

Pulchinski 1.3

$$\chi = \frac{1}{4\pi} \int_M d\tau d\sigma (-\tau)^{1/2} R$$

$$\text{Weyl: } (-\tau)^{1/2} R' = (-\tau)^{1/2} (R - 2 \nabla^2 w).$$

$$\Rightarrow \Delta \chi = \frac{1}{4\pi} \int_M d\tau d\sigma (-\tau)^{1/2} (-2 \nabla^2 w)$$

$$= -\frac{1}{2\pi} \int_M d\tau d\sigma (-\tau)^{1/2} \nabla^2 w.$$

$$\text{write } (-\tau)^{1/2} \nabla_a [\nabla^a w] = \partial_a [(-\tau)^{1/2} (\partial^a w)]$$

$$\Delta \chi = -\frac{1}{2\pi} \int_M d\sigma^a \partial_a [(-\tau)^{1/2} (\partial^a w)]$$

$\uparrow$

$$\text{cf form } \int d\vec{x} \, \vec{\nabla} \cdot \vec{F}(\vec{x})$$

$$= -\frac{1}{2\pi} \int_{\partial M} d\ell (-\tau)^{1/2} (\partial^a w)$$

$$= \oint_S d\vec{\ell} \cdot \vec{F}$$

$$\Rightarrow \Delta \chi = -\frac{1}{2\pi} \int_{\partial M} ds (\partial^a w) n_a$$

$s$  affine parameter,  
proper time, no unit  
normal.

Now we consider

$$\Delta_{\text{negh}} = \frac{1}{2\pi} \int_{\partial M} ds \, k, \quad k = t^a n_b \nabla_a t^b$$

$\left\{ \begin{array}{l} t^a \text{ unit tangent} \\ n_b \text{ unit normal} \end{array} \right.$

$$\nabla_a t^b = \partial_a t^b + \Gamma_{ac}^b t^c$$

$$\Gamma_{ac}^b = \frac{1}{2} g^{b\lambda} [ \partial_{a\lambda} g_{c\lambda} + \partial_{c\lambda} g_{a\lambda} - \partial_{ac} g_{\lambda\lambda} ]$$

under  $g_{ab} \rightarrow e^{2w} g_{ab}, \quad \Gamma \rightarrow \Gamma'$

$$\Gamma_{ac}^{b'} = \frac{1}{2} e^{-2w} g^{b\lambda} [ \partial_c [ e^{2w} g_{a\lambda} ] + \partial_a [ e^{2w} g_{c\lambda} ] - \partial_\lambda [ e^{2w} g_{ac} ] ]$$

$$= \frac{1}{2} e^{-2w} g^{b\lambda} \left\{ 2(\partial_c w) e^{2w} g_{a\lambda} + e^{2w} g_{a\lambda, c} + 2(\partial_a w) e^{2w} g_{c\lambda} + e^{2w} g_{c\lambda, a} - 2(\partial_\lambda w) e^{2w} g_{ac} - e^{2w} g_{ac, \lambda} \right\}$$

$$\Gamma_{ac}^{b'} = \Gamma_{ac}^b + g^{b\lambda} [ (\partial_c w) g_{a\lambda} + (\partial_a w) g_{c\lambda} - (\partial_\lambda w) g_{ac} ]$$

$$= \Gamma_{ac}^b + [ (\partial_c w) g_a^b + (\partial_a w) g_c^b - (\partial^b w) g_{ac} ]$$



$$\Rightarrow \nabla_a t^b \rightarrow \nabla_a t^b + g^b_a (\partial_c w) t^c + \cancel{g^b_a} (\partial_a w) t^b - (\partial^b w) t_a$$

$$\begin{aligned} \Rightarrow t^a \nabla_a t^b &\rightarrow t^a \nabla_a t^b + (\partial_c w) t^c t^b + (\partial_a w) t^a t^b - (\partial^b w) t_a t^a \\ &= t^a \nabla_a t^b + 2 (\partial_c w) t^c t^b - (\partial^b w) \quad (t_a t^a = 1) \end{aligned}$$

$$\Rightarrow n_b t^a \nabla_a t^b \rightarrow n_b t^a \nabla_a t^b \left( - (\partial^b w) n_b \right) \quad (t^b n_b = 0)$$

plugging this back into  $\frac{1}{2\pi} \int_{\partial M} ds k$ ,  $k = t^a n_b \nabla_a t^b$

$$\triangle_{\text{weyl}} \frac{1}{2\pi} \int_{\partial M} ds k = \frac{1}{2\pi} \int_{\partial M} ds (\Delta k)$$

$$= \left| \frac{1}{2\pi} \int_{\partial M} ds (-\partial^b w) n_b \right|$$

This matches  $\Delta \chi$  for  $\chi = \frac{1}{4\pi} \int d^2 \sigma (-\gamma)^{1/2} R$  we found earlier, except for a sign change, but we can just let  $\pm k \rightarrow \mp k$ .

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