

PHYS 4 Exam 6 Cheat Sheet (with L^AT_EX)
Angular Kinematics

$$\theta = \frac{S}{r}; \omega = \frac{d\theta}{dt}; \alpha = \frac{d^2\theta}{dt^2}; (1) \omega(t) = \omega_0 + \alpha t$$

$$(2) \theta = \theta_0 + \omega_0 t + \frac{1}{2}\alpha t^2; (3) \omega^2 = \omega_0^2 + 2\alpha\Delta\theta$$

$$v_t = \omega r; a_t = \alpha r; a_c = \omega r^2; T = \frac{2\pi}{\omega} = \frac{1}{f}$$

Electric Fields and Forces

$$e = 1.602 \times 10^{-19} \text{C}; \epsilon_0 = 8.85 \times 10^{-12} \frac{\text{C}^2}{\text{Nm}^2}$$

$$k = 8.99 \times 10^9 \frac{\text{Nm}^2}{\text{C}^2} = \frac{1}{4\pi\epsilon_0}$$

$$\vec{F} = \frac{kq_1q_2}{r^2}\hat{r} = \frac{kq_1q_2}{r^3}\vec{r}; \vec{E} = \frac{kq}{r^2}\hat{r} = \frac{kq}{r^3}\vec{r}; F = qE$$

In a diagram, the direction of an electric field is represented by the direction of its arrows, while the strength of the field is represented by the proximity of the lines.

$$\lambda = \frac{Q}{r}; \sigma = \frac{Q}{A}; \rho = \frac{Q}{V}$$

$$E = \int dE = \int \frac{k dq}{r^3}\vec{r} = \int \frac{k\lambda}{r^3}\vec{r}dr$$

$$\vec{E}_{ring}(z) = \frac{kqz}{(z^2 + R^2)^{3/2}}\hat{k}$$

For a rod of length L, measured at a distance d from the close end from the rod of charge Q.

$$\vec{E}_{axis} = -\frac{kQ}{d(d-L)}\hat{i}$$

For a rod of length L, measured perpendicular to the rod at a distance d from the close end from the rod of charge Q.

$$\vec{E} = k\lambda \left[\frac{1}{z} - \frac{1}{L^2 + z^2} \right] \hat{i} + \frac{k\lambda L}{z\sqrt{L^2 + z^2}} \hat{j}$$

$$V = k\lambda \ln \left(\frac{L + \sqrt{L^2 + d^2}}{d} \right)$$

$$\vec{E}_{arc} = \frac{k\lambda}{r} \begin{pmatrix} 2\sin(\frac{\theta}{2}) \\ 0 \end{pmatrix}$$

$$\vec{E}_{disc} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{z^2 + R^2}} \right)$$

For a spherical shell of radius R.

$$\vec{E} = \begin{cases} 0 & \text{if } r < R \\ \frac{kq}{r^2}\hat{r} & \text{otherwise} \end{cases}$$

If $r < R$, $\Delta V = 0$. If $r \rightarrow \infty$, $V = 0$.

Solid sphere of radius R.

$$\vec{E} = \begin{cases} \frac{kqr}{R^3} & \text{if } r < R \\ \frac{kq}{r^2}\hat{r} & \text{otherwise} \end{cases}$$

Gauss' Law

$$\Phi = \frac{q_{enc}}{\epsilon_0}; \Phi = \oint \vec{E} \cdot d\vec{A}$$

A must be a Gaussian surface. If \vec{E} is constant on the surface, it can be simplified to $\Phi = E * A$. Conductors in an electric field have $\vec{E} = 0$ inside. Electrons move to ensure this. Inside, $\Phi = 0$.

Electrical Potential Difference

Path independent. For $\vec{E}(x, y, z)$:

$$\Delta V = \frac{\Delta U}{q} = - \int_i^f \vec{E} \cdot d\vec{x} = \int_i^f dV$$

Electric field lines go from more positive to more negative voltage.

Equipotential surface (ES): Surface with same V .

Conductors have equipotential volumes and $\vec{E} = 0$

$$V = \frac{kq}{r} = \int \frac{k dQ}{r}; \vec{E} = -\nabla V$$

Capacitance (C)

Relationship between charged separated and potential difference. $Q = C * \Delta V$ To find capacitance:

1. Draw a picture
2. Determine direction of \vec{E}
3. Determine \vec{E} (Gauss' and determined distributions help), then $\Delta V = - \int \vec{E} \cdot d\vec{s}$
4. Calculate C with $C = \frac{Q}{\Delta V}$

For parallel plates, $C = \frac{A\epsilon_0}{d}$.

For cylindrical capacitor length L , $C = \frac{2\pi L\epsilon_0}{\ln(b/a)}$.

Concentric spheres of radii a and b , $C = 4\pi\epsilon_0 \frac{ab}{b-a}$.

Isolated sphere of radius R , $C = 4\pi\epsilon_0 R$.

Since $W = q\Delta V$, $\Delta U = \frac{1}{2}C*\Delta V^2 = \frac{q^2}{2C}$ (Electric Potential Energy)

$$u = \frac{1}{2}\epsilon_0 E^2 = \frac{U}{Vol}$$

A dielectric/material is in an electric field has a dielectric constant κ . In it, ϵ_0 is replaced with $\kappa\epsilon_0$. κ of metals is considered ∞ . $\kappa(vacuum) = 1$

If you put a dielectric in a capacitor, treat it like a network of capacitors in a creative alignment.

Add a dielectric to charged capacitor:

$$Q_\kappa = Q_0; V_\kappa < V_0; C_\kappa > C_0; U_\kappa < U_0$$

Add a dielectric to battery-connected capacitor:

$$V_\kappa = V_0; Q_\kappa > Q_0; C_\kappa > C_0; U_\kappa > U_0$$

Current

$$I = \frac{dq}{dt}$$

Ohm's Law: $V = IR$

Junction rule: For any point on a circuit, $I_{in} = I_{out}$

Stored charge at junction slows down I_{in} & speeds up I_{out}

Resistors

For a resistor of resistance R , length L , cross section A , resistivity ρ .

$$R = \frac{V}{I} = \rho \frac{L}{A}; P = IV = I^2 R = \frac{V^2}{R}; V = \frac{dW}{dq}$$

Current Density

For a cross-section \vec{A} , $dI = \vec{J} \cdot d\vec{A}$

$$\vec{J} = e * \vec{v}_d * n = \frac{\vec{E}}{\rho}$$

Circuits

Batteries keep ΔV constant

Long end of battery diagram is + side

Series Parallel

Capacitor $\frac{1}{C} = \sum \frac{1}{C_i}$ $C = \sum C_i$

Resistor $R = \sum R_i$ $\frac{1}{R} = \sum \frac{1}{R_i}$

RL circuit, current approaches equilibrium $\frac{\mathcal{E}}{R}$ (growth & decay resp).

$$I = \frac{\mathcal{E}}{R}(1 - e^{-t/\tau}); \tau = \frac{L}{R}; I = I_0 e^{-t/\tau}$$

Charge of charging capacitor (growth & decay):

$$q = C\mathcal{E}(1 - e^{-t/\tau}); q = q_0 e^{-t/RC}$$

Kirchoff's Loop Rule: The total voltage throughout a circuit must sum to zero. Ex: $\mathcal{E} = L \frac{dI}{dt} + IR$.

For an inductor, energy stored $U_B = \frac{1}{2}LI^2$.

Magnetism

$$\vec{F} = q(\vec{v} \times \vec{B})$$

$$\vec{F} = I\vec{L} \times \vec{B}$$

↑ For electric field B , charged wire of length L . ↑

Hall effect: electrons traveling through a strip (length l) of conductor in a magnetic field B generate an internal capacitor. Charge Q , current I .

$$Q = \frac{BI}{Vl}$$

For conductor moving at v through B , area A , electron drift v_d .

$$V = vBd; v_d = \frac{I}{neA}$$

Free charge moving in magnetic field centers source at radius r .

$$qvB = \frac{mv^2}{r}; f = \frac{qB}{2\pi m}$$

Energy density in a magnetic field $u_B = \frac{B^2}{2\mu_0}$.

Current generates a magnetic field. Current rounding area A N times makes a magnetic dipole.

$$\tau = NI\vec{A} \times \vec{B} = \vec{\mu} \times \vec{B}$$

$$U(\theta) = -\vec{\mu} \cdot \vec{B}$$

A distance R from a straight wire, $B = \frac{\mu_0 I}{2\pi R}$. Along wire b from wire a : $B = \frac{\mu_0 I_a}{2\pi d}$.

At center of arc of angle ϕ of radius R , $B = \frac{\mu_0 I \phi}{4\pi R}$.

Ampere's Law

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 I_{enc}$$

Inside an ideal (infinite) solenoid, $B = \mu_0 In = \mu_0 I \frac{N}{L}$.

Inside an ideal toroid, $B = \frac{\mu_0 IN}{2\pi R}$.

Along coil axis (coils create magnetic dipoles), $B = \frac{\mu_0 \vec{\mu}}{2\pi z^3}$.

Farraday's Law

In a conductor of area A , mag. field B .

$$\Phi_B = \int_S \vec{B} \cdot d\vec{A}; \mathcal{E} = \frac{d\Phi_B}{dt}$$

In a solenoid, inductance $L = \frac{N\Phi_B}{I}$.

For length l , $\frac{L}{l} = \mu_0 n^2 A$.

Creates self-opposing emf of magnitude $\mathcal{E} = -L \frac{dI}{dt}$.

Electric Dipoles

$$\vec{E} = \begin{cases} < 0 & \text{if } -\frac{d}{2} < z < \frac{d}{2} \\ > 0 & \text{otherwise} \end{cases}$$
$$= \frac{2kQd}{z^3 \left(1 - \frac{d^2}{4z^2}\right)^2} \hat{d}$$

ESs are \perp to \vec{p} . In an electric field:

$$\vec{p} = Q\vec{d}$$

$$\vec{\tau} = \vec{p} \times \vec{E}$$

$$U = -\vec{p} \cdot \vec{E}$$

$$W_{net} = \Delta K = -\Delta U$$

Maxwell's Equations

Gradient

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Integral

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$$

$$\oint \vec{E} \cdot d\vec{l} = -\int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{A}$$

$$\oint \vec{E} \cdot d\vec{A} = 0$$

$$\oint \vec{E} \cdot d\vec{l} = \mu_0 I_{enc} + \mu_0 \epsilon_0 \frac{\partial \Phi_E}{\partial t}$$