

# IIB\_branes

## IIB brane solutions

This is a quick reference for the standard half-BPS Type IIB brane solutions.

Unless stated otherwise, we use the **10d string-frame** metric, and split coordinates into

- worldvolume directions  $x^\mu$ ,  $\mu = 0, \dots, p$
- transverse directions  $y^i$ ,  $i = 1, \dots, 9 - p$

with flat line elements

$$ds^2(\mathbb{R}^{1,p}) = \eta_{\mu\nu} dx^\mu dx^\nu, \quad ds^2(\mathbb{R}^{9-p}) = \delta_{ij} dy^i dy^j.$$

## Dilaton and frames

We write the **10d dilaton** as a background value plus a fluctuation,

$$\Phi = \langle \Phi \rangle + \phi, \quad g_s = \langle e^\Phi \rangle = e^{\langle \Phi \rangle}, \quad e^\Phi = g_s e^\phi.$$

Unless otherwise stated, all solutions below give  $ds_{10, str}^2$  and  $\Phi$ .

The string-frame and Einstein-frame metrics are related by the Weyl rescaling

$$g_{MN}^{(E)} = e^{-\Phi/2} g_{MN}^{(str)}, \quad ds_{10, E}^2 = e^{-\Phi/2} ds_{10, str}^2.$$

In terms of  $g_s$  and  $\phi$  this is

$$ds_{10, E}^2 = g_s^{-1/2} e^{-\phi/2} ds_{10, str}^2, \quad ds_{10, str}^2 = g_s^{1/2} e^{\phi/2} ds_{10, E}^2.$$

For  $p < 7$  define  $r^2 = \delta_{ij} y^i y^j$  and the harmonic function

$$H(r) = 1 + \frac{Q_p}{r^{7-p}}.$$

For  $p = 7$  (codimension 2),  $H$  is replaced by a logarithm (see the D7 section).

We also use the shorthand volume form

$$\text{vol}_{1,p} = dx^0 \wedge dx^1 \wedge \cdots \wedge dx^p.$$

## Dp (template, $p$ odd)

For a Dp-brane with  $p = -1, 1, 3, 5$  (and formally  $p = 7, 9$  with caveats), the standard ansatz is

$$ds_{10, \text{str}}^2 = H^{-1/2} ds^2(\mathbb{R}^{1,p}) + H^{1/2} ds^2(\mathbb{R}^{9-p}),$$

$$e^\Phi = g_s H^{(3-p)/4}, \quad e^\phi = H^{(3-p)/4},$$

$$C_{p+1} = \frac{1}{g_s} (H^{-1} - 1) \text{vol}_{1,p}, \quad F_{p+2} = dC_{p+1}.$$

Using  $ds_{10,E}^2 = e^{-\Phi/2} ds_{10, \text{str}}^2$ , the Einstein-frame metric for a Dp in this normalization is

$$ds_{10,E}^2 = g_s^{-1/2} (H^{(p-7)/8} ds^2(\mathbb{R}^{1,p}) + H^{(p+1)/8} ds^2(\mathbb{R}^{9-p})).$$

The magnetic dual satisfies  $F_{8-p} = (-1)^{p*} {}_{10}F_{p+2}$ .

## F1 (fundamental string)

Worldvolume:  $\mathbb{R}^{1,1}$ .

$$ds_{10, \text{str}}^2 = H^{-1} ds^2(\mathbb{R}^{1,1}) + ds^2(\mathbb{R}^8),$$

$$e^{2\Phi} = g_s^2 H^{-1}, \quad e^{2\phi} = H^{-1},$$

$$B_2 = (H^{-1} - 1) \text{vol}_{1,1}, \quad H_3 = dB_2.$$

## NS5

Worldvolume:  $\mathbb{R}^{1,5}$ .

$$ds_{10, \text{str}}^2 = ds^2(\mathbb{R}^{1,5}) + H ds^2(\mathbb{R}^4),$$

$$e^{2\Phi} = g_s^2 H, \quad e^{2\phi} = H,$$

$$H_3 = *_\mathbb{R}^4 dH.$$

## D(-1) (D-instanton)

This is Euclidean and is most cleanly written in Einstein frame, where the metric is flat:

$$ds_{10,E}^2 = ds^2(\mathbb{R}^{10}).$$

The axio-dilaton varies with a harmonic function  $H(r) = 1 + \frac{Q_{-1}}{r^8}$ :

$$e^\Phi = g_s H, \quad e^\phi = H,$$

and the RR scalar  $C_0$  is sourced with

$$dC_0 = \pm d(H^{-1}).$$

Equivalently, in terms of the axio-dilaton

$$au = C_0 + ie^{-\Phi} = C_0 + ig_s^{-1}e^{-\phi}$$

one has a BPS relation between  $C_0$  and the dilaton profile.

## D1

Worldvolume:  $\mathbb{R}^{1,1}$ .

$$ds_{10,str}^2 = H^{-1/2} ds^2(\mathbb{R}^{1,1}) + H^{1/2} ds^2(\mathbb{R}^8),$$

$$e^\Phi = g_s H^{1/2}, \quad e^\phi = H^{1/2},$$

$$C_2 = \frac{1}{g_s} (H^{-1} - 1) \text{vol}_{1,1}, \quad F_3 = dC_2.$$

## D3

Worldvolume:  $\mathbb{R}^{1,3}$ .

$$ds_{10,str}^2 = H^{-1/2} ds^2(\mathbb{R}^{1,3}) + H^{1/2} ds^2(\mathbb{R}^6),$$

$$e^\Phi = g_s, \quad \phi = 0,$$

$$C_4 = \frac{1}{g_s} (H^{-1} - 1) \text{vol}_{1,3}.$$

The self-dual 5-form field strength is

$$F_5 = (1 + *_{10}) dC_4.$$

## D5

Worldvolume:  $\mathbb{R}^{1,5}$ .

$$ds_{10, str}^2 = H^{-1/2} ds^2(\mathbb{R}^{1,5}) + H^{1/2} ds^2(\mathbb{R}^4),$$

$$e^\Phi = g_s H^{-1/2}, \quad e^\phi = H^{-1/2}.$$

One can write either the electric 6-form potential or the magnetic 2-form potential:

$$C_6 = \frac{1}{g_s} (H^{-1} - 1) \text{vol}_{1,5}, \quad F_7 = dC_6,$$

with  $F_3 = - *_{10} F_7$ .

## D7

Worldvolume:  $\mathbb{R}^{1,7}$  and transverse space is 2d. A single isolated D7-brane has logarithmic behavior rather than a power-law harmonic function.

A common local form (away from the brane core) is to take

$$H(r) = h_0 + Q_7 \log \frac{r}{r_0}, \quad r^2 = (y^1)^2 + (y^2)^2,$$

and use the same warped metric ansatz

$$ds_{10, str}^2 = H^{-1/2} ds^2(\mathbb{R}^{1,7}) + H^{1/2} (dr^2 + r^2 d\theta^2),$$

$$e^\Phi = g_s H^{-1}, \quad e^\phi = H^{-1}.$$

The D7 is electrically charged under  $C_8$  (and magnetically under  $C_0$ ):

$$C_8 = \frac{1}{g_s} (H^{-1} - 1) \text{vol}_{1,7}, \quad F_9 = dC_8.$$

More precisely, the exact backreacted solution is naturally described in terms of the axio-dilaton  $\tau = C_0 + i e^{-\Phi}$  which is holomorphic in the transverse complex coordinate and has  $SL(2, \mathbb{Z})$  monodromy around the brane.

## D9

Space-filling: there is no transverse space. In practice one takes  $H = 1$ , so the fields are constant and the metric is flat.

$$ds_{10, str}^2 = ds^2(\mathbb{R}^{1,9}), \quad e^\Phi = g_s, \quad \phi = 0.$$

## (p,q) multiplets (generated by S-duality)

Type IIB has an  $SL(2, \mathbb{Z})$  duality acting on the axio-dilaton  $\tau$  and the two-form doublet  $(B_2, C_2)$ . Applying it to F1 and NS5 generates the full family of  $(p, q)$  strings and  $(p, q)$  fivebranes. In Einstein frame the metric warp factors take the same schematic form as above, while the dilaton and the NS/RR potentials mix into  $SL(2, \mathbb{Z})$  covariant combinations.