

Lecture 11:

Variational Quantum Algorithms and Quantum Chemistry

COMP3366

Quantum algorithms & computing architecture

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Objectives:

- **[O1] Concepts:** VQA, VQE, Ansatz, cost functions, quantum chemistry
- **[O3] Algorithm design:** Implementing VQE with Qiskit.

Previously in COMP3366:
“QAQA”

Max-Cut problem

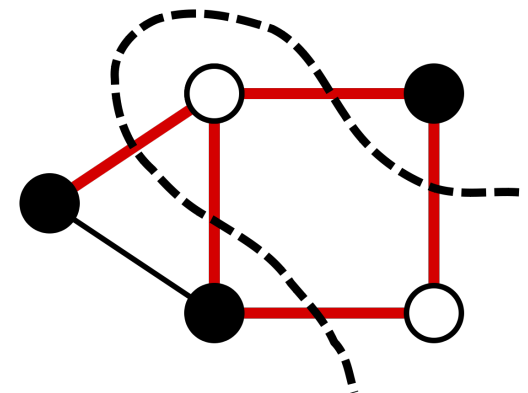
- Max-Cut:

For a graph $G = (V, E)$, split the set of vertices V of a graph into 2 disjoint subsets, such that the number of edges connecting the 2 subsets is maximized.

- Maximum cut as an optimization problem:

$$\max_s \frac{1}{2} \sum_{(i,j) \in E} (1 - s_i s_j) \quad s_i \in \{1, -1\}$$

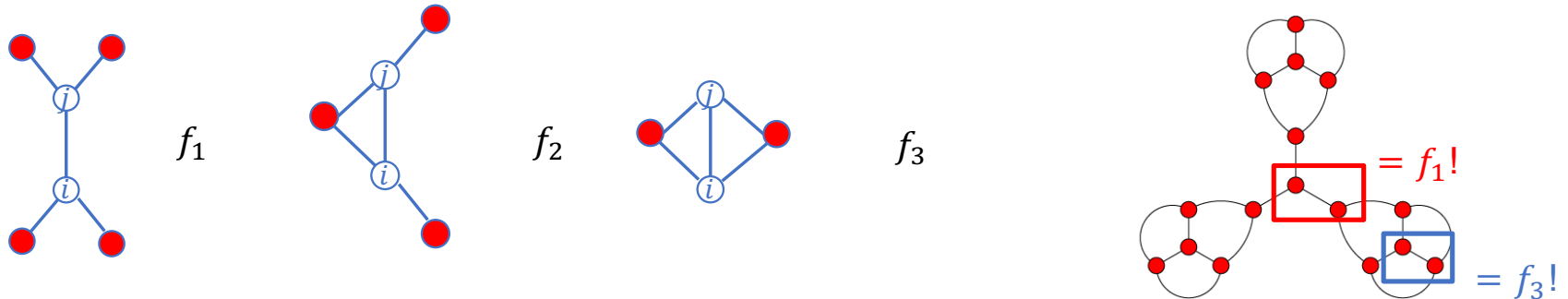
- Max-Cut is a prototypical **NP-complete problem**:
Easy to verify a solution, but extremely hard to solve!



5 edges in the maximum cut

Quantum Approximate Optimization Algorithm

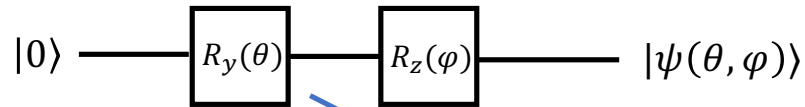
- QAOA is a **classical-quantum** variational algorithm for getting approximate solutions to the Max-Cut problem.
- For any 3-regular graph, a 6-qubit quantum computer is sufficient for executing a QAOA with depth $p = 1$ that yields a 0.6924-approximate solution. **The depth of quantum circuits depends only on p and not on n !**



- QAOA may be extended to other optimization methods and be applied, for instance, to finance and engineering.
It is a good proposal for demonstrating the power of **near-term** quantum computers.

Part I: Variational Quantum Algorithms

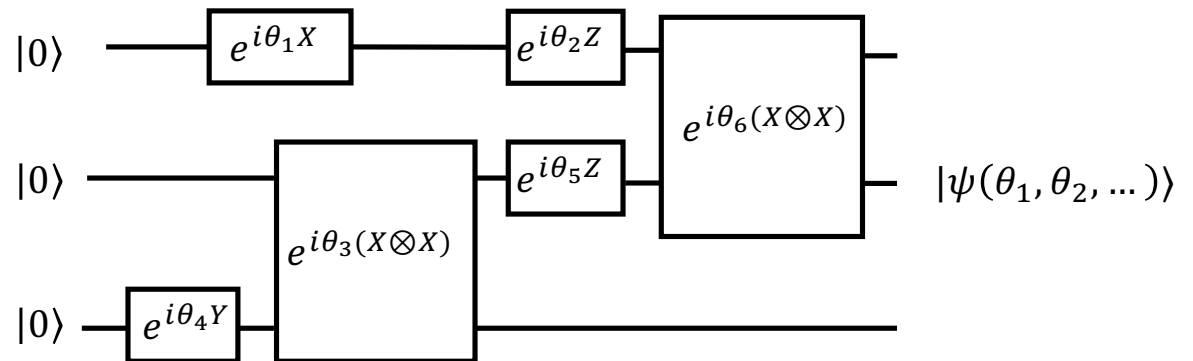
Variational states & circuits



- A single qubit state, generated by a variational circuit:

Gate types are fixed (X/Y/Z rotation); parameters vary

- A general state: $|\psi(\theta)\rangle = U(\theta)|0\rangle^{\otimes n}$ where $\theta = (\theta_1, \theta_2, \dots)$ is a vector of variable parameters



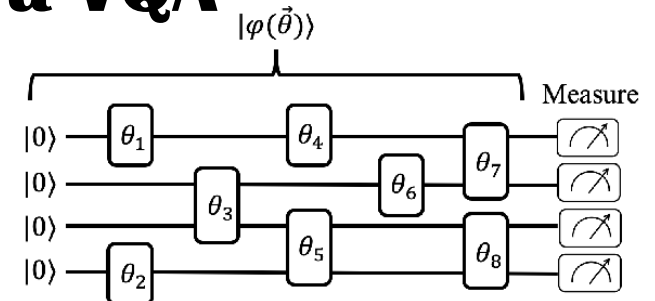
What is a VQA

- A **VQA (Variational Quantum Algorithm)** is an algorithm where a variational quantum circuit is trained step-by-step to minimize a cost function $\mathcal{C}(\theta)$.
- VQAs are **classical-quantum hybrid algorithms**, where classical and quantum computers join force to tackle a hard problem.
- In the near-term future, the classical-quantum hybrid is the most promising structure of quantum computing!



Standard procedure of a VQA

1. Select an **ansatz** of the circuit and initiate it in a random configuration $\vec{\theta}_0$
2. Run the circuit in the current configuration $\vec{\theta}$.
3. Evaluate the **cost function** $C(\vec{\theta})$ with the output state of the circuit.
4. If $C(\vec{\theta}) \leq C_{\text{term}}$, terminate and output $\vec{\theta}$; otherwise, go to Step 5.
5. Repeat Steps 2-3 for multiple times to get the values of C in a neighborhood of $\vec{\theta}$.
Compute the gradients $\frac{\partial C}{\partial \theta_1}, \frac{\partial C}{\partial \theta_2}, \dots$
6. Update $\vec{\theta} \rightarrow \vec{\theta}'$ using **gradient descent**. Go to Step 2.

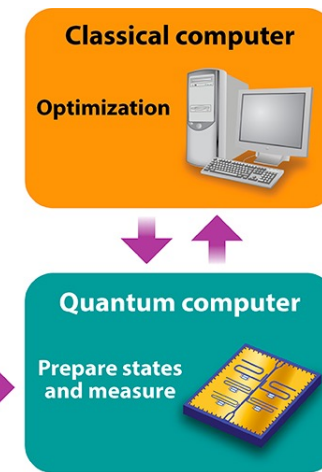


Quantum

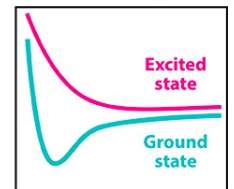


Encode

Classical



Molecular spectra



(Fig. aps.org)

Example: QAOA (Lecture 10)

- QAOA is a VQA.

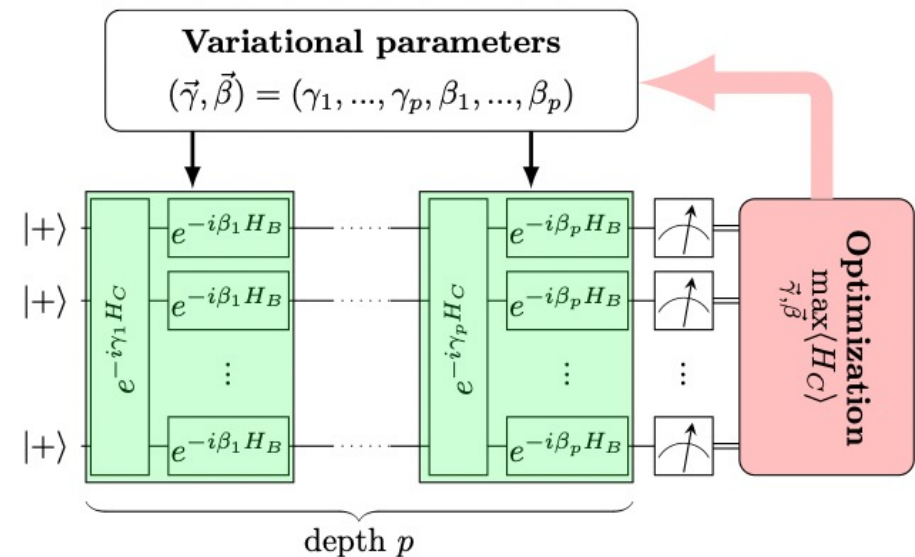
- Ansatz:

$$|\psi(\vec{\beta}, \vec{\gamma})\rangle = e^{-i\beta_p H_M} e^{-i\gamma_p H_C} \dots e^{-i\gamma_1 H_M} e^{-i\beta_1 H_C} |+\rangle^{\otimes n}$$

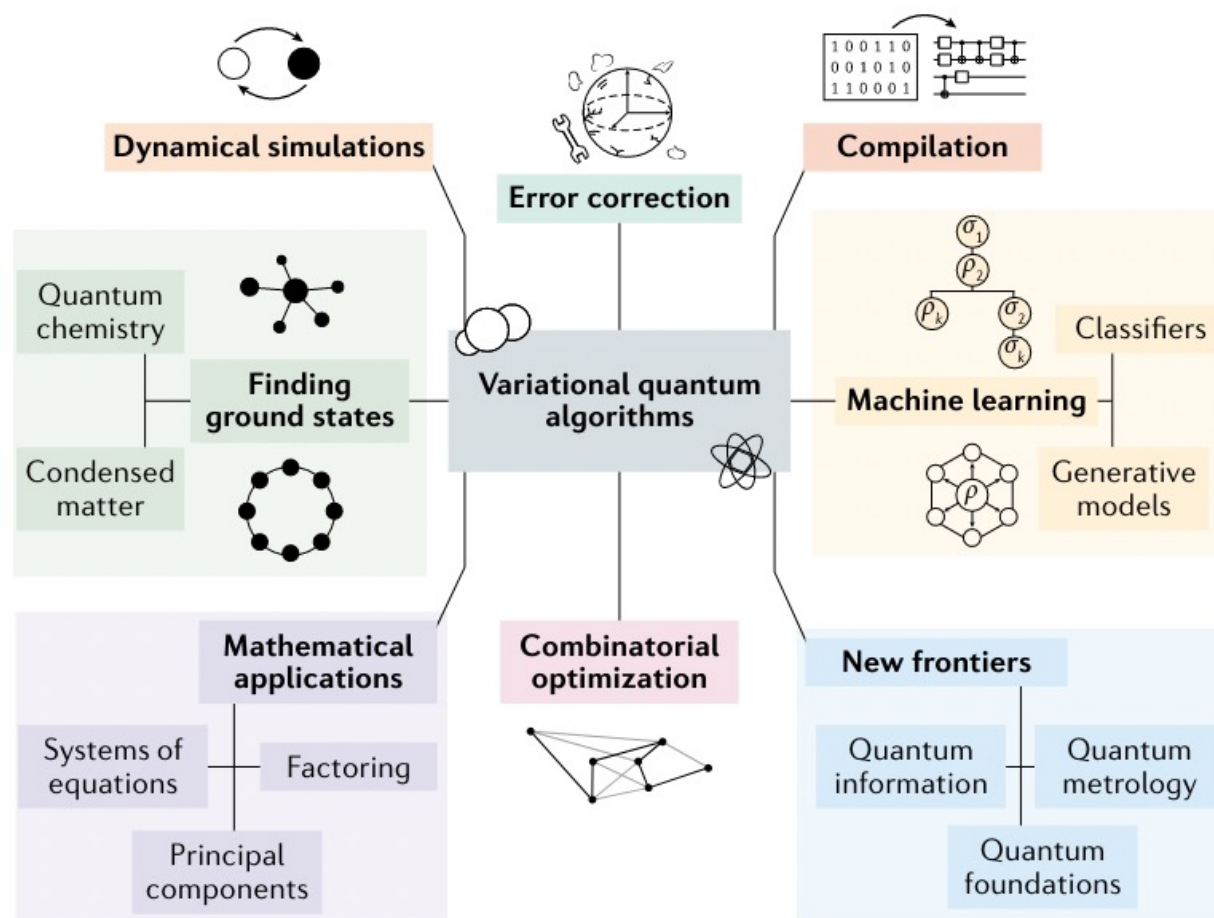
- Cost function:

$$C(\vec{\beta}, \vec{\gamma}) = \langle \psi(\vec{\beta}, \vec{\gamma}) | H_C | \psi(\vec{\beta}, \vec{\gamma}) \rangle$$

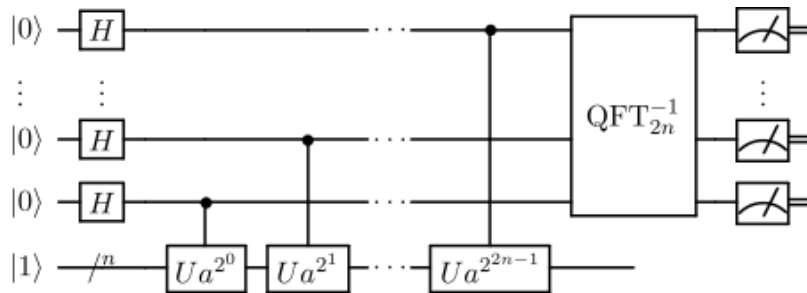
where H_C is the Hamiltonian that encodes a Max-Cut problem.



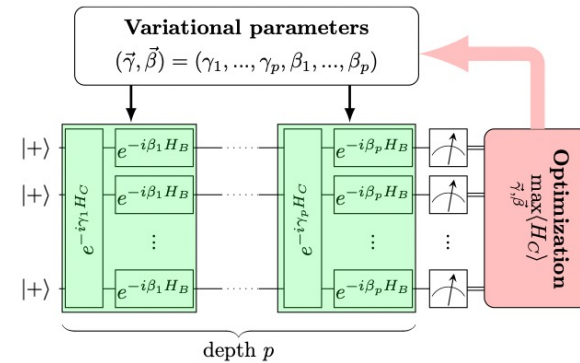
Applications of VQA



VQA vs Fully quantum algorithms



- “One-shot”.
- Fully coherent.
(No conversion of quantum-classical information until the end.)
- Large scale and deep circuits.
- Usually no ancillary variables (except for the input).



- Iterative.
- Low depth.
Circuit depth is brought down by introducing multiple iterations.
- Use classical computers as assistants.
- Involve ancillary variables.

*There are 2 crucial ingredients of a VQA:
The **cost function**, and the **ansatz**!*



Cost function

- A VQA is trained via **step-by-step reducing a cost function $C(\theta)$** .
- **The cost function** of a VQA is a function that maps the trainable parameters θ to real numbers:

$$C(\theta) = f(\{|\psi_k\rangle\}, \{O_k\}, U(\theta)) \in \mathbb{R}$$

- Here $\{|\psi_k\rangle\}$ are input states from a training set and $\{O_k\}$ is a set of observables.
- Example:

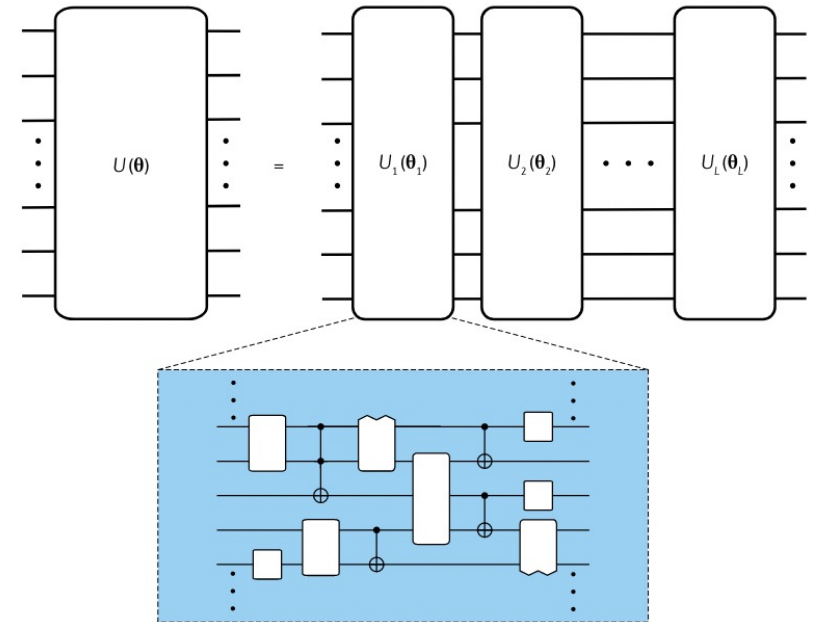
$$C(\theta) = \sum_k f_k \langle \psi_k | U(\theta)^\dagger O_k U(\theta) | \psi_k \rangle$$

Criteria for a good cost

- Usually, a cost function $C(\theta)$ should satisfy the following criteria:
 - 1. Trainability:**
It should be possible to efficiently optimize the parameters θ .
 - 2. Faithfulness:**
The minimum of $C(\theta)$ should correspond to the solution of the problem.
 - 3. Measurability:**
It should be efficiently estimable by performing measurements on a quantum computer and post-processing the classical data.

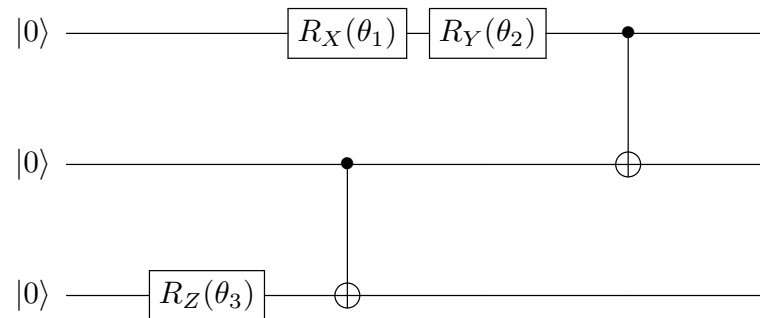
Ansatzes

- Usually, a variational circuit consists of a sequence of **circuit blocks**. Each block admits the same structure as specified by the ansatz.
- The choice of ansatzes is very important and may affect the accuracy of the solution as well as the efficiency of convergence.



Why does the choice of ansatz matter?

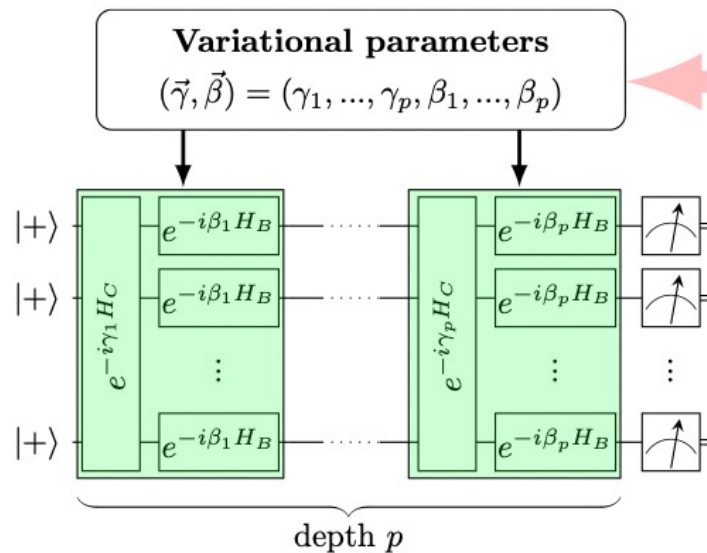
Is this a good ansatz?



- No. For example, the last two qubits are wasted!
- Then, is there any trick or common-sense when choosing the ansatz?
- Two common types of good ansatz: **problem-inspired** and **hardware-inspired**.

Problem-inspired ansatz

- Why “problem”-inspired:
The ansatz is designed based on information about the problem.
- Example: the quantum alternating operator ansatz in QAOA is a problem-inspired ansatz. It is based on the idea of simulated adiabatic quantum computing:

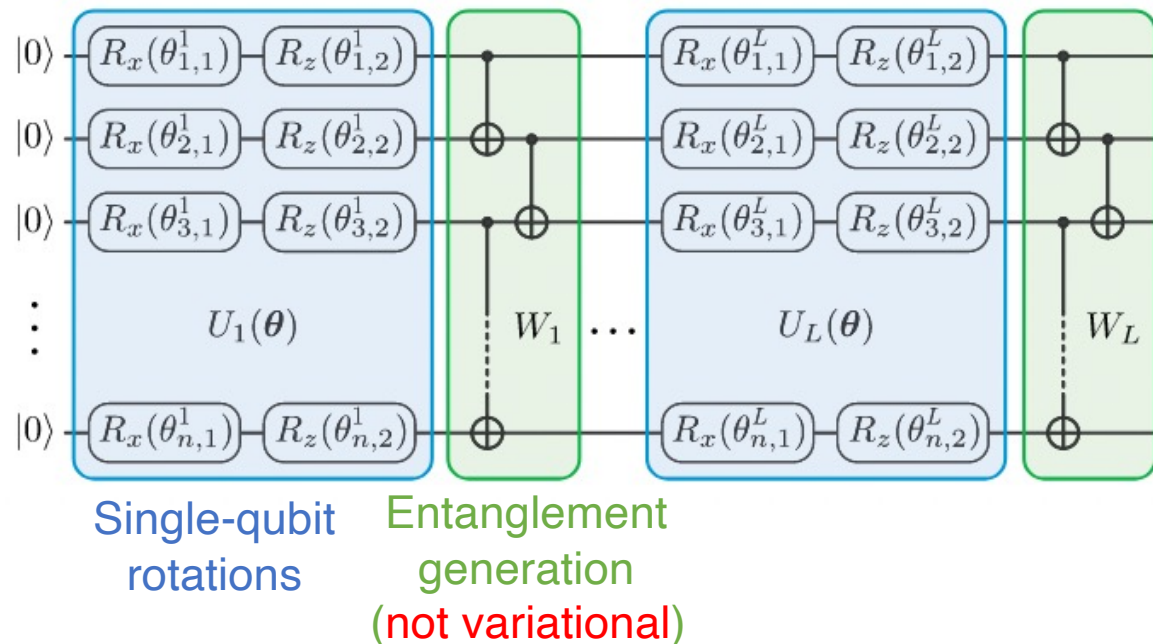


Hardware-efficient ansatz

- Problem-inspired ansatzes ensure **fast convergence to a satisfying solution**, but they could lead to a deep circuit or hard-to-implement gate sets.
- Besides, it might be hard to find such an inspiration for some problems.
- Alternatively, we can consider using hardware-efficient ansatzes.
- The choice of ansatz involves a trade-off:
 - Good and fast convergence to the solution → problem-inspired ansatzes
 - More implementable & “bullet-proof” → hardware-efficient ansatzes
- In the lucky case, you may find an ansatz with both merits.
(There are fair and smart partners, but not many ...)

Hardware-efficient ansatz

- A prototype of a hardware-efficient ansatz:



- Single-qubit rotations and CNOT are among the universal gate set and have high fidelity in most implementations.

Part II:

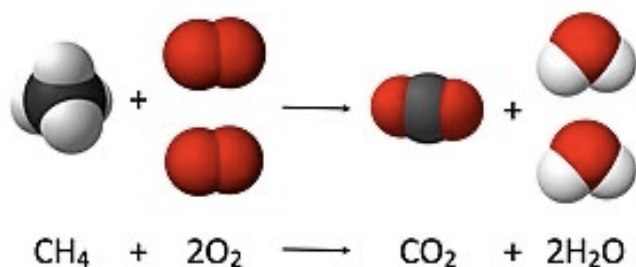
Variational quantum eigensolvers

Hamiltonians in quantum chemistry

- Every quantum system, e.g., a molecule or an atom, has an associated Hamiltonian H that **determines** its feature via Schrödinger's equation:

$$\frac{i\partial|\psi\rangle}{\partial t} = H|\psi\rangle.$$

- In quantum chemistry, many fundamental tasks (e.g., calculating the reaction rate) **boil down to finding the ground state** associated with the Hamiltonian of a molecule.



The ground state problem

- In general, H is a Hermitian operator that can be diagonalized as

$$H = \sum_i E_i |\psi_i\rangle\langle\psi_i|$$

with $E_0 \leq E_1 \leq \dots$. Here $\{|\psi_i\rangle\}$ are the energy eigenstates.

- Given a complex H , it is one of the most important questions in modern physics to find the ground state $|\psi_0\rangle$ as well as the (minimum) ground energy E_0 .
- **Goal:** We want to solve the ground state problem in quantum chemistry with a near-term quantum computer.
- **Question:** How to map “states of a molecule” to states of a quantum computer?

Hamiltonians in quantum chemistry

- Example: In quantum chemistry, we would like to know the ground state of a molecule, and the particles (electrons) are **fermions**.
- Hamiltonian: $H = \sum_j h_j a_j^\dagger a_j + \sum_{ij} h_{ij} a_j^\dagger a_i^\dagger a_i a_j$
- Here a_j, a_j^\dagger are the annihilation and creation operators of the j -th **orbital**.



- Fermions are indistinguishable whereas qubits are distinguishable. They obey **Pauli exclusion principle**:

$$\{a_i, a_j\} = 0, \{a_i^\dagger, a_j^\dagger\} = 0, \{a_i, a_j^\dagger\} = \delta_{i,j}$$

- To simulate Fermionic systems on a quantum computer, we need to construct **“simulated” annihilation and creation operators** satisfying the above relation.

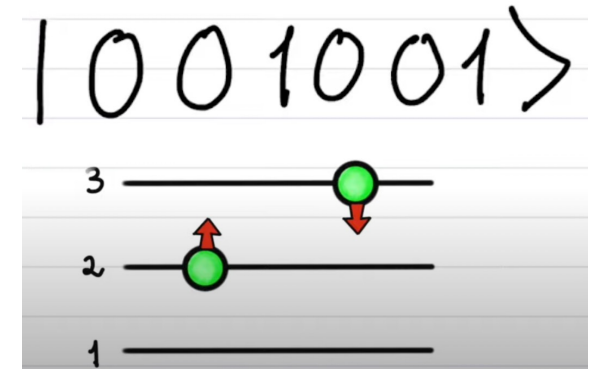
From Fermions to qubits

- Mapping **occupations** to qubit states:
Spin orbital \rightarrow qubits;
 $|0\rangle, |1\rangle \rightarrow$ unoccupied, occupied
- (Jordan-Wigner) Operators of qubit systems

$$\hat{a}_j^\dagger := \left(\prod_{m < j} Z_m \right) \sigma_{+,j}$$

where $\sigma_{+,j} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}_j = \frac{1}{2}(X_j + iY_j)$ obey Pauli exclusion principle: $\{\hat{a}_i, \hat{a}_j\} = 0, \{\hat{a}_i^\dagger, \hat{a}_j^\dagger\} = 0, \{\hat{a}_i, \hat{a}_j^\dagger\} = \delta_{i,j}$.

- Conclusion: We can map a fermionic Hamiltonian to a **Hamiltonian of multiple qubits** by replacing every a_j^\dagger (a_j) by \hat{a}_j^\dagger (\hat{a}_j).
- The number of **qubits** = the number of spin orbitals **> the number of fermions**.



2 electrons, 3 orbitals \rightarrow 6 qubits

Ground state via optimization

- Rayleigh-Ritz quotient formulation of the ground state energy:

$$E_0 = \min_{|\psi\rangle} \frac{\langle\psi|H|\psi\rangle}{\langle\psi|\psi\rangle}$$

- We can determine the ground state by solving the above optimization with H being the system's Hamiltonian.
- Idea:
 - Use a variational circuit to produce trial states $|\psi(\theta)\rangle$ with θ being variational parameter(s).
 - Measure the Hamiltonian for $|\psi(\theta)\rangle$ as the cost function $C_{VQE}(\theta) = \langle\psi(\theta)|H|\psi(\theta)\rangle$.
 - Update θ to make the cost function smaller.


Challenges

- We need to verify the following for $C_{VQE}(\theta) = \langle \psi(\theta) | H | \psi(\theta) \rangle$:

1. **Trainability:**

It should be possible to efficiently optimize the parameters θ .

2. **Faithfulness:**

If $|\psi_0\rangle$ is contained in the family $\{|\psi(\theta)\rangle\}_\theta$, then the faithfulness follows immediately from the definition. 

3. **Measurability:**

This is a highly non-trivial problem! 

It may not be possible to efficiently measure a generic H .

Measurement of energy

- Consider Hamiltonians of the following form:

$$H = \sum_{i,\alpha} h_{\alpha}^i \sigma_{\alpha}^i + \sum_{i,j,\alpha,\beta} h_{\alpha\beta}^{ij} \sigma_{\alpha}^i \otimes \sigma_{\beta}^j + \dots$$

where σ_{α}^i denotes a Pauli operator σ_{α} ($\sigma_{0,1,2,3} = I, X, Y, Z$) on the i th qubit and $h_{\alpha}^i \dots$ are coefficients.

- By linearity, we have $\langle H \rangle = \sum_{i,\alpha} h_{\alpha}^i \langle \sigma_{\alpha}^i \rangle + \sum_{i,j,\alpha,\beta} h_{\alpha\beta}^{ij} \langle \sigma_{\alpha}^i \otimes \sigma_{\beta}^j \rangle + \dots$
- Pauli operators can be efficiently estimated (by measurement in the bases $\{|0/1\rangle\}, \{|\pm\rangle\}, \{|\pm i\rangle\}$) \Rightarrow each $\langle \sigma_{\alpha}^i \rangle$ can be efficiently measured.

Energy measurability

- Any Hamiltonian can be expressed as

$$H = \sum_{i=1}^m c_i H_i$$

with each H_i being tensor product of Paulis (i.e., a Pauli string) and c_i being real coefficients.

- To measure $\langle H \rangle$, we can measure each $\langle H_i \rangle$ separately and evaluate $\langle H \rangle = \sum_i c_i \langle H_i \rangle$ on a classical computer.
- Complexity grows as $O(m)$: the state preparation must be repeated for several times to measure $\langle H_i \rangle$, for $i = 1, \dots, m$.
- Conclusion: $\langle H \rangle$ can be measured efficiently if H can be written as the sum of “not too many” Pauli strings.

Procedure of a VQE

- **VQE** (i.e., the **V**ariational **Q**uantum **E**igensolver) is a VQA that is tailor-made for the ground state problem.
- Input: $H = \sum_{i=1}^m c_i H_i$, each H_i being tensor product of single-qubit Paulis. Output: ground state of H and the ground energy (up to a small error).
- VQE routine:
 1. **For $i = 1, 2, \dots, m$, do:**
State preparation:
Prepare a parametrized quantum state $|\psi(\theta)\rangle$ via a variational circuit $U(\theta)|0\rangle^{\otimes n}$.
Energy estimation: Measure H_i for $|\psi(\theta)\rangle$.
 2. **Classical feedback:**
Evaluate $\langle H \rangle = \sum_i c_i \langle H_i \rangle$ and update θ using a classical optimization method.
 3. **Repeat** the above procedure until a convergence criteria is satisfied.

Example: finding the ground state of H_2

- Step 1: Setting up the equivalent Hamiltonian
- Configuration of the molecule (atom types, nuclei distance ...)

```
from pennylane import numpy as np

symbols = ["H", "H"]
coordinates = np.array([0.0, 0.0, -0.6614, 0.0, 0.0, 0.6614])
```

- Set the number of qubits (= 4)

A built-in function does the Jordan-Wigner transform for us, outputting a 4-qubit Hamiltonian that is equivalent to the original Fermionic Hamiltonian.

```
import pennylane as qml

H, qubits = qml.qchem.molecular_hamiltonian(symbols, coordinates)
print("Number of qubits = ", qubits)
print("The Hamiltonian is ", H)
```

Out:

```
Number of qubits = 4
The Hamiltonian is      (-0.2427450126096431) [Z2]
+ (-0.2427450126096431) [Z3]
+ (-0.04207255194723672) [I0]
+ (0.1777135822907777) [Z1]
+ (0.17771358229077772) [Z0]
+ (0.12293330449306411) [Z0 Z2]
+ (0.12293330449306411) [Z1 Z3]
+ (0.16768338855610976) [Z0 Z3]
+ (0.16768338855610976) [Z1 Z2]
+ (0.17059759276839878) [Z0 Z1]
+ (0.17627661394198885) [Z2 Z3]
+ (-0.04475008406304564) [Y0 Y1 X2 X3]
+ (-0.04475008406304564) [X0 X1 Y2 Y3]
+ (0.04475008406304564) [Y0 X1 X2 Y3]
+ (0.04475008406304564) [X0 Y1 Y2 X3]
```

- Step 2: Initialization of the variational circuit.
- We want to prepare a simple variational state of the form

$$|\Psi(\theta)\rangle = \cos\frac{\theta}{2} |1100\rangle - \sin\frac{\theta}{2} |0011\rangle$$

- There is only one variable θ .
- This is realized by a simple circuit:
- Set $\theta = 0$ as the initial configuration.

```
dev = qml.device('default.qubit', wires=4)

@qml.qnode(dev)
def circuit(phi):
    qml.PauliX(wires=0)
    qml.PauliX(wires=1)
    qml.DoubleExcitation(phi, wires=[0, 1, 2, 3])
    return qml.state()

circuit(0.1)
```


- Step 3: Update θ to minimize the energy.

```
# store the values of the cost function
energy = [cost_fn(theta)]

# store the values of the circuit parameter
angle = [theta]

max_iterations = 100
conv_tol = 1e-06

for n in range(max_iterations):
    theta, prev_energy = opt.step_and_cost(cost_fn, theta)

    energy.append(cost_fn(theta))
    angle.append(theta)

    conv = np.abs(energy[-1] - prev_energy)

    if n % 2 == 0:
        print(f"Step = {n}, Energy = {energy[-1]:.8f} Ha")

    if conv <= conv_tol:
        break

print("\n" f"Final value of the ground-state energy = {energy[-1]:.8f} Ha")
print("\n" f"Optimal value of the circuit parameter = {angle[-1]:.4f}")
```

```
Step = 0, Energy = -1.12799983 Ha
Step = 2, Energy = -1.13466246 Ha
Step = 4, Energy = -1.13590595 Ha
Step = 6, Energy = -1.13613667 Ha
Step = 8, Energy = -1.13617944 Ha
Step = 10, Energy = -1.13618736 Ha
Step = 12, Energy = -1.13618883 Ha
```

We can see that the energy goes down step-by-step.

- Step 4: Output the ground state & energy.

```
Final value of the ground-state energy = -1.13618883 Ha
```

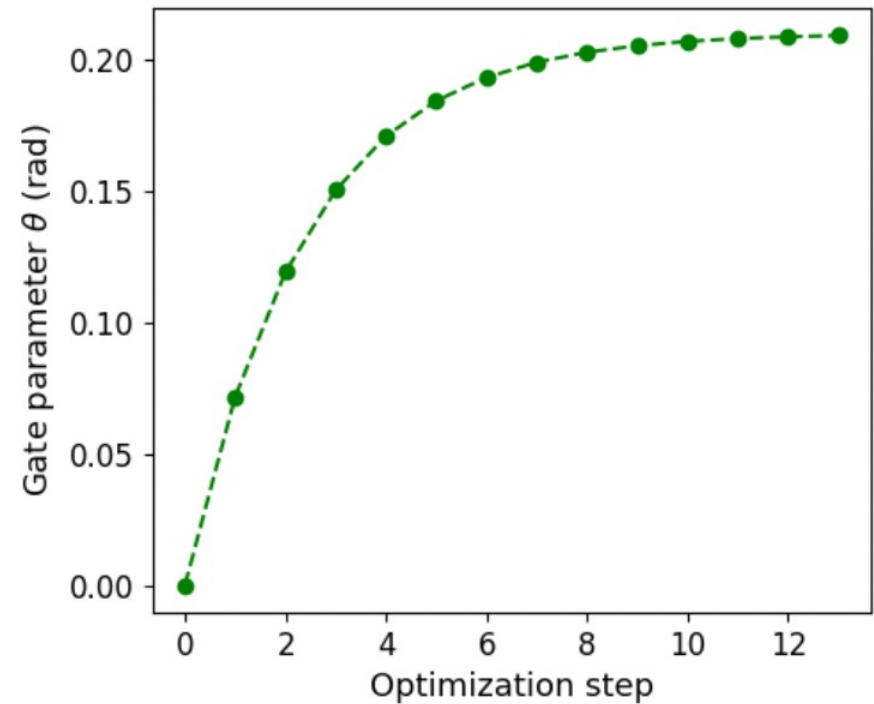
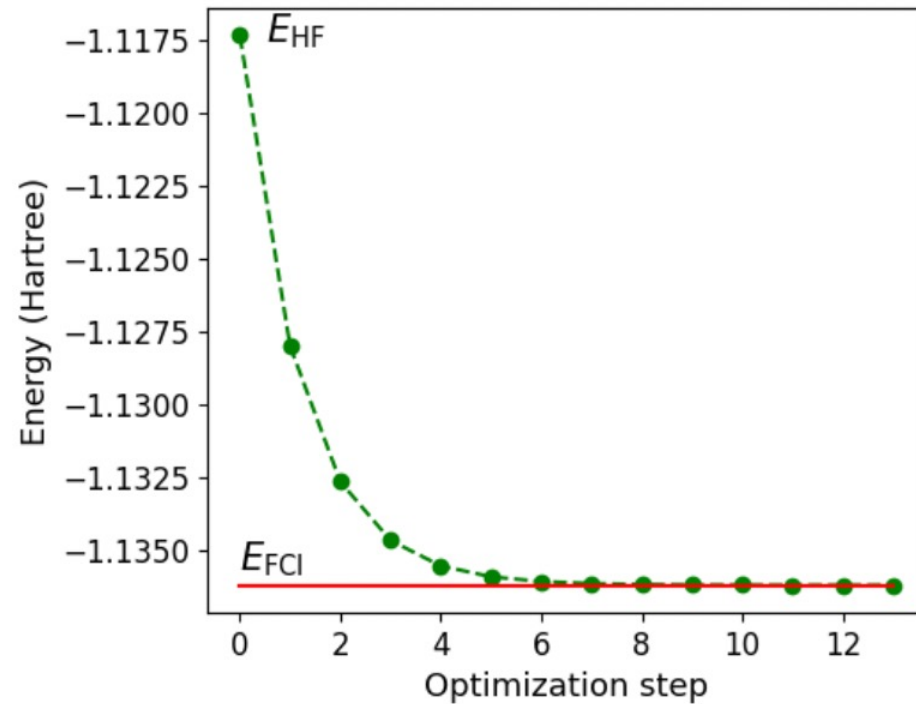
```
Optimal value of the circuit parameter = 0.2089
```

- The minimizer is identified

$$|\Psi(\theta^*)\rangle = 0.994 |1100\rangle - 0.104 |0011\rangle,$$

which can be mapped back to the ground state of the original (Fermionic) Hamiltonian of H_2 by reversing the Jordan-Wigner transform.

- Changes of E and θ during the iteration:



See https://pennylane.ai/qml/demos/tutorial_vqe.html for the complete code.

Summary

- VQA: cost function and ansatz.
- VQE: a variational solver for eigenstate problems.
- Application of VQE in quantum chemistry:
mapping spin orbitals to qubits, and energy measurement.

Homework

- **Review** the lecture slides; you may find the review questions in the next slides helpful.
- Read the Qiskit tutorial on VQE
<https://qiskit.org/textbook/ch-applications/vqe-molecules.html>.

Review questions

- What are the key ingredients of a VQA?
- What are the criteria for a good cost function?
- Give an example of a problem-inspired/hardware-efficient ansatz.
- When simulating a fermionic system, we often require more qubits than the number of fermions. Why?
- For the toy example of H_2 ground state, is the ansatz problem-inspired or hardware-efficient?