

# PHY 410 Homework 4 Solutions

1. K + K. Ch 4, #2 Earth-Sun distance  $d = 1.5 \cdot 10^{13} \text{ cm}$

Energy flux density at Earth from Sun:  $I = 0.136 \frac{\text{J}}{\text{s} \cdot \text{cm}^2}$

$$\begin{aligned} \text{a) } P &= I \cdot 4\pi d^2 = 0.136 \frac{\text{J}}{\text{s} \cdot \text{cm}^2} \cdot 4\pi \cdot (1.5 \cdot 10^{13} \text{ cm})^2 \\ &= \underline{\underline{3.8 \cdot 10^{26} \frac{\text{J}}{\text{s}}}} \end{aligned}$$

b) The energy flux density at the surface of the Sun is

$$\frac{P}{4\pi R_{\odot}^2} = \sigma_B T_{\odot}^4 \quad \odot = \text{Sun}$$

$$T_{\odot} = \left( \frac{P}{4\pi R_{\odot}^2 \sigma_B} \right)^{\frac{1}{4}} = \left( \frac{3.8 \cdot 10^{26} \frac{\text{J}}{\text{s}}}{4\pi \cdot (7 \cdot 10^{10} \text{ cm})^2 \cdot 5.67 \cdot 10^{-12} \frac{\text{J}}{\text{s} \cdot \text{cm}^2 \cdot \text{K}^4}} \right) = \underline{\underline{5700 \text{ K}}}$$

2. K + K. Ch 4, #5 Surface temperature of the Earth

The Earth receives energy from the Sun over an area of  $\pi R_E^2$ .  
The Earth emits black-body radiation over its total area  $4\pi R_E^2$ .

$$P_{\text{in}} = P_{\text{out}}$$

$$I_{\text{solar}} \cdot \pi R_E^2 = \sigma_B T_E^4 \cdot 4\pi R_E^2$$

I will use  $I_{\text{solar}}$  from the previous problem. But you can derive it from the given numbers.

$$T_E = \left( \frac{I_{\text{solar}}}{4\sigma_B} \right)^{\frac{1}{4}} = \left( \frac{0.136 \frac{\text{J}}{\text{s} \cdot \text{cm}^2}}{4 \cdot 5.67 \cdot 10^{-12} \frac{\text{J}}{\text{s} \cdot \text{cm}^2 \cdot \text{K}^4}} \right)^{\frac{1}{4}} = \underline{\underline{280 \text{ K}}}$$

3. K. + K. Ch. 4 # 8

$$\text{region 1: } J_u^1 = \sigma_B (T_u^4 - T_m^4)$$

$$\text{region 2: } J_u^2 = \sigma_B (T_m^4 - T_l^4)$$

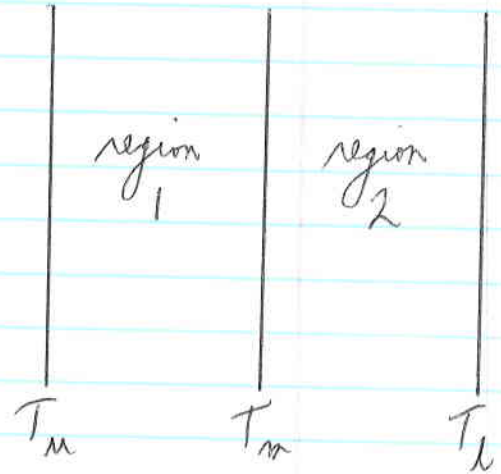
Under steady-state conditions,

$$J_u^1 = J_u^2$$

$$T_u^4 - T_m^4 = T_m^4 - T_l^4$$

$$2T_m^4 = T_u^4 + T_l^4$$

$$T_m = \left( \frac{T_u^4 + T_l^4}{2} \right)^{1/4}$$



$$J_u^1 = \sigma_B \left[ T_u^4 - \left( \frac{T_u^4 + T_l^4}{2} \right) \right] = \frac{1}{2} \sigma_B (T_u^4 - T_l^4)$$

This is half the energy flux we would get if the middle plane were not there.

4. K. + K. Ch 4, #1

Number of thermal photons  
in a cavity of volume  $V = L^3$ .

Method 1, following pages 93 + 94 in K. + K.

$$N = \sum_n \langle S_n \rangle = 2 \cdot \frac{1}{8} \int_0^\infty 4\pi n^2 dn \langle S_n \rangle$$

$$= \pi \int_0^\infty \frac{n^2 dn}{e^{h\omega_n/\tau} - 1} \quad \text{where } \omega_n = \frac{n\pi c}{L}$$

$$\text{Let } x = \frac{h\pi c}{L\tau} n \quad dx = \frac{h\pi c}{L\tau} dn$$

$$N = \pi \left( \frac{L\tau}{h\pi c} \right)^3 \int_0^\infty \frac{x^2 dx}{e^x - 1}$$

Mathematica says the integral  
is  $2 \zeta(3) = 2.404$

$$N = 2.404 \frac{V}{\pi^2} \left( \frac{\tau}{hc} \right)^3$$

Method 2, following my lecture on 2/2/18:

$$N = \sum_n \langle S_n \rangle = 2 \cdot \frac{V}{(2\pi)^3} \cdot \int_0^\infty 4\pi k^2 dk \cdot \frac{1}{e^{hck/\tau} - 1}$$

$$\text{let } x = \frac{hck}{\tau} \quad dx = \frac{hc}{\tau} dk$$

$$N = \frac{V}{\pi^2} \left( \frac{\tau}{hc} \right)^3 \int_0^\infty \frac{x^2 dx}{e^x - 1} = 2.404 \frac{V}{\pi^2} \left( \frac{\tau}{hc} \right)^3$$

5. Maxwell Relation  $F = U - \tau \sigma$   $N$  is fixed, so  $dN = 0$

$$dF = dU - d(\tau \sigma) = (\cancel{\tau d\sigma} - p dV) - (\cancel{\tau d\sigma} + \sigma d\tau)$$

$$dF = -\sigma d\tau - p dV \quad (1)$$

We also can write  $dF = \left(\frac{\partial F}{\partial \tau}\right)_V d\tau + \left(\frac{\partial F}{\partial V}\right)_\tau dV \quad (2)$

Compare (1) + (2):  $\sigma = -\left(\frac{\partial F}{\partial \tau}\right)_V \quad p = -\left(\frac{\partial F}{\partial V}\right)_\tau$

Crossed 2<sup>nd</sup> derivatives:  $\frac{\partial^2 F}{\partial V \partial \tau} = \frac{\partial^2 F}{\partial \tau \partial V}$

more explicitly:  $\left(\frac{\partial}{\partial V} \left(\frac{\partial F}{\partial \tau}\right)_V\right)_\tau = \left(\frac{\partial}{\partial \tau} \left(\frac{\partial F}{\partial V}\right)_\tau\right)_V$   
 $-\left(\frac{\partial \sigma}{\partial V}\right)_\tau = -\left(\frac{\partial p}{\partial \tau}\right)_V$

$$\underline{\underline{\left(\frac{\partial \sigma}{\partial V}\right)_\tau = \left(\frac{\partial p}{\partial \tau}\right)_V}}$$



$$6. \quad \langle s \rangle = \frac{1}{e^{\frac{k\omega}{\tau}} - 1} = \frac{1}{e^x - 1} \quad \text{where } x \equiv \frac{k\omega}{\tau}$$

a) High temperature expansion:  $x \ll 1$

$$e^x - 1 \approx x + \frac{x^2}{2} + \frac{x^3}{6} + \dots = x \left( 1 + \frac{x}{2} + \frac{x^2}{6} + \dots \right)$$

$$\begin{aligned} \langle s \rangle &= \frac{1}{x} \frac{1}{1 + \left( \frac{x}{2} + \frac{x^2}{6} \right)} = \frac{1}{x} \left[ 1 - \left( \frac{x}{2} + \frac{x^2}{6} \right) + \left( \frac{x}{2} + \frac{x^2}{6} \right)^2 - \dots \right] \\ &= \frac{1}{x} \left[ 1 - \frac{x}{2} - \frac{x^2}{6} + \frac{x^2}{4} + O(x^3) \right] \\ &= \frac{1}{x} \left[ 1 - \frac{x}{2} + \frac{x^2}{12} + O(x^3) \right] \end{aligned}$$

$$\langle s \rangle = \frac{\tau}{k\omega} \left[ 1 - \frac{1}{2} \frac{k\omega}{\tau} + \frac{1}{12} \left( \frac{k\omega}{\tau} \right)^2 + O\left( \frac{k\omega}{\tau} \right)^3 \right]$$

$$\langle s \rangle_{\text{asymptotic}} = \frac{\tau}{k\omega} \quad a = -\frac{1}{2} \quad b = \frac{1}{12}$$

$$\begin{aligned} b) \quad U &= \left( \langle s \rangle + \frac{1}{2} \right) k\omega = \langle s \rangle k\omega + \frac{1}{2} k\omega \\ &= \tau \left[ 1 - \frac{1}{2} \frac{k\omega}{\tau} + \frac{1}{12} \left( \frac{k\omega}{\tau} \right)^2 \right] + \frac{1}{2} k\omega \\ &= \tau - \cancel{\frac{1}{2} k\omega} + \frac{1}{12} \frac{(k\omega)^2}{\tau} + \dots + \cancel{\frac{1}{2} k\omega} \end{aligned}$$

$$U = \tau \left[ 1 + \frac{1}{12} \left( \frac{k\omega}{\tau} \right)^2 + \dots \right]$$