

# A Multi-Species, Multi-Fluid Model for Simulating Plasma Interpenetration

SIAM Conference on Computational Science and Engineering

Debojyoti Ghosh<sup>1</sup>, Thomas Chapman<sup>2</sup>, Richard Berger<sup>2</sup>,  
Andris Dimits<sup>3</sup>, Jeffrey W. Banks<sup>4</sup>



<sup>1</sup>Center for Applied Scientific Computing, LLNL

<sup>2</sup>Weapons and Complex Integration, LLNL

<sup>3</sup>Physics and Life Sciences, LLNL

<sup>4</sup>Mathematical Sciences, Rensselaer Polytechnic Institute

February 25 – March 1, 2019, Spokane, WA



LLNL-PRES-768123

This work was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under contract DE-AC52-07NA27344. Lawrence Livermore National Security, LLC

Lawrence Livermore  
National Laboratory

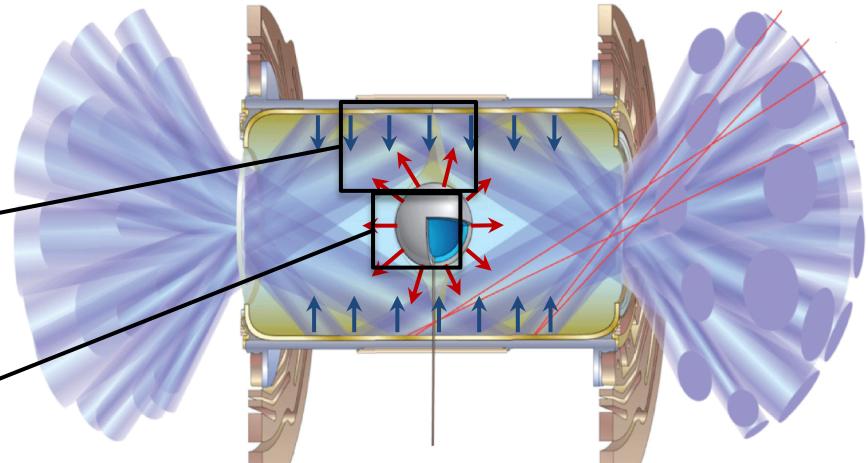
# Background and Motivation

## Inertial Confinement Fusion: Colliding plasmas from hohlraum wall and capsule

Interpenetration of plasma flows from capsule and hohlraum wall

- o Large range of  $Z$ :  $2 \leq Z \leq 60$
- o Supersonic flows ( $\Delta u \approx 10^8$  cm/s)

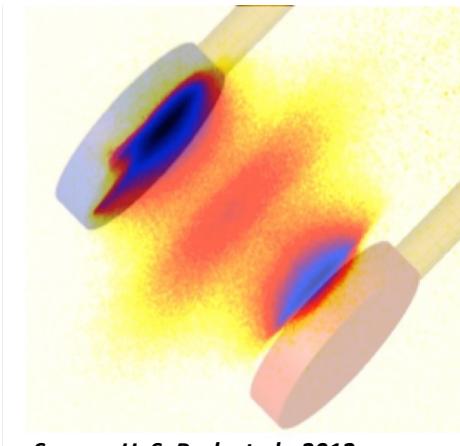
Species separation inside target capsule



Source: <https://csdl-images.computer.org/mags/cs/2014/06/figures/mcs20140600421.gif>

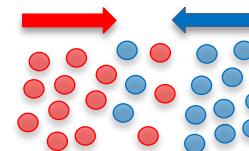
## High Energy Density Physics (HEDP) Experiments

Carbon plasma streams ablating off paddles hit by laser beams and colliding with each other

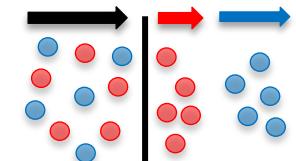


Source: H. S. Park et al., 2012

## Multifluid phenomena that we want to model



Interpenetrating plasmas



Plasma species separation

# Current simulation tools are *not sufficiently versatile*

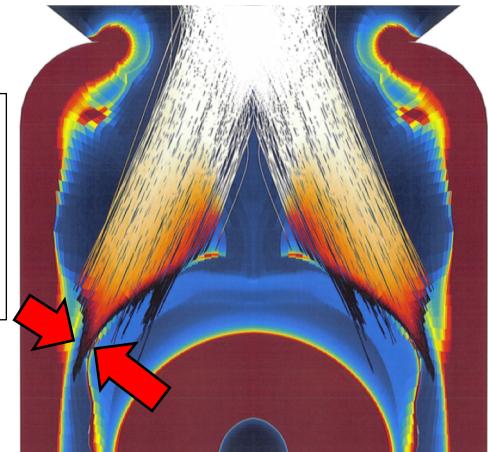
## Single-Fluid Multi-species Hydrodynamic Solvers

- Example: *HYDRA, LASNEX*
- Single velocity field insufficient to model multiple inter-penetrating fluids
- Unphysical shocks

Lack key physics

Simulation of plasma dynamics in hohlraum using *HYDRA*

Density pile-up predicted when plasma streams collide



## Collisional Kinetic Solvers

- Example: *LOKI, OSIRIS, PSC*
- High computational cost to simulate small volumes
- Impractical for experimental scales

Too Expensive

Current workarounds: species diffusion models

Needs: physics beyond single-fluid theory – multi-fluid, multi-species, with local kinetic effects

# Governing Equations: We solve the inviscid Euler equations for each ion species

$$\frac{\partial \rho_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha) = 0 \quad \boxed{\alpha = 1, \dots, n_s}$$

$$\frac{\partial \rho_\alpha \mathbf{u}_\alpha}{\partial t} + \nabla \cdot (P_\alpha + \rho_\alpha \mathbf{u}_\alpha \otimes \mathbf{u}_\alpha) = -Z_\alpha e n_\alpha \nabla \phi + \sum_{\beta \neq \alpha} \mathbf{R}_{\alpha,\beta} \quad \boxed{\text{Interaction between species}}$$

$$\frac{\partial \mathcal{E}_\alpha}{\partial t} + \nabla \cdot [(\mathcal{E}_\alpha + P_\alpha) \mathbf{u}_\alpha] = -Z_\alpha e n_\alpha \mathbf{u}_\alpha \cdot \nabla \phi + \sum_{\beta \neq \alpha} (\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_\alpha + Q_{\alpha,\beta})$$

**Assuming quasineutral, isothermal electrons\***

$$\nabla \phi = \frac{T_e}{n_e} \nabla n_e + \frac{1}{n_e} \sum_\alpha R_{e,\alpha}$$

*Electron momentum equation neglecting inertia terms and assuming*

$$P_e = n_e T_e$$

*Frictional drag*

$$\mathbf{R}_{\alpha,\beta} = m_\alpha n_\alpha \nu_{\alpha,\beta} (\mathbf{u}_\beta - \mathbf{u}_\alpha)$$

*Frictional heating and thermal equilibration*

$$Q_{\alpha,\beta} = Q_{\alpha,\beta}^{\text{fric}} + Q_{\alpha,\beta}^{\text{eq}}$$

$$Q_{\alpha,\beta}^{\text{fric}} = m_{\alpha,\beta} n_\alpha \nu_{\alpha,\beta} (\mathbf{u}_\beta - \mathbf{u}_\alpha)^2$$

$$Q_{\alpha,\beta}^{\text{eq}} = -3m_\alpha n_\alpha \frac{\nu_{\alpha,\beta}}{m_\alpha + m_\beta} (T_\alpha - T_\beta)$$

# Reformulated Governing Equations

*Ion Euler equations with isothermal, quasineutral  $e^-$*

Advective nature of electrostatic force



- Included **electron pressure** on LHS with hydrodynamic pressure
- Derived the **eigenstructure** for **characteristic-based discretization**

Effect of discretization error in dense species on **dynamics of sparse species**



Reformulation of electrostatic source terms to *avoid sums/differences of terms of disparate scales*

$$\frac{\partial \rho_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha) = 0,$$

$$\frac{\partial \rho_\alpha \mathbf{u}_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha \otimes \mathbf{u}_\alpha + P_\alpha^*) = Z_\alpha T_e n_e \nabla \left( \frac{n_\alpha}{n_e} \right) + \frac{Z_\alpha n_\alpha}{n_e} \sum_\beta \mathbf{R}_{e,\beta} + \mathbf{R}_{\alpha,e} + \sum_{\beta \neq \alpha} \mathbf{R}_{\alpha,\beta},$$

$$\begin{aligned} \frac{\partial \mathcal{E}_\alpha}{\partial t} + \nabla \cdot \{(\mathcal{E}_\alpha + P_\alpha^*) \mathbf{u}_\alpha\} &= Z_\alpha T_e n_e \nabla \left( \frac{\mathbf{u}_\alpha n_\alpha}{n_e} \right) + \frac{Z_\alpha n_\alpha}{n_e} \sum_\beta \mathbf{u}_\alpha \cdot \mathbf{R}_{e,\beta} + \sum_{\beta \neq \alpha} (\mathbf{R}_{\alpha,\beta} \cdot \mathbf{u}_\alpha + Q_{\alpha,\beta}) \\ &\quad + \mathbf{R}_{\alpha,e} \cdot \mathbf{u}_\alpha + Q_{\alpha,e}^{\text{eq}}, \end{aligned}$$

where  $P_\alpha^* = P_\alpha + Z_\alpha T_e n_\alpha$

Electron pressure

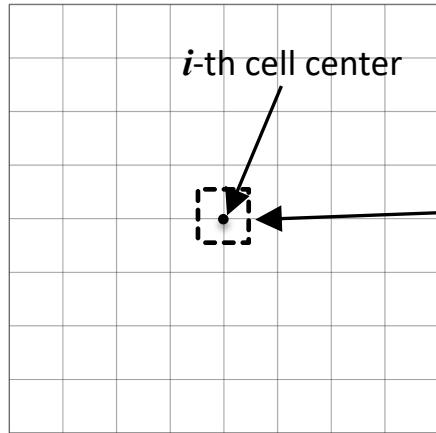
is the “augmented pressure” (hydro +  $e^-$ )

Wavespeeds (eigenvalues) :  $\mathbf{v}, \mathbf{v} \pm \sqrt{\frac{\gamma_\alpha P_\alpha^*}{\rho_\alpha}}$

# Summary of Numerical Method

## High-Order Conservative Finite-Difference/Finite-Volume Method

4<sup>th</sup> order finite-volume discretization (using the CHOMBO library) with AMR



Spatially-discretized ODE in time (integrated in time using 4<sup>th</sup> order Runge-Kutta method)

3D Domain  $\Omega \equiv \{ \mathbf{x} : 0 \leq \mathbf{x} \cdot \mathbf{e}_d \leq L_d, 1 \leq d \leq 3 \}$

discretized into computational cells

$$\omega_{\mathbf{i}} = \prod_{d=1}^3 \left[ \left( \mathbf{i} - \frac{1}{2} \mathbf{e}_d \right) h, \left( \mathbf{i} + \frac{1}{2} \mathbf{e}_d \right) h \right]$$

$\mathbf{i}$ : 3-dimensional integer index ( $i, j, k$ )  
 $h$ : grid spacing

$$\frac{\partial \bar{\mathbf{u}}_{\mathbf{i}}}{\partial t} = \frac{1}{h} \sum_{d=1}^3 \left( \langle \hat{\mathbf{F}}_{\mathbf{i} + \frac{1}{2} \mathbf{e}_d} \rangle - \langle \hat{\mathbf{F}}_{\mathbf{i} - \frac{1}{2} \mathbf{e}_d} \rangle \right)$$

Cell-averaged solution

Face-averaged fluxes

$$\mathbf{u} = \begin{bmatrix} \vdots \\ \rho_{\alpha} \\ \rho_{\alpha} \mathbf{v}_{\alpha} \\ \mathcal{E}_{\alpha} \\ \vdots \end{bmatrix}$$

Strong shocks and gradients  
O(1) to O(1e-14)



- Characteristic-based discretization
- 5<sup>th</sup>-order WENO scheme with Monotonicity-Preserving limiting

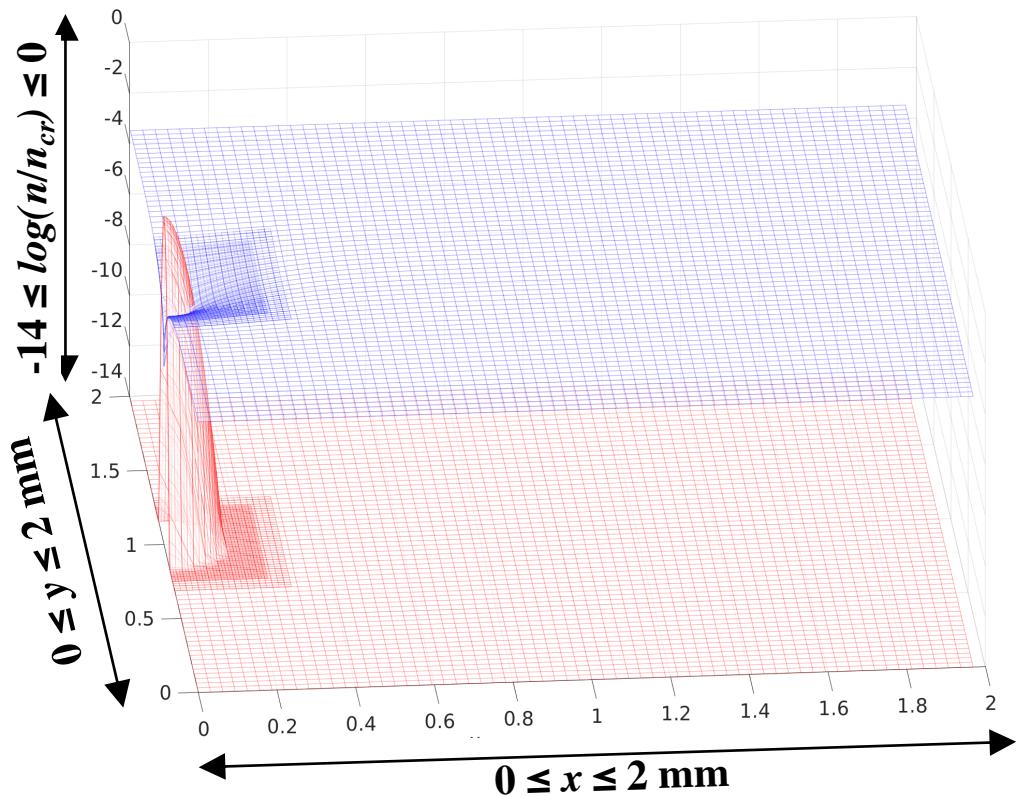
# Example: Single Species Expansion into Gas Fill with AMR – *Initial Setup*

## Expansion of a **carbon blob** in the presence of **helium gas fill** (2D)

- **Initial solution:** a carbon species piled up on one end (*Gaussian blob* density); gas fill present in the space everywhere else.
- **Boundary conditions:** Solid wall BCs along  $x$  and  $y$

### Reference quantities:

Mass: *proton mass* ( $1.6730\text{e-}24$  g);  
Number density:  $n_{crit}$  ( $9.0320\text{e+}21 \text{ cm}^{-3}$ );  
Length: 1 mm; Time:  $3.2314\text{e-}09$  s;  
Temperature: 1 keV ( $1.6022\text{e-}09$  ergs)



# Example: Single Species Expansion into Gas Fill with AMR – Solution Evolution

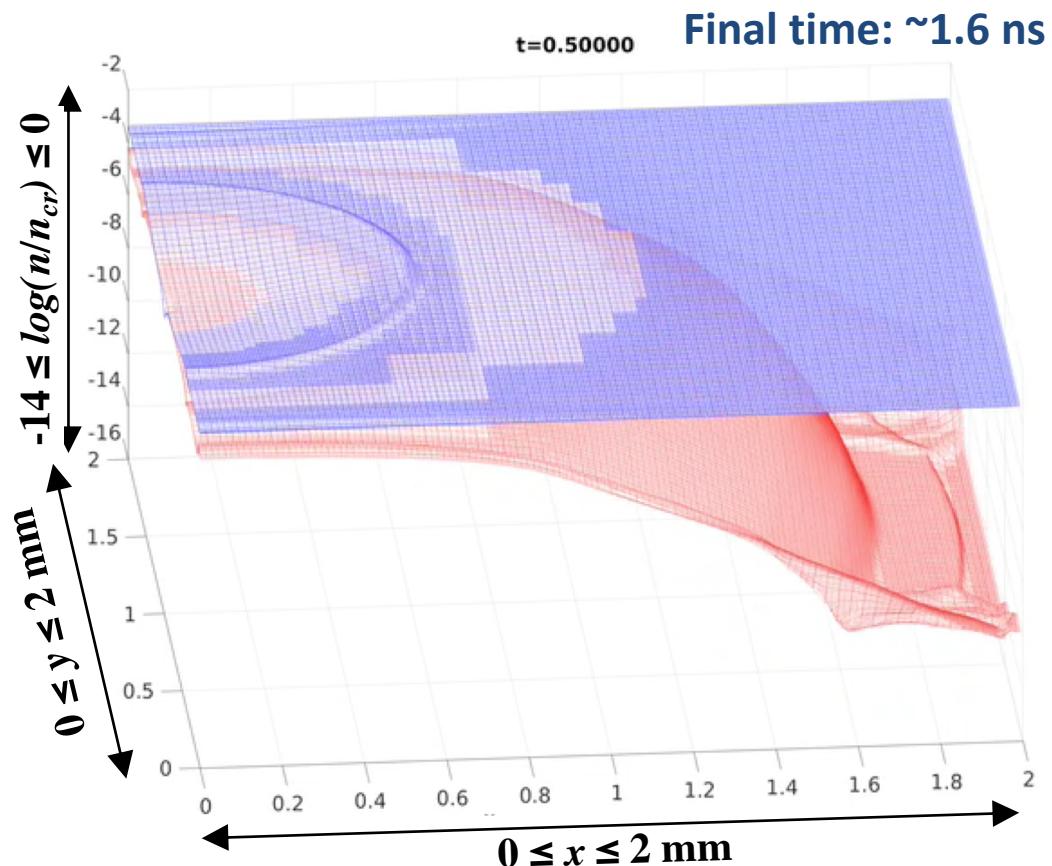
## Expansion of a **carbon blob** in the presence of **helium gas fill** (2D)

- **Initial solution:** a carbon species piled up on one end (*Gaussian blob* density); gas fill present in the space everywhere else.
- **Boundary conditions:** Solid wall BCs along  $x$  and  $y$

**AMR:** Refined mesh adaptively generated in regions of high gradients

### Reference quantities:

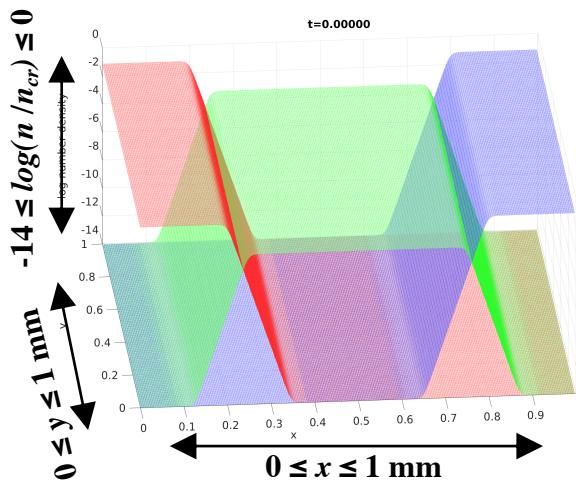
Mass: *proton mass* ( $1.6730\text{e-}24$  g);  
Number density:  $n_{crit}$  ( $9.0320\text{e+}21 \text{ cm}^{-3}$ );  
Length: 1 mm; Time:  $3.2314\text{e-}09$  s;  
Temperature: 1 keV ( $1.6022\text{e-}09$  ergs)



# Example: Two Species Interpenetration with Gas Fill Problem Setup

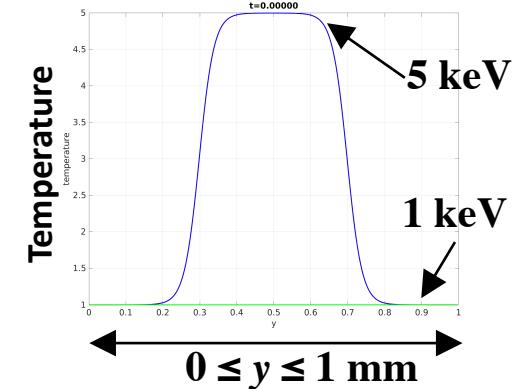
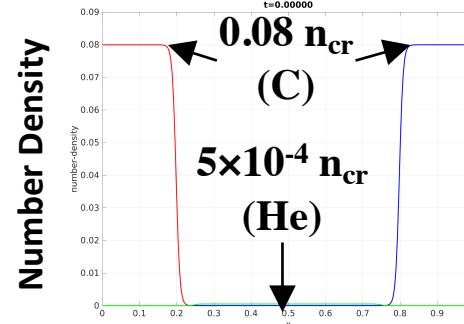
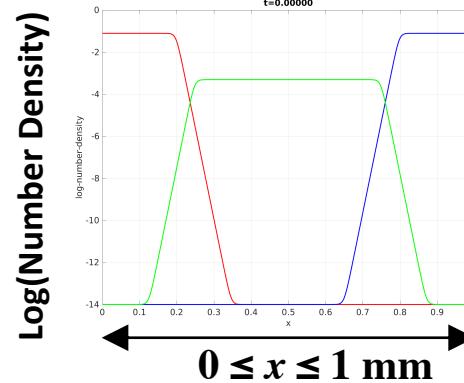
Interpenetration of **carbon** and **carbon** streams in the presence of **helium** gas fill (2D)

- Initial solution: two species piled up on either end (*smoothed slab density*); gas fill present in the space in between.
- Temperature variation along  $y$  – the plasmas are hotter in the center of the domain



Boundary conditions:

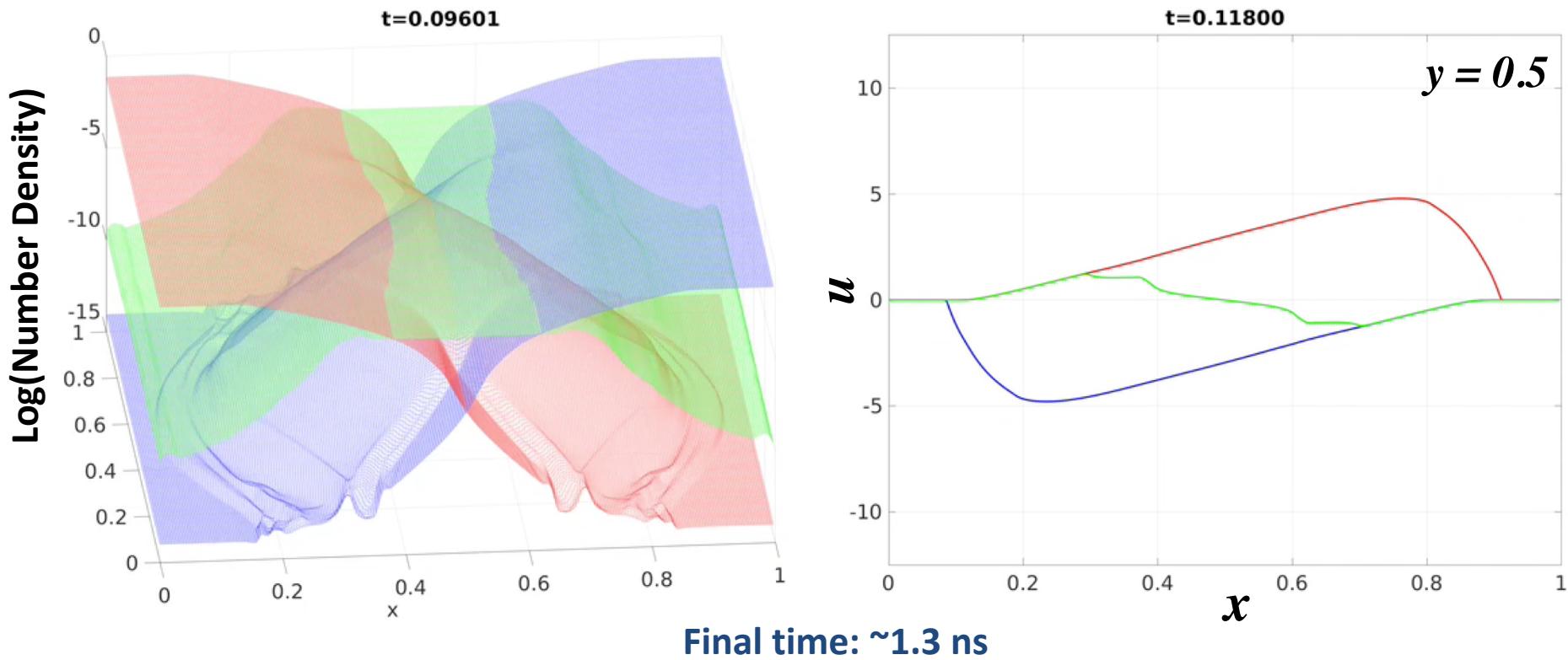
- Solid wall BCs along  $x$
- Periodic along  $y$



Reference quantities:

- Mass: proton mass ( $1.6730\text{e-}24 \text{ g}$ )
- Number density:  $n_{crit}$  ( $9.0320\text{e+}21 \text{ cm}^{-3}$ )
- Length: 1 mm
- Temperature: 1 keV ( $1.6022\text{e-}09 \text{ ergs}$ )

# Example: Two Species Interpenetration with Gas Fill



- **Species interaction** prevents one species from reaching the other end of the domain along  $x$
- The fill gas is pushed towards the center of the domain by the carbon streams.

## Reference quantities:

Number density:  $n_{crit}$  ( $9.0320e+21 \text{ cm}^{-3}$ );  
Length: 1 mm; Time:  $3.2314e-09$  s;  
Velocity:  $3.0946e+07 \text{ cm/s}$

# Stiffness of Collisional Terms for High-Z Species

## Ion-ion collisional interaction term

**interaction term:** Frictional force & heating, and thermal equilibration of *ion species 1* due to *ion species 2*

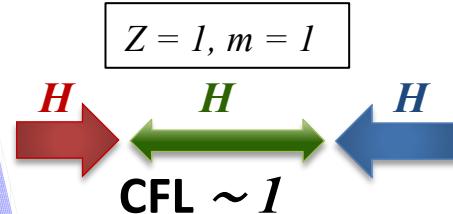
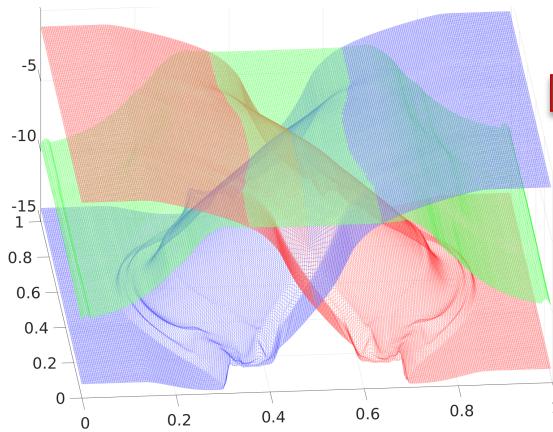
where

$$\nu_{12} = \left( \frac{4\sqrt{2\pi}}{3} \right) \frac{Z_1^2 Z_2^2 n_2 \Lambda}{m_1 m_{12}} \left[ r (u_1 - u_2)^2 + \frac{T_1}{m_1} + \frac{T_2}{m_2} \right]^{-\frac{3}{2}} \left( \frac{e^4 n_0 x_0}{T_0^2} \right)$$

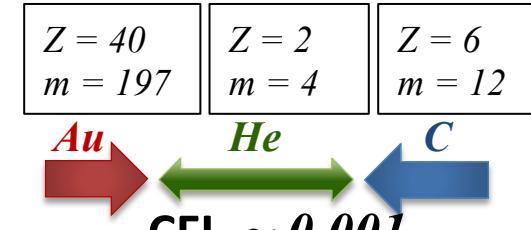
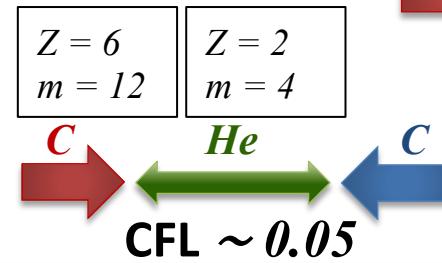
$$\left[ \begin{array}{c} 0 \\ m_1 n_1 \nu_{12} (u_2 - u_1) \\ m_1 n_1 \nu_{12} (u_2 - u_1) u_1 + m_{12} n_1 \nu_{12} (u_2 - u_1)^2 + 3 \frac{m_1 n_1 \nu_{12}}{m_1 + m_2} (T_2 - T_1) \end{array} \right]$$

Typically,  $n \propto (1/Z)$

$$\propto Z_1 Z_2 \sqrt{\frac{m_1 m_2}{m_1 + m_2}}$$



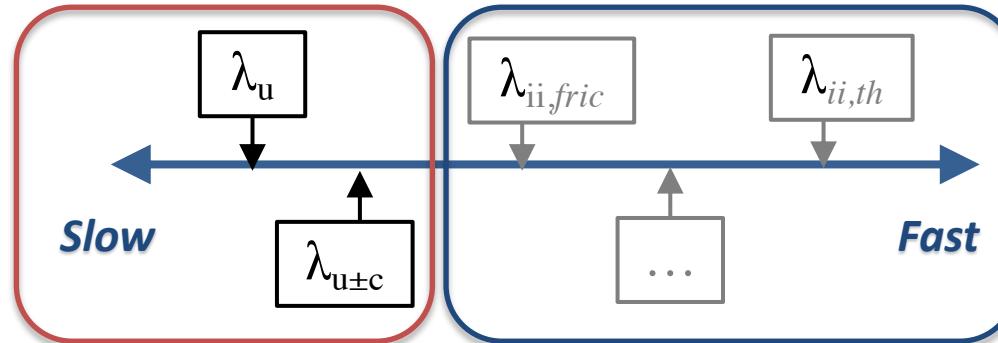
(CFL based on the hyperbolic term on the LHS)



# Implicit-Explicit (IMEX) Time Integration

Resolve scales of interest; Treat implicitly faster scales

Advective and acoustic time scales  
(nonstiff terms)



Explicit time integration    Implicit time integration

Collisional time scales – ion-ion and  $e^-$ -ion friction & thermal equilibration (stiff terms)

## ODE in time

Resulting from spatial discretization of PDE

$$\frac{d\mathbf{y}}{dt} = \mathcal{R}(\mathbf{y})$$

IMEX time integration: *partition RHS*

$$\mathcal{R}(\mathbf{y}) = \mathcal{R}_{\text{stiff}}(\mathbf{y}) + \mathcal{R}_{\text{nonstiff}}(\mathbf{y})$$

Linear stability constraint  
on time step

$$\Delta t \left( \lambda \left[ \frac{d\mathcal{R}_{\text{nonstiff}}(\mathbf{y})}{d\mathbf{y}} \right] \right) \in \{z : |R(z)| \leq 1\}$$

Time step constrained by eigenvalues (time scales) of *nonstiff component of RHS*

# Additive Runge-Kutta (ARK) Time Integrators

## Multistage, high-order, conservative IMEX methods

*Butcher tableaux representation*

$$\begin{array}{c|ccccc}
 0 & 0 & & & & \text{Explicit RK} \\
 c_2 & a_{21} & 0 & & & \\
 \vdots & \vdots & \ddots & & & \\
 c_s & a_{s1} & \cdots & a_{s,s-1} & 0 & \\
 \hline
 b_1 & \cdots & \cdots & & b_s &
 \end{array}
 \quad +
 \quad
 \begin{array}{c|ccccc}
 0 & 0 & & & & \text{DIRK} \\
 \tilde{c}_2 & \tilde{a}_{21} & \gamma & & & \\
 \vdots & \vdots & \ddots & & \gamma & \\
 \tilde{c}_s & \tilde{a}_{s1} & \cdots & \tilde{a}_{s,s-1} & \gamma & \\
 \hline
 b_1 & \cdots & \cdots & & b_s &
 \end{array}$$

**Time step:** From  $t_n$  to  $t_{n+1} = t_n + \Delta t$        $s \rightarrow$  number of stages

*Stage solutions*

$$\mathbf{y}^{(i)} = \mathbf{y}_n + \Delta t \sum_{j=1}^{i-1} a_{ij} \mathcal{R}_{\text{nonstiff}} \left( \mathbf{y}^{(j)} \right) + \Delta t \sum_{j=1}^i \tilde{a}_{ij} \mathcal{R}_{\text{stiff}} \left( \mathbf{y}^{(j)} \right), \quad i = 1, \dots, s$$

$$\mathbf{y}_{n+1} = \mathbf{y}_n + \Delta t \sum_{i=1}^s b_i \mathcal{R} \left( \mathbf{y}^{(i)} \right) \quad \textit{Step completion}$$

Kennedy & Carpenter, J. Comput. Phys., 2003



# Implicit Stage Solution

Requires solving nonlinear system of equations

Rearranging the stage solution expression:

$$\underbrace{\frac{1}{\Delta t \tilde{a}_{ii}} \mathbf{y}^{(i)} - \mathcal{R}_{\text{stiff}}(\mathbf{y}^{(i)}) - \left[ \mathbf{y}_n + \Delta t \sum_{j=1}^{i-1} \left\{ a_{ij} \mathcal{R}_{\text{nonstiff}}(\mathbf{y}^{(j)}) + \tilde{a}_{ij} \mathcal{R}_{\text{stiff}}(\mathbf{y}^{(j)}) \right\} \right]}_{\mathcal{F}(y) = 0} = 0$$

**Jacobian-free Newton-Krylov** method (*Knoll & Keyes, J. Comput. Phys., 2004*):

Initial guess:  $y_0 \equiv \mathbf{y}_0^{(i)} = \mathbf{y}^{(i-1)}$

Newton update:  $y_{k+1} = y_k - \mathcal{J}(y_k)^{-1} \mathcal{F}(y_k)$

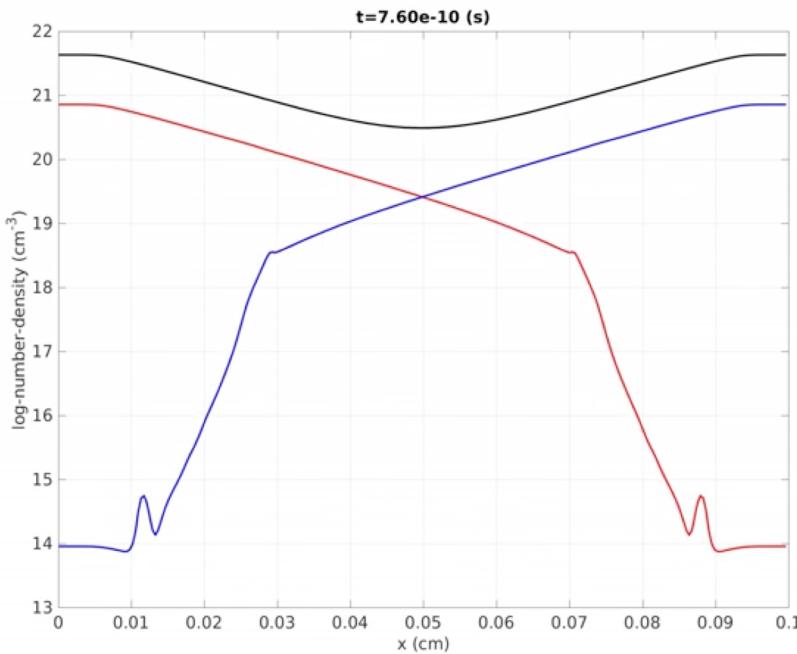
GMRES solver  
(preconditioned)  
 $\mathcal{J}(y_k) \Delta y = \mathcal{F}(y_k)$

Action of the Jacobian on a vector approximated by *directional derivative*

$$\mathcal{J}(y_k) x = \left. \frac{d\mathcal{F}(y)}{dy} \right|_{y_k} x \approx \frac{1}{\epsilon} [\mathcal{F}(y_k + \epsilon x) - \mathcal{F}(y_k)]$$

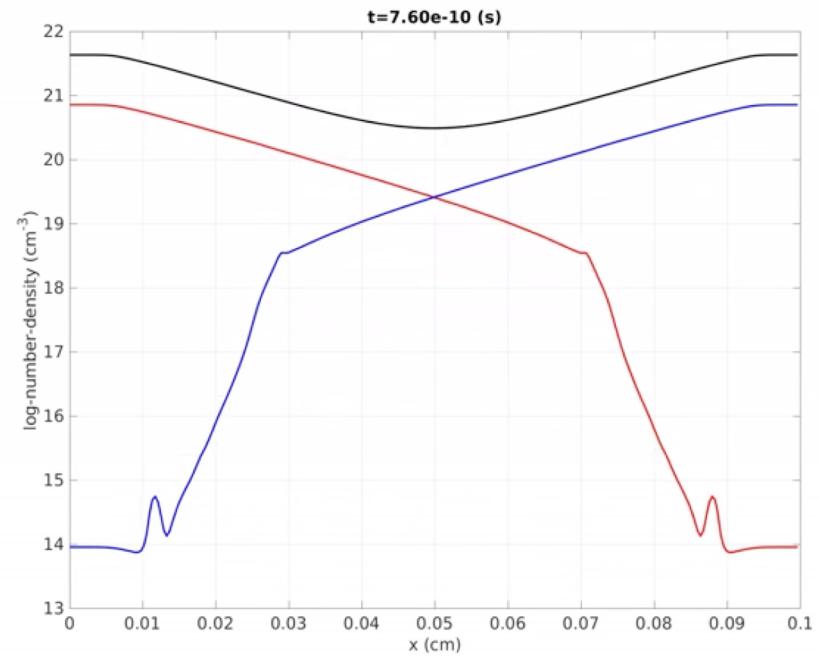
# Work-in-Progress: Preliminary Example

## 1D Carbon–Carbon Interpenetration



**ARK4 (6-stage, 4<sup>th</sup> order, 5 implicit stages)**  
**CFL = 0.5**

**$\sim 1200$  time steps, 1 hr 40 mins**



**RK4 (4-stage, 4<sup>th</sup> order)**  
**CFL = 0.05**  
 **$\sim 12,000$  time steps, 7 hrs 15 mins**

No preconditioner implemented yet; however, note *implicit term is block diagonal (no spatial derivatives)*

*Good agreement between IMEX and explicit solutions; need to verify convergence*

# Conclusions and Future Work

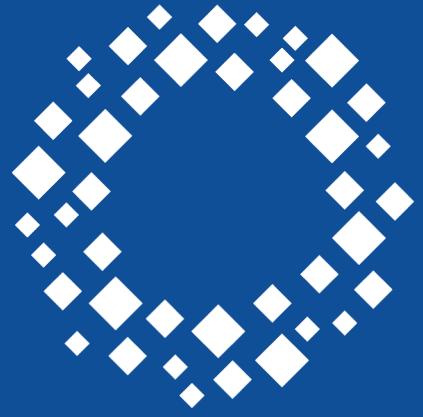
## Summary

### EUCLID: EUlerian Code for pLasma Interpenetration Dynamics

- Developed a **3D, parallel, AMR-capable multifluid flow solver**
- Implemented the *quasineutral, isothermal electron model* as a computationally tractable electron model for our target applications.
- *Verified EUCLID for accuracy and convergence* (benchmark cases, manufactured solutions)
- *Simulated flows motivated by laboratory astrophysics experiments and ICF hohlraums.*

## Current and Future Work

- Implementation of **IMEX time integrators**
  - Newton's method convergence difficulties for large time steps
  - Implement an efficient preconditioner
- Conduct **simulations of plasma interpenetration experiments** (e.g. Ross *et al.*, 2013, Le Pape *et al.*, ongoing)
- Investigate *higher-fidelity electron models*, for example, adding an electron energy equation.
- Add *source terms to energy equations to simulate heating*



# CASC

Center for Applied  
Scientific Computing



**Lawrence Livermore  
National Laboratory**

#### **Disclaimer**

This document was prepared as an account of work sponsored by an agency of the United States government. Neither the United States government nor Lawrence Livermore National Security, LLC, nor any of their employees makes any warranty, expressed or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States government or Lawrence Livermore National Security, LLC. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States government or Lawrence Livermore National Security, LLC, and shall not be used for advertising or product endorsement purposes.

# Multifluid vs. Single Fluid Simulations - How do the solutions differ?

Interpenetration of **two hydrogen streams** in the presence of **hydrogen gas fill** (2D)

