

# Worksheet #8

PHYS 4C: Waves and Thermodynamics

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## 1 Problem 1

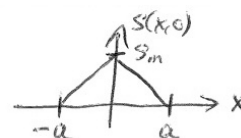
The speed of sound in steel is 5941 m/s. Steel has a density of around 7900 kg/m<sup>3</sup> (depends somewhat on the alloy content).

- Based on this information, what is the bulk modulus of steel?
- A steel rod of cross-sectional area  $A$  is placed on the  $x$ -axis. The following sound pulse is sent through the steel rod in the  $+x$  direction:

$$f(x) = s(x, 0) = s_m(1 - |x|/a), \text{ if } |x| < a.$$

$$f(x) = s(x, 0) = 0, \text{ if } |x| \geq a.$$

Determine the total energy of this pulse.



Sound pulse graph

- Now suppose a sinusoidal sound wave with frequency 440 Hz is sent through steel with a sound level of 100 dB. Calculate  $s_m$  and  $\Delta p_m$  for this sound wave.

### 1.1 Solution (a)

The bulk modulus is used as part of an equation for the velocity.

$$v = \sqrt{\frac{B}{\rho}} \quad (1)$$

This can be solved for the bulk modulus.

$$B = \rho v^2 \quad (2)$$

We know all the values necessary, so we can solve this equation.

$$B = (7900 \text{ kg/m}^3)(5941 \text{ m/s})^2 = \boxed{2.788 \times 10^{11} \text{ kg/m} \cdot \text{s}^2} \quad (3)$$

## 1.2 Solution (b)

If this is the initial point, there is no motion, so  $\frac{ds}{dt} = 0$  at all points. There is an equation for the potential energy.

$$dU = \frac{1}{2}B \left( \frac{\partial s}{\partial x} \right)^2 A dx \quad (4)$$

At time  $t = 0$ , the partial derivative of  $s$  is calculatable, but it would have to be divided into two cases ( $x \geq 0$  and  $x < 0$ ). For the case of  $x \geq 0$ :

$$\frac{\partial s}{\partial x} = \frac{\partial}{\partial x} \left( s_m \left( 1 - \frac{x}{a} \right) \right) = -\frac{s_m}{a} \quad (5)$$

For the case of  $x < 0$ :

$$\frac{\partial s}{\partial x} = \frac{\partial}{\partial x} \left( s_m \left( 1 + \frac{x}{a} \right) \right) = \frac{s_m}{a} \quad (6)$$

Since the negative is the only thing differentiating these, it is safe to say that  $\left( \frac{\partial s}{\partial x} \right)^2$  is identical for both.

These in turn can be plugged into our equation for the potential energy.

$$dU = \frac{1}{2}B \frac{s_m^2}{a^2} A dx \quad (7)$$

This in turn can be integrated between  $-a$  and  $a$ . Technically, we would be integrating between  $-\infty$  and  $\infty$ , but since  $s(x, 0) = 0$  everywhere outside the range of  $(-a, a)$ , all the other spots would result in zero to begin with.

$$\int_{-\infty}^{\infty} dU = \int_{-\infty}^{\infty} \frac{1}{2}BA \frac{s_m^2}{a^2} dx \quad (8)$$

$$U = \int_{-a}^a \frac{1}{2}BA \frac{s_m^2}{a^2} dx = \frac{1}{2}BA \frac{s_m^2}{a^2} \int_{-a}^a dx \quad (9)$$

$$= \frac{1}{2}BA \frac{s_m^2}{a^2} [x]_{-a}^a = \frac{1}{2}BA \frac{s_m^2}{a^2} (a - (-a)) = \boxed{BA \frac{s_m^2}{a}} \quad (10)$$

### 1.3 Solution (c)

We know the sound level. This can be used to calculate the intensity.

$$\beta = 100 \text{ dB} = (10 \text{ dB}) \log_{10} \frac{I}{I_0} \quad (11)$$

$$10 = \log_{10} \frac{I}{I_0} \quad (12)$$

$$10^{10} = \frac{I}{10^{-12} \text{ W/m}^2} \quad (13)$$

$$10^{10} * 10^{-12} = 10^{-2} \text{ W/m}^2 = I \quad (14)$$

Next, we know the bulk modulus ( $B$ ), the speed of sound ( $v$ ), and the frequency ( $f$ ). The latter can be turned into the angular speed ( $\omega$ ) by arithmetic.

$$\omega = 2\pi f \quad (15)$$

We have an equation for the intensity that includes all of these values and the displacement amplitude, the latter of which can be solved for.

$$I = \frac{1}{2} \rho v \omega^2 s_m^2 \quad (16)$$

$$s_m = \sqrt{\frac{2I}{\rho v \omega^2}} = \sqrt{\frac{2I}{\rho v (2\pi f)^2}} = \sqrt{\frac{I}{2\rho v (\pi f)^2}} \quad (17)$$

$$= \sqrt{\frac{10^{-2} \text{ W/m}^2}{2(7900 \text{ kg/m}^3)(5941 \text{ m/s})(\pi * 440 \text{ Hz})^2}} \quad (18)$$

$$= \boxed{7.47 \times 10^{-9} \text{ m}} \quad (19)$$

The pressure amplitude is calculatable similarly.

$$\Delta p_m = (v \rho \omega) s_m \quad (20)$$

$$= (5941 \text{ m/s})(7900 \text{ kg/m}^3)(2\pi * 440 \text{ Hz})(7.47 \times 10^{-9} \text{ m}) \quad (21)$$

$$= \boxed{968.85 \text{ Pa}} \quad (22)$$

## 2 Problem 2

A half-open organ pipe is tuned to A(440) (i.e., the fundamental frequency is 440 Hz). Air has a density of  $1.21 \text{ kg/m}^3$  and a speed of sound of  $343 \text{ m/s}$  at  $20^\circ\text{C}$ .

- What is the length of the pipe?
- What is the maximum kinetic energy density (per unit volume) at the open end of the pipe if  $s_m = 2.0 \mu\text{m}$ ? At the closed end?
- If the ambient temperature were raised from  $20^\circ\text{C}$  to  $40^\circ\text{C}$ , what would be the new fundamental frequency of the pipe (ignore changes in the length of the pipe due to the temperature change)?

### 2.1 Solution (a)

This is calculatable from the formula for the  $n$ th frequency of a sound wave in what is in this case a half-open tube.

$$f = \frac{nv}{4L} \quad (23)$$

We should solve for  $L$ . Since it is the fundamental frequency, the value of  $n$  would be 1.

$$L = \frac{nv}{4f} = \frac{1 * 343 \text{ m/s}}{4 * 440 \text{ Hz}} = \boxed{0.195 \text{ m}} \quad (24)$$

### 2.2 Solution (b)

The lecture gives us an equation for the kinetic energy density. We have a known equation for the sound wave.

$$\rho_K = \frac{1}{2}\rho \left( \frac{\partial s}{\partial t} \right)^2 \quad (25)$$

$$s(x, t) = s_m \cos(kx - \omega t) \quad (26)$$

Our value for  $\omega$  would be  $2\pi$  times the frequency of the wave, which we do have. At this point, we can take the derivative of  $s(x, t)$  with respect to  $t$ .

$$\frac{\partial s}{\partial t} = -s_m \omega \sin(kx - \omega t) \quad (27)$$

Put that into the equation for the kinetic energy density. This will allow us to find the kinetic energy density.

$$\rho_K = \frac{1}{2}\rho s_m^2 \omega^2 \cos^2(kx - \omega t) \quad (28)$$

The kinetic energy density will be at its maximum when  $\cos(kx - \omega t) = 1$  and consequently  $\cos^2(kx - \omega t) = 1$ . We can substitute those in here. Not that this would be the maximum kinetic energy density at the open end, where there would be an antinode.

$$\rho_{K;\text{max};\text{open}} = \frac{1}{2}\rho s_m^2 \omega^2 \quad (29)$$

$$= \frac{1.21 \text{ kg/m}^3}{2} (2.0 \times 10^{-6} \text{ m})^2 (2\pi * 440 \text{ Hz})^2 \quad (30)$$

$$= \boxed{1.85 \times 10^{-5} \text{ J/m}^3} \quad (31)$$

On the other hand, at the closed end, there would be a node rather than an antinode. At a node, the kinetic energy is always zero, so at the closed end, the maximum kinetic energy density would be  $\boxed{0}$ .

## 2.3 Solution (c)

There is a formula for the speed of sound in air, equal to the square root of the adiabatic index times the universal gas constant times the temperature divided by the molar mass.

$$v = \sqrt{\frac{\gamma R T}{M}} \quad (32)$$

We can rearrange this to isolate the speed of sound and the temperature.

$$\frac{v}{\sqrt{T}} = \sqrt{\frac{\gamma R}{M}} \quad (33)$$

In this case, everything on the left side of the equation will be constant, so we can use a comparison of fractions. From there, we can solve for the speed of sound at 40°C relative to 20°C.

$$\frac{v_{20^\circ\text{C}}}{\sqrt{T_{20^\circ\text{C}}}} = \frac{v_{40^\circ\text{C}}}{\sqrt{T_{40^\circ\text{C}}}} \quad (34)$$

$$v_{40^\circ\text{C}} = v_{20^\circ\text{C}} \sqrt{\frac{T_{40^\circ\text{C}}}{T_{20^\circ\text{C}}}} \quad (35)$$

We can reuse our formula for the frequency in a half-open tube. This would still have the fundamental frequency ( $n = 1$ ). The temperatures we use should be in Kelvin.

$$f_{40^\circ\text{C}} = \frac{nv}{4L} = \frac{v_{20^\circ\text{C}}}{4 * \frac{v_{20^\circ\text{C}}}{4f_{20^\circ\text{C}}}} \cdot \sqrt{\frac{T_{40^\circ\text{C}}}{T_{20^\circ\text{C}}}} = f_{20^\circ\text{C}} \sqrt{\frac{T_{40^\circ\text{C}}}{T_{20^\circ\text{C}}}} \quad (36)$$

$$= 440 \text{ Hz} \cdot \sqrt{\frac{313.15\text{K}}{293.15\text{K}}} = \boxed{454.8 \text{ Hz}} \quad (37)$$