

Conclusion to Last Lecture

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In notes.

Non-Linear Conservation Laws

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Remember the transport equation for $p(x,t)$
for $-\infty < x < \infty, t > 0$

$$p_t + q_x = 0$$

where q = flux function.

For the traffic equation

$$q = cp$$

$$\frac{\text{mass}}{\text{time}} = \frac{\text{length}}{\text{time}} \cdot \frac{\text{mass}}{\text{length}}$$

Previously we set $c = c(x,t)$ which could be thought of as the speed limit.

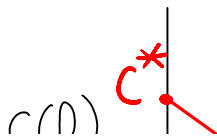
A much more realistic model is

$$c = c(\rho)$$

This will yield a non-linear DE for the density.

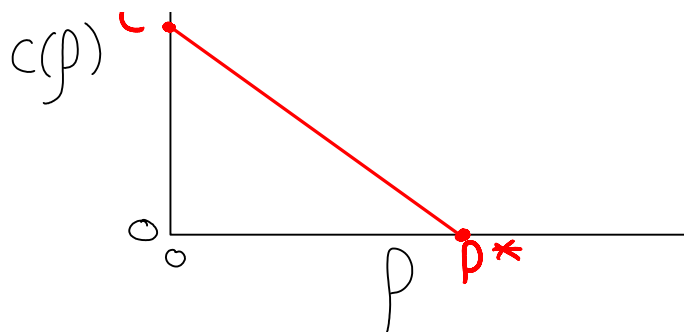
Choose simplest model: a linear model for $c(\rho)$

c linear $\rightarrow q$ quadratic \rightarrow non-linear



c^* = open highway speed

ρ^* = jam density

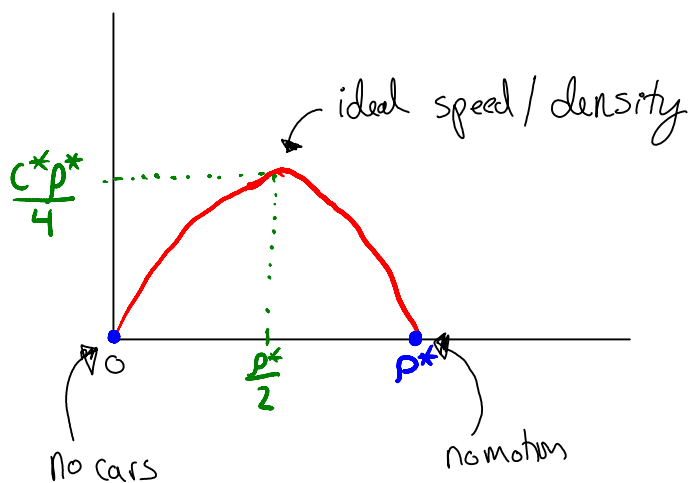


$\rho^* = \text{jam density}$

We will limit density $0 \leq \rho \leq \rho^*$

$$c(\rho) = c^* \left(1 - \frac{\rho}{\rho^*}\right)$$

So $q(\rho) = c\rho = c^* \left(1 - \frac{\rho}{\rho^*}\right)\rho$



Our traffic equation becomes

$$\rho_t + [q(\rho)]_x = 0$$

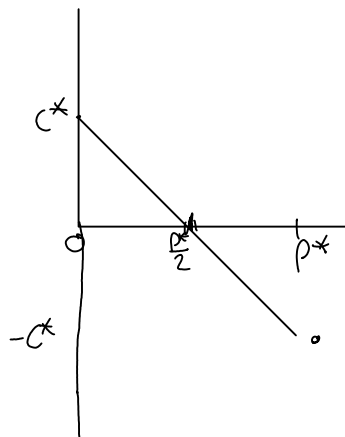
This yields

$$\rho_t + q'(\rho)\rho_x = 0$$

Let $q'(\rho) = v(\rho) = c^* \left(1 - 2\frac{\rho}{\rho^*}\right)$

$$\text{Let } g'(p) = v(p) = c^* \left(1 - 2 \frac{p}{p^*}\right)$$

$$\rho_t + v(p) \rho_x = 0$$



$c \equiv$ phase velocity (speed of a car)

$v \equiv$ group velocity (speed of a line of constant density)

can be negative, related to fact that if car put on its breaks, break lights propagate backwards.

Let's simplify our lives,

set $c^* = 1$, $p^* = 1$ (like gaussian units)

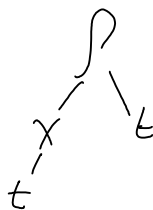
$$\text{D.E.: } \rho_t + v(p) \rho_x = 0 \quad -\infty < x < \infty$$

$$v(p) = 1 - 2p \quad 0 < t$$

$$\text{I.C.: } \rho(x, 0) = f(x) \quad -\infty < x < \infty$$

Know initial condition

Solve using characteristics. Look for $x(t)$, a characteristic curve on which p is constant.



$$\frac{dp}{dt} = \frac{\partial p}{\partial t} + \frac{dx}{dt} \frac{\partial p}{\partial x} = 0$$

But from traffic equation

$$p_t + V(p)p_x = 0,$$

This implies

$$\boxed{\frac{dx}{dt} = V(p)}$$

and also on this curve

$$\boxed{\frac{dp}{dt} = 0}$$

a function of p and x on right would go here

\Downarrow
 $p = \text{const. on characteristics}$

The characteristics are: the solutions to

$$\frac{dx}{dt} = V(p)$$

← can do this because p is constant on characteristic

$$x = V(p)t + \xi$$

$$\xi = x - V(p)t$$

Now

$$p(\xi) = A(\xi)$$

so convert by fact to

$$p(x, t) = A(x - V(p)t) \quad \text{but}$$

$$p(x, 0) = A(x) = f(x), \quad \text{so } A = f \text{ and}$$

$$\boxed{p(x, t) = f(x - V(p)t)}$$

Good News: I have a solution.

Bad News: The solution is implicit.

Let's check this. Certainly

I.C. $\rho(x, 0) = f(x)$ ✓

What about the P.E.?

$$\begin{array}{c} f \\ | \\ \xi \\ \swarrow \downarrow \searrow \\ x \quad p \quad t \\ \swarrow \searrow \\ x \quad t \end{array} = x - v(p)t$$

$$\rho(x, t) = f(\xi) = f(x - v(p)t)$$

$$\frac{\partial \rho}{\partial x} = \frac{\partial}{\partial x} [f(\xi)] = f'(\xi) \xi_x \text{ but } \xi_x = \frac{\partial}{\partial x} (x - v(p)t) = 1 - v'(p)p_x t$$

so

$$\rho_x = f'(\xi) (1 - v'(p)p_x t)$$

and we

$$\rho_x (1 - f'(\xi) v'(p) t) = f'(\xi)$$

$$\rho_x = \frac{f'(\xi)}{1 + f'(\xi) v'(p) t}$$

$$\frac{\partial p}{\partial t} = \frac{\partial}{\partial t} [f(\xi)] = f'(\xi) \xi_t$$

$$\text{but } \xi_t = \frac{\partial}{\partial t} (x - v(p)t) = 0 - v'(p)p_t t - v(p)$$

$$p_t = f'(\xi) (-v'(p)p_t t - v(p))$$

$$p_t (1 + f'(\xi) v'(p) t) = -v(p)$$

$$p_t = \frac{-f'(\xi) v(p)}{1 + f'(\xi) v'(p) t} = -v(p) p_x$$

$$\boxed{p_t + v(p) p_x = 0} \quad \checkmark \quad \text{☺}$$

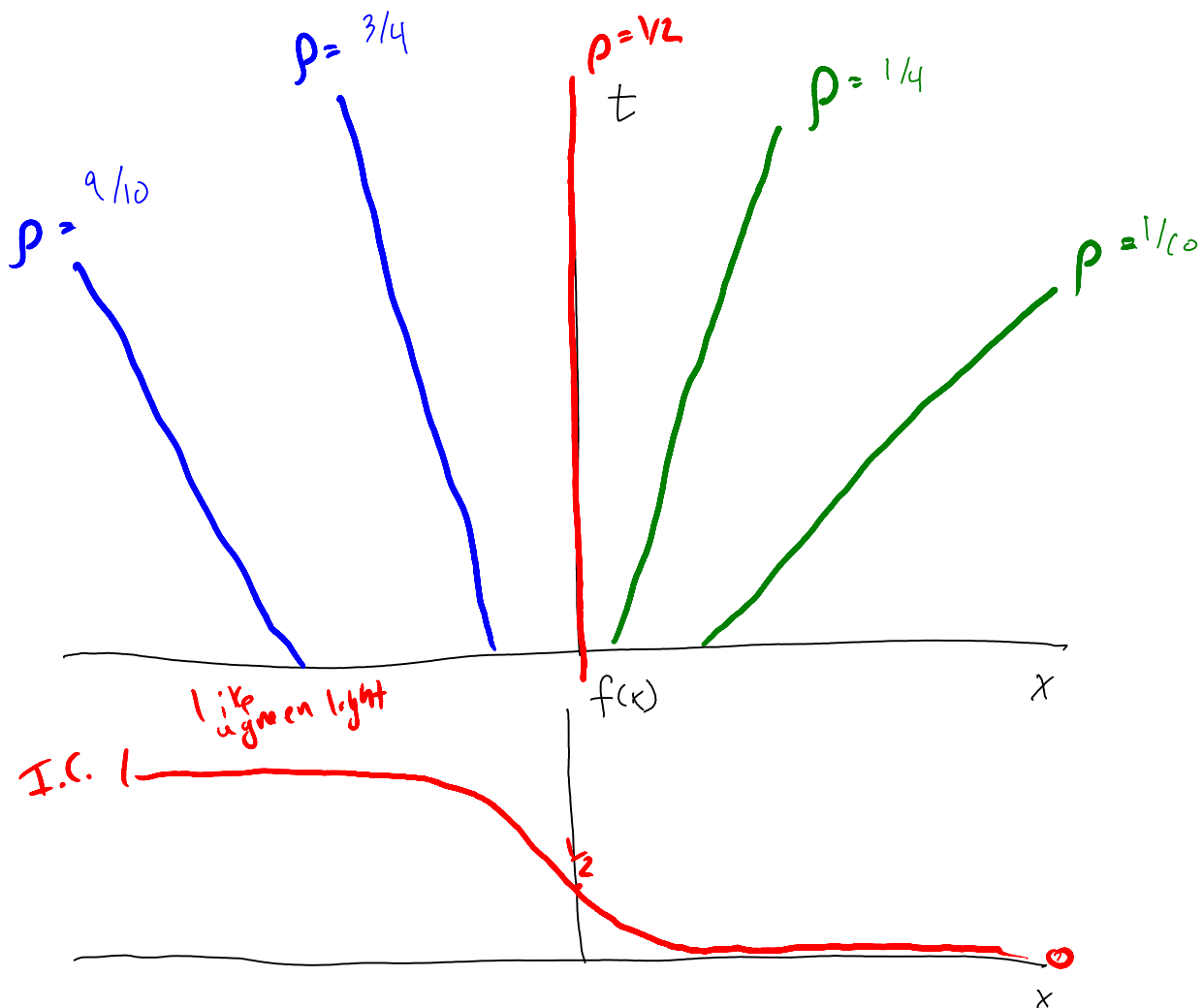
Example Solutions

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Maple Worksheet online

"The Green Light Problem"

Characteristics (lines of constant density)



Spreads out as time increases, \rightarrow linear constant (in HW)

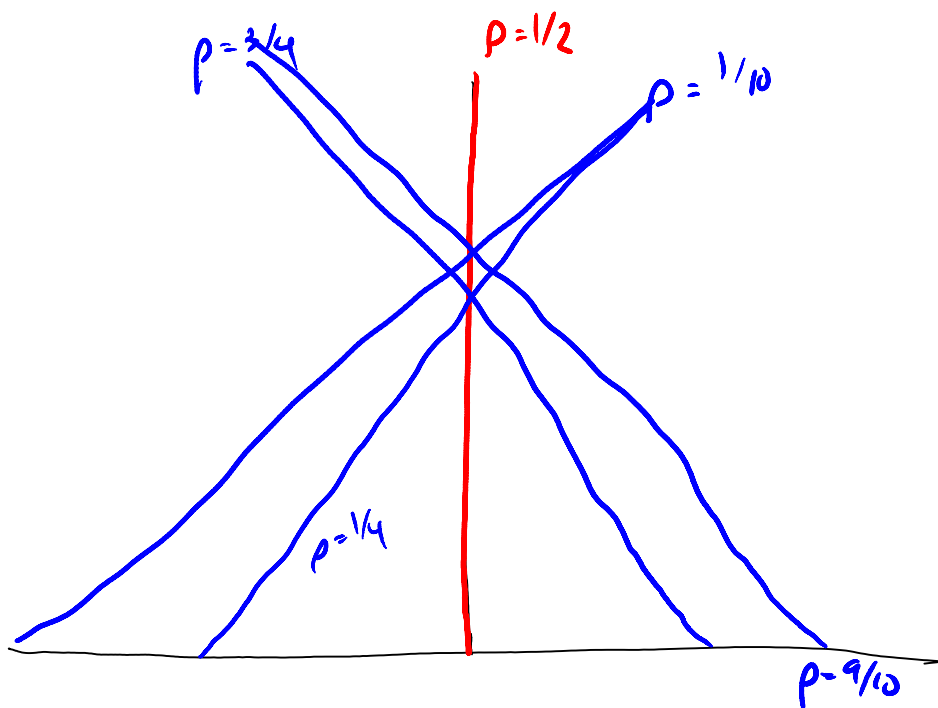
Called a rarefaction



Density is spreading.

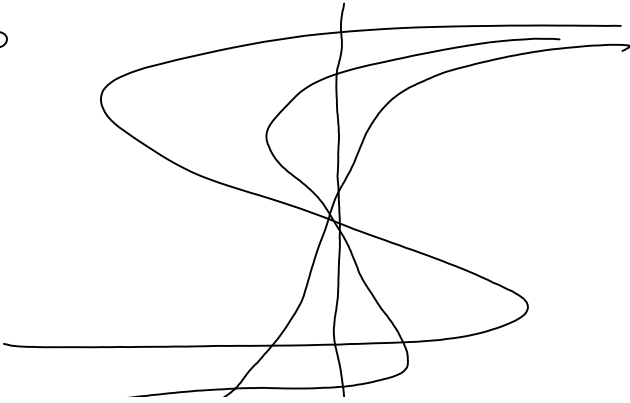
Example 2

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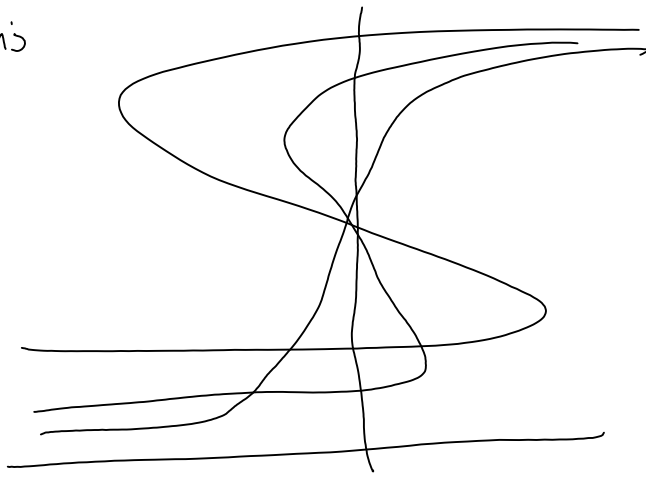
Characteristics cross?

sols



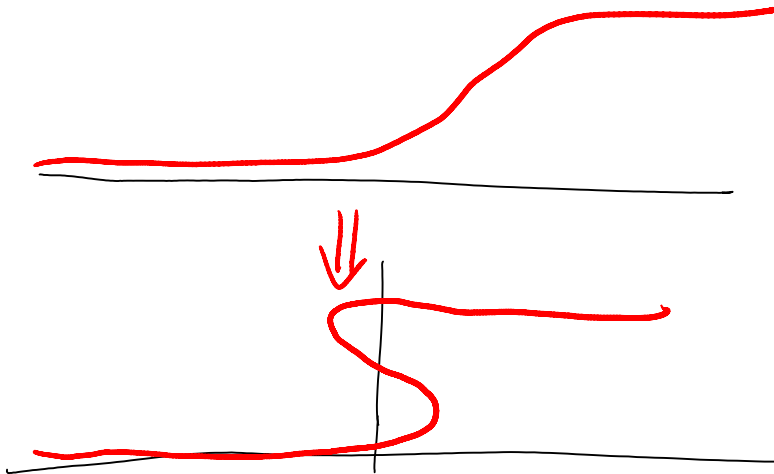
multi-valued
solutions

Solns



multi-valued
solutions

This is compression - solutions become
multi-valued. (huh?)



we broke
our physical
model. How do
we find a
single
valued solution?

How do you find a single valued solution?

- Look for the unique solution w/ a jump
that conserves mass.

This is called a **SHOCK**

Jump is called a SHOCK.

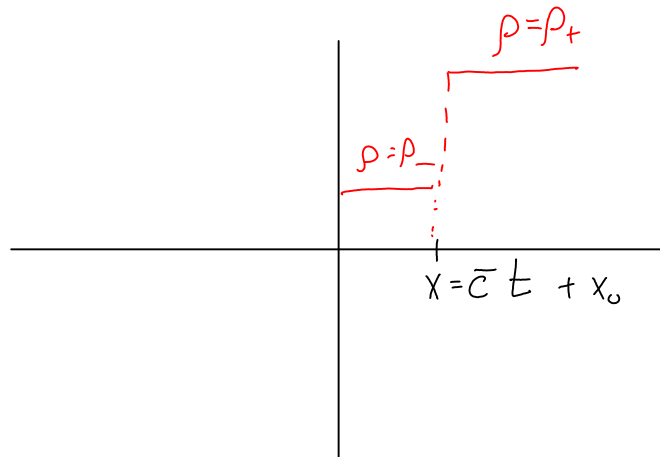
How to do this?

Shock Example

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Ex. Look for a piecewise constant solution



(look for a traveling wave solution)

$$\text{Let } \rho = \bar{\rho}(x - \bar{c}t) = \bar{\rho}(\xi)$$

$$\rho_t + \frac{\partial}{\partial x} (Q(\rho)) = 0$$

$$\text{Note } \rho_t = \bar{\rho}_\xi \dot{\xi}_t = -\bar{c} \bar{\rho}_\xi$$

$$\text{and } \frac{\partial}{\partial x} = \frac{\partial}{\partial \xi} \frac{\partial \xi}{\partial x} = \frac{\partial}{\partial \xi}$$

The equation for $\bar{\rho}$ is

$$-\bar{c} \bar{\rho}_\xi + \frac{\partial}{\partial \xi} [Q(\rho)] = 0$$

$$\bar{p}_3 = \frac{\partial}{\partial \xi} [Q(\rho)]$$

Integrate from $\int_{\xi=-\xi}^{\xi=\xi}$

$$\int_{-\xi}^{\xi} \bar{p}_3 d\xi = \int_{-\xi}^{\xi} \frac{\partial}{\partial \xi} Q(\rho) d\xi$$

$$\bar{p}_3 \Big|_{-\xi}^{\xi} = Q(\rho) \Big|_{-\xi}^{\xi}$$

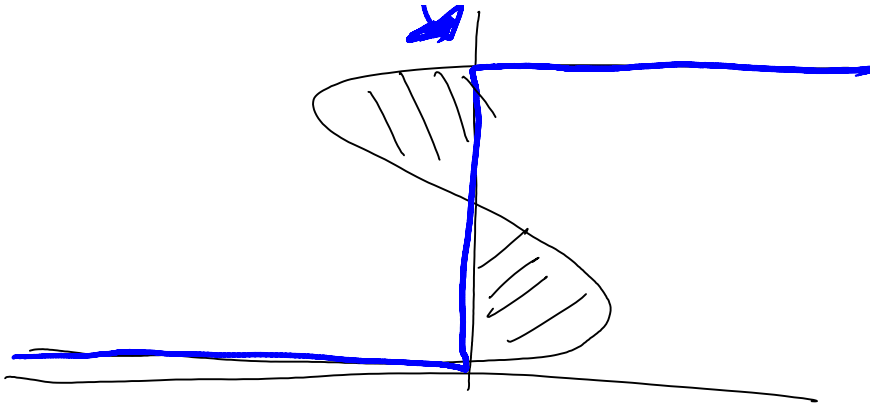
$$\bar{c}[\rho_1, \rho_2] = Q(\rho_+) - Q(\rho_-)$$

$$\bar{c} = \frac{Q(\rho^+) - Q(\rho^-)}{\rho^+ - \rho^-} \quad \text{Shock speed}$$

Unique speed at which you can move such that the shock can move at.

From last prob





equal area rule

Choose the shock that cuts off equal areas.

Moral of the story:

non-linear systems are nasty

Sometimes you've got to add more physics

One non-linear part in FLW, no shocks.