

Polchinski 1.3

$$x = \frac{1}{4\pi} \int_M d\tau d\sigma (-r)^{1/2} R$$

$$\text{Weyl: } (-r')^{1/2} R' = (-r)^{1/2} (R - 2\bar{v}^2 w).$$

$$\Rightarrow \Delta x = \frac{1}{4\pi} \int_M d\tau d\sigma (-r)^{1/2} (-2\bar{v}^2 w)$$

$$= -\frac{1}{2\pi} \int_M d\tau d\sigma (-r)^{1/2} \bar{v}^2 w$$

$$\text{write } (-r)^{1/2} \partial_a [\bar{v}^a w] = \partial_a [(-r)^{1/2} (\partial^a w)]$$

$$\Delta x = -\frac{1}{2\pi} \int_M d\sigma^a \partial_a [(-r)^{1/2} (\partial^a w)]$$

↑
M

$$= -\frac{1}{2\pi} \oint_M d\vec{\ell} (-r)^{1/2} (\partial^a w) \vec{\ell}^a = \oint_S d\vec{\ell} \cdot \vec{F}$$

$$\Rightarrow \Delta x = -\frac{1}{2\pi} \oint_M d\sigma \partial_a (\partial^a w) n_a$$

[s] affine parameter,
proper time, no unit

no norm,

$$\int_M d\sigma \partial_a (\partial^a w) = \int_M d\sigma \partial_a \partial^a w$$

time = distance / speed, $\partial_a \partial^a w$

Now we consider

$$\Delta_{\text{wedge}} \xrightarrow[2\pi]{1} \int ds \ k, \quad k = t^a n_b \nabla a^b$$

$\left. \begin{array}{l} t^a \text{ unit tangent} \\ n_b \text{ unit normal} \end{array} \right\}$

$$\nabla a^b = \partial a^b + \Gamma_{ac}^b t^c$$

$$\Gamma_{ac}^b = \frac{1}{2} g^{b\lambda} [g_{a\lambda,c} + g_{c\lambda,a} - g_{ac,\lambda}]$$

$$\text{under } g_{ab} \xrightarrow{e^{2w}} g_{ab}, \quad \Gamma \rightarrow \Gamma'$$

$$\begin{aligned} \Gamma_{ac}^{b'} &= \frac{1}{2} e^{-2w} g^{b\lambda} [\partial_c [e^{2w} g_{a\lambda}] + \partial_a [e^{2w} g_{c\lambda}] \\ &\quad - \partial_\lambda [e^{2w} g_{ac}]] \end{aligned}$$

$$\begin{aligned} &= \frac{1}{2} e^{-2w} g^{b\lambda} \left\{ 2(\partial_c w) e^{2w} g_{a\lambda} + e^{2w} g_{a\lambda,c} + \right. \\ &\quad \left. - 2(\partial_a w) e^{2w} g_{c\lambda} + e^{2w} g_{c\lambda,a} - \right. \\ &\quad \left. - 2(\partial_\lambda w) e^{2w} g_{ac} - e^{2w} g_{ac,\lambda} \right\} \end{aligned}$$

$$\Gamma_{ac}^{b'} = \Gamma_{ac}^b + g^{b\lambda} [(\partial_c w) g_{a\lambda} + (\partial_a w) g_{c\lambda} - (\partial_\lambda w) g_{ac}]$$

$$= \Gamma_{ac}^b + [(\partial_c w) g_a^b + (\partial_a w) g_c^b - (\partial^b w) g_{ac}]$$

$$\Rightarrow \nabla_a t^b \rightarrow \nabla_a t^b + g^b_a (\partial_c w) t^c + \cancel{(\partial_a w) t^b} - (\partial^b w) t_a$$

$$\Rightarrow t^a \nabla_a t^b \rightarrow t^a \nabla_a t^b + (\partial_c w) t^c t^b + (\partial_a w) t^a t^b - (\partial^b w) t_a t^a$$

$$= t^a \nabla_a t^b + 2 (\partial_c w) t^c t^b - (\partial^b w) \quad (t_a t^a = 1)$$

$$\Rightarrow n_b t^a \nabla_a t^b \rightarrow n_b t^a \nabla_a t^b \left\{ - (\partial^b w) n_b \right\} \quad (t^b n_b = 0)$$

Plugging this back into $\frac{1}{2\pi} \int_{\partial M} ds k$, $k = t^a n_b \nabla_a t^b$

$$\Delta_{\text{west}} \frac{1}{2\pi} \int_{\partial M} ds k = \underbrace{\frac{1}{2\pi} \int_{\partial M} ds (\Delta k)}_{= \left| \frac{1}{2\pi} \int_{\partial M} ds (-\partial^b w) n_b \right|}$$

This matches Δx for $x = \frac{1}{4\pi} \int d\sigma^a (-\tau)^{1/2} R$ we found earlier, except for a sign change, but we can just let ~~$\tau \rightarrow -\tau$~~ $k \rightarrow \pm k$.

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