

Name grading key

Homework Assignment #13 due in class Friday December 2

Staple this cover sheet in front of your solutions.

Write the requested answers on this sheet, and do the detailed solutions on your own paper.

[71] Problem 8.4 ★

Answer: Write the equation of motion.  $\mu \ddot{\mathbf{r}} + \nabla U = 0$  :: 1 point

[72] Problem 8.6 ★

Answer: No answer required here. Check the calculations of  $l_1$  and  $l_2$ . :: 1 point

[73] Problem 8.12 ★★

Answer: Write the equations for  $\omega_\theta$  and  $\omega_r$ .  $\omega_\theta = \omega_r = \mu(Gm_1m_2)^2 / l^{3/2}$  :: 2 points

[74] Problem 8.15 ★

Answer: By what percent would you expect the "constant" to vary?  
percent variation  $\sim 0.1$  :: 1 point

[75] Problem 8.16 ★★

Answer: Write the equations for x and y. :: 2 points

$$x(\varphi) = c \cos \varphi / (1 + \varepsilon \cos \varphi) \quad \text{and} \quad y(\varphi) = c \sin \varphi / (1 + \varepsilon \cos \varphi)$$

$$(x + \varepsilon a)^2 / a^2 + y^2 / b^2 = 1 \quad \text{where } a \text{ and } b \text{ are as given in Eq. 8.52.}$$

# Homework Assignment #13

## [71] Problem 8.4

The Lagrangian is  $\mathcal{L} = \frac{1}{2}M\dot{\vec{R}}^2 + \frac{1}{2}\mu\dot{\vec{r}}^2 - U(r)$   
 and Lagrange's equation is  $\frac{d}{dt}\left(\frac{\partial \mathcal{L}}{\partial \dot{q}}\right) - \frac{\partial \mathcal{L}}{\partial q} = 0$ .

$$\text{For } q = x, \quad \frac{d}{dt}(\mu\dot{x}) + \frac{\partial U}{\partial x} = \mu\ddot{x} + \frac{\partial U}{\partial x} = 0$$

$$\text{For } q = y, \quad \mu\ddot{y} + \frac{\partial U}{\partial y} = 0$$

$$\text{For } q = z, \quad \mu\ddot{z} + \frac{\partial U}{\partial z} = 0.$$

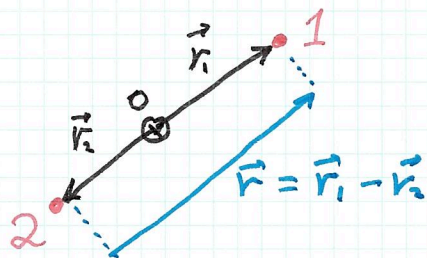
Thus

$$\mu\ddot{\vec{r}} + \nabla U = 0$$

which is the same as a particle of mass  $\mu$   
 in potential energy  $U(\vec{r})$ .

[92] Problem 8.6

The CM frame of reference



$$\vec{r}_1 = \frac{m_2}{M} \vec{r}$$

$$\vec{r}_2 = -\frac{m_1}{M} \vec{r}$$

$$\vec{\ell}_1 = \vec{r}_1 \times m_1 \dot{\vec{r}}_1 = m_1 \left(\frac{m_2}{M}\right)^2 \vec{r} \times \dot{\vec{r}} = \frac{m_2}{M} \mu \vec{r} \times \dot{\vec{r}}$$

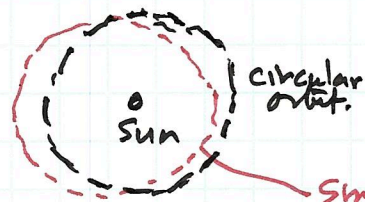
$$\vec{\ell}_2 = \vec{r}_2 \times m_2 \dot{\vec{r}}_2 = m_2 \left(\frac{m_1}{M}\right)^2 \vec{r} \times \dot{\vec{r}} = -\frac{m_1}{M} \mu \vec{r} \times \dot{\vec{r}}$$

$$\vec{L} = \vec{\ell}_1 + \vec{\ell}_2 = \mu \vec{r} \times \dot{\vec{r}}$$

Thus  $\vec{\ell}_1 = \frac{m_2}{M} \vec{L}$  and  $\vec{\ell}_2 = -\frac{m_1}{M} \vec{L}$ , as claimed.



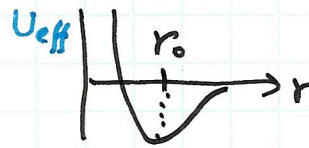
## [73] Problem 8.12



Small deviation from a circular orbit.

(a) For the circular orbit

$$U_{\text{eff}}(r) = -\frac{Gm_1 m_2}{r} + \frac{l^2}{2\mu r^2}$$

$U_{\text{eff}}$  

$$U'_{\text{eff}}(r) = \frac{Gm_1 m_2}{r^2} - \frac{l^2}{\mu r^3}$$

$$U'_{\text{eff}}(r_0) = 0 \Rightarrow \boxed{r_0 = \frac{l^2}{\mu Gm_1 m_2}}$$

(b) Stability of the circular orbit.

Consider  $r(t) = r_0 + \epsilon(t)$  where  $\epsilon$  is small.

$$U''_{\text{eff}}(r_0) = \frac{-2Gm_1 m_2}{r_0^3} + \frac{3l^2}{\mu r_0^4} = \frac{1}{r_0^3} \left[ -2Gm_1 m_2 + \frac{3l^2}{\mu} \frac{\mu Gm_1 m_2}{l^2} \right]$$

$$= \frac{Gm_1 m_2}{r_0^3} \text{ which is positive; so the circular orbit is stable.}$$

Small radial oscillations.

$\mu \ddot{r} + \frac{\partial U_{\text{eff}}}{\partial r} = 0$  where  $r(t) = r_0 + \epsilon(t)$  small

$$U_{\text{eff}}(r_0 + \epsilon) \approx U_{\text{eff}}(r_0) + \underbrace{\epsilon U'_{\text{eff}}(r_0)}_{=0} + \frac{1}{2} \epsilon^2 \underbrace{U''_{\text{eff}}(r_0)}_{Gm_1 m_2 / r_0^3}$$

$$\text{So } \mu \ddot{\epsilon} + \epsilon \frac{Gm_1 m_2}{r_0^3} = 0$$

$$\ddot{\epsilon} + \omega_R^2 \epsilon = 0 \Rightarrow \text{radial oscillations have } \omega_R^2 = \frac{Gm_1 m_2}{\mu r_0^3}$$

$$\text{The angular velocity is } \omega_\phi = \frac{d\phi}{dt} = \frac{l}{\mu r_0^2}.$$

$$\frac{\omega_R}{\omega_\phi} = \sqrt{\frac{Gm_1 m_2}{\mu r_0^3}} \frac{\mu r_0^2}{l} = \sqrt{\frac{Gm_1 m_2}{\mu r_0^3}} \frac{\mu r_0^2}{\sqrt{r_0 \mu Gm_1 m_2}} = \underline{\underline{1}}$$

∴ The period of radial oscillations is equal to the period of revolutions of the planet.

[74] Problem 8.15

Kepler's third law states  $\tau^2 = C a^3$  for any planet, where  $C$  is a constant. It is approximately true, and  $C \approx \frac{4\pi^2}{GM_{\text{sun}}}$ .

But more accurately,  $C = \frac{4\pi^2}{G(M_{\text{sun}} + m)}$  where  $m = \text{mass of the planet}$ .

The planet masses range from Jupiter,  $M_J = 2 \times 10^{27} \text{ kg}$  to much smaller masses (e.g., Earth,  $m_E = 6 \times 10^{24} \text{ kg}$ ).

So the variations of  $C$  are of order

$$\begin{aligned} \frac{C_E}{C_J} &= \frac{M_S + m_J}{M_S + m_E} \approx 1 + \frac{m_J}{M_S} = 1 + \frac{2 \times 10^{27} \text{ kg}}{2 \times 10^{30} \text{ kg}} \\ &= 1.001 \end{aligned}$$

Variation  $\sim 0.001 = 0.1 \text{ percent}$

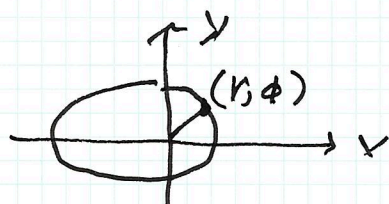


# [75] Problem 8.16

In polar coordinates, the equation for a Keplerian orbit is  $r(\phi) = \frac{C}{1 + \epsilon \cos \phi}$  where  $C > 0$  and  $\epsilon \geq 0$ .

Consider a bounded orbit:  $0 \leq \epsilon < 1$ .

Write the orbit equation in Cartesian coordinates  $(x, y)$ .



$$x = r \cos \phi \quad \text{and} \quad y = r \sin \phi$$

We have  $r = \frac{C}{1 + \epsilon x/r}$

- $r + \epsilon x = C$
- $r = C - \epsilon x$
- $r^2 = x^2 + y^2 = (C - \epsilon x)^2$   
 $= C^2 - 2C\epsilon x + \epsilon^2 x^2$
- $(1 - \epsilon^2)x^2 + 2C\epsilon x + y^2 = C^2$
- $(1 - \epsilon^2)\left[x + \frac{C\epsilon}{1 - \epsilon^2}\right]^2 + y^2$   
 $= C^2 + \frac{C^2 \epsilon^2}{1 - \epsilon^2} = \frac{C^2}{1 - \epsilon^2}$

• Thus

$$\frac{\left(x + \frac{C\epsilon}{1 - \epsilon^2}\right)^2}{C^2/(1 - \epsilon^2)^2} + \frac{y^2}{C^2/(1 - \epsilon^2)} = 1$$

Result

$$\frac{(x+d)^2}{a^2} + \frac{y^2}{b^2} = 1$$

where

$$d = \frac{C\epsilon}{1 - \epsilon^2} \quad \text{or} \quad d = \epsilon a$$

$$a^2 = \frac{C^2}{(1 - \epsilon^2)^2} \quad \text{or} \quad a = \frac{C}{1 - \epsilon^2}$$

$$b^2 = \frac{C^2}{1 - \epsilon^2} \quad \text{or} \quad b = \frac{C}{\sqrt{1 - \epsilon^2}}$$