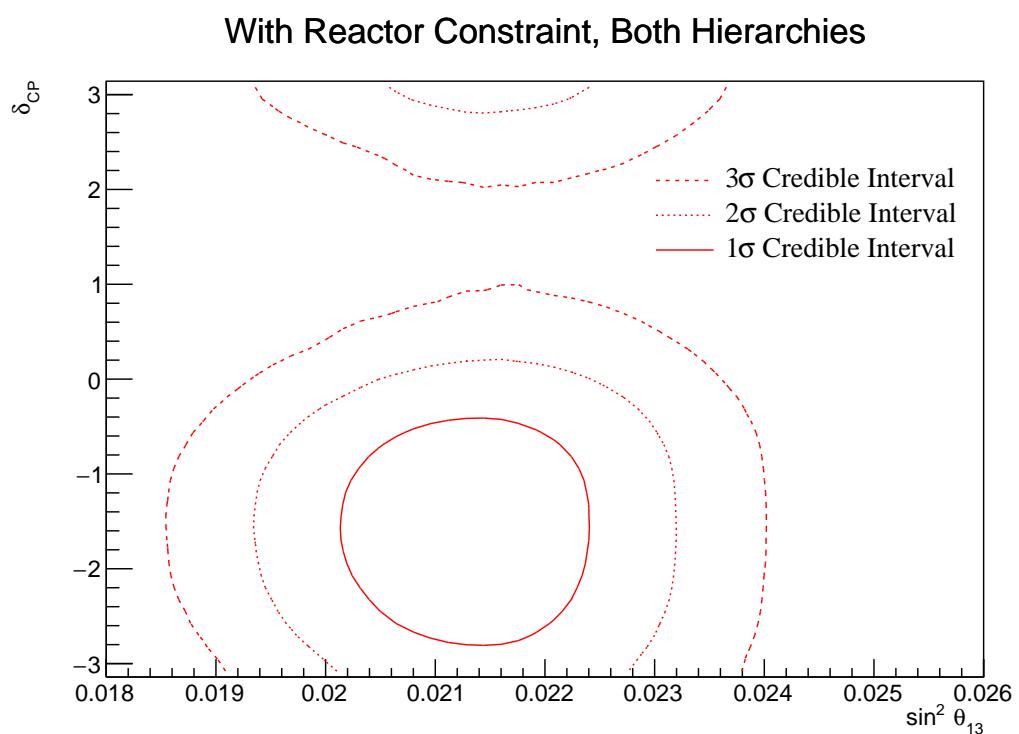


Figure 8.29: The two-dimensional posterior probability density distribution in $\delta_{CP} - \sin^2(\theta_{13})$, marginalised over both hierarchies, from the joint beam and atmospheric fit where the reactor constraint is applied.

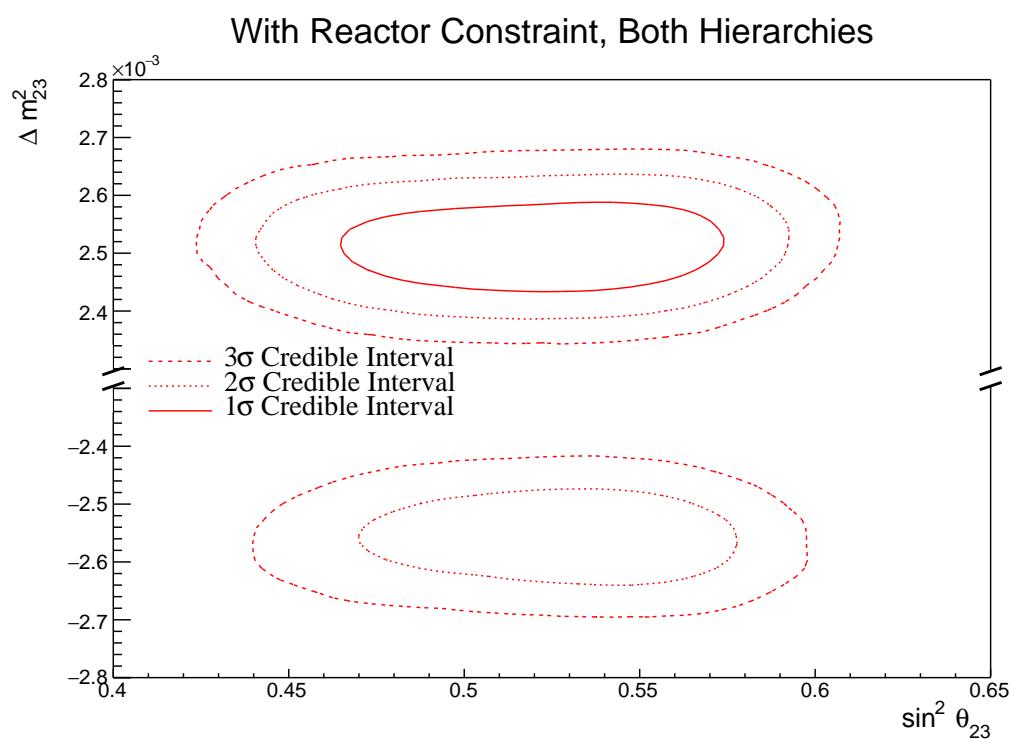


Figure 8.30: The two-dimensional posterior probability density distribution in $\Delta m_{32}^2 - \sin^2(\theta_{23})$, marginalised over both hierarchies, from the joint beam and atmospheric fit where the reactor constraint is applied.

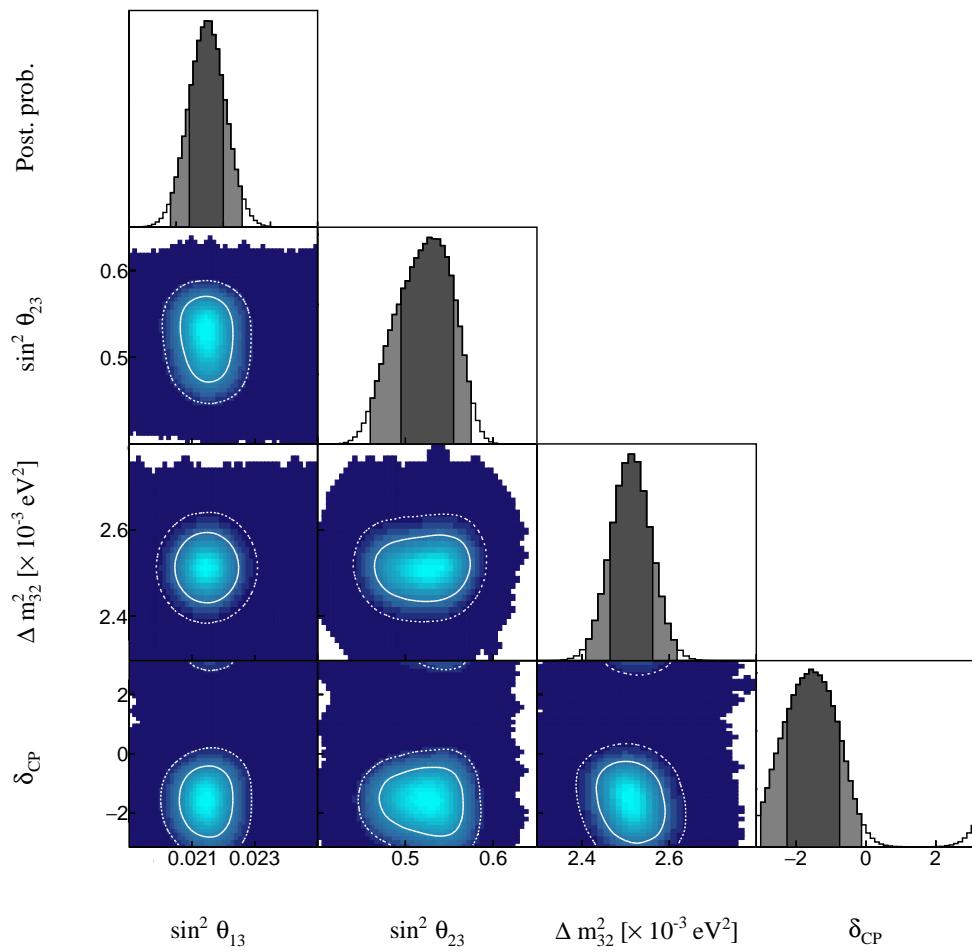


Figure 8.31: The posterior probability density distribution from the joint beam and atmospheric fit where the reactor constraint is applied. The distribution is given for each two-dimensional permutation of the oscillation parameters of interest. The one-dimensional distribution of each parameter is also given.

3257 8.3.6 Comparison to Latest T2K Sensitivities without Reactor 3258 Constraint

3259 The benefits of the joint beam and atmospheric analysis can be determined by
 3260 comparing the sensitivities to the beam-only analysis. This section presents those
 3261 comparisons for sensitivities built using the Asimov A oscillation parameters
 3262 defined in Table 2.2 and the post-BANFF systematic tune. The reactor constraint
 3263 is not applied within either of the fits used in these comparisons.

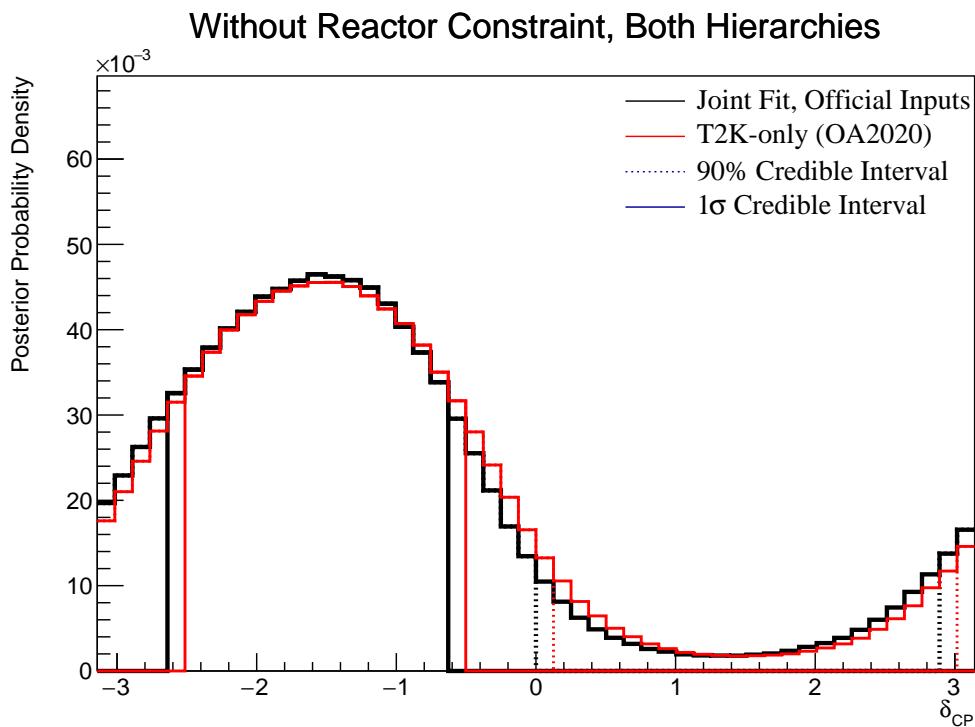


Figure 8.32: The one-dimensional posterior probability density distribution in δ_{CP} compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

3264 The sensitivity, marginalised over both hierarchies, to δ_{CP} from the joint
 3265 beam-atmospheric and beam-only fits is presented in Figure 8.32. As expected
 3266 from the likelihood scans (Figure 8.4), the sensitivity to δ_{CP} is not significantly
 3267 increased. This is because the known oscillation parameter value ($\delta_{CP} = -1.601$)
 3268 lies at the position where the beam samples dominate the sensitivity compared
 3269 to the SK samples.

The sensitivity to Δm_{32}^2 of the joint beam-atmospheric fit is illustrated in Figure 8.33, where the posterior distribution has been marginalised over both hierarchies. The 1σ credible interval of the joint beam and atmospheric fit is entirely contained within the NH region. This shows the significant increase in the ability of the fit to determine the correct mass hierarchy, as compared to the beam-only analysis. This is further evidenced by the fact that the 90% credible intervals from the joint fit are also tighter in the IH region as compared to the beam-only analysis. The Bayes factor for mass hierarchy determination for the beam-only and joint beam and atmospheric are $B(\text{NH}/\text{IH}) = 1.91$ and $B(\text{NH}/\text{IH}) = 3.67$, respectively. According to Jeffrey's scale (Table 4.1), the beam-only analysis represents a weak preference for the NH hypothesis whereas the joint fit returns a substantial preference for the NH hypothesis. To summarise, the joint beam-atmospheric fit has a substantial preference for the correct hierarchy without the requirement of external constraints.

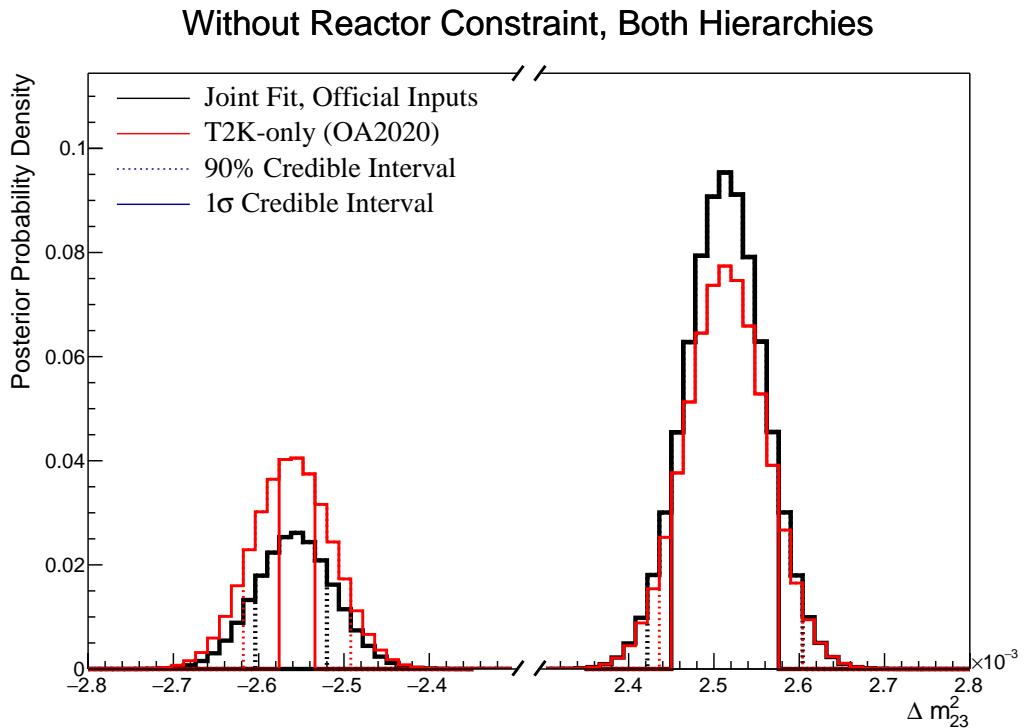


Figure 8.33: The one-dimensional posterior probability density distribution in Δm_{32}^2 compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

The sensitivity to $\sin^2(\theta_{23})$, marginalised over both hierarchies, for both the beam-only and joint beam and atmospheric analysis are presented in Figure 8.34. The peak of the posterior distribution from the joint analysis is more aligned with the known value of $\sin^2(\theta_{23}) = 0.528$ as compared to the beam-only analysis. This indicates that the marginalisation effects from other oscillation parameters ($\sin^2(\theta_{13}) - \sin^2(\theta_{23})$) presented in Figure 8.22) are less prevalent in the projection of this parameter. The Bayes factors for the beam-only and joint beam-atmospheric fit are $B(\text{UO}/\text{LO}) = 1.56$ and $B(\text{UO}/\text{LO}) = 1.74$, respectively. Therefore, the joint beam-atmospheric fit does prefer the UO more strongly than the beam-only analysis, albeit slightly.

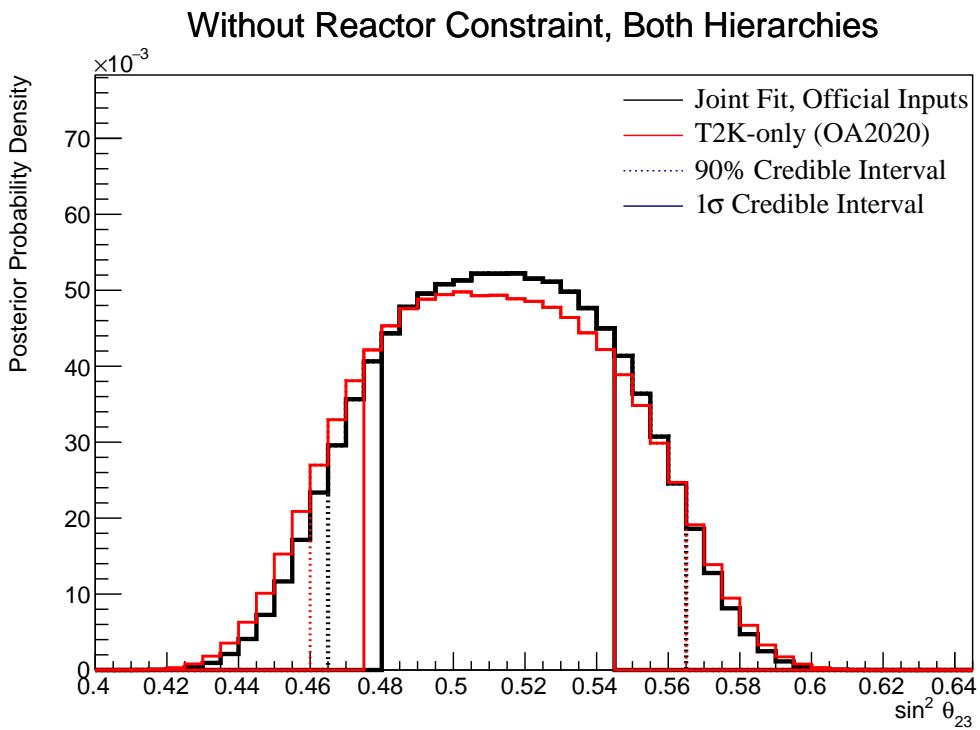


Figure 8.34: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

Whilst the beam-only and joint beam-atmospheric fits have similar sensitivity to δ_{CP} and $\sin^2(\theta_{23})$ when projected in one-dimension, the benefit of the joint analysis becomes more obvious when the sensitivities are presented in two-

dimensions. The sensitivity of the two fits to the appearance parameters ($\delta_{CP} - \sin^2(\theta_{13})$) are illustrated in Figure 8.35.

The width of the 99% joint fit credible interval in $\sin^2(\theta_{13})$ is squeezed in the region of $\delta_{CP} \sim 0$ compared to the beam-only analysis. This is the same behaviour that is seen in the appearance likelihood scans presented in Figure 8.2. The 1σ and 90% also exhibit slightly tighter constraints on δ_{CP} . This is most prevalent in the region of $\delta_{CP} \sim 0$ and $\sin^2(\theta_{13}) \sim 0.03$. Whilst the atmospheric samples do not have significant sensitivity to $\sin^2(\theta_{13})$ (as shown in Figure 8.1), they aid in breaking the degeneracy between the oscillation parameters allowing for tighter constraints.

The sensitivity to the disappearance parameters $\sin^2(\theta_{23}) - \Delta m_{32}^2$, marginalised over both hierarchies, is presented in Figure 8.36 for both the beam-only and joint beam-atmospheric fits. Whilst the one-dimensional sensitivity comparisons considered so far show the improvements of the joint fit, the two-dimensional projection really shows the benefit of adding the atmospheric samples to the beam samples. The area contained within the IH credible intervals is drastically reduced in the joint fit. This follows from the better determination of the mass hierarchy seen in the Bayes factor comparisons. The 1σ joint fit credible interval in the IH region more strongly favours the UO as compared to the beam-only fit. Even in the NH region, the width of the credible intervals in $\sin^2(\theta_{23})$ decrease, albeit to a smaller extent.

The change in sensitivity to $\delta_{CP} - \Delta m_{32}^2$ is illustrated in Figure 8.37. As expected, the contours presented within the IH region are much smaller in the joint fit due to the increased sensitivity to mass hierarchy determination. This culminates in a region around $\delta_{CP} \sim \pi/2$ which is excluded at 3σ . This behaviour is not present within the beam-only analysis. Consistent with the previous observations, the area contained within the IH credible intervals is significantly reduced in comparison to the beam-only analysis.

The sensitivity to Δm_{32}^2 and $\sin^2(\theta_{23})$, as a function of $\sin^2(\theta_{13})$, is presented in Figure 8.38 and Figure 8.39, respectively. These sensitivities are marginalised

Without Reactor Constraint, Both Hierarchies

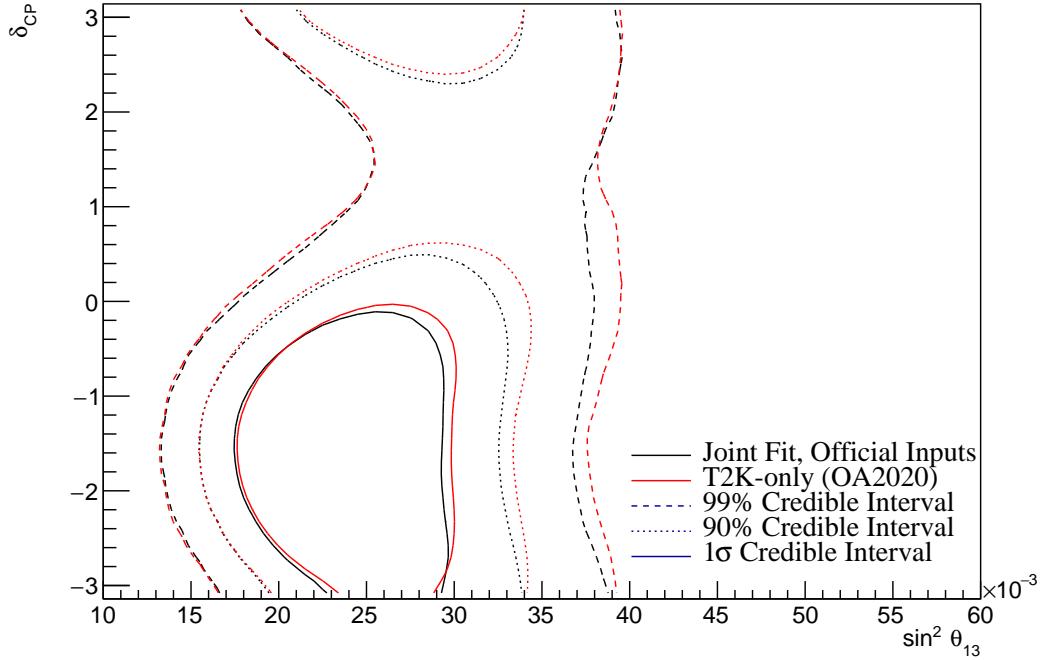


Figure 8.35: The two-dimensional posterior probability density distribution in δ_{CP} – $\sin^2(\theta_{13})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

over both hierarchies. As expected from the previous observations, the Δm_{32}^2 contours within IH region of the joint fit are much smaller than the beam-only analysis. Notably, the joint fit IH 1σ credible intervals exclude the region around the reactor constraint. This is not a bias from the fit as the known value for Δm_{32}^2 is in the NH region. This does suggest that the application of the reactor constraint would further increase the preference for NH in the joint fit as compared to its effect on the beam-only analysis.

The beam-only and joint beam-atmospheric fits have a slightly different contour shape between the $\sin^2(\theta_{13})$ and $\sin^2(\theta_{23})$ parameters, as illustrated by Figure 8.39. The joint analysis disfavours the wrong octant hypothesis more strongly in the region of high $\sin^2(\theta_{13})$. This suggests that the application of the reactor constraint will favour the UO more strongly in the joint analysis compared to the beam-only analysis.

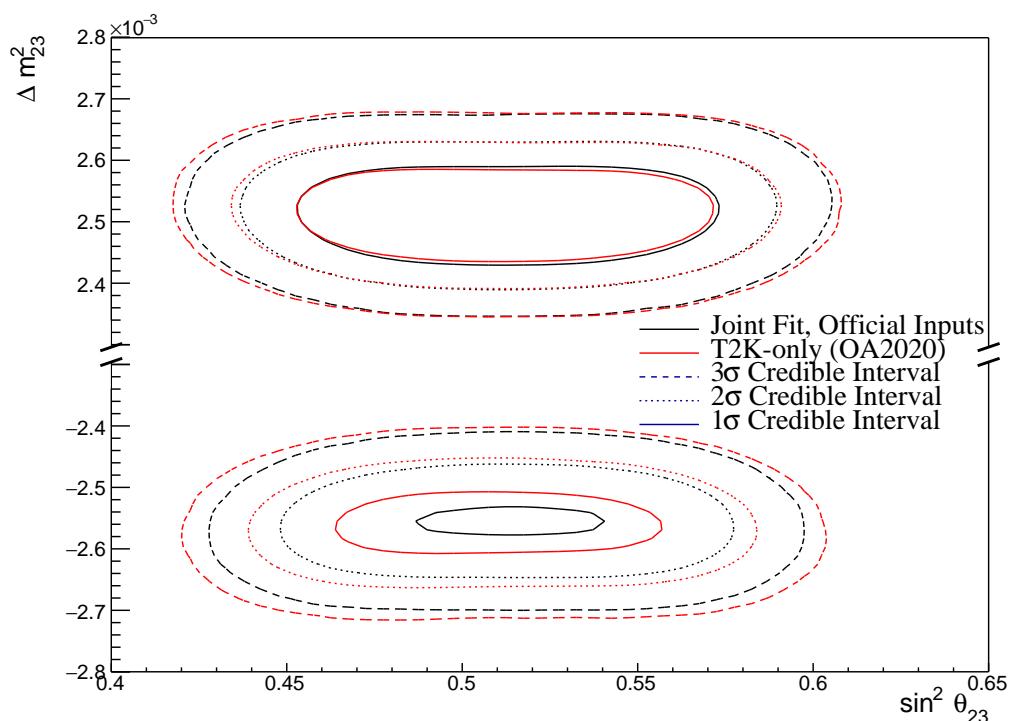


Figure 8.36: The two-dimensional posterior probability density distribution in Δm_{32}^2 – $\sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

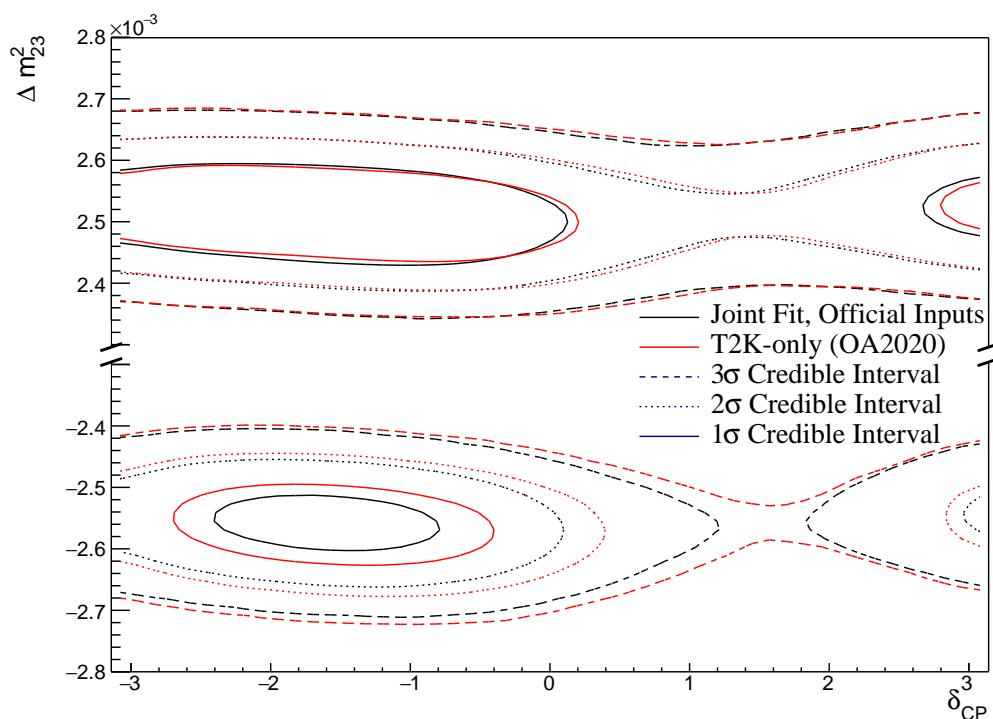


Figure 8.37: The two-dimensional posterior probability density distribution in $\Delta m_{32}^2 - \Delta_{CP}$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

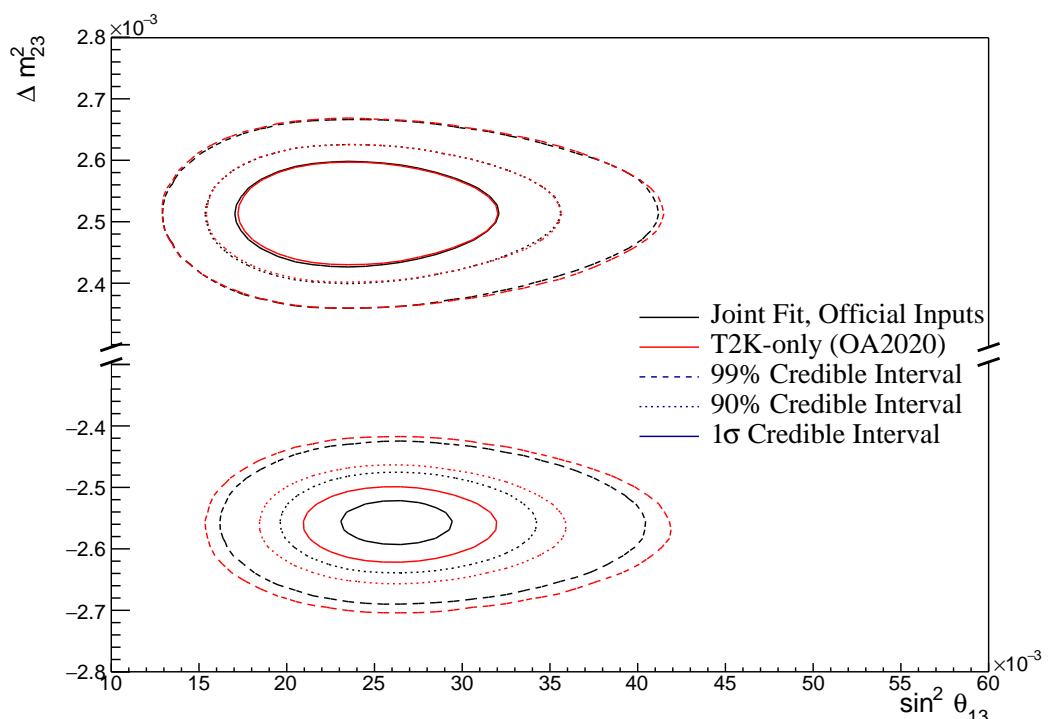


Figure 8.38: The two-dimensional posterior probability density distribution in Δm_{32}^2 – $\sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

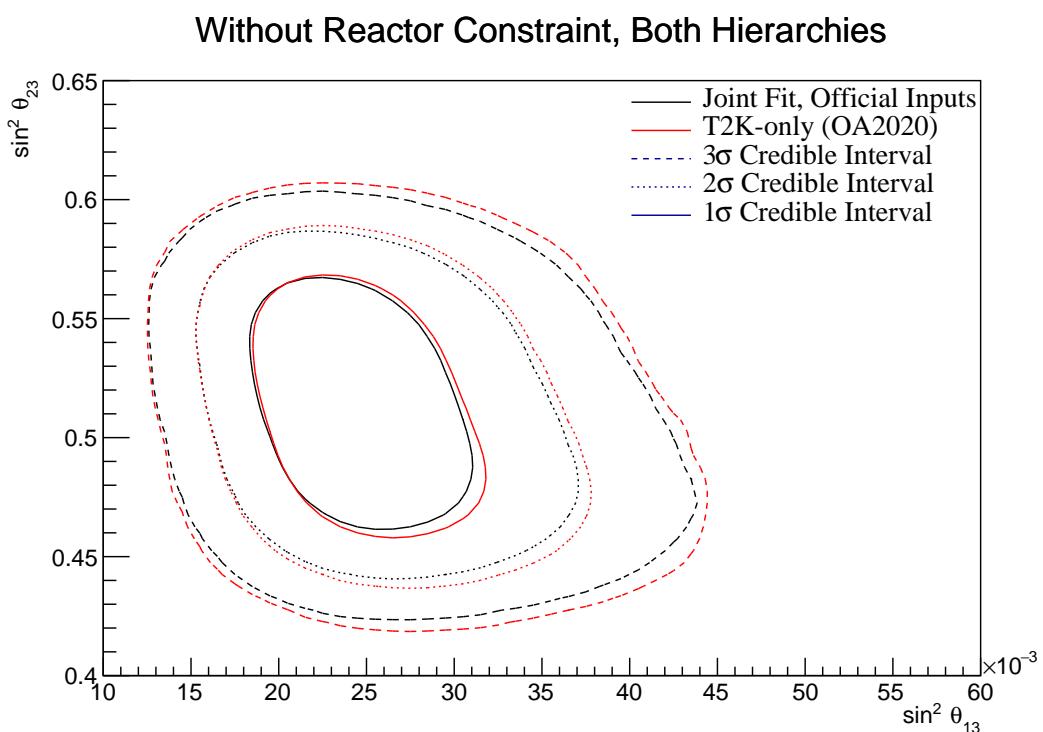


Figure 8.39: The two-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$ – $\sin^2(\theta_{13})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

3340 8.3.7 Comparison to Latest T2K Sensitivities with Reactor Con- 3341 straint

3342 The comparison between the beam-only and joint beam-atmospheric fits are
 3343 compared in subsection 8.3.6. Those comparisons were made with the reactor
 3344 constraint not applied to either of the fits. This section illustrates the com-
 3345 parison when the reactor constraint is applied. As shown in Figure 8.38, the
 3346 application of the reactor constraint is expected to significantly increase the
 3347 joint fit's preference for the NH hypothesis, as compared to the beam-only
 3348 analysis. Figure 8.40 illustrates the sensitivities of the two fits to the disappearance
 3349 parameters ($\sin^2(\theta_{23}) - \Delta m_{32}^2$) marginalised over both hierarchies and with the
 3350 reactor constraint applied. This plot clearly illustrates the benefit of the joint
 3351 beam and atmospheric analysis. The 1σ credible interval in the IH region is
 3352 entirely removed in the joint analysis, illustrating the improved NH preference.

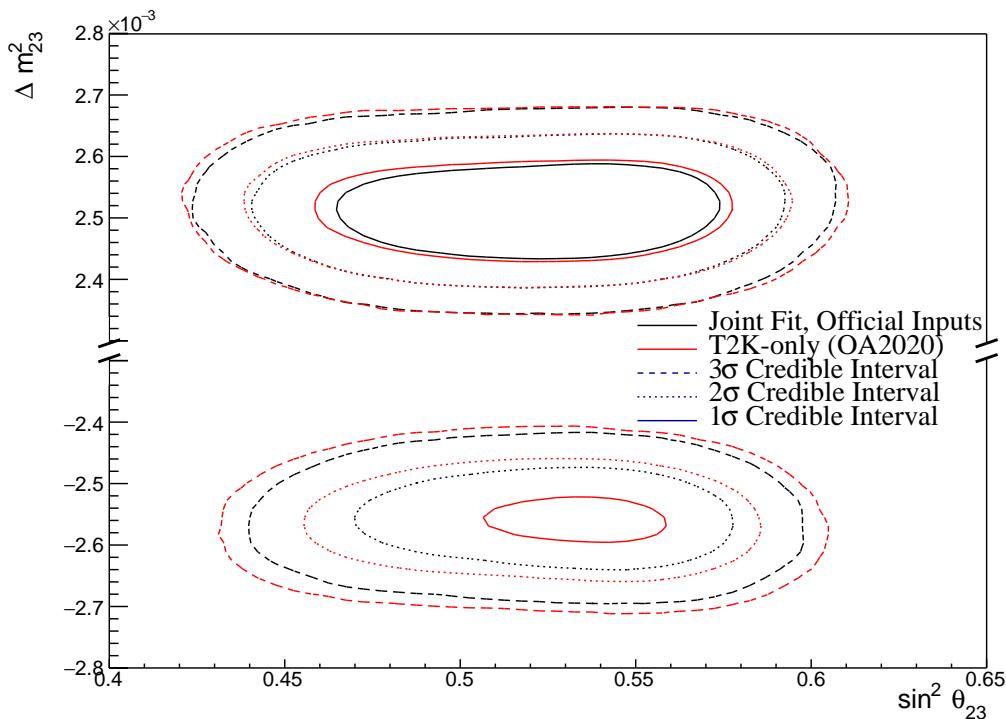


Figure 8.40: The two-dimensional posterior probability density distribution in $\Delta m_{32}^2 - \sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is applied in both fits. The distributions are marginalised over both hierarchies.

3353 The credible intervals of the joint fit are also tighter in the $\sin^2(\theta_{23})$ dimension
3354 than the beam-only analysis in both mass hierarchy regions. This shows that
3355 beyond the ability of the joint fit to prefer the NH more strongly than the beam-
3356 only analysis, the precision to which it can measure $\sin^2(\theta_{23})$ is also improved.
3357 The Bayes factor for NH preference is calculated as $B(\text{NH}/\text{IH}) = 7.29$ and
3358 $B(\text{NH}/\text{IH}) = 3.41$ for the joint beam-atmospheric and beam-only analysis,
3359 respectively. Whilst both present a significant preference for the NH hypothesis
3360 (Table 4.1), the joint fit's preference is much stronger. A similar conclusion can be
3361 made regarding the Bayes factors for UO preference which are $B(\text{UO}/\text{LO}) = 2.86$
3362 and $B(\text{UO}/\text{LO}) = 2.67$ for the joint beam-atmospheric and beam-only analysis,
3363 respectively. Both of these represent a mild preference for the UO but there is
3364 a stronger preference observed in the joint analysis.

3365 The sensitivity of the beam-only and joint beam-atmospheric analyses, to the
3366 appearance parameters ($\delta_{CP} - \sin^2(\theta_{13})$), are compared in Figure 8.41. These
3367 results are marginalised over both hierarchies and include the reactor constraint
3368 on $\sin^2(\theta_{13})$. For this particular set of known oscillation parameters (AsimovA
3369 defined in Table 2.2), the beam-only analysis dominates the sensitivity. The
3370 joint fit does slightly increase the sensitivity to δ_{CP} but it does not change any
3371 conclusions that would be made.

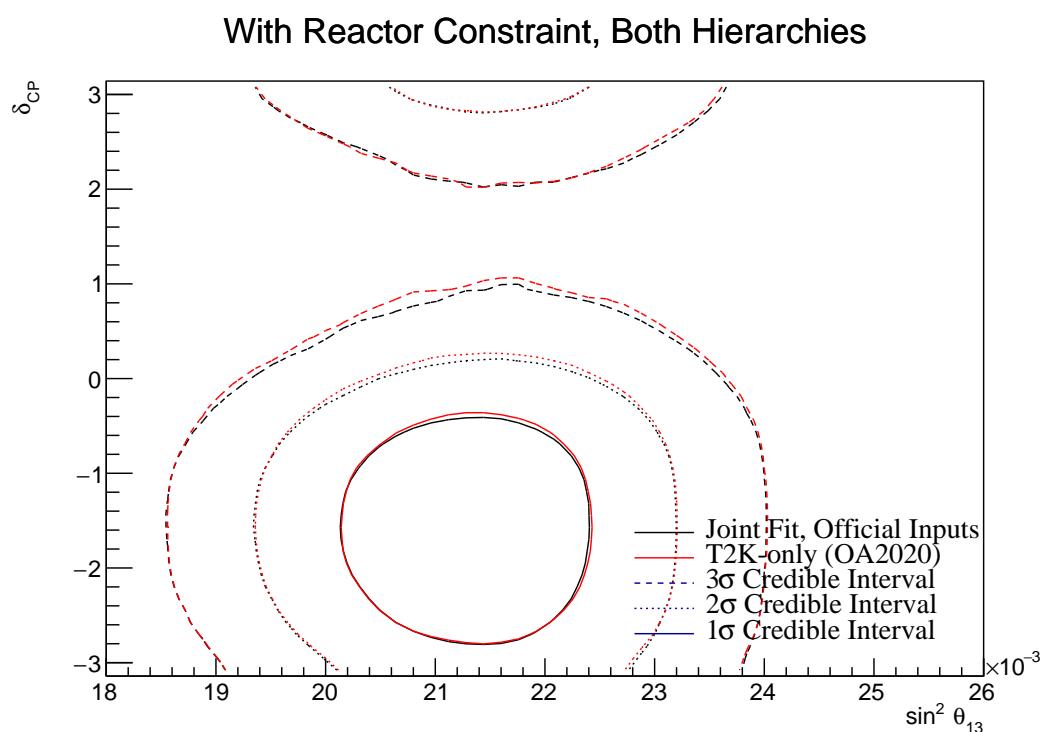


Figure 8.41: The two-dimensional posterior probability density distribution in δ_{CP} – $\sin^2(\theta_{13})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is applied in both fits. The distributions are marginalised over both hierarchies.

3372 8.3.8 Effect of Asimov Parameter Set

3373 Figure 8.4 and Figure 8.5 show that the choice of the parameter set at which the
 3374 Asimov data is made can affect the conclusion. ‘AsimovA’ oscillation parameters
 3375 are defined at a region of δ_{CP} which is dominated by the T2K experiment. This
 3376 explains why the addition of the atmospheric samples does not significantly in-
 3377 crease the sensitivity to δ_{CP} , as illustrated in subsection 8.3.6 and subsection 8.3.7.
 3378 This section presents the sensitivities when ‘AsimovB’ oscillation parameters,
 3379 as defined in Table 2.2, are assumed (alongside the post-BANFF tune) when
 3380 building the Asimov data.

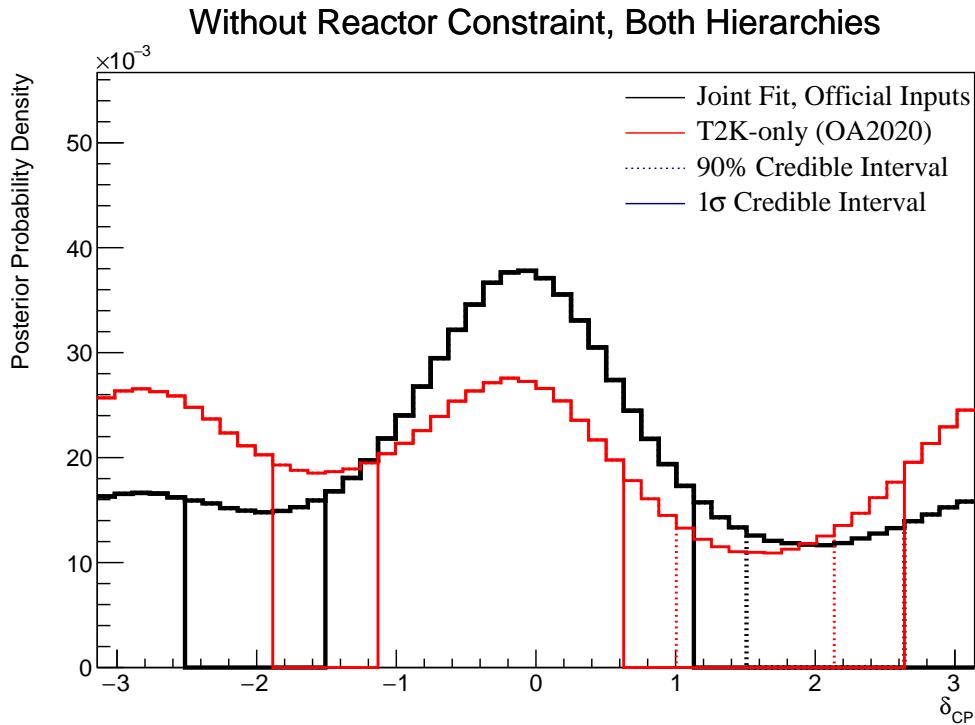


Figure 8.42: The one-dimensional posterior probability density distribution in δ_{CP} compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

3381 The sensitivity to δ_{CP} for the joint beam and atmospheric fit is presented
 3382 in Figure 8.42. The results are compared to those from the beam-only analysis
 3383 in [181]. The reactor constraint is not applied in either of the fits. The known
 3384 oscillation parameter value is $\delta_{CP} = 0$. The shape of the posterior distribution

from the joint analysis is more peaked at $\delta_{CP} = 0$ as compared to the beam-only analysis which has approximately the same posterior probability density at $\delta_{CP} = 0$ and $\delta_{CP} = \pi$. This shows the ability of the joint analysis to better determine the correct phase of δ_{CP} if the true value was CP-conserving. The 1σ credible intervals and the position of the highest posterior probability density are given in Table 8.9.

Parameter	Interval	HPD
δ_{CP} , (BH)	$[-\pi, -2.51], [-1.51, 1.31]$	-0.06
δ_{CP} , (NH)	$[-1.13, 1.63]$	0.06
δ_{CP} , (IH)	$[-3.02, -1.88], [-1.76, 0.13]$	-0.44
Δm_{32}^2 (BH) [$\times 10^{-3}$ eV 2]	$[-2.60, -2.49], [2.46, 2.59]$	2.51
Δm_{32}^2 (NH) [$\times 10^{-3}$ eV 2]	$[2.47, 2.56]$	2.52
Δm_{32}^2 (IH) [$\times 10^{-3}$ eV 2]	$[-2.61, -2.52]$	-2.57
$\sin^2(\theta_{23})$ (BH)	$[0.43, 0.48], [0.55, 0.59]$	0.45
$\sin^2(\theta_{23})$ (NH)	$[0.43, 0.49], [0.55, 0.58]$	0.45
$\sin^2(\theta_{23})$ (IH)	$[0.44, 0.48], [0.54, 0.59]$	0.57

Table 8.9: The position of the highest posterior probability density (HPD) and width of the 1σ credible interval for the SK atmospheric-only fit. The reactor constraint is not applied. The values are presented by which hierarchy hypothesis is assumed: marginalised over both hierarchies (BH), normal hierarchy only (NH) and inverted hierarchy only (IH).

Naively, if just the 1σ credible interval were considered without observing the shape of the distribution, it would appear that the joint analysis would have a worse sensitivity to δ_{CP} due to the larger interval around δ_{CP} . The 1σ credible interval for the beam-only analysis is given as the range $\delta_{CP} = [-\pi, -1.88], [-1.13, 0.63]$ and $[2.64, \pi]$ which contains 56% of all values of δ_{CP} . The joint beam and atmospheric analysis contains 52% of all δ_{CP} values within the 1σ credible interval. Therefore, if the area within the 1σ credible interval were to be compared between the two fits, the joint analysis would be shown to have better precision.

This contradiction stems from the methodology in which the credible interval is calculated. The technique used in this analysis (documented in subsection 4.3.2)

3401 fills the credible interval by selecting bins in order of magnitude until 68% of the
 3402 posterior density is contained. If instead, the credible interval was calculated
 3403 by expanding around the highest posterior probability, the benefits of the joint
 3404 fit would be more obvious. In the case where the shape of the posterior was
 3405 Gaussian, these two techniques would be equivalent.

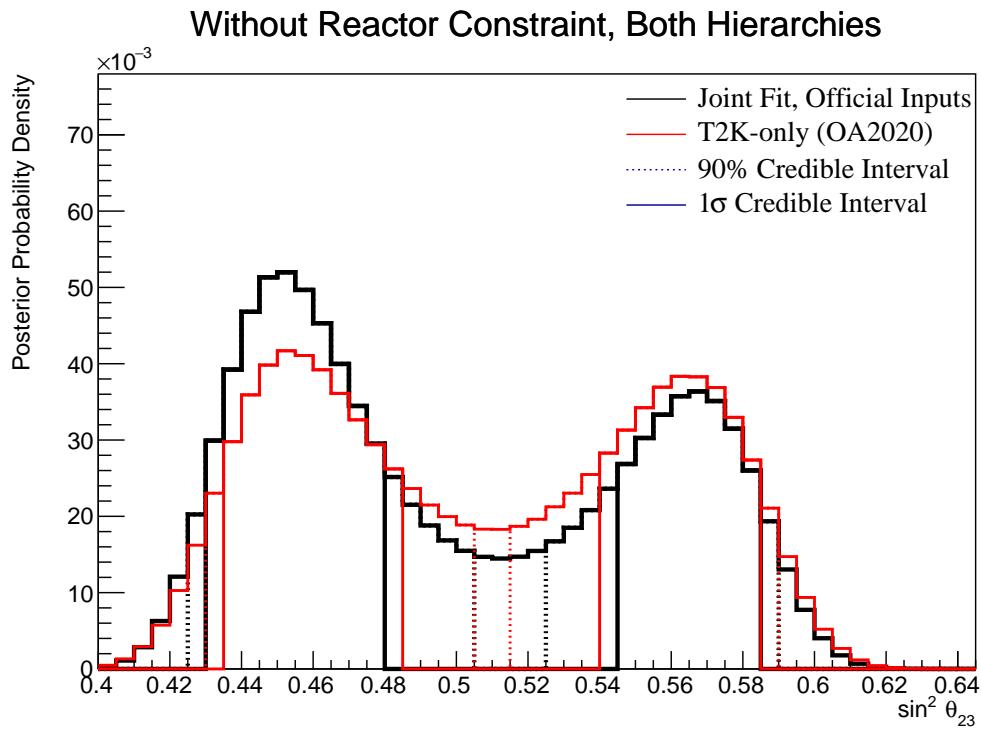


Figure 8.43: The one-dimensional posterior probability density distribution in $\sin^2(\theta_{23})$ compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

3406 The sensitivity of the joint beam and atmospheric fit to $\sin^2(\theta_{23})$ is presented
 3407 in Figure 8.43. The sensitivity is compared to that of the beam-only analysis
 3408 in [181]. The reactor constraint is not applied in either of the fits being com-
 3409 pared. The Asimov parameter value is $\sin^2(\theta_{23}) = 0.45$ and the sensitivities are
 3410 marginalised over both hierarchies. Clearly, the joint beam and atmospheric
 3411 fit has a much larger probability density in the region surrounding the known
 3412 oscillation parameters. This shows the better octant determination of the joint
 3413 analysis compared to the beam-only fit. The ratio of the posterior density at

³⁴¹⁴ the peak of the lower octant to the peak of the upper octant from the joint fit is
³⁴¹⁵ 1.43 compared to 1.09 from the beam-only analysis. This shows further support
³⁴¹⁶ for the joint analysis in correctly selecting the lower octant, which is the correct
³⁴¹⁷ hypothesis given the known oscillation parameters.

	LO ($\sin^2 \theta_{23} < 0.5$)	UO ($\sin^2 \theta_{23} > 0.5$)	Sum
NH ($\Delta m_{32}^2 > 0$)	0.35	0.24	0.59
IH ($\Delta m_{32}^2 < 0$)	0.19	0.22	0.41
Sum	0.54	0.46	1.00

Table 8.10: The distribution of steps in a joint beam and atmospheric fit, presented as the fraction of steps in the upper (UO) and lower (LO) octants and the normal (NH) and inverted (IH) hierarchies. The reactor constraint is not applied. The Bayes factors are calculated as $B(\text{NH}/\text{IH}) = 1.43$ and $B(\text{LO}/\text{UO}) = 1.19$.

³⁴¹⁸ The distribution of steps, split by hierarchy and octant hypothesis, is presented
³⁴¹⁹ in Table 8.10. The Bayes factor for hierarchy and octant determination are
³⁴²⁰ $B(\text{NH}/\text{IH}) = 1.43$ and $B(\text{LO}/\text{UO}) = 1.19$, respectively. The octant Bayes factor
³⁴²¹ is now presented as LO/UO as the known oscillation parameter is contained
³⁴²² within the lower octant. These values compare to $B(\text{NH}/\text{IH}) = 1.08$ and
³⁴²³ $B(\text{LO}/\text{UO}) = 0.91$ from the beam-only analysis. This shows additional evidence
³⁴²⁴ of the joint analysis's preference for selecting the correct octant and hierarchy
³⁴²⁵ hypothesis. Comparisons to the AsimovA Bayes factors presented in Table 8.6
³⁴²⁶ show how the preference for the correct octant and hierarchy depend on the
³⁴²⁷ true value of δ_{CP} and $\sin^2(\theta_{23})$.

³⁴²⁸ The sensitivity of the beam-only and joint beam-atmospheric analysis to
³⁴²⁹ Δm_{32}^2 is given in Figure 8.44. Both of the results are marginalised over both
³⁴³⁰ hierarchies and the reactor constraint is not applied in either analysis. The joint
³⁴³¹ analysis has a stronger preference for the correct hierarchy (NH) which is shown
³⁴³² by the higher Bayes factor ($B(\text{NH}/\text{IH}) = 1.43$) compared to the beam-only
³⁴³³ analysis ($B(\text{NH}/\text{IH}) = 1.08$).

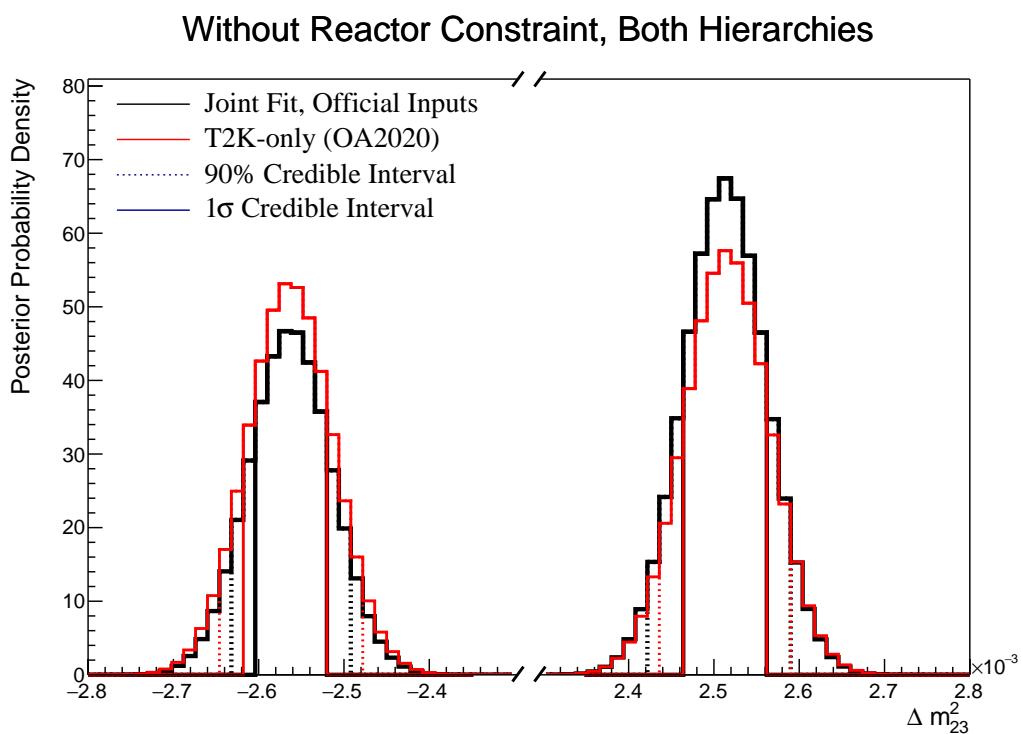


Figure 8.44: The one-dimensional posterior probability density distribution in Δm_{32}^2 compared between the joint beam-atmospheric fit (Black) and the latest T2K sensitivities (Red) [181]. The reactor constraint is not applied in either fit. The distributions are marginalised over both hierarchies.

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Conclusions and Outlook

Appendices

A

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Atmospheric Sample Spectra

3439 This appendix documents the interaction mode breakdown of all the atmospheric
3440 samples used within the analysis. The generated tune of the model parameters
3441 and the Asimov A oscillation parameter set (defined in Table 2.2) are assumed.
3442 The livetime of SK-IV is taken to be 3244.4 days.

3443 A.1 Binning

3444 The lepton momentum and cosine zenith binning edges for the atmospheric
3445 samples used within this analysis are defined in Table A.1.

3446 A.2 Fully Contained Sub-GeV Samples

3447 The interaction mode breakdown of the fully contained Sub-GeV samples are
3448 shown in Figure A.1 and Figure A.2, for the samples with enriched CC0 π and
3449 CC1 π^\pm respectively.

3450 The CC0 π sample are dominated by CCQE events ($\sim 70\%$) with smaller
3451 contributions of 2p2h ($\sim 12\%$) and CC1 π ($\sim 10\%$) components. The energy peaks
3452 around 300 MeV, which is slightly below that of the T2K samples but still has
3453 significant contribution upto 1 GeV which overlaps the T2K sample energy range.

Sample	$\cos(\theta_Z)$ Bins	Momentum Bin Edges ($\log_{10}(P)$ MeV)
SubGeV-elike-0dcy	10	2.0, 2.4, 2.6, 2.8, 3.0, 3.2
SubGeV-elike-1dcy	1	2.0, 2.4, 2.6, 2.8, 3.0, 3.2
SubGeV-mulike-0dcy	10	2.0, 2.4, 2.6, 2.8, 3.0, 3.2
SubGeV-mulike-1dcy	10	2.0, 2.4, 2.6, 2.8, 3.0, 3.2
SubGeV-mulike-2dcy	1	2.0, 2.4, 2.6, 2.8, 3.0, 3.2
SubGeV-pi0like	1	2.0, 2.2, 2.4, 2.6, 2.8, 3.2
MultiGeV-elike-nue	10	3.0, 3.4, 3.7, 4.0, 5.0
MultiGeV-elike-nuebar	10	3.0, 3.4, 3.7, 4.0, 5.0
MultiGeV-mulike	10	3.0, 3.4, 5.0
MultiRing-elike-nue	10	3.0, 3.4, 3.7, 5.0
MultiRing-elike-nuebar	10	3.0, 3.4, 3.7, 5.0
MultiRing-mulike	10	2.0, 3.124, 3.4, 3.7, 5.0
MultiRing-Other1	10	3.0, 3.4, 3.7, 4.0, 5.0
PC-Stop	10	2.0, 3.4, 5.0
PC-Through	10	2.0, 3.124, 3.4, 3.7, 5.0
Upmu-Stop	10	3.2, 3.4, 3.7, 8.0
Upmu-Through-Showering	10	2.0, 8.0
Upmu-Through-NonShowering	10	2.0, 8.0

Table A.1: The reconstructed cosine zenith and lepton momentum binning assigned to the atmospheric samples. The “ $\cos(\theta_Z)$ Bins” column illustrates the number of bins uniformly distributed over the $-1.0 \leq \cos(\theta_Z) \leq 1.0$ region for fully and partially contained samples and $-1.0 \leq \cos(\theta_Z) \leq 0.0$ region for up- μ samples.

3454 The one-ring CC1 π samples, where the pion is tagged via its decay electron,
 3455 are dominated by CC1 π events ($\sim 75\%$) with a small contribution of CCM π
 3456 ($\sim 10\%$). The two-ring pion sample is mostly dominated by the NC1 π^0 via
 3457 resonances, and has several equally-sized contributions from CCQE, NC1 π^\pm via
 3458 resonances, and NC coherent pion production, where the π^0 likely comes from
 3459 nucleon and π^\pm final state interactions in the nucleus.

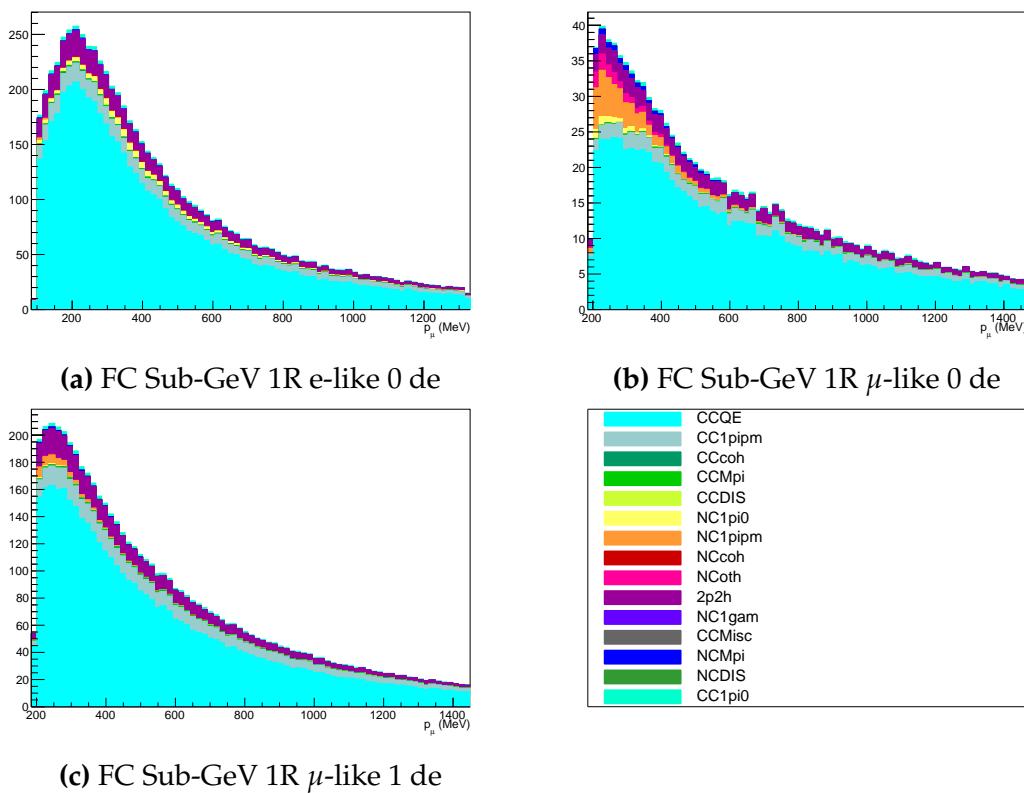


Figure A.1: Breakdown by interaction mode of the FC Sub-GeV atmospheric samples targeting CC 0π events.

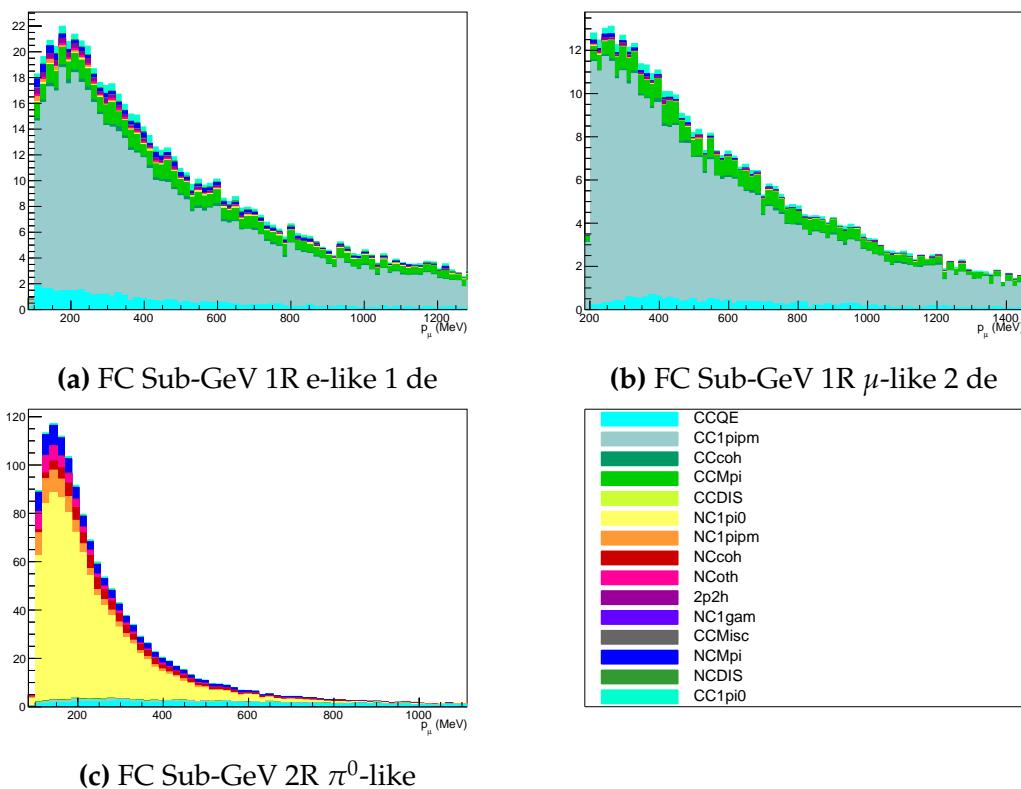


Figure A.2: Breakdown by interaction mode of the FC Sub-GeV atmospheric samples targeting single pion events.

3460 A.3 Fully Contained Multi-GeV Samples

3461 The interaction mode breakdown of fully contained multi-GeV samples is high-
 3462 lighted in Figure A.3. Due to the event selection applied in SK which targets π^+
 3463 and π^- separation, the ν_e sample mainly consists of events with pions (single pion
 3464 production or multi-pion/DIS interactions). The pion separation is explained in
 3465 Section section 6.1. This reasoning also explains the significant CCQE contribution
 3466 of the $\bar{\nu}_e$ sample. The muon-like sample is dominated by CCQE interactions with
 3467 $\sim 10 - 15\%$ 2p2h and CC1 π contribution of events.

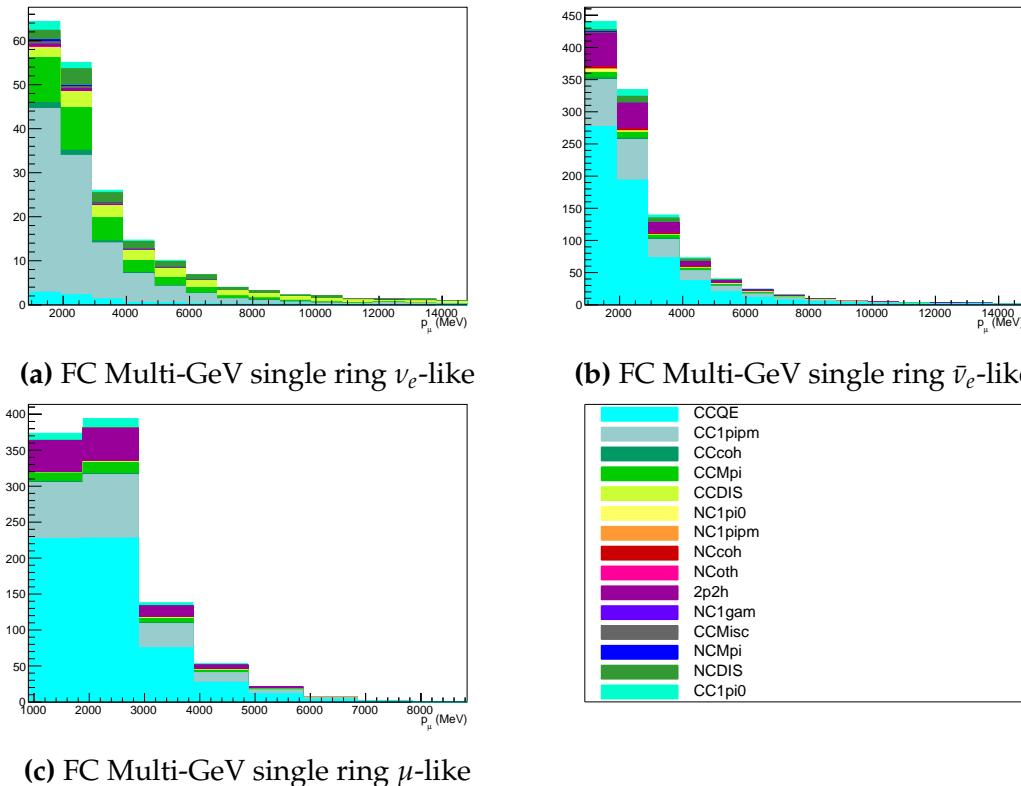


Figure A.3: Breakdown by interaction mode of the FC Multi-GeV single ring atmospheric samples.

3468 A.4 Fully Contained Multi-Ring Samples

3469 The interaction mode breakdown of fully contained multi-ring events is shown
 3470 in Figure A.4. These samples see more interaction modes contributing in general,
 3471 and there is a much larger contribution from multi-pion and DIS interaction
 3472 modes, compared to the other samples.

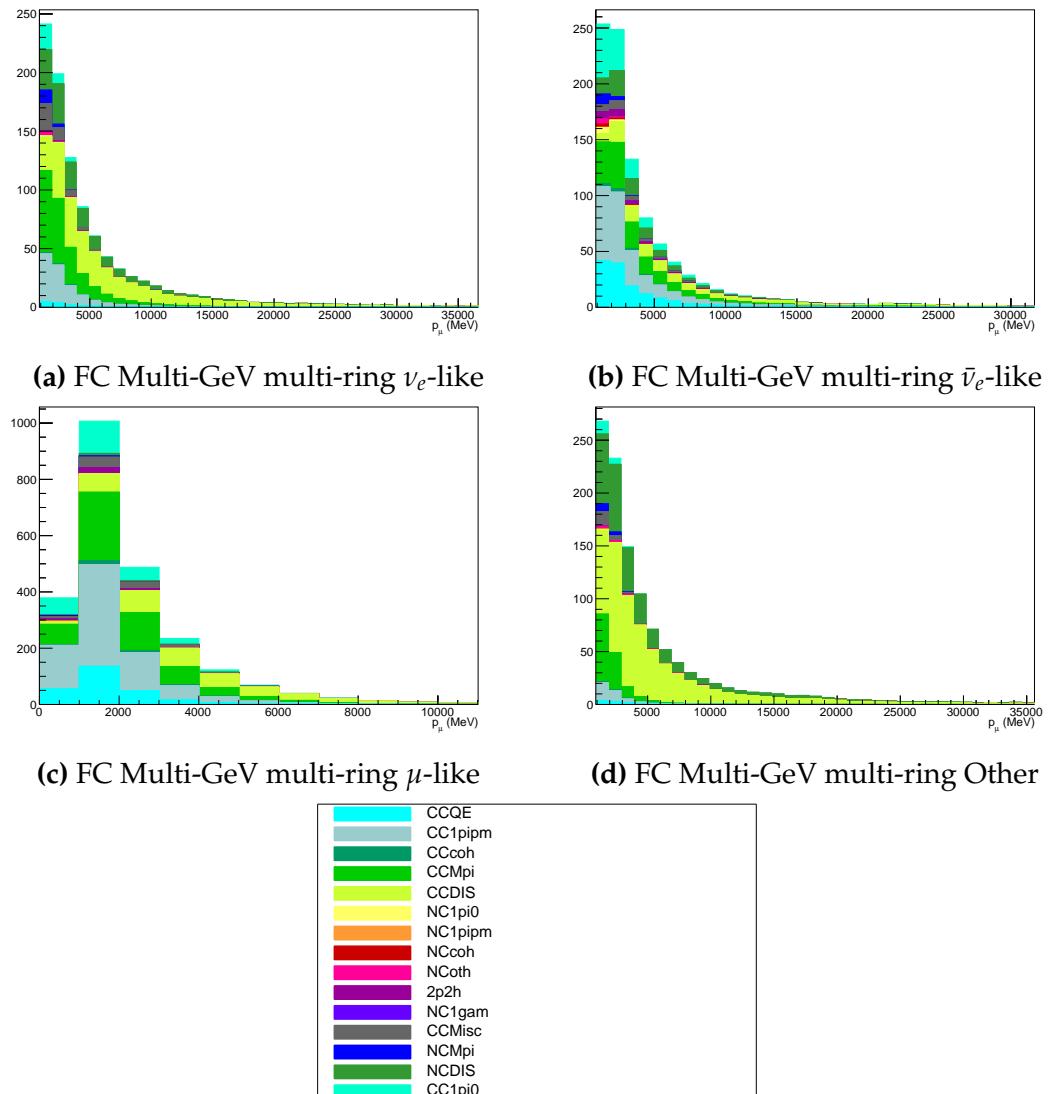


Figure A.4: Breakdown by interaction mode of the FC Multi-GeV multi-ring atmospheric samples.

3473 A.5 Partially Contained Samples

3474 The breakdown for partially contained samples is highlighted in Figure A.5.
 3475 As with the multi-ring samples, there is no dominating interaction mode. The
 3476 neutrino energies of events in this sample extend into the tens of GeV and become
 3477 dominated by DIS interaction modes in the high energy limit.

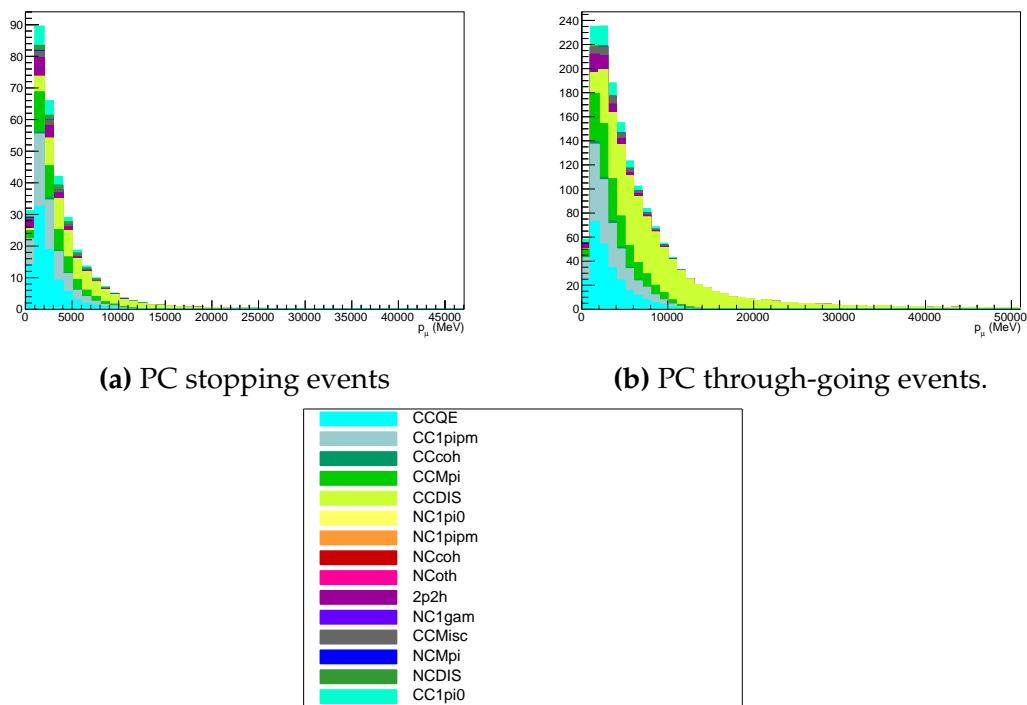


Figure A.5: Breakdown by interaction mode of the PC atmospheric samples.

3478 A.6 Upward-Going Muon Samples

3479 The breakdown for upward-going muons is illustrated in Figure A.6. These
 3480 samples are significantly dominated by DIS interactions with energies extending
 3481 up into the hundreds of GeV.

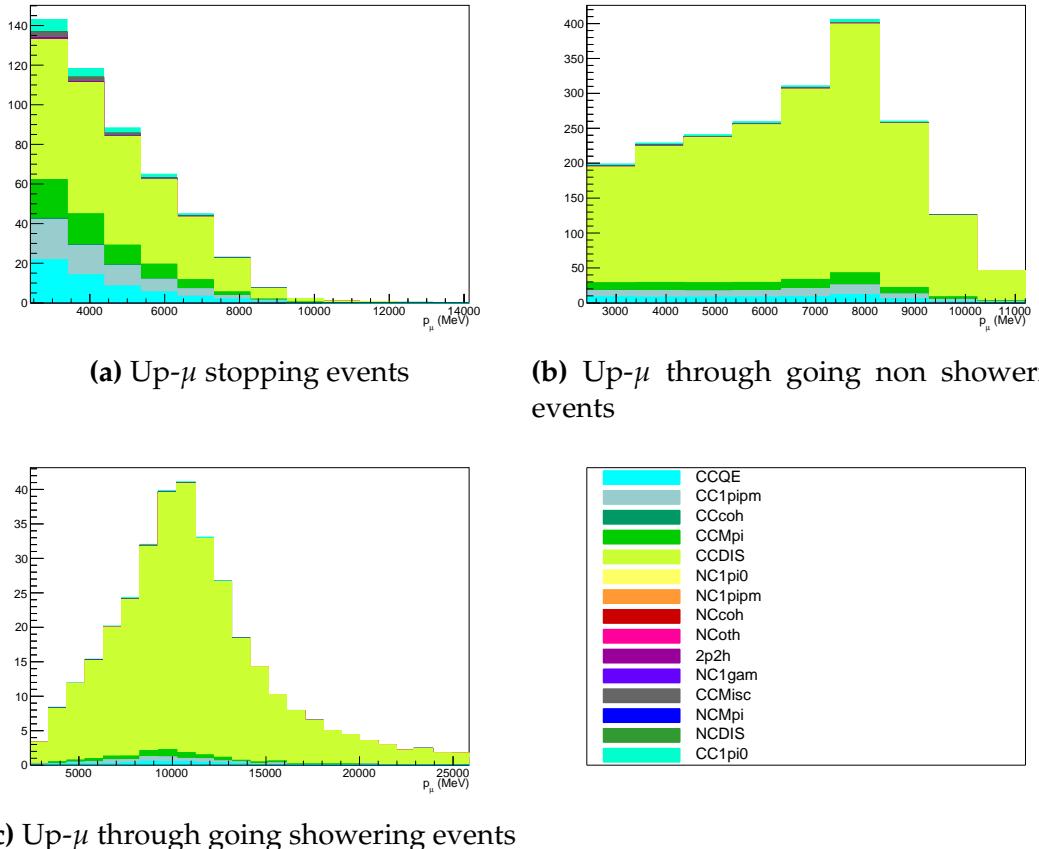


Figure A.6: Breakdown by interaction mode of the atmospheric upward going muon samples.

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