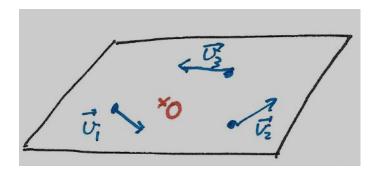
# Section 3.5

# Angular Momentum for Several Particles

#### Read Section 3.5.

Consider N particles, with masses  $m_{\alpha}$ , velocities  $\boldsymbol{v}_{\alpha}$  and positions  $\boldsymbol{r}_{\alpha}$ .

$$\{ \alpha = 1 \ 2 \ 3 \dots N \}$$



The total angular momentum is

$$\mathbf{L} = \sum_{\alpha=1}^{\mathbf{N}} \boldsymbol{\ell}_{\alpha}$$

$$\mathbf{L} = \sum_{\alpha=1}^{\mathbf{N}} m_{\alpha} \mathbf{r}_{\alpha} \times \mathbf{v}_{\alpha}$$

Remember:  $\ell$  and L are defined w.r.t. O.

### The importance of L ...

[I] If the internal forces are *central*, thendL/dt = the external torque.

[I] For an isolated system with central internal forces, dL/dt = 0; i.e., then L is a constant of the motion.

# Figure 3.8

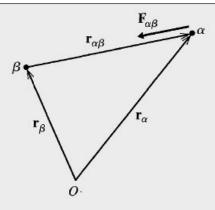


Figure 3.8 The vector  $\mathbf{r}_{\alpha\beta} = (\mathbf{r}_{\alpha} - \mathbf{r}_{\beta})$  points to particle  $\alpha$  from particle  $\beta$ . If the force  $\mathbf{F}_{\alpha\beta}$  is central (points along the line joining  $\alpha$  and  $\beta$ ), then  $\mathbf{r}_{\alpha\beta}$  and  $\mathbf{F}_{\alpha\beta}$  are collinear and their cross product is zero.

# What is this figure telling us?

For an isolated system with central forces, the total angular momentum is constant.

Consider a system with two isolated particles and a *central* force.

 $F_{\alpha\beta}$  = the force on  $\alpha$  exerted by  $\beta$ 

Consider an arbitrary origin:

$$r_{\alpha}$$
 = position vector of  $\alpha$ 

$$r_{\beta}$$
 = position vector of  $\beta$ 

The total angular momentum is

$$\mathbf{L} = \boldsymbol{\ell}_{\alpha} + \boldsymbol{\ell}_{\beta}$$

$$= \mathbf{m}_{\alpha} \mathbf{r}_{\alpha} \times \mathbf{v}_{\alpha} + \mathbf{m}_{\beta} \mathbf{r}_{\beta} \times \mathbf{v}_{\beta}$$

and so

$$dL/dt = \mathbf{r}_{\alpha} \times \mathbf{F}_{\alpha\beta} + \mathbf{r}_{\beta} \times \mathbf{F}_{\beta\alpha}$$

$$= (r_{\alpha} - r_{\beta}) \times F_{\alpha\beta} = r_{\alpha\beta} \times F_{\alpha\beta} = 0$$

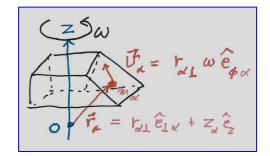
#### Moment of Inertia I

- You learned this in PHY 183.
- I "Moment of inertia" is a property of a solid body.
- I.e., there is a continuum mass density,  $\rho(\mathbf{r})$ .
- I Often we'll have  $\rho(r) = constant$ , called uniform mass density.
- I Moment of inertia is defined with respect to an axis of rotation, which might be a symmetry axis of the body (but not necessarily).

# <u>Definition of the moment of inertia</u>

(This is Taylor's Problem 3.30; more complete discussion in Chapter 10.)

Consider rotation about the z axis,



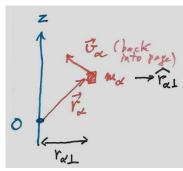
Divide the body into N small parts (treat them as particles) with masses  $m_{\alpha}$  {  $\alpha$  = 1 2 3 ... N }; then take the continuum limit N  $\rightarrow \infty$  and  $m_{\alpha} \rightarrow 0$  with M constant.

What is the total angular momentum?

# Definition of the moment of inertia

Consider rotation about the z axis.

Let  $m_{_{\alpha}}$  be one particle in the system



What is the angular momentum?

$$\vec{l}_{d} = m_{d} \vec{r}_{d} \times \vec{v}_{d}$$

$$= m_{d} (r_{d\perp} \hat{e}_{d\perp} + Z_{d} \hat{e}_{z}) \times (r_{d\perp} \omega \hat{e}_{d} \omega)$$

$$= m_{d} (r_{d\perp}^{2} \omega \hat{e}_{z} + m_{d} Z_{d} r_{d\perp} \omega (-\hat{e}_{d\perp}))$$

For uniform density, 
$$\beta = \frac{M}{\int_{B} r_{\perp}^{2} dV}$$

$$I = M \frac{\int_{B} r_{\perp}^{2} dV}{\int_{B} dV}$$

#### Example 3.3

# A lump of putty collides with a turntable

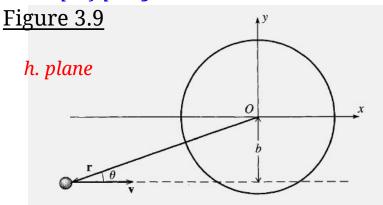


Figure 3.9 A lump of putty of mass m is thrown with velocity  $\mathbf{v}$  at a stationary turntable. The putty's line of approach passes within the distance b of the table's center O.

After the collision the lump of putty sticks to the turntable (i.e., it's an *inelastic collision*).

Note the *impact parameter* b defined in the picture.

The problem is to calculate the final angular velocity of the turntable with the lump of putty stuck to it.

# § The principle is conservation of angular momentum.

§ Before the collision (anyplace on the dashed line)

$$\vec{L} = m \left( * \hat{e}_x + y \hat{e}_y \right) * \left( v_x \hat{e}_x \right)$$

$$= m \left( -b \right) v_o \left( -\hat{e}_z \right) = m b v_o \hat{e}_z$$

§ After the collision:

$$L_{Z} = I\omega + mR^{2}\omega$$

$$= (I + mR^{2})\omega$$

$$\omega = \frac{mb \sigma_{0}}{I + mR^{2}}$$

For a disk, 
$$I = \frac{1}{2}MR^2 \Rightarrow \omega = \frac{bv_0}{R^2(\frac{M}{2m} + 1)}$$

Angular momentum vector of a rigid body that rotates about an axis through the center of mass position

$$\int_{B}^{Z} \int_{CoM} \int$$

=  $I_{\text{for this axis}} \omega e_{\mathbf{z}}$  because we specified that the CoM lies on the axis of rotation;  $\oint e_{\mathbf{z}} d\phi = 0$ 

### Example 3.4

# A Sliding and Spinning Dumbbell

See Figure 3.10. Kick the sphere on the left—an *impulsive force*—as shown. Calculate the motion.

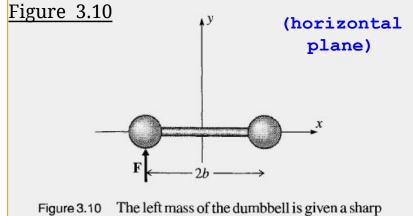


Figure 3.10 I he left mass of the dumbbell is given a snarp tap in the y direction.

$$M = 2m$$
 and  $I = 2 m b^2$ 

#### **Principles:**

The impulsive force causes the center of mass to accelerate briefly; impulse  $F \Delta t = \kappa$ ;

$$d\mathbf{P}/dt = \mathbf{F}^{\mathbf{ext}} \Rightarrow \Delta \mathbf{P} = \mathbf{F} \Delta t = \kappa \mathbf{e}_{\mathbf{v}}$$

The impulsive *torque* causes a brief angular acceleration;

$$dL_z/dt = N_z = -bF$$
  $\Rightarrow$   $\Delta L_z = -bF \Delta t = -b\kappa$ 

After the impulse, the momentum and angular momentum are constant ( $\exists$  no force and  $\exists$  no torque);

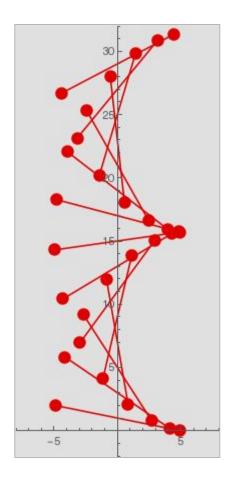
$$v_{CM} = P/M = \kappa / (2m)$$

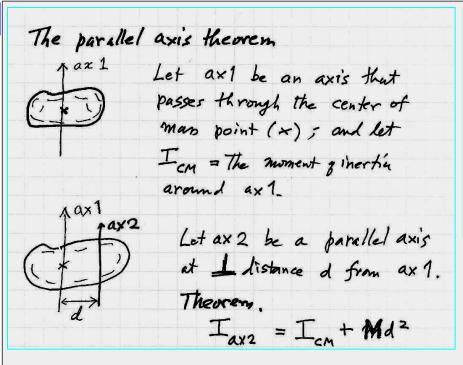
and

$$\omega = L/I = -b\kappa/I$$
 so  $\omega = -v_{CM}/b$ 

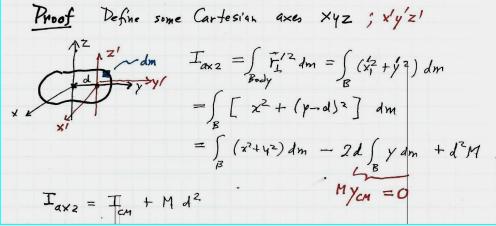
# Describe the motion after the impulsive force:

- the center of mass moves along the y axis;
- the two spheres revolve round the center of mass;
- $\mathbf{v}_{cm} = -\mathbf{b} \ \omega$





Chapter 10 (PHY 422) – the inertia tensor



Homework Assignment #6
due in class Friday, October 14
[27] Problem 3.16 \*
[28] Problem 3.20 \*\*
[29] Problem 3.22 \*\*
[30] Problem 3.27 \*\*
[30x] Problem 3.35 \*\*

Use the cover sheet.