Game Physics

GPR350, Fall 2019 Daniel S. Buckstein

Angular Dynamics Week 9

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Angular Dynamics

- Force, torque and Newton's 2nd law
- Conversion between force and torque
- Moment of inertia
 - Inertia tensor
 - Frame of reference
 - Food for thought

Cool stuff:

http://physics-help.info/physicsguide/appendices/si_units.shtml https://people.eecs.berkeley.edu/~jfc/mirtich/massProps.html

Recap: Linear Motion

 Position of a particle is determined by integration (e.g. explicit Euler method):

$$x_{t+dt} = x_t + \frac{dx}{dt}dt = x_t + v_t dt$$

$$v_{t+dt} = v_t + \frac{dv}{dt}dt = v_t + a_t dt$$

- x: position; dx/dt: derivative of position
- v: velocity; dv/dt: derivative of velocity (acceleration)

Recap: Linear Motion

 Acceleration is determined by rearranging Newton's 2nd law:

$$F = ma$$

$$a = m^{-1}F = \frac{F}{m}$$

- m: mass (kg)
- a: acceleration (m/s²)
- F: net force on particle (Newtons: N = kg m/s²)

Recap: Linear Motion

Force is also related to linear momentum:

$$\frac{dp}{dt} = \frac{d}{dt}(mv) = m\frac{dv}{dt}$$

$$= ma = F$$

$$\frac{dp}{dt} = \frac{d}{dt}(mv) = m\frac{dv}{dt}$$

- m: mass (kg, assuming constant)
- v: linear velocity (m/s)
- p: linear momentum (kg m/s)

 Rotation/orientation is solved using similar methods:

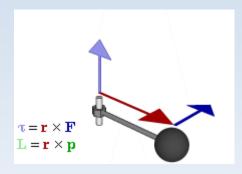
$$q_{t+dt} = q_t + \frac{dq}{dt}dt = q_t + \omega_t q_t \frac{1}{2}dt$$

$$\omega_{t+dt} = \omega_t + \frac{d\omega}{dt}dt = \omega_t + \alpha_t dt$$

- q: orientation (quaternion)
- omega: angular velocity; alpha: angular acceleration

- How do we calculate angular acceleration???
- *Torque*: "twist force"

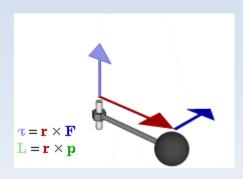
$$\tau = r \times F$$



- r: "moment arm" or "lever arm" vector (m)
- F: force (N)
- tau: torque (Newton-meters: Nm)

- How do we calculate angular acceleration???
- Angular momentum to start with:

$$L = r \times p$$



- r: "moment arm" or "lever arm" vector (m)
- p: linear momentum (kg m/s)
- L: angular momentum (kg m²/s)

- How do we calculate angular acceleration???
- Angular momentum: also calculated with respect to angular velocity:

$$L = I\omega$$

- I: "moment of inertia" (kg m²)
- omega: angular velocity (angle/s)
- L: angular momentum (kg m²/s)

Is torque related to angular momentum?

$$L = I\omega$$

$$\frac{dL}{dt} = \frac{d}{dt}(I\omega) = I\frac{d\omega}{dt}$$

$$= I\alpha = \tau$$

- I: moment of inertia (kg m², assuming constant)
- alpha: angular acceleration (rad/s²)
- tau: net torque on particle (Nm = kg m/s²)

- Angular acceleration is determined using the rotational version of Newton's 2nd law:
 - Linear:

$$F = ma$$
$$a = m^{-1}F$$

Angular:

$$\tau = I\alpha$$
$$\alpha = I^{-1}\tau$$

- Can convert between linear and angular model using the original torque formula
 - Linear to angular:

$$\tau = r \times F$$

Angular to linear:

$$F = \tau \times r$$

Moment of Inertia

- Problem: in vector physics, moment of inertia is not necessarily just a scalar...
- Effectively represents "resistance to angular change over the distribution of mass"
- Mass as a uniform scale rather than a scalar:

$$\vec{F} = \boldsymbol{m}_{3 \times 3} \ \vec{a} = \begin{bmatrix} m \\ m \end{bmatrix} \begin{bmatrix} \vec{a}_x \\ \vec{a}_y \\ \vec{a}_z \end{bmatrix}$$

Moment of Inertia

 Recovering acceleration from force would use the inverse of this uniform scale:

$$\vec{a} = m_{3\times 3}^{-1} \vec{F} = \begin{bmatrix} m^{-1} \\ m^{-1} \\ m^{-1} \end{bmatrix} \vec{F}$$

• We might call this a "mass tensor"

Moment of Inertia

 This rule is used for angular since inertia is not necessarily uniform across a rigid body:

$$\vec{\tau} = \mathbf{I}_{3\times3} \; \vec{\alpha}$$
$$\vec{\alpha} = \mathbf{I}_{3\times3}^{-1} \; \vec{\tau}$$

- The moment of inertia in matrix form is called an "inertia tensor"
- Pre-defined for some shapes, or calculate in general form

- Examples (uniform distribution):
 - Solid sphere of radius r and mass m

$$I = \begin{bmatrix} \frac{2}{5}mr^2 & & \\ & \frac{2}{5}mr^2 & \\ & & \frac{2}{5}mr^2 \end{bmatrix}$$

Hollow sphere of radius r and mass m

$$I = \begin{bmatrix} \frac{2}{3}mr^2 & & & \\ & \frac{2}{3}mr^2 & & \\ & & \frac{2}{3}mr^2 \end{bmatrix}$$

- Examples (non-uniform distribution):
 - Solid box* of width w, height h, depth d and mass m

$$I = \begin{bmatrix} \frac{1}{12}m(h^2 + d^2) & & \\ & \frac{1}{12}m(d^2 + w^2) & \\ & & \frac{1}{12}m(w^2 + h^2) \end{bmatrix}$$

Hollow box** of width w, height h, depth d and mass m

$$I = \begin{bmatrix} \frac{5}{3}m(h^2 + d^2) \\ & \frac{5}{3}m(d^2 + w^2) \\ & \frac{5}{3}m(w^2 + h^2) \end{bmatrix}$$

https://physics.stackexchange.com/questions/105229/tensor-of-inertia-of-a-hollow-cube/105234 https://en.wikipedia.org/wiki/List_of_moments_of_inertia

*cube:
$$I = \frac{1}{6}ms^2$$

- Examples (non-uniform distribution):
 - Solid cylinder* of radius r, height h and mass m

$$I = \begin{bmatrix} \frac{1}{12}m(3r^2 + h^2) & & \\ & \frac{1}{12}m(3r^2 + h^2) & \\ & & \frac{1}{2}mr^2 \end{bmatrix}$$

Solid cone* of radius r, height h and mass m about apex

$$I = \begin{bmatrix} \frac{3}{5}mh^2 + \frac{3}{20}mr^2 \\ & \frac{3}{5}mh^2 + \frac{3}{20}mr^2 \\ & \frac{3}{10}mr^2 \end{bmatrix}$$

*axis parallel to third local basis

- Beware of rods: inverse cannot be "real"!
 - Rod* of length I and mass m spinning about end

$$I = \begin{bmatrix} \frac{1}{3}ml^2 \\ \frac{1}{3}ml^2 \\ 0 \end{bmatrix}$$

nope 🕾

Rod* of length I and mass m spinning about center

$$I = \begin{bmatrix} \frac{1}{12}ml^2 & & \\ & \frac{1}{12}ml^2 & \\ & & 0 \end{bmatrix}$$

*axis parallel to third local basis

- General tensor for a rigid system of particles:
- For a list of n particles, each with mass m_k and position x_k , calculate the total mass M and center of mass c_m :

$$M = \sum_{k=0}^{n-1} m_k$$

$$c_m = \frac{1}{M} \sum_{k=0}^{n-1} m_k x_k$$

Moment arm can be calculated at any point:

$$r=x-c_m$$

Center of mass c_m is our "reference point"

Inertia tensor from a set of masses:

$$I = \begin{bmatrix} I_{\chi\chi} & I_{y\chi} & I_{z\chi} \\ I_{\chi y} & I_{yy} & I_{zy} \\ I_{\chi z} & I_{yz} & I_{zz} \end{bmatrix}$$

$$I_{xy} = I_{yx}$$
, $I_{yz} = I_{zy}$, $I_{zx} = I_{xz}$

$$r_k = x_k - c_m = (a_k, b_k, c_k)$$

Calculate the terms:

$$I_{xx} = \sum_{k=0}^{n-1} m_k (b_k^2 + c_k^2)$$

$$I_{yy} = \sum_{k=0}^{n-1} m_k (c_k^2 + a_k^2)$$

$$I_{zz} = \sum_{k=0}^{n-1} m_k (a_k^2 + b_k^2)$$

$$I_{xy} = -\sum_{k=0}^{n-1} m_k a_k b_k$$

$$I_{yz} = -\sum_{k=0}^{n-1} m_k b_k c_k$$

$$I_{ZX} = -\sum_{k=0}^{n-1} m_k c_k a_k$$

• The inertia tensor is inverted and used to solve angular acceleration:

$$\vec{\tau} = I\vec{\alpha}$$
$$\vec{\alpha} = I^{-1}\vec{\tau}$$

- The final problem: *frame of reference*
- The vectors are in world space, while the tensor (and its inverse) is local to the object...

$$\vec{\alpha}_{\text{world}} = I_{\text{local}}^{-1} \vec{\tau}_{\text{world}}$$

nope ⊗

- Perform change of basis: bring torque into local space, apply tensor, move back to world
- $\frac{\text{world}}{R_{\text{local}}}$ is the object's world orientation:

$$\vec{\alpha}_{\text{world}} = \begin{pmatrix} \text{world} R_{\text{local}} I_{\text{local}}^{-1} & \text{local} R_{\text{world}} \end{pmatrix} \vec{\tau}_{\text{world}}$$

$$\vec{\alpha}_{\text{world}} = \begin{pmatrix} \text{world} R_{\text{local}} I_{\text{local}}^{-1} & \text{world} R_{\text{local}}^{-1} \end{pmatrix} \vec{\tau}_{\text{world}}$$

- Perform change of basis: bring torque into local space, apply tensor, move back to world
- Therefore, every update we calculate a "modified" inertia tensor (inverse) to use:

$$I_t^{-1} = R I^{-1} R^{-1}$$

The final angular acceleration update:

$$\alpha_t = I_t^{-1} \tau_t$$

• *Pro tip*: know the behavior of uniform scales:

$$S = \begin{bmatrix} S & & \\ & S & \\ & & S \end{bmatrix}$$

 What happens when you multiply it with any other 3x3 matrix 'A'?

$$SA \equiv sA = As \equiv AS$$

A uniform scale matrix commutes!

 If your inertia tensor is a uniform scale (e.g. sphere, cube), can cancel out change of basis:

$$I_t^{-1} = R I^{-1} R^{-1}$$

 $\equiv R R^{-1} I^{-1}$

Technically this also happens for mass...

$$a = R \ m_{3\times 3}^{-1} \ R^{-1} \ F$$
 $a \equiv R \ R^{-1} \ m_{3\times 3}^{-1} \ F$
 $a \equiv m_t^{-1} F$
 $m_t^{-1} = m^{-1}$ (constant scalar)

 Force application will also be off if we do not also update the center of mass (which is local)

$$r_{\text{fworld}} = x_{\text{fworld}} - c_{m_{\text{local}}}$$
 $c_{m_{\text{world}}} = {^{\text{world}}} R_{\text{local}} c_{m_{\text{local}}}$

All operations now solvable in world space:

$$r_{
m fworld} = x_{
m fworld} - c_{m_{
m world}}$$
 $au_{
m world} = r_{
m fworld} imes F_{
m world}$
 $lpha_{
m world} = I_{
m world}^{-1} au_{
m world}$

 We now have a complete set of dynamics formulas in the same frame of reference:

Force generators:

$$r_{f} = x_{f} - c_{m}$$

$$\tau = r_{f} \times F$$

$$F_{t} = \sum F$$

$$\tau_{t} = \sum \tau$$

$$a_{t} = m_{t}^{-1}F_{t} = m^{-1}F_{t}$$

$$\alpha_{t} = I_{t}^{-1}\tau_{t} = R I^{-1} R^{-1} \tau_{t}$$

Integration (Euler):

$$x_{t+dt} = x_t + v_t dt$$

$$v_{t+dt} = v_t + a_t dt$$

$$q_{t+dt} = q_t + \omega_t q_t \frac{dt}{2}$$

$$q'_{t+dt} = \hat{q}_{t+dt}$$

$$\omega_{t+dt} = \omega_t + \alpha_t dt$$

Food for thought: relationship with velocity:

$$\omega = \frac{r \times v}{|r|^2}$$

$$I = m|r|^2$$

Therefore,

$$L = I\omega = m|r|^{2} \frac{r \times v}{|r|^{2}}$$
$$= m(r \times v) = r \times mv = r \times p$$

- Final food for thought:
- Minimize the number of loops
 - E.g. calculate total mass and center of mass simultaneously
 - E.g. calculate and store inertia tensor (and inverse) once per object
- Don't forget to normalize orientation after integration!

- Challenge!
- Prove the tensor formulas for spheres and cubes using a collection of "particles"
- I.e. read discrete mesh data
- Box: calculation should yield solid or hollow box tensor
- Sphere: higher resolution; should be closer to solid or hollow tensor

The end.

Questions? Comments? Concerns?

