



# Moiré heterostructures as a condensed-matter quantum simulator

Dante M. Kennes<sup>1,2,12</sup> , Martin Claassen<sup>3,4,12</sup> , Lede Xian<sup>2,5,12</sup> , Antoine Georges<sup>3,6,7,8</sup>, Andrew J. Millis<sup>3,9</sup>, James Hone<sup>10</sup> , Cory R. Dean<sup>9</sup>, D. N. Basov<sup>9</sup> , Abhay N. Pasupathy<sup>9</sup> and Angel Rubio<sup>2,3,11</sup>

**Twisted van der Waals heterostructures have latterly received prominent attention for their many remarkable experimental properties and the promise that they hold for realizing elusive states of matter in the laboratory. We propose that these systems can, in fact, be used as a robust quantum simulation platform that enables the study of strongly correlated physics and topology in quantum materials. Among the features that make these materials a versatile toolbox are the tunability of their properties through readily accessible external parameters such as gating, straining, packing and twist angle; the feasibility to realize and control a large number of fundamental many-body quantum models relevant in the field of condensed-matter physics; and finally, the availability of experimental readout protocols that directly map their rich phase diagrams in and out of equilibrium. This general framework makes it possible to robustly realize and functionalize new phases of matter in a modular fashion, thus broadening the landscape of accessible physics and holding promise for future technological applications.**

The inseparability of groups of particles into products of their individual states underpins the most counterintuitive predictions of quantum mechanics. In condensed-matter physics, recent focus has been placed on finding and understanding quantum materials — systems in which a delicate balance between the crystal lattice and strong electronic interactions lead to new inseparable, collective behaviours with exotic properties. Encompassing rich phenomenologies ranging from superconductivity at high temperatures to topologically ordered states of matter with fractionalized excitations, numerous materials have recently come under intense scrutiny for the realization of quantum phases on demand<sup>1</sup>.

A major impediment to the systematic study of correlated electron physics is the comparative lack of tunability of conventional chemical compounds. The main means of control are pressure, strain and doping. This limited range of control knobs severely limits experimental exploration of the wider phase diagram, essential to guide the discovery of novel quantum phases and to controlled studies of quantum criticality, unconventional phase transitions, or realizations of exotic topological states of matter that are both of fundamental interest and could have applications in emerging quantum technologies such as quantum computing and simulation.

Another impediment lies in the difficulty of solving these paradigmatic quantum Hamiltonians via present computational approaches. As an alternative, in the early 80s, Feynman proposed quantum simulation as a paradigm to turn the problem upside down<sup>2</sup>: implementing clean models using real physical systems, to ‘simulate’ the ground state, thermodynamic behaviour or non-equilibrium dynamics of such models. If the implementation is well controlled, the results can be used to disentangle the more

complex material’s phenomena for which the original model was set up, as well as to guide the way towards stabilizing and controlling new and exotic phases of matter in real quantum materials.

Much progress has been achieved via studying ultracold gases of bosonic and fermionic atoms, which can be confined in optical lattices to realize lattice models of condensed-matter physics in a controlled manner<sup>3</sup>. While these systems are highly controllable and can faithfully realize certain quantum Hamiltonians, it remains an experimental challenge to engineer Hamiltonians with tunable long-range interactions and especially to access low-temperature emergent long-range ordered phases, such as the intriguing *d*-wave state of the *t*–*t'* repulsive Hubbard model<sup>4</sup>.

This Review identifies moiré heterostructures of van der Waals materials as an alternative and complementary condensed-matter approach to realize a large set of highly controllable quantum Hamiltonians. While such systems do not typically afford the high level of isolation and precise tunability that cold atomic systems do, the large degree of tunability control available through readily accessible experimental knobs allows for the exploration of phase diagrams of vast and novel sets of many-body Hamiltonians in and out of equilibrium and at very low temperatures. In particular, we review how different choices of two-dimensional (2D) heterostructure compositions at different twist angles realize a wide range of effective low-energy electronic tight-binding Hamiltonians with various geometries, dimensionality, frustration, spin–orbit coupling as well as interactions beyond the Hubbard limit, including partially screened and unscreened Coulomb interactions. This allows for flexible interaction engineering as well as control of band structure and topology.

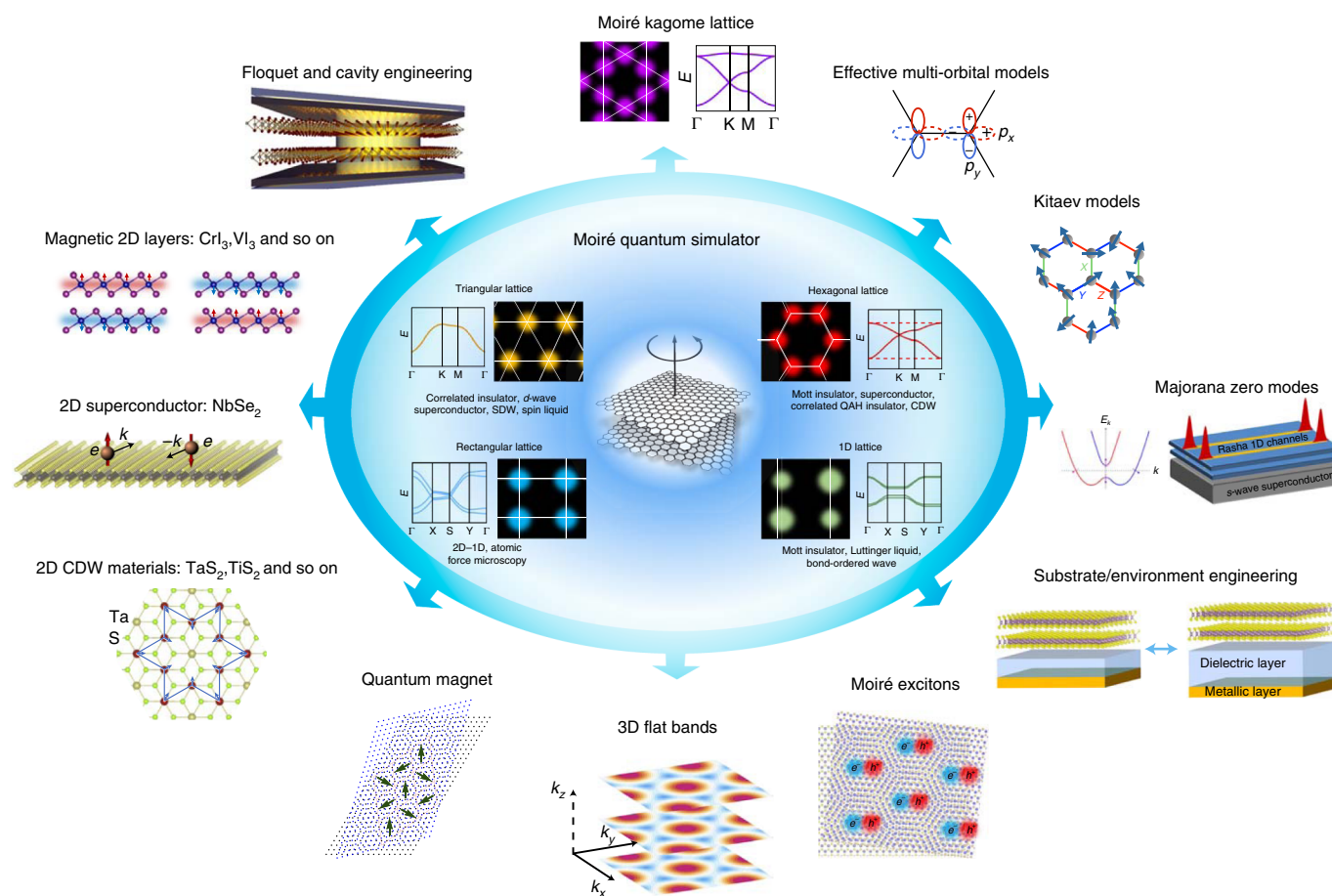
<sup>1</sup>Institut für Theorie der Statistischen Physik, RWTH Aachen University and JARA-Fundamentals of Future Information Technology, Aachen, Germany.

<sup>2</sup>Max Planck Institute for the Structure and Dynamics of Matter and Center Free-Electron Laser Science, Hamburg, Germany. <sup>3</sup>Center for Computational Quantum Physics, Flatiron Institute, New York, NY, USA. <sup>4</sup>Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA, USA.

<sup>5</sup>Songshan Lake Materials Laboratory, Guangdong, China. <sup>6</sup>Collège de France, Paris, France. <sup>7</sup>CPHT, CNRS, École Polytechnique, IP Paris, Palaiseau, France.

<sup>8</sup>DQMP, Université de Genève, Geneva, Switzerland. <sup>9</sup>Department of Physics, Columbia University, New York, NY, USA. <sup>10</sup>Department of Mechanical Engineering, Columbia University, New York, NY, USA. <sup>11</sup>Nano-Bio Spectroscopy Group, Universidad del País Vasco UPV/EHU, San Sebastian, Spain.

<sup>12</sup>These authors contributed equally: Dante M. Kennes, Martin Claassen, Lede Xian. ✉e-mail: [Dante.Kennes@rwth-aachen.de](mailto:Dante.Kennes@rwth-aachen.de); [apn2108@columbia.edu](mailto:apn2108@columbia.edu); [Angel.Rubio@mpsd.mpg.de](mailto:Angel.Rubio@mpsd.mpg.de)



**Fig. 1 | Moiré quantum simulator.** Stacking sheets of van der Waals heterostructures with a twist gives rise to a plethora of effective low-energy Hamiltonians. Some of these realizations that were studied in the literature are given in the inner circle of the figure with the corresponding lattice, band structure and potential phases of matter that can be realized. For the red band structure, the dashed lines indicate the flat bands additionally present in the case of twisted bilayer  $\text{MoS}_2$  compared with TBG. In the outer circle, we outline a perspective on possible future directions of moiré quantum simulators with many intriguing developments to be expected.  $k$ ,  $k_x$ ,  $k_y$  and  $k_z$  denote momentum and momentum in the  $x$ ,  $y$  and  $z$  directions, respectively.  $E$  and  $E_k$  refer to band energy.

## Realizing model quantum Hamiltonians in van der Waals heterostructures

Moiré heterostructures of van der Waals materials exploit quantum interference to quench the effective kinetic energy scales, which permits both driving the system at low energies to an interaction-dominated regime, and drastically enhancing anisotropies or reducing symmetries of the monolayer. Conceptually, the physical mechanism can be understood straightforwardly in analogy to classical moiré interference patterns (Fig. 1, centre): if two identical patterns with a discrete translation symmetry in two dimensions are superimposed at an angle, the combined motif retains a periodicity at much longer wavelengths for an infinite set of small rational angles.

Similarly, a heterostructure of two or more monolayers with commensurate lattice constants stacked at a twist retains spatial periodicity for small commensurate angles, albeit with a much larger moiré unit cell. This moiré superlattice can span many unit cells and defines a new crystal structure with a mini Brillouin zone. Within this moiré Brillouin zone, the folded bands lead to a multitude of crossings, which are subsequently split via interlayer hybridization<sup>5</sup>. Crucially, lattice relaxation enhances these avoided crossings and typically segregates sets of almost-dispersionless bands on millielectronvolt energy scales<sup>6</sup>, which can be addressed individually via gate-tuning the chemical potential.

A complementary view of the physics of a moiré band follows from considering its real space description. A localized electron in a moiré band can be viewed as occupying a virtual ‘moiré orbital’ that can extend over hundreds of atoms, effectively spreading over the entire moiré unit cell. Importantly, the shape of moiré orbitals changes with the size of the moiré unit cell so that the electronic interactions become a function of the twist angle as well, decreasing at a slower rate than the kinetic energy scales for decreasing angle. This permits tuning a priori weakly interacting electronic systems into regimes dominated by electronic correlations in a controlled fashion.

A first demonstration of the above phenomenology was realized in materials with hexagonal structure and multiple atomic orbitals. The moiré band width for twisted bilayers of graphene does not behave monotonically as a function of twist angle but exhibits a series of ‘magic’ angles<sup>5</sup> at which the moiré bands near charge neutrality become almost dispersionless over a large fraction of the moiré Brillouin zone.

More broadly, as reductions of the electronic band width equivalently enhance the role of competing energy scales, this microscopic structural knob to selectively quench the kinetic energy scales in 2D materials opens up possibilities to selectively engineer heterostructures with properties that are dominated via many-body electronic interactions, ultrastrong spin–orbit interactions from heavier

**Table 1 | Overview of possible quantum Hamiltonians, materials realizations and phases in twisted moiré heterostructures**

Lattice	Model	Possible materials realizations	Correlated phases
<b>Twisted heterostructures of weakly correlated van der Waals monolayers</b>			
Honeycomb	Two-orbital extended Hubbard model <sup>18</sup> ; fragile topological insulator <sup>107</sup>	TBG (BN substrate, with/without twist)	Mott insulation <sup>8</sup> ; superconductivity <sup>7</sup> ; correlated QAH insulator <sup>22,23</sup>
	Two-orbital extended Hubbard model	Twisted double bilayer graphene	Ferromagnetic insulator superconductivity <sup>12,13</sup> ; triplet pairing <sup>108</sup>
	Asymmetric $p_x$ - $p_y$ Hubbard model <sup>29,30</sup>	Twisted bilayer MoS <sub>2</sub> , MoSe <sub>2</sub>	Nematic (anti)ferromagnets <sup>29</sup>
	Domain wall networks	Small-angle TBG with domain reconstruction <sup>57,58,66,109</sup>	
Triangular	Hubbard model (with/without strong SOC)	Twisted bilayer WS <sub>2</sub> , WSe <sub>2</sub> (ref. <sup>34</sup> ); twisted WS <sub>2</sub> /WSe <sub>2</sub> heterostructures <sup>35,36</sup> ; twisted double bilayers of WSe <sub>2</sub> (ref. <sup>110</sup> )	Correlated insulator <sup>34</sup> ; superconductivity?; Wigner crystals <sup>35</sup>
	Doped multi-orbital Hubbard models	Twisted heterostructures of MoS <sub>2</sub> , WS <sub>2</sub> , WSe <sub>2</sub>	Moiré excitons <sup>100,101,111</sup>
	Multi-orbital Kanamori models	Twisted bilayer BN	SDW; $d$ -wave superconductivity <sup>32</sup>
Rectangular	1D ionic Hubbard model 1D-2D crossover	Twisted bilayer GeSe	Luttinger liquid; Mott insulator; bond density waves <sup>43</sup>
	Inverted band insulator, strong SOC	Twisted bilayer WTe <sub>2</sub>	Quantum spin Hall insulator, fractional Chern/topological insulator
Any	Hofstadter models	TBG or TMDs in strong magnetic fields	Fractional Chern insulator <sup>112</sup>
Kagome	Kagome Heisenberg model	TBD	Z <sub>2</sub> QSL; U(1) QSL; quantum chiral spin liquid; valence bond crystal
Decorated kagome	Hubbard model (putative)	Twisted bilayer MoS <sub>2</sub> , MoSe <sub>2</sub>	TBD
3D	Flat-band Hubbard-Kanamori models	Twisted multilayer 'staircase'	TBD
<b>Proximity effects</b>			
Honeycomb, triangular	Proximity-induced Rashba SOC	TBG on WS <sub>2</sub> , WSe <sub>2</sub> substrate <sup>113</sup>	Correlated QSH insulator
Rectangular	Proximity-induced superconductivity	Superconductor, twisted bilayer GeSe, TMD 'sandwich' heterostructure	1D Kitaev superconductor; Majorana bound states
<b>Twisted heterostructures of correlated monolayers</b>			
	Moiré ferromagnet <sup>70</sup>	Twisted odd-multilayer CrI <sub>3</sub>	Moiré domain wall; ferromagnets
	Moiré Kitaev model	Twisted multilayer $\alpha$ -RuCl <sub>3</sub>	Kitaev QSL; stripe order; Majorana fermions
	TBD	Twisted bilayer TaSe <sub>2</sub>	TBD
	TBD	Twisted bilayer NbSe <sub>2</sub>	TBD

TBD, to be discovered; QSL, quantum spin liquid; SOC, spin-orbit coupling.

elements, electron–lattice interactions or electron–photon interactions, permitting realization of a wide range of novel correlated or topological phenomena. Here we provide a perspective on how different universal classes of quantum many-body Hamiltonians can be engineered and controlled using moiré heterostructures

### The moiré quantum simulator

Table 1 summarizes some of the realizable lattice structures, associated model Hamiltonians and featured quantum phases that can be achieved within a 2D twisted van der Waals heterostructure framework. In our discussion, we will concentrate first on those systems realized or proposed in the literature, summarized in the inner circle of Fig. 1. Then, we will provide a perspective on future potential control and simulation possibilities, shown in the outer circle of Fig. 1.

**Honeycomb lattice.** So far, most experimental and theoretical research has concentrated on twisted bilayer graphene (TBG) (see refs. <sup>7,8</sup>, and a review<sup>9</sup> for an overview on TBG), with extensions

to triple- and quadruple-layer generalizations<sup>10–16</sup>. TBG realizes a moiré superstructure with the relevant low-energy degrees of freedom effectively again forming a honeycomb lattice, although at millielectronvolt energy scales and imbued with an additional orbital degree of freedom reflecting the two layers of the original system<sup>17,18</sup>.

With superconductivity<sup>7,19–21</sup>, correlated insulators<sup>8,20,21</sup> and the quantum anomalous Hall (QAH) effect<sup>22,23</sup> already realized experimentally, an intriguing possibility indicated by recent theoretical analyses is that repulsive interactions favour topological  $d+id$  instead of nodal  $d$ -wave pairing<sup>24–26</sup>. Such topological superconducting states are potentially relevant to topological quantum computing<sup>27</sup>, and can be harnessed and controlled via tailored laser pulses<sup>28</sup>.

Expanding the catalogue of engineered lattice structures, we next consider twisting two monolayers of MoS<sub>2</sub> (refs. <sup>29–31</sup>). In these structures, families of flat bands emerge at the band edges, the first two of which realizes a single orbital Hubbard model while the second realizes a more exotic strongly asymmetric  $p_x$ - $p_y$  Hubbard model<sup>29,30</sup>, both on the honeycomb lattice (see inner circle of Fig. 1). The latter emerges at small twist angle, where the overall band width

of these families of flat bands is tuned into the millielectronvolt regime. Due to destructive interference within the strongly asymmetric  $p_x$ - $p_y$  Hubbard model itself (meaning that the two orbitals have very different hopping amplitudes), attached at the bottom and top of these flat bands one can find bands with even much lower dispersion. The asymmetric  $p_x$ - $p_y$  Hubbard model is therefore another lattice Hamiltonian that can effectively be engineered in this solid-state framework, and also controlled by external parameters, such as strain or fields. This could provide interesting insights into highly degenerate systems and the interplay of magnetism with these types of band structure. Similar phenomena are expected to emerge in other twisted transition metal dichalcogenide (TMD) homobilayers, such as twisted bilayer  $\text{MoSe}_2$  and  $\text{WS}_2$ .

**Triangular lattices.** A natural extension of TBG involves two twisted sheets of boron nitride (BN), which is often called white graphene for its structural similarity and large bandgap. When twisting two BN sheets, the consequences are even more dramatic; while in graphene, the quenching of the kinetic energy scales relies on being closed to a set of magic angles, in BN three families of flat band emerge, where the flattening of the bands is a monotonous function of twist angle. The effective lattice structure that is engineered at low enough twist angle is effectively triangular<sup>32,33</sup>.

A similar engineered lattice structure is obtained by twisting two sheets of tungsten diselenide  $\text{WSe}_2$  (ref. <sup>34</sup>) or other TMD heterostructures<sup>35–37</sup>. In all these cases, a superstructure emerges that effectively confines the low-energy physics onto a triangular lattice. In addition, the heavier transition metal elements in such TMD systems impose substantial spin–orbit coupling onto the low-energy degrees of freedom, which can in turn induce interesting topological properties.

Fundamentally, triangular lattices are prototypical model systems to study the role of geometric frustration on electronic and magnetic orders. Using twisted materials one might thus simulate effectively the properties of these theoretically challenging models, relevant also for understanding many other condensed-matter systems, such as bis(ethylenedithio)tetrathiafulvalene organic compounds and their rich phase diagrams<sup>38–40</sup>. Debated correlated states of matter that can be realized this way range from  $120^\circ$  Neel ordered states to exotic, topological forms of superconductivity and quantum chiral spin liquids<sup>41</sup>.

Indeed, in the case of TMD homobilayers, features consistent with superconducting phases (with further experimental evidence for a truly superconducting state still highly desirable) have already been seen in close proximity to the insulating phase<sup>34</sup>. An advantage of these materials is that the angular precision needed for the emergence of new quantum phases is not as stringent as in graphene, so samples are easier to produce. However, the quality of the constituent materials tends to be poorer, and might need to be improved to allow a well-controlled realization of the model Hamiltonians and phenomena described in Table 1 and Fig. 1.

**Rectangular lattices.** While all of the examples above build on stacking monolayers with a three-fold ( $120^\circ$ ) rotational symmetric lattice structure, this is certainly not the only interesting scenario that one can consider<sup>42</sup>. To illustrate this, we next turn our attention to twisted bilayer germanium selenide ( $\text{GeSe}$ )<sup>43</sup> or similarly germanium sulfide ( $\text{GeS}$ ). Monolayer  $\text{GeSe}$  has a rectangular unit cell with a  $180^\circ$  rotational symmetry. When two sheets of  $\text{GeSe}$  are twisted, a similar quenching of the kinetic energy scales can be found as in the above discussed cases; however, the effective lattice structure realized by such a system is a square lattice at small angles of  $8^\circ$ – $15^\circ$  while at lower angles the hopping in one of the principal directions is reduced much more strongly than the corresponding hopping along the second principal direction. At about  $6^\circ$ , the effective lattice system realized is entirely dispersionless along one direction

while it has a residual (though small due to the band narrowing at small angles) dispersion along the second one (see inner circle of Fig. 1). Thus, twist angle provides a control knob to tune the system from effectively 2D to effectively 1D low-energy physics, providing new insights into this fundamental and poorly understood problem.

One-dimensional quantum systems exhibit strong collective effects. Bilayer  $\text{GeSe}$  at a small twist angle should thus provide an experimental example to study this emergent quantum regime, where Luttinger liquid, Mott insulating and bond-ordered wave states dominate the stage. Furthermore, sizable spin–orbit coupling of the low-energy bands emerges at larger twist angles, suggesting that interesting spin-momentum locked states with possibly nematic properties could be engineered<sup>44,45</sup>.

## Experimental realizations and readout of the moiré quantum simulator

Next we outline some experimental progress on the realization and readout of controlled moiré structures.

**Experimental realizations.** Moiré structures on layered materials have been observed by microscopy for at least the past three decades. In 3D crystals, they can appear naturally during growth due to the presence of screw dislocations, or can be produced by mechanical damage of the crystal structure. These structures remained mostly curiosities until the advent of graphene, when it was realized that moiré structures are an effective way of tuning band structure (strong peaks in the local density of states were observed when twisting graphene on bulk graphite)<sup>46</sup>. Such moiré patterns are commonly seen in chemical vapour deposition-grown multilayer graphene and TMD films, and many experiments to understand the properties of these patterns were performed in the past decade<sup>47,48</sup>. In these samples, the twist angles generated were largely uncontrolled, and often showed disorder on submicrometre lengthscales, making the general applicability of these samples to other experiments limited.

The first steps on the road to making moiré heterostructures was established by mechanically stacking two monolayers<sup>49</sup> of graphene and hexagonal BN (hBN) to make high-quality graphene layers for transport measurements. Due to the natural lattice mismatch between graphene and hBN, such structures naturally featured moiré textures<sup>47</sup>. The most successful technique to realize twisted bilayer systems is the ‘tear and stack’ method and its variants<sup>50</sup>. This technique essentially separates a single monolayer into two pieces, which are then naturally aligned with each other. One of the pieces is picked up, rotated by the desired angle and stacked with the second piece to create the desired moiré structure. This simple and elegant technique has so far proved to be the most productive method, but has its limitations in reproducibility and uniformity of the moiré patterns produced (other techniques such as pushing with an atomic force microscopy tip<sup>51</sup> or introducing heterostrain<sup>52</sup> have been developed). The development of techniques to measure the wavelength and uniformity of the Moiré patterns *in situ* are promising developments in the search to create the high-quality and precise moiré structures that we are discussing here.

The in-plane bonding in 2D materials can arise dominantly from covalent bonding or alternatively from an ionic one. For the former, exfoliation and twisting sheets of materials should work well, while for the latter it will be difficult to do the same. In addition, there is a large class whose bonding is intermediate between ionic and covalent. For those, exfoliation and twisting techniques need to be tested experimentally.

**Readout of the ground-state properties.** Electrical transport has been the technique of choice to probe the ground-state quantum properties of moiré structures. The first experiments in this direction were the measurements of quantum Hall effects in graphene on



hBN, which revealed the influence of the moiré potential via the formation of Hofstadter butterfly patterns in the magnetotransport<sup>53</sup>.

A dramatic development was the discovery of collective quantum phases in TBG near the so-called magic angle<sup>5</sup>. After the original discovery of insulating<sup>8</sup> and superconducting phases<sup>7</sup>, new orbitally magnetized phases with a quantized Hall effect, indicating the development of topological phases that spontaneously break time-reversal symmetry, have also been observed<sup>22,20</sup>. Similar collective and topological phases have been seen in the case of ABC graphene on hBN<sup>11</sup>, in which case the moiré pattern between the ABC graphene and hBN provides additional flattening of the band structure that already displays a peak in the density of states at charge neutrality.

More recently, the methods applied to graphene have also been applied to TMD-based heterostructures. The first experiments indicating the development of interaction-driven insulating phases have been performed in both TMD heterobilayers and homobilayers<sup>34,36</sup>.

**Spectroscopic readout.** Moiré structures present two great advantages for experiments compared with traditional solid-state materials. The first is that the size of the ‘lattice’ is expanded from the atomic scale to several nanometres, allowing optical and near-field techniques to directly probe the local spectroscopic properties of the system, in close analogy with the quantum gas microscope invented in the case of cold atomic gases. The second advantage is in their intrinsic 2D nature, which makes it possible to measure the properties of the entire sample with surface-sensitive spectroscopic probes such as scanning tunnelling microscopy or angle-resolved photoemission. This does come with its disadvantages however, as many traditional spectroscopic probes applied to solid-state systems (neutron spectroscopy, for example) produce signals that are proportional to the volume of the sample being measured, and are thus not easily applicable to these materials.

Scanning tunnelling microscopy has been applied extensively to study the band structure in twisted graphene structures<sup>46,47,54–56</sup>, to visualize internal edge states that exist at domain boundaries in small-angle twisted bilayer and four-layer graphene systems, and to relate these states to the topological properties of the individual domains<sup>57,58</sup>. Touching on many-body correlation effects, prominent results obtained so far are a measurement of the insulating gap at half-filling, evidence that this insulating state does not break translational symmetry but does break rotational symmetry near half-filling<sup>46,54,56</sup> and the discovery of an insulating phase at charge neutrality in four-layer rhombohedral graphene produced by twisting two bilayers<sup>58</sup>.

Furthermore, van der Waals heterostructures display nearly all optical phenomena found in solids, including plasmonic oscillations of free electrons characteristic of metals, light emission/lasing and excitons encountered in semiconductors, and intense phonon resonances typical of insulators<sup>59,60</sup>. Therefore, optical spectroscopies allow one to reconstruct the role of moiré superlattice potentials in the electronic structure of twisted multilayers as well as in their electron and lattice dynamics. For example, infrared spectroscopy on a graphene/BN heterostructure shows that the moiré superlattice potential is dominated by a pseudospin-mixing component analogous to a spatially varying pseudomagnetic field<sup>61</sup>, and photoluminescence has uncovered the systematic evolution of light-emitting properties associated with interlayer excitons in TMD heterobilayers<sup>62</sup>.

Conventional diffraction-limited optics suffers from the spatial resolution restricted by the wavelength of light and therefore provides only area-averaged information on the electromagnetic response of van der Waals heterostructures. Modern nano-optical methods<sup>63,64</sup> allow one to overcome this limitation and enable optical enquiry at the lengthscale commensurate with moiré periodicities. In addition, nano-optical methods give access to spectroscopy

and imaging of hybrid light–matter modes known as polaritons<sup>59</sup>. Moiré structures also allow the setting up of a highly accurate tool enabling moiré metrology of energy landscapes in 2D van der Waals structures<sup>65</sup>.

We conclude this subsection by noting that functionalized atomic force microscope tips enable multiple nanoscale contrasts in addition to nano-infrared studies, including nano-Raman, magnetic force microscopy, Kelvin probe and piezo force microscopy<sup>66</sup>. All these different methods provide complimentary insight carrying separate messages about the studied phenomena<sup>65</sup>. Co-located visualization of contrasts obtained with multiple functionalized scanning probes offers an opportunity for multimessenger imaging of the interplay between electronic, magnetic and lattice effect at the lengths scale of moiré domains<sup>67</sup>.

### What comes next?

An important question is which lattice structure might be realizable using moiré physics in the future. Some promising future avenues of research are illustrated in Table 1 and in the outer circle of Fig. 1 — while extensive, these lists are not exhaustive and we are confident that other low-energy model realizations will emerge as more heterostructure ‘building blocks’ become experimentally viable. We concentrate on a few highly sought after candidates next.

### Quantum LEGO with correlated or topological monolayers.

Beyond utilizing a priori weakly correlated monolayers as over-viewed above, a promising approach concerns constructing twisted heterostructures out of monolayers that by themselves already exhibit correlation-driven quantum phases. We present a short perspective of prototype 2D materials that should be studied in the future.

- (1) In niobium diselenide ( $\text{NbSe}_2$ ), superconducting behaviour can be found even when exfoliated down to the monolayer (at temperatures of the order of 3 K)<sup>68</sup>. Two such sheets might provide a solid-state-based in-road into the question of how the superconducting state emerges within a conventional Bardeen–Cooper–Schrieffer picture and how this state crosses over to a strong coupling regime.
- (2) In the 2D van der Waals material chromium(III) iodide ( $\text{CrI}_3$ ), quantum magnetic properties can be realized that depend on the number of layers and relative stacking order<sup>69</sup>. Twisted heterostructures made of  $\text{CrI}_3$  could permit accessing antiferromagnetically and ferromagnetically aligned domains in a flexible fashion using moiré lattice engineering<sup>70,71</sup>.
- (3)  $\text{RuCl}_3$  realizes a Kitaev–Heisenberg magnet with stripe order<sup>72</sup>. Layering on a substrate at a twist, or twisting layers of  $\text{RuCl}_3$ , can provide a method to tune the magnetic interactions and putatively tune the system towards a quantum spin liquid, as well as study the effects of doping a proximal spin liquid in a controlled manner.
- (4)  $\text{TaSe}_2$  realizes a correlation-driven insulating state that is characterized by a star-of-David charge density wave (CDW) configuration at low temperatures that yields a rather large unit cell (Fig. 1, outer circle)<sup>73</sup>. Again, this correlation-driven phase is present even in the untwisted case. Twisting two of these structures atop each other raises the interesting questions of whether and how superstructures made of these larger star-of-David CDWs can be realized and what are their manifestation in the electronic properties of the material.
- (5) Beyond van der Waals materials, the moiré approach to modify electronic structures is also at play in correlated oxides. These are epitaxial structures and therefore cannot be rotated easily. However a moiré pattern in films of the prototypical magnetoresistive oxide  $\text{La}_{0.67}\text{Sr}_{0.33}\text{MnO}_3$  can be epitaxially grown on  $\text{LaAlO}_3$  substrates. The net effect is that both the electronic

conductivity and ferromagnetism are modulated by this moiré engineering over mesoscopic scales<sup>67</sup>. This opens up yet another route in the combinatory space of chemical compositions to use in moiré systems.

- (6) 1T' WTe<sub>2</sub> monolayers realize a 2D topological insulator due to a spin-orbit-driven topological band inversion<sup>74</sup>. While layering two WTe<sub>2</sub> monolayers would naively lead to a Z<sub>2</sub> trivial insulating state, the superlattice potential at finite twist angles can in principle induce a series of stacking-controlled topological transitions. A remarkable consequence of such a device would be the possibility to study the combination of non-trivial band topology and electronic correlations in moiré bands with quenched dispersion, potentially giving rise to experimentally realizable manifestations of fractional Chern insulators or fractional topological insulator phases.

**Highly controllable geometrically frustrated lattices.** One long-sought research goal regards the controlled realization of kagome or Kitaev quantum magnets<sup>75</sup>. In conventional candidate systems such as herbertsmithite, RuCl<sub>3</sub>, iridates or metal-organic complexes, disorder-free realizations of the low-energy magnetic Hamiltonian or suppression of unwanted and longer-ranged spin exchange interactions are major challenges. A moiré realization of the kagome lattice or of strongly spin-orbit-coupled multi-orbital moiré models in the strong-interaction limit would permit tuning magnetic interactions in a controlled manner, possibly allowing a realization of a quantum spin liquid state. This would elevate moiré heterostructures as a platform to simulate an experimentally elusive phase of matter in a highly controllable setting, with a potentially worthwhile starting candidate being twisted layers of RuCl<sub>3</sub>.

Another direction recently put forward in the realization of an effective kagome lattice, is twisted MoS<sub>2</sub>, which was discussed above in the context of the asymmetric  $p_x$ - $p_y$  Hubbard model. Beyond this at very small angles, the next family of bands, which can be accessed by further doping the system, would effectively realize a multi-orbital generalization of a kagome lattice<sup>29,30</sup>. Future studies need to address whether such small angles and large degree of doping can be realized experimentally.

**Proximity effects and spin-orbit coupling.** One key advantage of 2D systems is the possibility to induce effects via proximity to an engineered substrate. To avoid chemical modifications of the moiré system by the interface through strong chemical bonds, which would destroy the low-energy moiré bands of the original heterostructure, one can employ 2D heterostacking. Many interesting effects are expected in flat-band systems, when superconducting or strongly spin-orbit-coupled substrates are used to imprint some of their properties on the 2D system under scrutiny. For instance, large induced spin-orbit coupling<sup>76</sup> might reveal novel topological superconducting states with applications to quantum computing. Similarly, proximitization with superconducting substrates could permit a controlled study of Majorana wires and lattices<sup>77–79</sup>, for example, via using a quasi-1D moiré material such as twisted GeSe. An important corollary of exploiting a moiré superstructure would be the delocalization of bound Majorana states on scales larger than the moiré unit cell, promoting their length scale from ångströms to nanometres. Potential topological phenomena include the demonstration of a high-temperature QAH effect and Majorana modes as well as antiferromagnetic topological insulators, which open a route towards topological spintronics.

**From 2D to 3D.** Another intriguing future research direction concerns extending the idea of simulating quantum models from the 1D or 2D (discussed so far) to the 3D realm using moiré systems. To this end, a new exfoliation technique was recently demonstrated,

to yield large-area atomically thin layers that can be stacked in any desired order and orientation to generate a whole new class of artificial materials, including stacking thick twisted materials<sup>80</sup>. We next discuss explicitly two examples, with obvious intriguing extensions possible.

- (1) Alternating twist. When stacking many layers atop each other, we want to define the twist angle of each layer to be measured with respect to the bottom one. Stacking 2D materials at alternating twist then means that the twist angle alternates between two values: zero and  $\alpha$ . If  $\alpha$  is small, in-plane localized sites will emerge by the moiré interference physics, and these sites will lie directly atop each other in the out-of-plane direction. The in-plane flat-band states will become dispersive along the out-of-plane direction. On slight doping, the system thus becomes dominantly a 1D metal with small residual in-plane coupling by the residual dispersion of the flat bands, similar to what has been discussed for bulk TaS<sub>2</sub> (ref. <sup>81</sup>). Such a 1D metal is susceptible to instabilities (such as CDW, spin density wave (SDW) or excitonic instabilities) as a 1D system is always perfectly nested. This will gap out the system along the out-of-plane direction and elevate the relevance of the small dispersion along the in-plane direction. A similar mechanism was recently observed to give rise to the 3D quantum Hall effect in bulk ZrTe<sub>5</sub> (ref. <sup>82</sup>) and similar fascinating surprises might be expected in the alternating twist configuration we propose here.
- (2) Continuous twist. Another intriguing idea is to stack layer by layer with constant twist angle  $\alpha_i$  between adjacent layers, such that the above defined twist angle with respect to the bottom layer increases linearly,  $\alpha = \alpha_0 i$  modulo 360°, where  $i$  is the layer index. Calculations to characterize such a system are challenging as the unit cell with respect to the in-plane direction grows huge or a quasi-crystal structure with no unit cell emerges when more and more layers are added. Here the generalized Bloch band theory developed in ref. <sup>83</sup> might be useful. However, the general physics of emerging flat bands should transfer to the third direction, providing an inroad to realize 3D quantum models with tunable ratio of potential and kinetic energy scales as discussed above for the 2D case.

**Moiré heterostructures out of equilibrium.** So far, we have addressed equilibrium properties of twisted systems and we have illustrated the rich phenomena that one can realize via engineering a variety of low-energy quantum many-body Hamiltonians. With increasingly sensitive pump-probe experiments available, another enticing possibility is to drive such systems out of equilibrium via short laser pulses or via embedding the material in a quantum cavity<sup>84–88</sup>. As the moiré potential lowers the relevant kinetic energy scales to millielectronvolts, these lie within reach of the typical energy scales of light-matter interactions, permitting outsized transient modifications of the electronic dynamics. Combining control of the twist angle with non-equilibrium perturbations thus holds great promise to offer a new platform with opportunities to realize novel correlated non-equilibrium states of matter.

- (1) Floquet engineering. In TBG, the sensitivity of correlated phases to millielectronvolt changes of the dispersion of the moiré bands as a function of deviation from the magic angle suggests that rich effects can be achieved already via dressing the electronic bands by weak light fields. In the presence of an optical pump, a photon-mediated renormalization of the hopping matrix elements between moiré unit cells can selectively modify the low-energy band structure and flatten its dispersion.
- (2) Cavity engineering. Analogously, cavity engineering recently attracted great interest as a means to exploit strong light-mat-

ter coupling to change electronic phases. The comparatively low energy scales of moiré bands and correlated phases in TBG suggests that coupling to a cavity photon mode can induce an outsized effect on the interacting electronic state. One such application is in photo-induced and cavity-induced superconductivity. New predictions for quantum phenomena in these cavities include the enhancement of superconductivity using the coupling between the vacuum fluctuations of the cavity and plasmonic modes of the environment<sup>87–90</sup>, the study of strongly correlated Bose–Fermi mixtures of degenerate electrons and dipolar excitons<sup>91–95</sup>.

- (3) Optically generated synthetic magnetic fields. The strong spin–orbit coupling in TMDs offers new possibilities to generate opto-magnetic quantum properties using dynamic fields. For instance, one can optically generate real magnetic fields from spin–valley polarization in TMD heterobilayers. Under circularly polarized excitation at the bandgap of one of the TMD monolayers in a heterobilayer, the initially large spin–orbit splitting can be further increased, leading to a high degree of spin–valley polarization of the hole. This, along with the lack of spin–valley polarization of the electron, will result in a net spin polarization, and thus, a real magnetic moment with the potential to induce novel non-equilibrium quantum phases.

**Excitons at will.** Moiré platforms also provide a basic infrastructure for achieving different interlayer exciton phases and excitonic lattices (Fig. 1). In particular, bilayers of TMDs were identified (even without twist) as intriguing candidates because the long lifetime of these charge-separated interlayer excitons<sup>96–98</sup> should facilitate their condensation<sup>99</sup>. Adding twist angle allows the realization of moiré exciton lattices<sup>62,100–103</sup>, which can feature topological exciton bands that support chiral excitonic edge states<sup>104</sup>. The high degree of control permits study of the localized–delocalized crossover of excitons in moiré lattices as well as a strain-induced crossover from 2D to 1D exciton physics<sup>105</sup>. Minimizing the kinetic energy to increase interactions is key to achieving robust condensation. For example, twisted WSe<sub>2</sub> homobilayers have tunable flat electronic bands over a range of twist angles (see above), with alternatives including experimentally accessible heterobilayers as well as strain and pressure engineering. Two such band-engineered TMDs separated by a uniform spacer layer (hBN)<sup>101</sup> will probably provide the best avenue to achieve exciton condensation in zero applied field and to then tune the properties of the condensate by gating, field, pressure and strain. Future strategies to enhance the tendencies of excitons to condensate beyond the use of twist angle to achieve flat bands and thus more prominent interactions could also include the use of cavity structures to enhance light–matter interaction (Fig. 1). Here, pump–probe nano-optical methods can both address and interrogate distinct regions in moiré heterostructures<sup>67</sup>.

**Skyrmions.** Skyrmions — vortex-like magnetic configurations — can arise as a spontaneously formed hexagonal lattice in the ground state of chiral magnets, or as quasiparticle excitations on top of a ferromagnetic ground state. The energetic stability of skyrmions is a consequence of an anti-symmetric exchange (also called Dzyaloshinskii–Moriya) interaction and they are interesting due to their non-trivial topological structure protecting them against perturbations and rendering them highly attractive as potential memory devices. One candidate system for which skyrmions have been predicted is monolayer CrI<sub>3</sub> (ref. 106). In the realm of twisted magnetic van der Waals heterostructures, moiré skyrmion lattices are favoured due to the long-period moiré pattern, which can be used to tune the periodicity and shape of the magnetization texture. In addition, an external magnetic field can be used to switch the vorticity and location of skyrmions.

## Conclusion

Here we have identified future directions of research into the blossoming fields of twisted van der Waals heterostructures with a particular emphasis on the high degree of control achievable in these materials. Controlled engineering of quantum Hamiltonians as tunable low-energy theories of these structures is at the forefront of this fascinating research field. This idea might be viewed as a condensed-matter realization of a quantum simulator in the future, which allows direct access to fundamental properties of quantum many-body systems, such as the delicate emergent physics of novel collective phases of matter, which have been experimentally elusive so far. In this sense, twisted van der Waals heterostructures might be the long-sought-after remedy to access the wider phase space, which should include all kinds of fundamentally and technologically relevant phases of matter. One of the major experimental challenges to be overcome next is to devise readout techniques of the magnetic properties in these systems as well as to push the precision of angle-resolved photoemission spectroscopy into the regime where it can resolve the physics within the small Brillouin zone of twisted heterostructures reliably. Both of these advances would provide cornerstones allowing the interrogation of the low-energy physics of the system with regards to its interacting properties more completely. Considering the tremendous combinatorial space of van der Waals materials to choose from, extensions to higher dimensions and the vibrant field of non-equilibrium control, including Floquet engineering and twistronics applied to these novel platforms, it is likely that only a small fraction of the potential phenomena achievable in twisted van der Waals heterostructures have been experimentally realized so far. This newly emerging research field is likely to yield many more exciting developments.

Received: 23 May 2020; Accepted: 14 December 2020;

Published online: 01 February 2021

## References

1. Basov, D. N., Averitt, R. D. & Hsieh, D. Towards properties on demand in quantum materials. *Nat. Mater.* **16**, 1077–1088 (2017).
2. Feynman, R. P. Simulating physics with computers. *Int. J. Theor. Phys.* **21**, 467–488 (1982).
3. Bloch, I., Dalibard, J. & Nascimbene, S. Quantum simulations with ultracold quantum gases. *Nat. Phys.* **8**, 267–276 (2012).
4. Zheng, B.-X. et al. Stripe order in the underdoped region of the two-dimensional Hubbard model. *Science* **358**, 1155–1160 (2017).
5. Bistritzer, R. & MacDonald, A. H. Moiré bands in twisted double-layer graphene. *Proc. Natl Acad. Sci. USA* **108**, 12233–12237 (2011).
6. Yoo, H. et al. Atomic and electronic reconstruction at the van der Waals interface in twisted bilayer graphene. *Nat. Mater.* **18**, 448–453 (2019).
7. Cao, Y. et al. Unconventional superconductivity in magic-angle graphene superlattices. *Nature* **556**, 43–50 (2018).
8. Cao, Y. et al. Correlated insulator behaviour at half-filling in magic-angle graphene superlattices. *Nature* **556**, 80–84 (2018).
9. Balents, L., Dean, C. R., Efetov, D. K. & Young, A. F. Superconductivity and strong correlations in moiré flat bands. *Nat. Phys.* **16**, 725–733 (2020).
10. Chen, G. et al. Signatures of tunable superconductivity in a trilayer graphene moiré superlattice. *Nature* **572**, 215–219 (2019).
11. Chen, G. et al. Evidence of a gate-tunable Mott insulator in a trilayer graphene moiré superlattice. *Nat. Phys.* **15**, 237–241 (2019).
12. Liu, X. et al. Tunable spin-polarized correlated states in twisted double bilayer graphene. *Nature* **583**, 221–225 (2020).
13. Shen, C. et al. Correlated states in twisted double bilayer graphene. *Nat. Phys.* **16**, 520–525 (2020).
14. Cao, Y. et al. Tunable correlated states and spin-polarized phases in twisted bilayer–bilayer graphene. *Nature* **583**, 215–220 (2020).
15. Burg, G. W. et al. Correlated insulating states in twisted double bilayer graphene. *Phys. Rev. Lett.* **123**, 197702 (2019).
16. He, M. et al. Symmetry breaking in twisted double bilayer graphene. *Nat. Phys.* **17**, 26–30 (2021).
17. Yuan, N. F. Q. & Fu, L. Model for the metal–insulator transition in graphene superlattices and beyond. *Phys. Rev. B* **98**, 045103 (2018).
18. Koshino, M. et al. Maximally localized Wannier orbitals and the extended Hubbard model for twisted bilayer graphene. *Phys. Rev. X* **8**, 031087 (2018).
19. Yankowitz, M. et al. Tuning superconductivity in twisted bilayer graphene. *Science* **363**, 1059–1064 (2019).



20. Lu, X. et al. Superconductors, orbital magnets and correlated states in magic-angle bilayer graphene. *Nature* **574**, 653–657 (2019).
21. Stepanov, P. et al. Untying the insulating and superconducting orders in magic-angle graphene. *Nature* **583**, 375–378 (2020).
22. Sharpe, A. L. et al. Emergent ferromagnetism near three-quarters filling in twisted bilayer graphene. *Science* **365**, 605–608 (2019).
23. Serlin, M. et al. Intrinsic quantized anomalous Hall effect in a moiré heterostructure. *Science* **367**, 900–903 (2020).
24. Xu, C. & Balents, L. Topological superconductivity in twisted multilayer graphene. *Phys. Rev. Lett.* **121**, 087001 (2018).
25. Liu, C.-C., Zhang, L.-D., Chen, W.-Q. & Yang, F. Chiral spin density wave and  $d + id$  superconductivity in the magic-angle-twisted bilayer graphene. *Phys. Rev. Lett.* **121**, 217001 (2018).
26. Kennes, D. M., Lischner, J. & Karrasch, C. Strong correlations and  $d + id$  superconductivity in twisted bilayer graphene. *Phys. Rev. B* **98**, 241407 (2018).
27. Nayak, C., Simon, S. H., Stern, A., Freedman, M. & Sarma, S. D. Non-abelian anyons and topological quantum computation. *Rev. Mod. Phys.* **80**, 1083 (2007).
28. Claassen, M., Kennes, D. M., Zingl, M., Sentef, M. A. & Rubio, A. Universal optical control of chiral superconductors and Majorana modes. *Nat. Phys.* **15**, 766–770 (2018).
29. Xian, L. et al. Realization of nearly dispersionless bands with strong orbital anisotropy from destructive interference in twisted bilayer MoS<sub>2</sub>. Preprint at <https://arxiv.org/abs/2004.02964> (2020).
- A study of twisted bilayer MoS<sub>2</sub> shown to feature strongly asymmetric  $p_x$ - $p_y$  multi-orbital low-energy electronic bands.**
30. Angeli, M. & MacDonald, A.  $\gamma$ -valley transition-metal-dichalcogenide moiré bands. Preprint at <https://arxiv.org/abs/2008.01735> (2020).
31. Naik, M. H. & Jain, M. Ultraflatbands and shear solitons in moiré patterns of twisted bilayer transition metal dichalcogenides. *Phys. Rev. Lett.* **121**, 266401 (2018).
32. Xian, L., Kennes, D. M., Tancogne-Dejean, N., Altarelli, M. & Rubio, A. Multiflat bands and strong correlations in twisted bilayer boron nitride: doping-induced correlated insulator and superconductor. *Nano Lett.* **19**, 4934–4940 (2019).
33. Ni, G. X. et al. Soliton superlattices in twisted hexagonal boron nitride. *Nat. Commun.* **10**, 4360 (2019).
34. Wang, L. et al. Correlated electronic phases in twisted bilayer transition metal dichalcogenides. *Nat. Mater.* **19**, 861–866 (2020).
35. Regan, E. C. et al. Mott and generalized Wigner crystal states in WS<sub>2</sub>/WS<sub>2</sub> moiré superlattices. *Nature* **579**, 359–363 (2020).
36. Tang, Y. et al. Simulation of Hubbard model physics in WS<sub>2</sub>/WS<sub>2</sub> moiré superlattices. *Nature* **579**, 353–358 (2020).
37. Wu, F., Lovorn, T., Tutuc, E. & MacDonald, A. H. Hubbard model physics in transition metal dichalcogenide moiré bands. *Phys. Rev. Lett.* **121**, 026402 (2018).
38. Shimizu, Y., Miyagawa, K., Kanoda, K., Maesato, M. & Saito, G. Spin liquid state in an organic Mott insulator with a triangular lattice. *Phys. Rev. Lett.* **91**, 107001 (2003).
39. Kagawa, F., Miyagawa, K. & Kanoda, K. Unconventional critical behaviour in a quasi-two-dimensional organic conductor. *Nature* **436**, 534–537 (2005).
40. Tajima, N., Sugawara, S., Tamura, M., Nishio, Y. & Kajita, K. Electronic phases in an organic conductor  $\alpha$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub>: ultra narrow gap semiconductor, superconductor, metal, and charge-ordered insulator. *J. Phys. Soc. Jpn* **75**, 051010–051010 (2006).
41. Szasz, A., Motruk, J., Zalete, M. P. & Moore, J. E. Chiral spin liquid phase of the triangular lattice hubbard model: a density matrix renormalization group study. *Phys. Rev. X* **10**, 021042 (2020).
42. Kariyado, T. & Vishwanath, A. Flat band in twisted bilayer Bravais lattices. *Phys. Rev. Res.* **1**, 033076 (2019).
43. Kennes, D. M., Xian, L., Claassen, M. & Rubio, A. One-dimensional flat bands in twisted bilayer germanium selenide. *Nat. Commun.* **11**, 1124 (2020).
44. Harter, J. W., Zhao, Z. Y., Yan, J.-Q., Mandrus, D. G. & Hsieh, D. A parity-breaking electronic nematic phase transition in the spin-orbit coupled metal Cd<sub>2</sub>Re<sub>2</sub>O<sub>7</sub>. *Science* **356**, 295–299 (2017).
45. Harter, J. W. et al. Evidence of an improper displacive phase transition in Cd<sub>2</sub>Re<sub>2</sub>O<sub>7</sub> via time-resolved coherent phonon spectroscopy. *Phys. Rev. Lett.* **120**, 047601 (2018).
46. Jiang, Y. et al. Charge order and broken rotational symmetry in magic-angle twisted bilayer graphene. *Nature* **573**, 91–95 (2019).
47. Li, G. et al. Observation of van Hove singularities in twisted graphene layers. *Nat. Phys.* **6**, 109–113 (2010).
48. Wong, D. et al. Local spectroscopy of moiré-induced electronic structure in gate-tunable twisted bilayer graphene. *Phys. Rev. B* **92**, 155409 (2015).
49. Dean, C. R. et al. Boron nitride substrates for high-quality graphene electronics. *Nat. Nanotechnol.* **5**, 722–726 (2010).
50. Kim, K. et al. Tunable moiré bands and strong correlations in small-twist-angle bilayer graphene. *Proc. Natl Acad. Sci. USA* **114**, 3364–3369 (2017).
51. Finney, N. R. et al. Tunable crystal symmetry in graphene-boron nitride heterostructures with coexisting moiré superlattices. *Nat. Nanotechnol.* **14**, 1029–1034 (2019).
52. Edelberg, D., Kumar, H., Shenoy, V., Ochoa, H. & Pasupathy, A. N. Tunable strain soliton networks confine electrons in van der Waals materials. *Nat. Phys.* **16**, 1097–1102 (2020).
53. Dean, C. R. et al. Hofstadter's butterfly and the fractal quantum Hall effect in moiré superlattices. *Nature* **497**, 598–602 (2013).
54. Kerelsky, A. et al. Maximized electron interactions at the magic angle in twisted bilayer graphene. *Nature* **572**, 95–100 (2019).
55. Xie, Y. et al. Spectroscopic signatures of many-body correlations in magic-angle twisted bilayer graphene. *Nature* **572**, 101–105 (2019).
56. Choi, Y. et al. Electronic correlations in twisted bilayer graphene near the magic angle. *Nat. Phys.* **15**, 1174–1180 (2019).
57. Huang, S. et al. Topologically protected helical states in minimally twisted bilayer graphene. *Phys. Rev. Lett.* **121**, 037702 (2018).
58. Kerelsky, A. et al. Moiré-less correlations in ABCA graphene. Preprint at <https://arxiv.org/abs/1911.00007> (2019).
59. Sunko, S. S. et al. Photonic crystals for nano-light in moiré graphene superlattices. *Science* **362**, 1153–1156 (2018).
60. Low, T. et al. Polaritons in layered two-dimensional materials. *Nat. Mater.* **16**, 182–194 (2017).
61. Shi, Z. et al. Gate-dependent pseudospin mixing in graphene/boron nitride moiré superlattices. *Nat. Phys.* **10**, 743–747 (2014).
62. Seyler, K. L. et al. Signatures of moiré-trapped valley excitons in MoSe<sub>2</sub>/WS<sub>2</sub> heterobilayers. *Nature* **567**, 66–70 (2019).
63. Liu, M., Sternbach, A. J. & Basov, D. N. Nanoscale electrodynamics of strongly correlated quantum materials. *Rep. Prog. Phys.* **80**, 014501 (2016).
64. Hesp, N. C. H. et al. Collective excitations in twisted bilayer graphene close to the magic angle. Preprint at <https://arxiv.org/abs/1910.07893> (2019).
65. Halbertal, D. et al. Moiré metrology of energy landscapes in van der Waals heterostructures. *Nat. Commun.* **12**, 242 (2021).
66. McGilly, L. J. et al. Visualization of moiré superlattices. *Nat. Nanotechnol.* **15**, 580–584 (2020).
67. McLeod, A. S. et al. Multi-messenger nanoprobe of hidden magnetism in a strained manganite. *Nat. Mater.* **19**, 397–404 (2020).
68. Xi, X. et al. Ising pairing in superconducting NbSe<sub>2</sub> atomic layers. *Nat. Phys.* **12**, 139–143 (2016).
69. Soriano, D., Katsnelson, M. I. & Fernández-Rossier, J. Magnetic two-dimensional chromium trihalides: a theoretical perspective. *Nano Lett.* **20**, 6225–6234 (2020).
70. Hejazi, K., Luo, Z.-X. & Balents, L. Noncollinear phases in moiré magnets. *Proc. Natl Acad. Sci. USA* **117**, 10721–10726 (2020).
71. Wang, C., Gao, Y., Lv, H., Xu, X. & Xiao, D. Stacking domain wall magnons in twisted van der Waals magnets. *Phys. Rev. Lett.* **125**, 247201 (2020).
72. Banerjee, A. et al. Proximate Kitaev quantum spin liquid behaviour in a honeycomb magnet. *Nat. Mater.* **15**, 733–740 (2016).
73. Chen, Y. et al. Strong correlations and orbital texture in single-layer 1T-TaSe<sub>2</sub>. *Nat. Phys.* **16**, 218–224 (2020).
74. Qian, X., Liu, J., Fu, L. & Li, J. Quantum spin Hall effect in two-dimensional transition metal dichalcogenides. *Science* **346**, 1344–1347 (2014).
75. Balents, L. Spin liquids in frustrated magnets. *Nature* **464**, 199–208 (2010).
76. Island, J. et al. Spin-orbit-driven band inversion in bilayer graphene by the van der Waals proximity effect. *Nature* **571**, 85–89 (2019).
77. Fu, L. & Kane, C. L. Superconducting proximity effect and Majorana fermions at the surface of a topological insulator. *Phys. Rev. Lett.* **100**, 096407 (2008).
78. Lutchyn, R. M., Sau, J. D. & Das Sarma, S. Majorana fermions and a topological phase transition in semiconductor-superconductor heterostructures. *Phys. Rev. Lett.* **105**, 077001 (2010).
79. Oreg, Y., Refael, G. & von Oppen, F. Helical liquids and Majorana bound states in quantum wires. *Phys. Rev. Lett.* **105**, 177002 (2010).
80. Liu, F. et al. Disassembling 2D van der Waals crystals into macroscopic monolayers and reassembling into artificial lattices. *Science* **367**, 903–906 (2020).
81. Darancet, P., Millis, A. J. & Marianetti, C. A. Three-dimensional metallic and two-dimensional insulating behavior in octahedral tantalum dichalcogenides. *Phys. Rev. B* **90**, 045134 (2014).
82. Tang, F. et al. Three-dimensional quantum Hall effect and metal-insulator transition in ZrTe<sub>5</sub>. *Nature* **569**, 537–541 (2019).
83. Wu, F., Zhang, R.-X. & Sarma, S. D. Three-dimensional topological twistrionics. *Phys. Rev. Res.* **2**, 022010 (2020).
84. Ruggenthaler, M., Tancogne-Dejean, N., Flick, J., Appel, H. & Rubio, A. From a quantum-electrodynamical light-matter description to novel spectroscopies. *Nat. Rev. Chem.* **2**, 0118 (2018).
85. Hübener, H. et al. Engineering quantum materials with chiral optical cavities. *Nat. Mater.* <https://doi.org/10.1038/s41563-020-00801-7> (2020).
86. Mazza, G. & Georges, A. Superradiant quantum materials. *Phys. Rev. Lett.* **122**, 017401 (2019).



87. Sentef, M. A., Ruggenthaler, M. & Rubio, A. Cavity quantum-electrodynamical polaritonically enhanced electron–phonon coupling and its influence on superconductivity. *Sci. Adv.* **4**, 6969 (2018).
88. Curtis, J. B., Raines, Z. M., Allocca, A. A., Hafezi, M. & Galitski, V. M. Cavity quantum Eliashberg enhancement of superconductivity. *Phys. Rev. Lett.* **122**, 167002 (2018).
89. Schlawin, F., Cavalleri, A. & Jaksch, D. Cavity-mediated electron–photon superconductivity. *Phys. Rev. Lett.* **122**, 133602 (2019).
90. Thomas, A. et al. Exploring superconductivity under strong coupling with the vacuum electromagnetic field. Preprint at <https://arxiv.org/abs/1911.01459> (2019).
91. Shimazaki, Y. et al. Signatures of tunable superconductivity in a trilayer graphene moiré superlattice. *Nature* **572**, 215–219 (2019).
92. Paravicini-Bagliani, G. L. et al. Magneto-transport controlled by Landau polariton states. *Nat. Phys.* **15**, 186–190 (2018).
93. Li, X. et al. Observation of Dicke cooperativity in magnetic interactions. *Science* **361**, 794–797 (2018).
94. Sun, Z. et al. Optical control of room-temperature valley polaritons. *Nat. Photon.* **11**, 491–496 (2017).
95. Latini, S., Ronca, E., Giovannini, U. D., Hübener, H. & Rubio, A. Cavity control of excitons in two dimensional materials. *Nano Lett.* **19**, 3473–3479 (2019).
96. Hong, X. et al. Ultrafast charge transfer in atomically thin  $\text{MoS}_2/\text{WS}_2$  heterostructures. *Nat. Nanotechnol.* **9**, 682–686 (2014).
97. Zhu, H. et al. Interfacial charge transfer circumventing momentum mismatch at two-dimensional van der Waals heterojunctions. *Nano Lett.* **17**, 3591–3598 (2017).
98. Jauregui, L. A. et al. Electrical control of interlayer exciton dynamics in atomically thin heterostructures. *Science* **366**, 870–875 (2019).
99. Wang, Z. et al. Evidence of high-temperature exciton condensation in two-dimensional atomic double layers. *Nature* **574**, 76–80 (2019).
100. Jin, C. et al. Observation of moiré excitons in  $\text{WSe}_2/\text{WS}_2$  heterostructure superlattices. *Nature* **567**, 76–80 (2019).
101. Shimazaki, Y. et al. Strongly correlated electrons and hybrid excitons in a moiré heterostructure. *Nature* **580**, 472–477 (2020).
102. Tran, K. et al. Evidence for moiré excitons in van der Waals heterostructures. *Nature* **567**, 71–75 (2019).
103. Alexeev, E. M. et al. Resonantly hybridized excitons in moiré superlattices in van der Waals heterostructures. *Nature* **567**, 81–86 (2019).
104. Wu, F., Lovorn, T. & MacDonald, A. H. Topological exciton bands in moiré heterojunctions. *Phys. Rev. Lett.* **118**, 147401 (2017).
105. Bai, Y. et al. Excitons in strain-induced one-dimensional moiré potentials at transition metal dichalcogenide heterojunctions. *Nat. Mater.* **19**, 1068–1073 (2020).
106. Tong, Q., Liu, F., Xiao, J. & Yao, W. Skyrmions in the moiré of van der Waals 2D magnets. *Nano Lett.* **18**, 7194–7199 (2018).
107. Po, H. C., Zou, L., Vishwanath, A. & Senthil, T. Origin of Mott insulating behavior and superconductivity in twisted bilayer graphene. *Phys. Rev. X* **8**, 031089 (2018).
108. Lee, J. Y. et al. Theory of correlated insulating behaviour and spin-triplet superconductivity in twisted double bilayer graphene. *Nat. Commun.* **10**, 5333 (2019).
109. Efimkin, D. K. & MacDonald, A. H. Helical network model for twisted bilayer graphene. *Phys. Rev. B* **98**, 035404 (2018).
110. An, L. et al. Interaction effects and superconductivity signatures in twisted double-bilayer  $\text{WSe}_2$ . *Nanoscale Horiz.* **5**, 1309–1316 (2020).
111. Wang, J. et al. Diffusivity reveals three distinct phases of interlayer excitons in  $\text{MoSe}_2/\text{WSe}_2$  heterobilayers. Preprint at <https://arxiv.org/abs/2001.03812> (2020).
112. Spanton, E. M. et al. Observation of fractional Chern insulators in a van der Waals heterostructure. *Science* **360**, 62–66 (2018).
113. Arora, H. S. et al. Superconductivity in metallic twisted bilayer graphene stabilized by  $\text{WSe}_2$ . *Nature* **583**, 379–384 (2020).

## Acknowledgements

This work is supported by the European Research Council (ERC-2015-AdG-694097) and Grupos Consolidados (IT1249-19). M.C., A.J.M., A.G. and A.R. are supported by the Flatiron Institute, a division of the Simons Foundation. We acknowledge funding by the Deutsche Forschungsgemeinschaft (DFG) under Germany's Excellence Strategy - Cluster of Excellence Matter and Light for Quantum Computing (ML4Q) EXC 2004/1 - 390534769, within the Priority Program SPP 2244 '2DMP' and Advanced Imaging of Matter (AIM) EXC 2056 - 390715994 and funding by the Deutsche Forschungsgemeinschaft (DFG) through RTG 1995 and RTG 2247. Support by the Max Planck Institute - New York City Center for Non-Equilibrium Quantum Phenomena is acknowledged. Work at Columbia is supported as part of Programmable Quantum Materials, an Energy Frontier Research Center funded by the US Department of Energy (DOE), Office of Science, Basic Energy Sciences (BES), under award DE-SC0019443. D.N.B. is a Vannevar Bush Faculty Fellow ONR-VB: N00014-19-1-2630. A.N.P. acknowledges support from the Air Force Office of Scientific Research via grant FA9550-16-1-0601.

## Competing interests

The authors declare no competing interests.

## Additional information

Correspondence should be addressed to D.M.K., A.N.P. or A.R.

Reprints and permissions information is available at [www.nature.com/reprints](http://www.nature.com/reprints).

**Publisher's note** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

© Springer Nature Limited 2021