

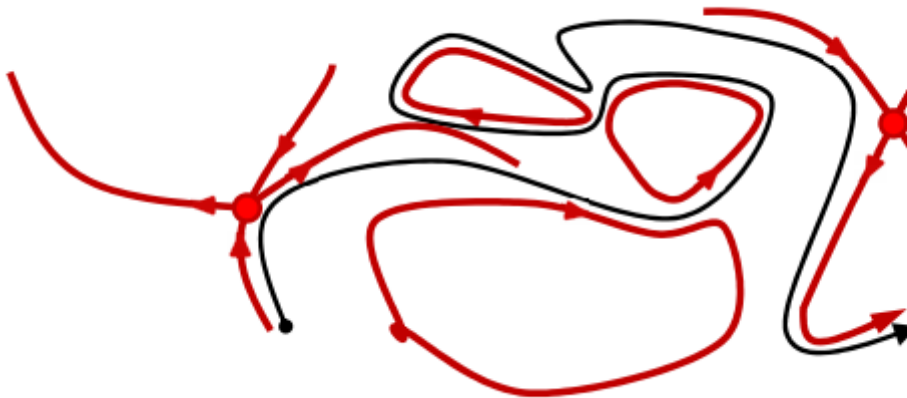
Symmetry Reduced Exact Coherent Structures in Plane Couette Flow

May 4, 2015

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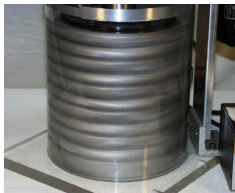
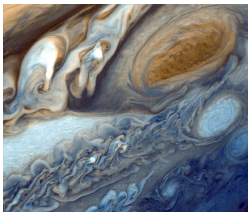


Turbulence



What is it, and why do we care?

- Turbulence is an irregular, chaotic, and dissipative flow phenomenon, with a vast number of (coupled) length scales.



- Shows up in many important systems - atmospheric and geophysical flows, turbine exhaust, stellar flows...

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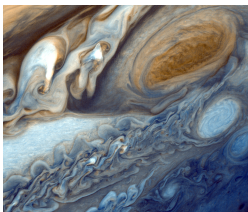
Future Work

Turbulence



What is it, and why do we care?

- ▶ Turbulence is an irregular, chaotic, and dissipative flow phenomenon, with a vast number of (coupled) length scales.



- ▶ Shows up in many important systems - atmospheric and geophysical flows, turbine exhaust, stellar flows...
- ▶ Direct numerical simulation (DNS) is very time consuming - need a simpler model!

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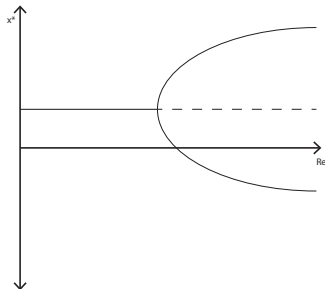
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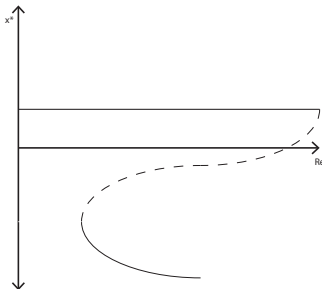


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- Of particular interest are transitions to turbulence from smooth (laminar) flow



Supercritical Transition



Subcritical Transition

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- Supercritical transitions are well described by linear stability analysis (inner-cylinder rotation in Taylor Couette flow)...

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- ▶ Supercritical transitions are well described by linear stability analysis (inner-cylinder rotation in Taylor Couette flow)...
- ▶ Subcritical transitions are sudden and unpredictable (pipe flow, plane Couette flow)



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Navier-Stokes

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The Navier-Stokes equation is the linear momentum conservation statement for Newtonian, incompressible fluids:

$$\left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{v},$$

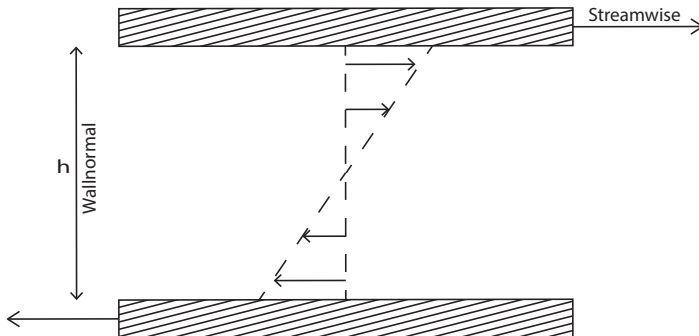
with the appropriate boundary conditions.

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Base Flow



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Plane Couette



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If the Reynolds number is not small, the damping term is about the same size as the rest of the terms, and we have to solve the full Navier-Stokes equations (for Couette flow)

$$\left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = \frac{1}{Re} \nabla^2 \mathbf{v},$$

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- Mean field approaches, such as the Reynolds Averaged Navier Stokes have had a great deal of success

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- ▶ Mean field approaches, such as the Reynolds Averaged Navier Stokes have had a great deal of success
- ▶ RANS works by time averaging out the turbulent perturbations to some mean flow, resulting in a new equation of motion for the mean flow...

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- ▶ Mean field approaches, such as the Reynolds Averaged Navier Stokes have had a great deal of success
- ▶ RANS works by time averaging out the turbulent perturbations to some mean flow, resulting in a new equation of motion for the mean flow...
- ▶ ... but dynamical properties of turbulence are lost in the time-averaging process

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Exact Coherent Structures



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- Instead of time averaging the perturbation away, try to analyze its dynamics

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- ▶ Instead of time averaging the perturbation away, try to analyze its dynamics
- ▶ Imagine the perturbation as living in some infinite dimensional space of all possible velocity fields (the 'state space')

Laminar Basin of Attraction



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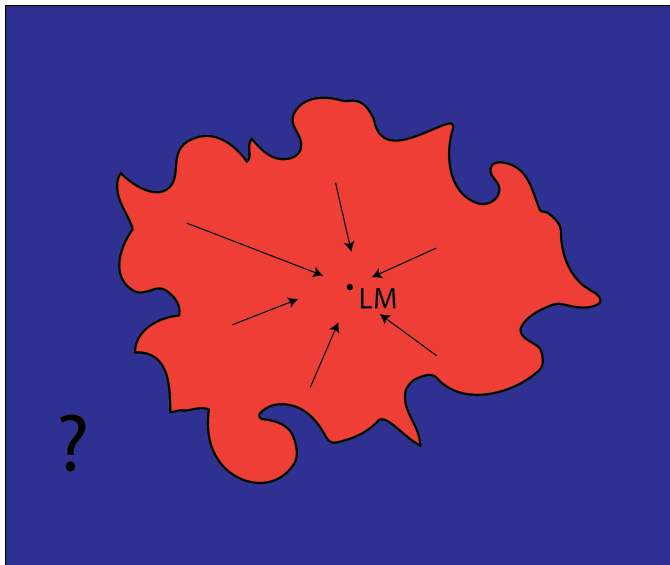
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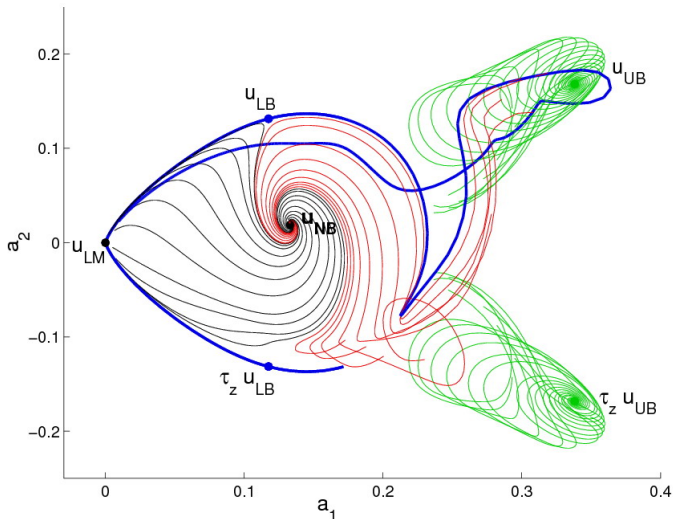
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Turbulence

Equilibria



Halcrow et. al (2008)

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- The full Navier-Stokes equations are invariant ($f_T(\sigma \mathbf{u}) = \sigma f_T(\mathbf{u})$) under any spatial rotation and translation

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- ▶ The full Navier-Stokes equations are invariant ($f_T(\sigma \mathbf{u}) = \sigma f_T(\mathbf{u})$) under any spatial rotation and translation
- ▶ Boundary conditions of plane Couette flow reduce the number of invariant symmetries to the group generated by the following transformations:

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- ▶ The full Navier-Stokes equations are invariant ($f_T(\sigma \mathbf{u}) = \sigma f_T(\mathbf{u})$) under any spatial rotation and translation
- ▶ Boundary conditions of plane Couette flow reduce the number of invariant symmetries to the group generated by the following transformations:
 - ▶ $\sigma_{xy} \rightarrow$ Rotation by π about the z axis

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- ▶ The full Navier-Stokes equations are invariant ($f_T(\sigma \mathbf{u}) = \sigma f_T(\mathbf{u})$) under any spatial rotation and translation
- ▶ Boundary conditions of plane Couette flow reduce the number of invariant symmetries to the group generated by the following transformations:
 - ▶ $\sigma_{xy} \rightarrow$ Rotation by π about the z axis
 - ▶ $\sigma_z \rightarrow$ Mirroring across the z axis

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- ▶ The full Navier-Stokes equations are invariant ($f_T(\sigma \mathbf{u}) = \sigma f_T(\mathbf{u})$) under any spatial rotation and translation
- ▶ Boundary conditions of plane Couette flow reduce the number of invariant symmetries to the group generated by the following transformations:
 - ▶ $\sigma_{xy} \rightarrow$ Rotation by π about the z axis
 - ▶ $\sigma_z \rightarrow$ Mirroring across the z axis
 - ▶ $\tau(l_x, l_z) \rightarrow$ Translation by (l_x, l_z) in the xz plane



- Integration and root finding in a symmetric subspace are faster - fewer computations required

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- ▶ Integration and root finding in a symmetric subspace are faster - fewer computations required
- ▶ Coherent structures can also exist in different inertial reference frames \rightarrow travelling waves

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- ▶ Integration and root finding in a symmetric subspace are faster - fewer computations required
- ▶ Coherent structures can also exist in different inertial reference frames \rightarrow travelling waves
- ▶ Symmetries can restrict the allowed directions of these travelling waves (Viswanath, 2007).



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- ▶ Integration and root finding in a symmetric subspace are faster - fewer computations required
- ▶ Coherent structures can also exist in different inertial reference frames \rightarrow travelling waves
- ▶ Symmetries can restrict the allowed directions of these travelling waves (Viswanath, 2007).
- ▶ Symmetric fields are also not necessarily representative of turbulence

Symmetric and Asymmetric Fields



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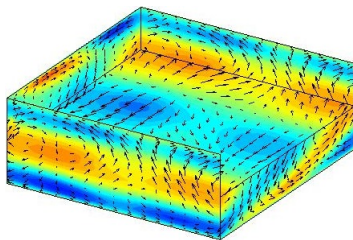
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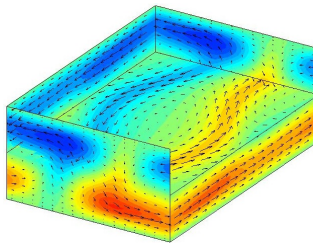
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A random, asymmetric field



A highly symmetric periodic orbit

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- I use Channelflow, a simulation library developed at UNH (Gibson, 2014) to analyze plane Couette flows.

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- ▶ I use Channelflow, a simulation library developed at UNH (Gibson, 2014) to analyze plane Couette flows.
- ▶ In addition to Navier-Stokes integration, Channelflow can also find coherent structures (with a good initial guess)

The Simulation

Generating a good guess



- How do we hunt through the immense phase space? (> 60,000 dimensions)

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Generating a good guess



- ▶ How do we hunt through the immense phase space? ($> 60,000$ dimensions)
- ▶ Start from a random initial condition

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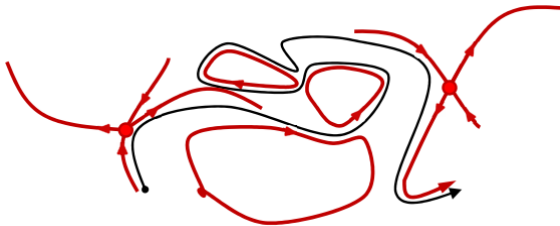
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The Simulation

Generating a good guess



- ▶ How do we hunt through the immense phase space? ($> 60,000$ dimensions)
- ▶ Start from a random initial condition
- ▶ Integrate flow forward in time



Borrero (2014)

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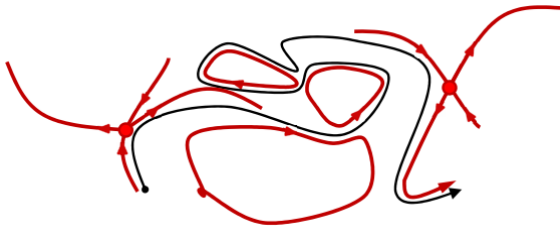
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Generating a good guess



- ▶ How do we hunt through the immense phase space? ($> 60,000$ dimensions)
- ▶ Start from a random initial condition
- ▶ Integrate flow forward in time



Borrero (2014)

- ▶ Look at recurrence graphs

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- Can compare the flow at time t to time $t + T$ by calculating $d(\mathbf{u}(t), \mathbf{u}(t + T))$

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- ▶ Can compare the flow at time t to time $t + T$ by calculating $d(\mathbf{u}(t), \mathbf{u}(t + T))$
- ▶ Use the physically motivated energy norm

$$|\mathbf{u}|^2 = \frac{1}{V} \int_V \mathbf{u} \cdot \mathbf{u} dV,$$

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- ▶ Can compare the flow at time t to time $t + T$ by calculating $d(\mathbf{u}(t), \mathbf{u}(t + T))$
- ▶ Use the physically motivated energy norm

$$|\mathbf{u}|^2 = \frac{1}{V} \int_V \mathbf{u} \cdot \mathbf{u} dV,$$

- ▶ Minima of these plots should indicate the presence of a periodic orbit

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Recurrence Plots



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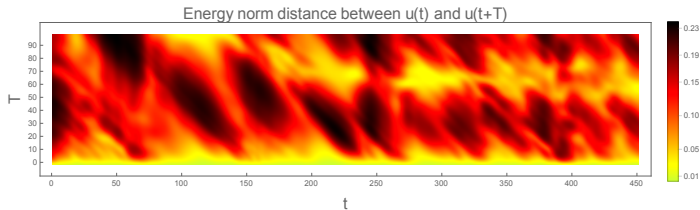
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- Once a promising candidate is found, it can be passed on to the root finder, to solve

$$\mathbf{u} - \sigma f_T(\mathbf{u}) = 0,$$

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- Once a promising candidate is found, it can be passed on to the root finder, to solve

$$\mathbf{u} - \sigma f_T(\mathbf{u}) = 0,$$

- where σ encodes the travelling wave velocity.

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- Once a promising candidate is found, it can be passed on to the root finder, to solve

$$\mathbf{u} - \sigma f_T(\mathbf{u}) = 0,$$

- where σ encodes the travelling wave velocity.
- The root finder uses the Newton method, with the step in 60,000 dimensions calculated using the generalized minimum residual method.

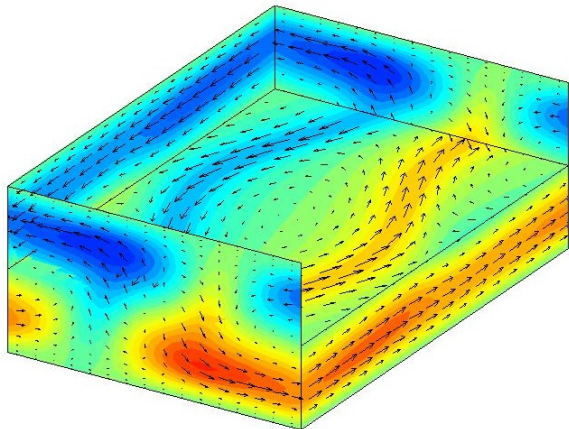
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P85.47



This orbit has period $T = 85.47$, and has symmetries generated by $\{e, \sigma_{xy}, \sigma_z \tau_{xz}\}$



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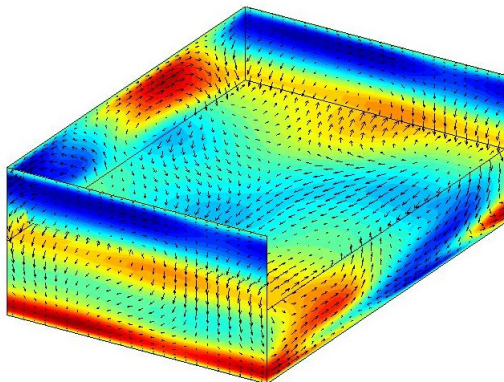
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P8.32



This orbit has period $T = 8.32$, and has symmetries generated by $\{e, \sigma_{xy}\tau_{xz}\}$.



- Even though it's in a reduced symmetry subspace that allows travelling waves, it has no relative velocity.

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- Coherent structures cannot exist for all Reynolds numbers
 - $Re = 0$ means only the laminar state can exist

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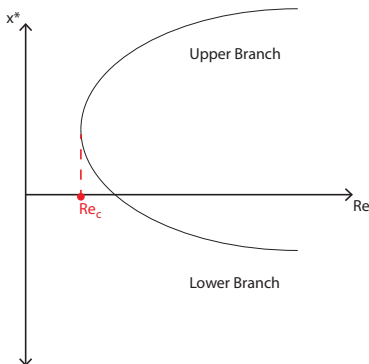
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- ▶ Coherent structures cannot exist for all Reynolds numbers
- $Re = 0$ means only the laminar state can exist
- ▶ Coherent structures will arise from some high-dimensional bifurcation



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- Near bifurcation, period will be low - is P8.32 near a bifurcation?



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- ▶ Near bifurcation, period will be low - is P8.32 near a bifurcation?
- ▶ Can try to continue this solution for lower Reynolds numbers to locate the bifurcation

P8.32



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Couette Flow

The Problem

Transitions to Turbulence

The System

Navier-Stokes

Bulk Flow

Statistical Methods

Exact Coherent Structures

Symmetry

The Simulation

Recurrence Graphs

Root Finding

23 Current Results

Future Work

- ▶ Near bifurcation, period will be low - is P8.32 near a bifurcation?
- ▶ Can try to continue this solution for lower Reynolds numbers to locate the bifurcation
- ▶ After bifurcation, coherent structures tend to branch off in pairs - can this orbit's pair be located?

Future Work



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24 Future Work

- Try to find travelling coherent structures in the reduced symmetry space

Future Work



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24 Future Work

- ▶ Try to find travelling coherent structures in the reduced symmetry space
- ▶ Try to find coherent structures in spaces with even less symmetry than P8.32

Future Work



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- ▶ Try to find travelling coherent structures in the reduced symmetry space
- ▶ Try to find coherent structures in spaces with even less symmetry than P8.32
- ▶ Find the bifurcation for P8.32 and continue the solution to find the other branch

24 Future Work

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