

# Fractional Marcus-Hush-Chidsey-Yakopcic current-voltage model for redox-based resistive memory devices

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April 14, 2023

We propose a circuit-level model combining the Marcus-Hush-Chidsey electron current equation and the Yakopcic equation for the state variable for describing resistive switching memory devices of the structure metal–ionic conductor–metal. We extend the dynamics of the state variable originally described by a first-order time derivative by introducing a fractional derivative with an arbitrary order between zero and one. We show that the extended model fits with great fidelity the current-voltage characteristic data obtained on a Si electrochemical metallization memory device with Ag-Cu alloy

## 1 Introduction

Substantial research efforts have been dedicated to the development of electrically-controlled resistive switching in metal-insulator-metal (MIM) devices or memristors, going from new materials discovery<sup>1–7</sup> to modelling and simulation<sup>8–10</sup>, and design and applications<sup>3,11–13</sup>. With both memory and logic capabilities combined at the hardware level, in addition to long retention times<sup>11</sup> and high switching rates<sup>14</sup> at relatively low energy consumption<sup>1,15</sup>, these devices are favorably seen as the next-generation building blocks for nonvolatile memories and neuromorphic computing applications<sup>11,12</sup>. In a typical memristor, the resistive switching is based on the electrically-stimulated change of cell resistance usually driven by internal ion redistribution, which actually depends not only on the applied excitation but also on the past history of the excitation<sup>6</sup>. Physical mechanisms associated with these reversible transitions have been attributed to different effects including valence change<sup>16</sup>, electrochemical metallization<sup>17</sup>, and phase change effects<sup>18</sup>. They can be either abrupt (binary) or gradual (analogue), and evolve at different timescales, leading to rich and complex device behaviors in this seemingly simple device structure of just three layers<sup>19</sup>. Furthermore, with the wide range of diversity in memristors materials and their morphologies, operating mechanisms, and manufacturing technologies there is an urgent need for the development of a general model capable of capturing accurately and effectively their complex nonlinear dynamics. This is crucial not only for the characterization and comparison between different memristor devices, but also for the investigation of larger scale memristor-based circuits and hybrid hardware architectures, and also to explore similar behaviors observed for instance in biological synapse systems<sup>20</sup>. While models at different size scales and thus with different degrees of physical details and computational complexity have been developed for memristors, including but not limited to ab initio<sup>21</sup>, kinetic Monte Carlo, and finite element method models<sup>22</sup>, in this work we focus on the circuit-level (compact) current-voltage behavior of memristors. From this point of view, memristors are generally described by

$$i = G(v, x)v, \quad (1)$$

$$\dot{x} = f(x, v), \quad (2)$$

where  $i = i(t)$  is the current through the device,  $v = v(t)$  is the applied voltage, and  $x = x(t)$  corresponds to a state variable or a group of state variables that quantify the internal dynamics of the device. These are, for example, width of doping region, concentration of vacancies in the gap region, and tunneling barrier width<sup>8</sup>. State variables can not be observed from external electrical behavior<sup>24</sup>. Eq. (1) follows the  $i$ - $v$  curve of the resistive device in question with  $G(v, x)$  being the generalized conductance, whereas Eq. (2) describes the dynamics of its internal state  $x$  based on its prehistory<sup>25</sup>. The actual state of a memristor can only be determined by solving Eqs. (1) and (2) self-consistently. Memristive systems as featured in terms of Eqs. (1) and (2) are known to possess a pinched hysteresis loop at the origin in the  $i$ - $v$  plane in the response to any periodic voltage source<sup>26</sup>. Being versatile and modular enough it is the Yakopcic model<sup>27–29</sup> which is most often used to simulate the nonlinear  $i$ - $v$  characteristic of wide range of memristors in response to sinusoidal and repetitive sweeping inputs. The model takes into account electron transmission effects, voltage threshold for state variable motion, and nonlinear velocity function for oxygen vacancies or dopant drift, considered to be the most relevant internal state information<sup>29</sup>. It follows on the steps of Strukov et al. work<sup>30</sup>, and describes the memristor as two resistors in series characterized by electron transmission equations so that<sup>29</sup>

$$I(t) = h_1(v)x + h_2(v)(1 - x). \quad (3)$$

Here,  $h_1$  is used to model the behavior in the low-resistance state of the device, and  $h_2$  captures its behavior<sup>Xiv:2302.09407v1 [cond-mat.mtrl-sci]</sup>

18 Feb 2023 2  $i$  or in the high-resistance state. The two electron transmission equations are weighted and mixed by the state variable  $x$  which is set to take values between zero and one<sup>25</sup>. In memristive devices, it is the

rate of change of the state variable  $x$  that is explicitly determined (2), and is given in the Yakopcic memristor model by the product of the two composite functions  $g(v)$  and  $f(x)$  such that<sup>29</sup>: ]