

# Physics 410 -- Useful Formulas #1

## 0. Physical Constants

$$N_A = 6.02 \times 10^{23} \text{ mol}^{-1} \quad k_B = 1.38 \times 10^{-23} \text{ J/K} \quad N_A k_B = 8.31 \text{ J/(mol} \cdot \text{K)}$$

$$\hbar = 1.055 \times 10^{-34} \text{ J-s} \quad e = 1.60 \times 10^{-19} \text{ C} \quad m_p = 1.67 \times 10^{-27} \text{ kg}$$

$$\text{Kittel and Kroemer notation: } \tau = k_B T, \quad \sigma = S / k_B$$

## I. Probability and statistics, and other mathematical formulas:

$$\text{mean value and variance: } \bar{X} \equiv \langle X \rangle = \sum_s X(s) P(s), \quad \langle (\Delta X)^2 \rangle = \langle X^2 \rangle - \langle X \rangle^2$$

$$\text{where } P(s) \text{ is a normalized probability distribution: } \sum_s P(s) = 1$$

$$\text{binomial distribution: } (p + q)^N = \sum_{n=0}^N \frac{N!}{n!(N-n)!} p^n q^{N-n}$$

$$\text{geometric series: } \sum_{n=0}^N x^n = \frac{1 - x^{N+1}}{1 - x}, \quad \sum_{n=0}^{\infty} x^n = \frac{1}{1 - x}, \quad \text{for } |x| < 1$$

$$\text{Stirling's approximation: } \ln(n!) \approx \frac{1}{2} \ln(2\pi) + (n + \frac{1}{2}) \ln(n) - n$$

$$\text{binomial multiplicity for large N: } g(N, s) = \frac{N!}{(\frac{N}{2} + s)!(\frac{N}{2} - s)!} \approx \left(\frac{2}{\pi N}\right)^{1/2} 2^N e^{-2s^2 / N}$$

$$\text{Gaussian integrals: } \int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi} \quad \int_{-\infty}^{\infty} x^2 e^{-x^2} dx = \frac{\sqrt{\pi}}{2}$$

$$\text{Normalized Gaussian probability distribution: } P(x) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{(x-\bar{x})^2}{2\sigma^2}}$$

$$\text{Taylor series: } f(x_0 + \Delta x) = f(x_0) + f'(x_0)\Delta x + \frac{f''(x_0)}{2!}(\Delta x)^2 + \dots$$

$$\text{examples: } \ln(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} - \dots \quad e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots$$

## II. Microcanonical ensemble: independent variables $U, V, N$

multiplicity function:  $g(U, V, N)$ ; entropy:  $\sigma(U, V, N) = \ln g(U, V, N)$

$$\text{temperature: } \frac{1}{\tau} = \left( \frac{\partial \sigma}{\partial U} \right)_{V, N} \quad \text{pressure: } p = \tau \left( \frac{\partial \sigma}{\partial V} \right)_{U, N}$$

Alternative formulation with independent variables  $\sigma, V, N$

$$\text{temperature: } \tau = \left( \frac{\partial U}{\partial \sigma} \right)_{V, N} \quad \text{pressure: } p = - \left( \frac{\partial U}{\partial V} \right)_{\sigma, N}$$

## III. Canonical ensemble: independent variables $\tau, V, N$

$$\text{Partition function: } Z = \sum e^{-\frac{\epsilon_s}{\tau}}, \text{ Canonical distribution function: } P_s = \frac{e^{-\frac{\epsilon_s}{\tau}}}{Z}$$

The numerator of  $P_s$  is called the "Boltzmann factor"

Partition function for a system of  $N$  identical subsystems or particles:

$$\text{Distinguishable: } Z_N = (Z_1)^N \quad \text{indistinguishable, Classical limit: } Z_N = \frac{(Z_1)^N}{N!}$$

$$\text{Mean Energy: } U = \tau^2 \frac{\partial(\ln Z)}{\partial \tau} = - \frac{\partial(\ln Z)}{\partial \beta} \quad \text{where } \beta = \frac{1}{\tau}$$

$$\text{Helmholtz free energy: } F = U - \tau \sigma = -\tau \ln Z$$

$$\text{entropy: } \sigma = - \left( \frac{\partial F}{\partial \tau} \right)_{V, N} \quad \text{pressure: } p = - \left( \frac{\partial F}{\partial V} \right)_{\tau, N}$$

$$\text{heat capacity: } C_V = \left( \frac{\partial U}{\partial \tau} \right)_V = \tau \left( \frac{\partial \sigma}{\partial \tau} \right)_V$$

## IV. Thermodynamic Identity for systems with fixed $N$ : $dU = \tau d\sigma - p dV = T dS - p dV$

For reversible processes:  $dQ = \tau d\sigma$ ,  $dW = p dV$ , so  $dU = dQ - dW$  for constant  $N$

Compare 1st Law of Thermodynamics:  $\Delta U = Q - W$ ,  $W$  is work done by the system.