



西安电子科技大学
XIDIAN UNIVERSITY

B39HF High Frequency Circuits

Lecture 1 Introduction

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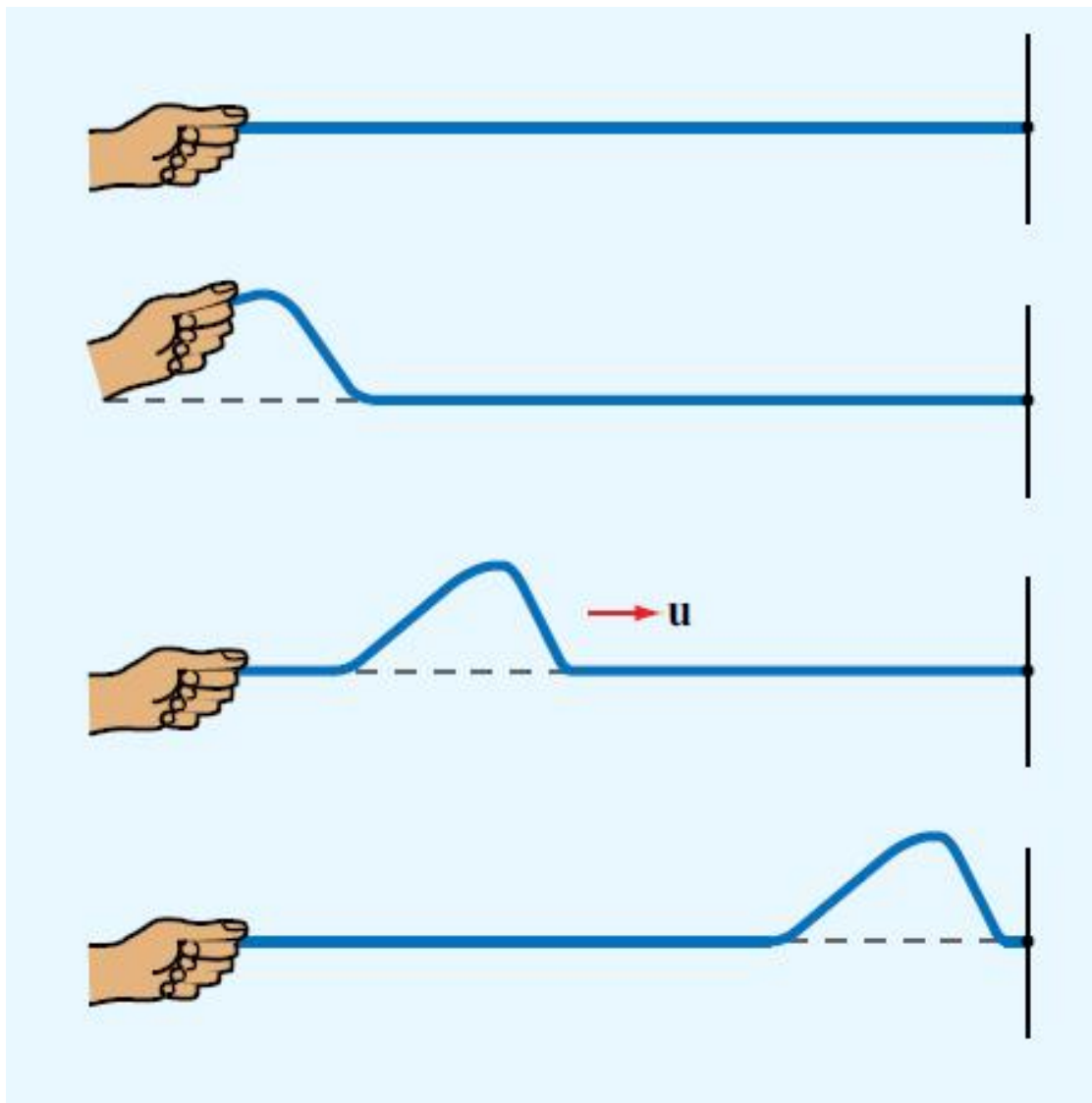
- Development of Electromagnetics
- Dimensions, Units, and Notation
- The Nature of Electromagnetism
- Traveling Waves
- The Electromagnetic Spectrum
- Review of Complex Numbers
- Review of Phasors



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A one-dimensional wave traveling on a string





- **Two-dimensional wave:** propagates out across a surface, like the ripples on a pond, and its disturbance can be described by two space variables.
- **Three-dimensional wave:** propagates through a volume and its disturbance may be a function of all three space variables. may take on many different shapes; include *plane waves*, *cylindrical waves*, and *spherical waves*.

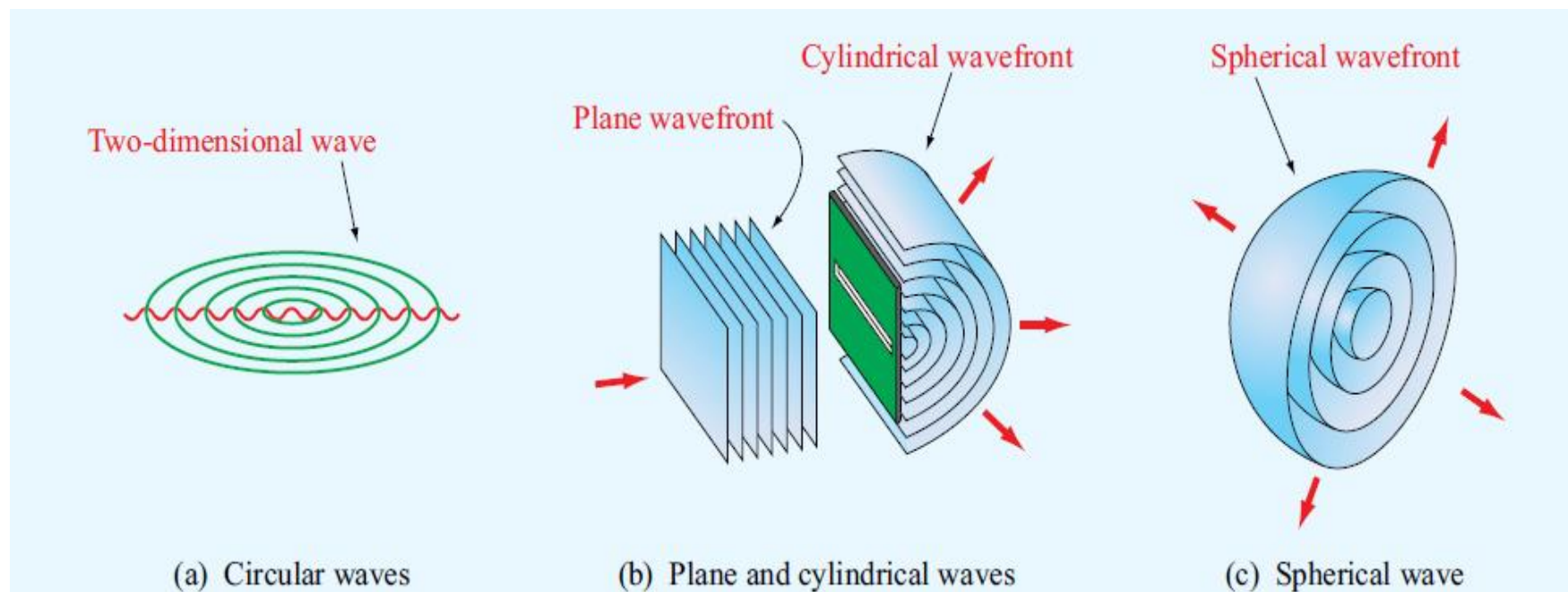
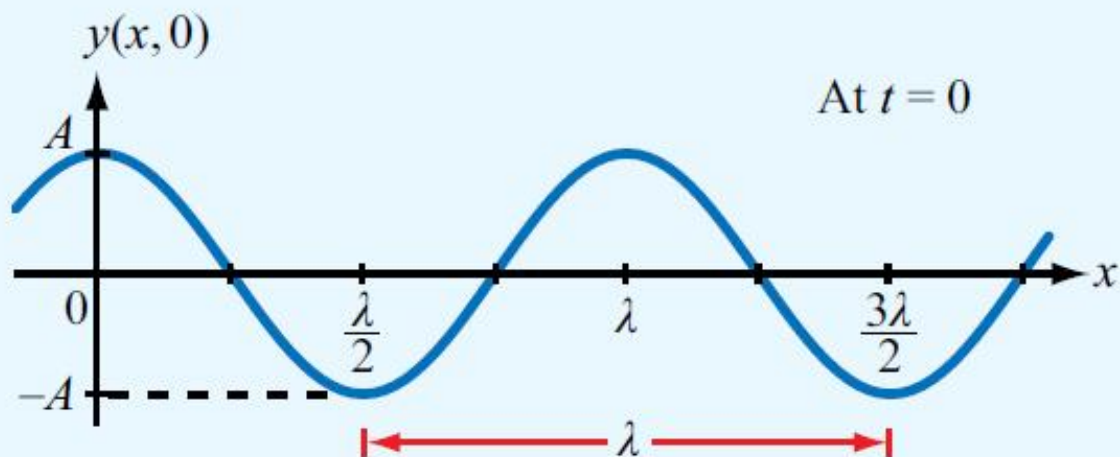
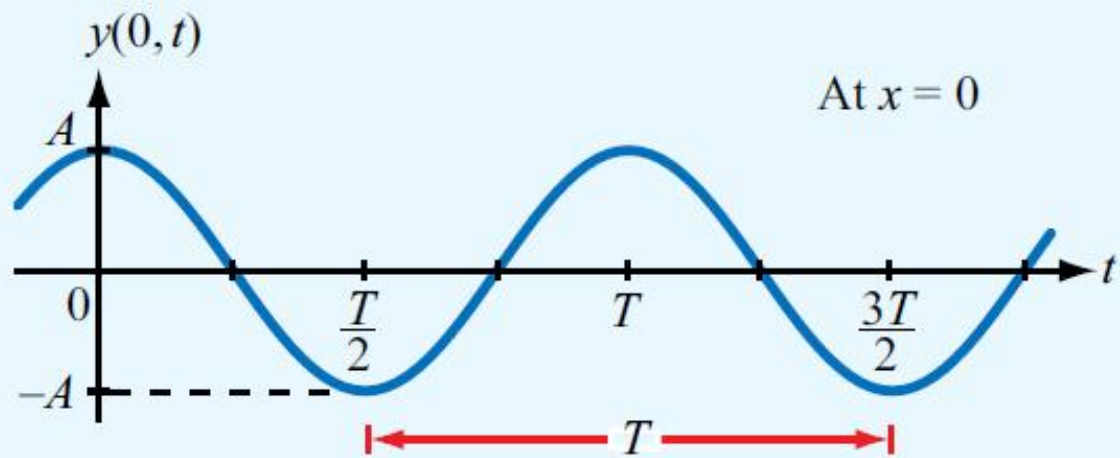


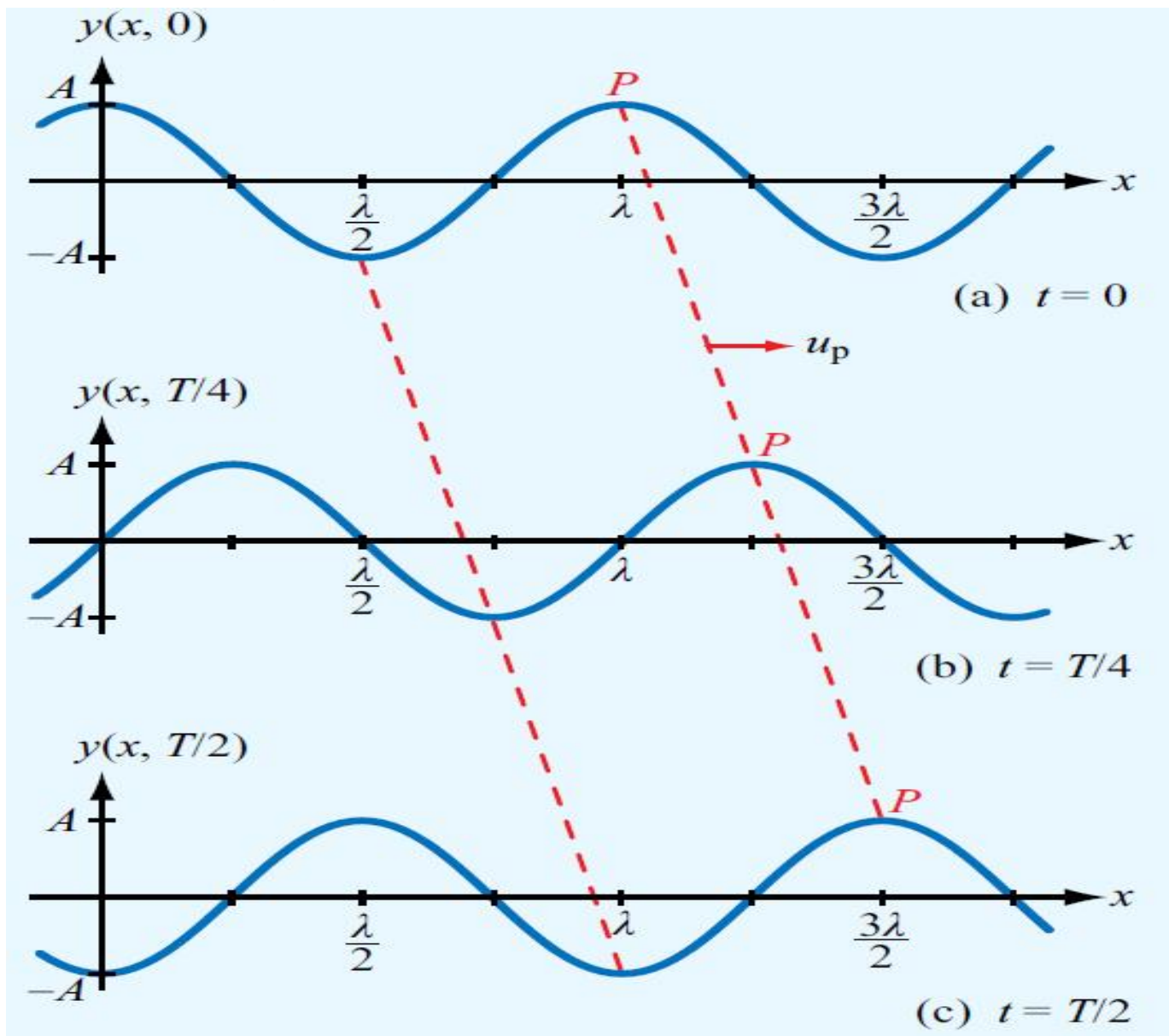
Figure 1-11 Examples of two-dimensional and three-dimensional waves: (a) circular waves on a pond, (b) a plane light wave exciting a cylindrical light wave through the use of a long narrow slit in an opaque screen, and (c) a sliced section of a spherical wave.



(a) $y(x, t)$ versus x at $t = 0$



(b) $y(x, t)$ versus t at $x = 0$





Sinusoidal Waves in a Lossless Medium

Phase of the Wave $\longrightarrow \phi(x, t)$

$$y(x, t) = A \cos\left(\frac{2\pi t}{T} - \frac{2\pi x}{\lambda} + \phi_o\right)$$

Amplitude $\nearrow A$

Time Period $\nearrow T$

Spatial Wavelength $\nearrow \lambda$

Reference Phase $\nearrow \phi_o$

Let us first analyze the simple case when $\phi_o = 0$:

$$y(x, t) = A \cos\left(\frac{2\pi t}{T} - \frac{2\pi x}{\lambda}\right) \quad (\text{m})$$



At the peaks of the wave pattern, the phase $\phi(x, t)$ is equal to zero or multiples of 2π radians. Thus,

$$\phi(x, t) = \frac{2\pi t}{T} - \frac{2\pi x}{\lambda} = 2n\pi, \quad n = 0, 1, 2, \dots$$

Had we chosen any other fixed height of the wave, say y_0 , and monitored its movement as a function of t and x , this again would have been equivalent to setting the phase $\phi(x, t)$ constant such that

$$y(x, t) = y_0 = A \cos \left(\frac{2\pi t}{T} - \frac{2\pi x}{\lambda} \right)$$
$$\frac{2\pi t}{T} - \frac{2\pi x}{\lambda} = \cos^{-1} \left(\frac{y_0}{A} \right) = \text{constant}$$
$$\downarrow$$
$$\frac{2\pi}{T} - \frac{2\pi}{\lambda} \frac{dx}{dt} = 0$$

which gives the *phase velocity* u_p as

$$u_p = \frac{dx}{dt} = \frac{\lambda}{T} \quad (\text{m/s})$$



Phase Velocity

$$u_p = \frac{\lambda}{T}$$

A traveling wave is characterized by a spatial wavelength λ , a time period T and a phase velocity $u_p = \lambda/T$.



phase velocity

$$u_p = \frac{dx}{dt} = \frac{\lambda}{T} \quad (\text{m/s})$$

The phase velocity, also called the *propagation velocity*, is *the velocity of the wave pattern* as it moves across the water surface.

$$f = \frac{1}{T} \quad (\text{Hz})$$

$$u_p = f\lambda \quad (\text{m/s})$$

with $y(x, t) = A \cos\left(\frac{2\pi t}{T} - \frac{2\pi x}{\lambda}\right)$ (m)



$$y(x, t) = A \cos\left(2\pi f t - \frac{2\pi}{\lambda} x\right) \\ = A \cos(\omega t - \beta x),$$

(wave moving along +x direction)

$$y(x, t) = A \cos(\omega t + \beta x).$$

(wave moving along -x direction)

$$\omega = 2\pi f \quad (\text{rad/s})$$

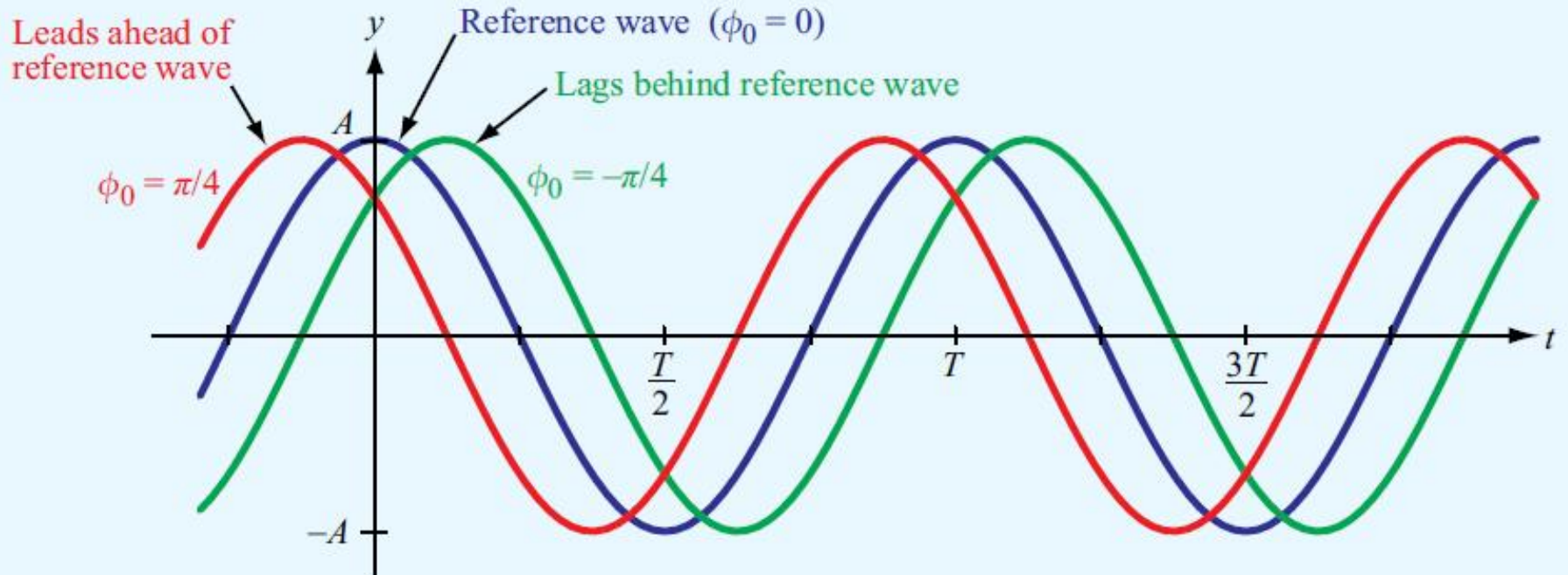
$$\beta = \frac{2\pi}{\lambda} \quad (\text{rad/m})$$



Direction of Travel

Determined by inspection of both t and x terms in phase of the wave

- If one sign is positive and the other negative, then the wave is traveling in the positive x -direction
- If both signs are positive or both are negative, then the wave is traveling in the negative x -direction



The constant reference ϕ_0 has no influence on either velocity or the direction of wave propagation



Sinusoidal Wave in a Lossless Medium

► Alternate Form

$$y(x, t) = A \cos(\omega t - \beta x + \phi_o)$$

Angular Velocity

Phase Constant or
Wavenumber



Sinusoidal Waves in a Lossy Medium

If a wave is traveling in the x direction in a *lossy medium*, its amplitude decreases as $e^{-\alpha x}$. This factor is called the *attenuation factor*, and α is called the *attenuation constant* of the medium and its unit is neper per meter (Np/m).

Attenuation constant (Neper/meter)

$$y(x, t) = A \underbrace{e^{-\alpha x}}_{\text{Attenuation Factor}} \cos(\omega t - \beta x + \phi_o)$$



Sinusoidal Waves in a Lossy Medium

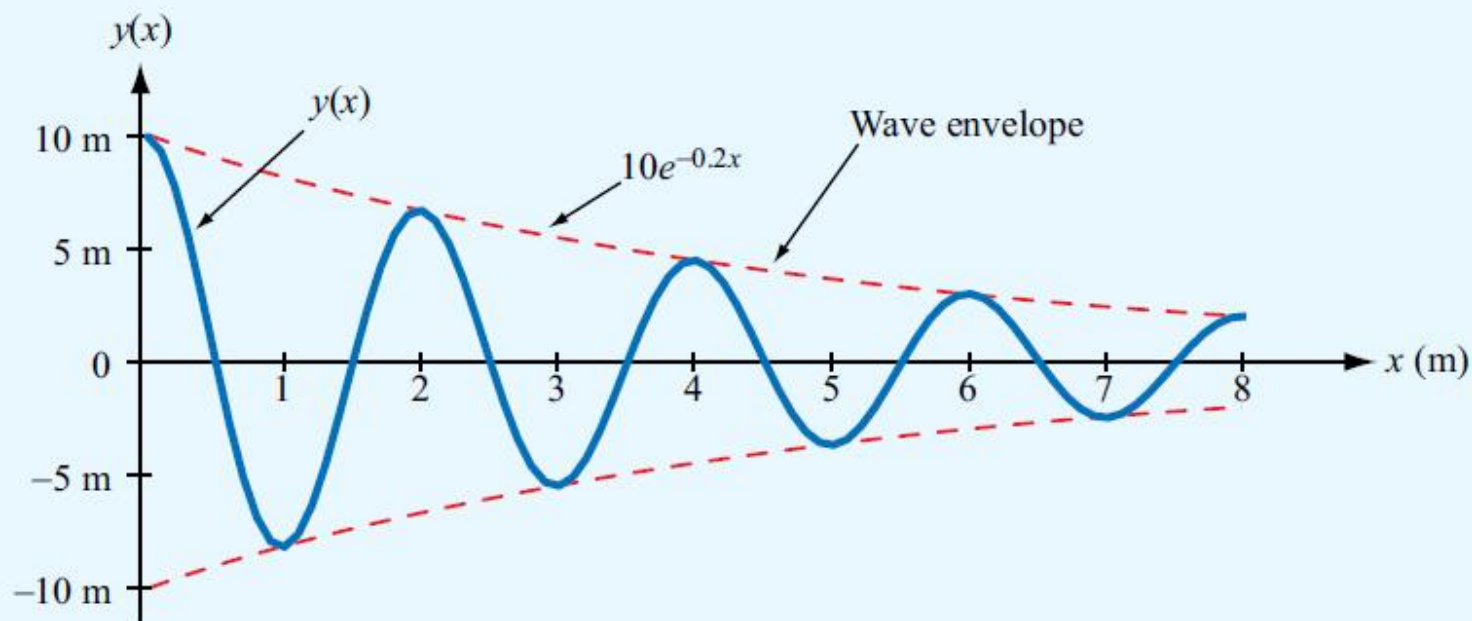


Figure 1-15 Plot of $y(x) = (10e^{-0.2x} \cos \pi x)$ meters. Note that the envelope is bounded between the curve given by $10e^{-0.2x}$ and its mirror image.



Example

The electric field of a traveling electromagnetic wave is given by

$$E(z, t) = 10 \cos(\pi \times 10^7 t + \pi z/15 + \pi/6) \text{ (V/m)}.$$

Determine

- (a) the direction of wave propagation
- (b) The wave frequency f
- (c) its wavelength λ
- (d) its phase velocity u_p .

Answer: (a) $-z$ direction, (b) $f = 5 \text{ MHz}$, (c) $\lambda = 30 \text{ m}$,
(d) $u_p = 1.5 \times 10^8 \text{ m/s}$.



Sinusoidal Waves in a Lossy Medium

Example

A laser beam of light propagating through the atmosphere is characterized by an electric field given by

$$E(x, t) = 150e^{-0.03x} \cos(3 \times 10^{15}t - 10^7x) \quad (\text{V/m}),$$

where x is the distance from the source in meters. The attenuation is due to absorption by atmospheric gases. Determine

- (a) the direction of wave travel,
- (b) the wave velocity, and
- (c) the wave amplitude at a distance of 200 m.



Sinusoidal Waves in a Lossy Medium

Solution: (a) Since the coefficients of t and x in the argument of the cosine function have opposite signs, the wave must be traveling in the $+x$ direction.

(b)

$$u_p = \frac{\omega}{\beta} = \frac{3 \times 10^{15}}{10^7} = 3 \times 10^8 \text{ m/s},$$

which is equal to c , the velocity of light in free space.

(c) At $x = 200$ m, the amplitude of $E(x, t)$ is

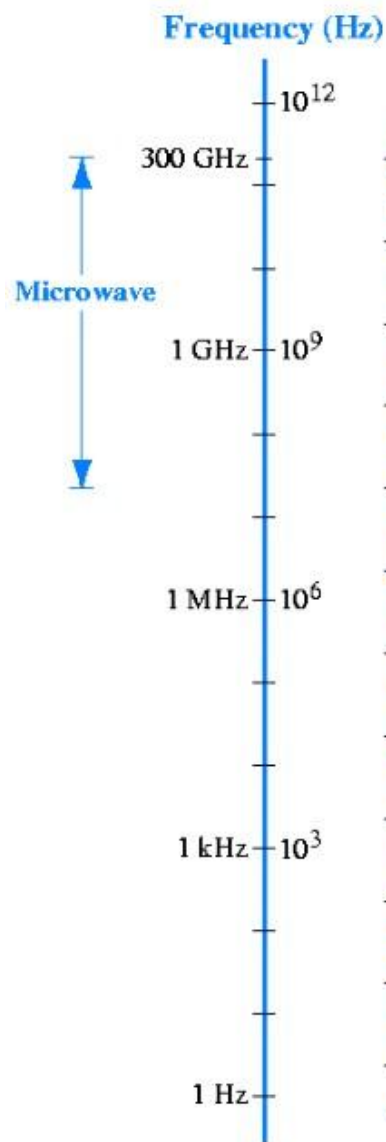
$$150e^{-0.03 \times 200} = 0.37 \quad (\text{V/m}).$$



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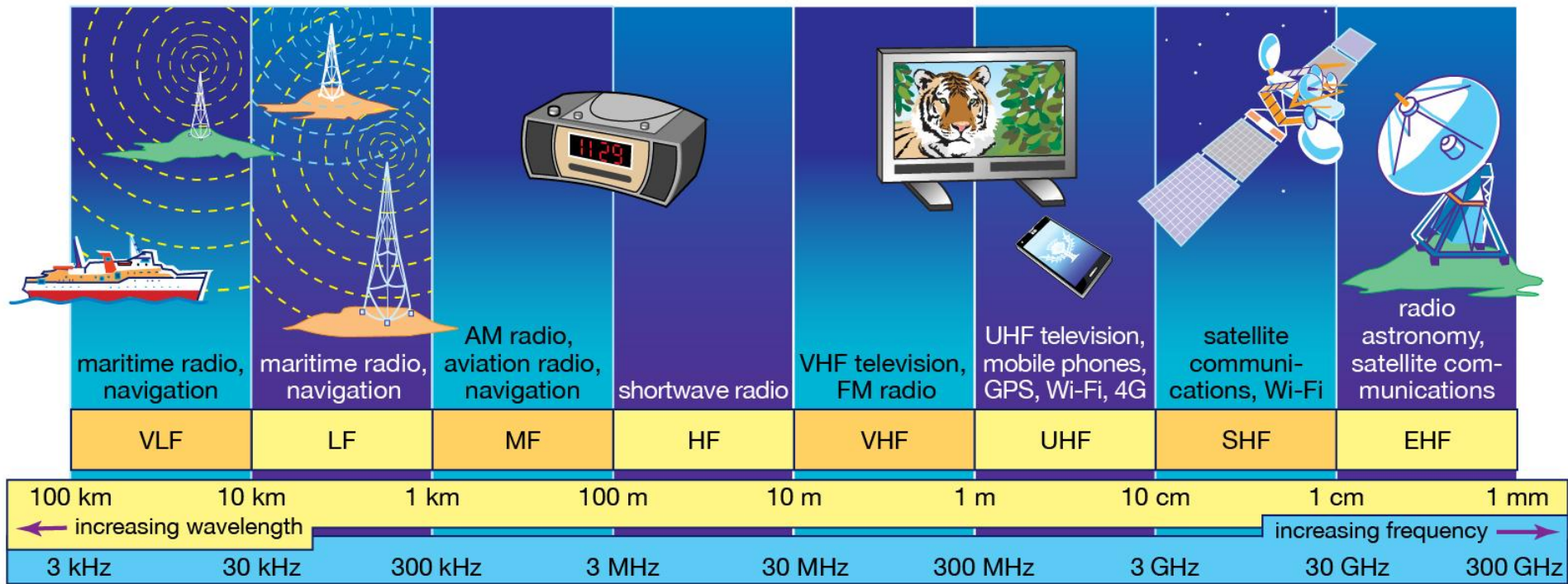
The Electromagnetic Spectrum



Band	Applications
Extremely High Frequency EHF (30 - 300 GHz)	Radar, advanced communication systems, remote sensing, radio astronomy
Super High Frequency SHF (3 - 30 GHz)	Radar, satellite communication systems, aircraft navigation, radio astronomy, remote sensing
Ultra High Frequency UHF (300 MHz - 3 GHz)	TV broadcasting, radar, radio astronomy, microwave ovens, cellular telephone
Very High Frequency VHF (30 - 300 MHz)	TV and FM broadcasting, mobile radio communication, air traffic control
High Frequency HF (3 - 30 MHz)	Short wave broadcasting
Medium Frequency MF (300 kHz - 3 MHz)	AM broadcasting
Low Frequency LF (30 - 300 kHz)	Radio beacons, weather broadcast stations for air navigation
Very Low Frequency VLF (3 - 30 kHz)	Navigation and position location
Ultra Low Frequency ULF (300 Hz - 3 kHz)	Audio signals on telephone
Super Low Frequency SLF (30 - 300 Hz)	Ionospheric sensing, electric power distribution, submarine communication
Extremely Low Frequency ELF (3 - 30 Hz)	Detection of buried metal objects
$f < 3$ Hz	Magnetotelluric sensing of the earth's structure



The Electromagnetic Spectrum



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Any *complex number* z can be expressed in *rectangular form* as

$$z = x + jy$$

where x and y are the *real* (Re) and *imaginary* (Im) parts of z , respectively, and $j = \sqrt{-1}$. That is,

$$x = \text{Re}(z), y = \text{Im}(z).$$

Alternatively, z may be cast in *polar form* as

$$z = |z| e^{j\theta} = |z| \angle \theta$$

where $|z|$ is the magnitude of z , θ is its phase angle, and $\angle \theta$ is a useful shorthand representation for $e^{j\theta}$. Applying *Euler's identity*,

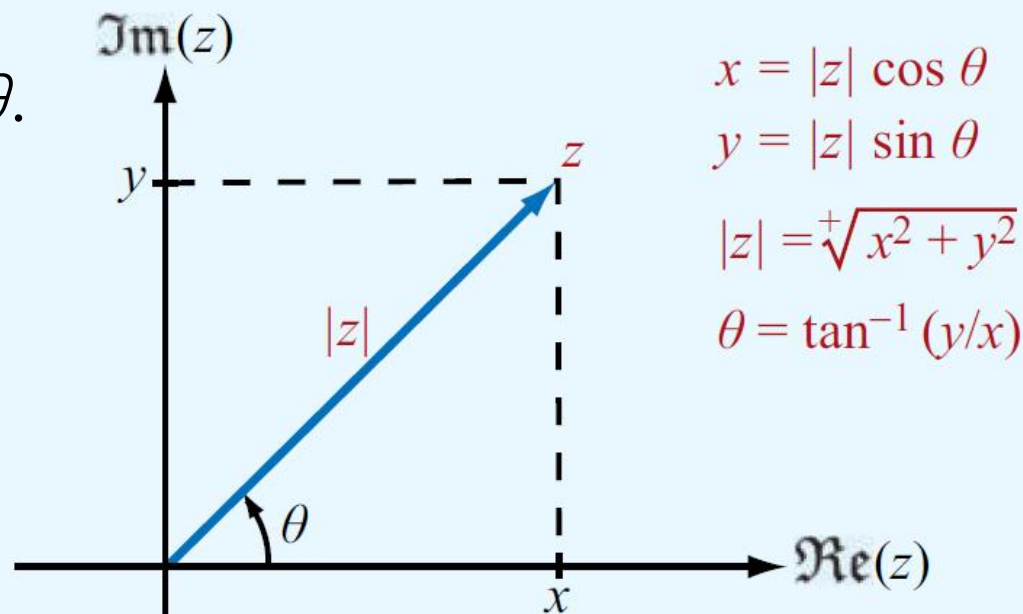
$$e^{j\theta} = \cos \theta + j \sin \theta$$



Review of Complex Numbers

$$z = |z| e^{j\theta} = |z| \cos \theta + j |z| \sin \theta.$$

The magnitude $|z|$ is equal to the positive square root of the product of z and its complex conjugate: $|z| = \sqrt{zz^*}$



The **complex conjugate** of z , denoted with a star superscript (or asterisk), is obtained by replacing j (wherever it appears) with $-j$, so that

$$z^* = (x + jy)^* = x - jy = |z| e^{-j\theta} = |z| \angle -\theta$$



Example

Given two complex numbers $V = 3 - j4$, $I = -(2 + j3)$,

(a) express V and I in polar form, and find (b) VI , (c) VI^* , (d) V/I ,
and (e) \sqrt{I}



Solution

$$(a) \quad |V| = \sqrt[4]{VV^*} = \sqrt[4]{(3-j4)(3+j4)}$$

$$= \sqrt[4]{9+16} = 5,$$

$$\theta_v = \tan^{-1}(-4/3) = -53.1^\circ$$

$$V = |V| e^{j\theta_v} = 5e^{-j53.1^\circ} = |5| \angle -53.1^\circ$$

$$|I| = \sqrt[4]{2^2 + 3^2} = \sqrt[4]{13} = 3.61$$

$$(\quad VI = 5e^{-j53.1^\circ} \times 3.61e^{j236.3^\circ}$$

$$b) \quad = 18.03e^{j(236.3^\circ - 53.1^\circ)}$$

$$= 18.03e^{j183.2^\circ}$$

$$(\quad VI^* = 5e^{-j53.1^\circ} \times 3.61e^{-j236.3^\circ} = 18.03e^{-j289.4^\circ} = 18.03e^{j70.6^\circ}$$

c)

$$(d) \quad \frac{V}{I} = \frac{5e^{-j53.1^\circ}}{3.61e^{j236.3^\circ}} = 1.39e^{-j289.4^\circ} = 1.39e^{j70.6^\circ}$$

$$(\quad \sqrt{I} = \sqrt{3.61e^{j236.3^\circ}} = \pm\sqrt{3.61}e^{j236.3^\circ/2} = \pm 1.90e^{j118.15^\circ}$$

e)

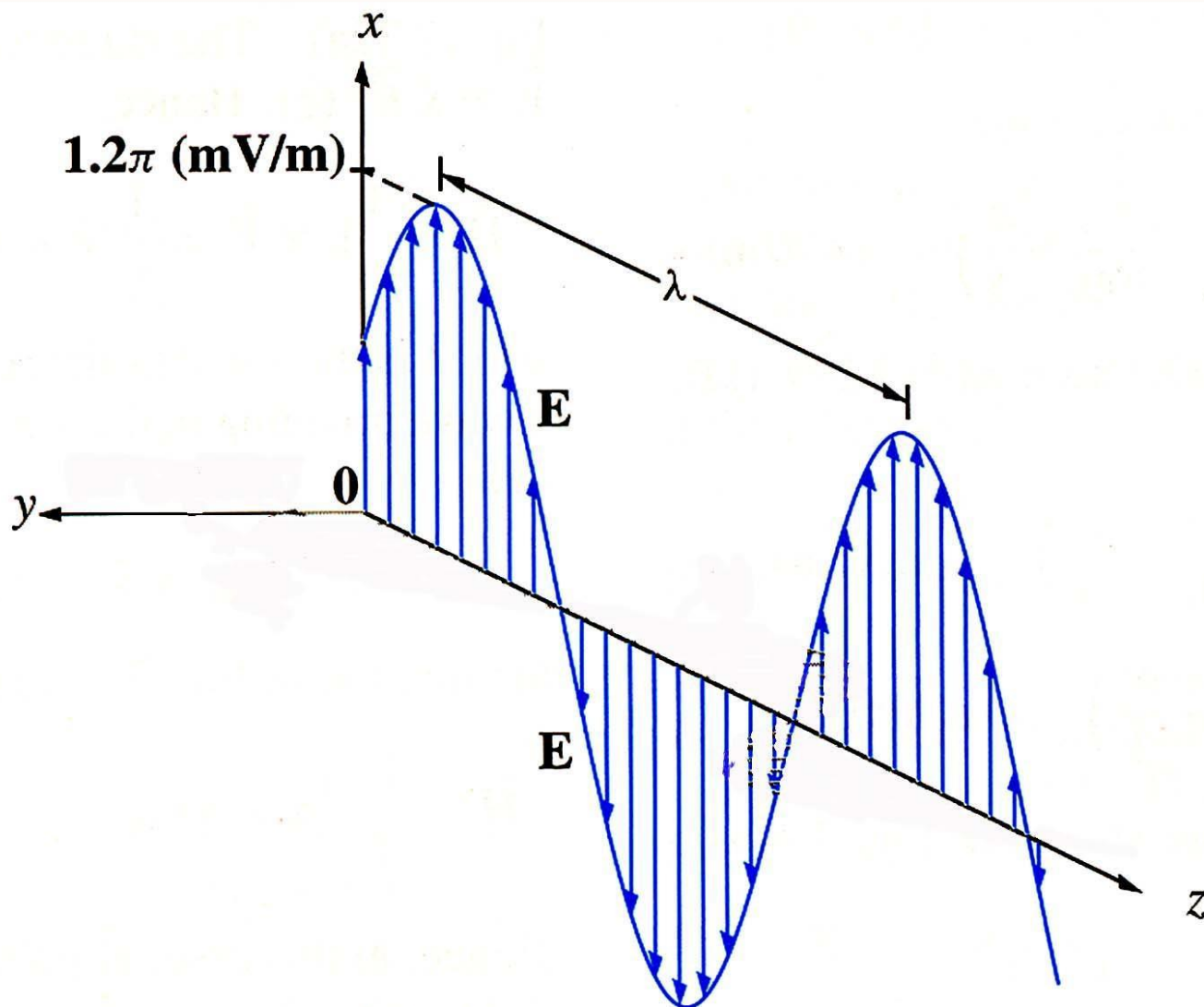
Since $I = (-2 - j3)$ is in the third quadrant in the complex plane

$$\theta_I = +180^\circ + \tan^{-1}(3/2) = 236.3^\circ$$

$$I = 3.61 \angle 236.3^\circ$$



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$$\mathbf{E}(x, y, z; t) = \text{Re}[\underbrace{\tilde{\mathbf{E}}(x, y, z)}_{\text{Vector Phasor}} e^{j\omega t}]$$

Instantaneous

Vector Phasor



- Phasors are useful when analyzing periodic functions in linear systems
- A simple cosine can be represented as a constant in the phasor domain
- Complicated time domain functions can be represented as an addition of sin and cosine functions (Fourier series)
- Analyzing the frequency content of the wave
- Using a cosine as a reference:



The simple RC circuit shown in **Fig. 1-20** contains a sinusoidally time-varying voltage source given by

$$v_s(t) = V_0 \sin(\omega t + \phi_0)$$

Application of Kirchhoff's voltage law gives the following loop **equation**:

$$Ri(t) + \frac{1}{C} \int i(t) dt = v_s(t)$$

Our objective is to obtain an expression for the current $i(t)$.

We can do this by solving **equation** in the time domain, which is somewhat cumbersome because the forcing function $v_s(t)$ is a sinusoid.

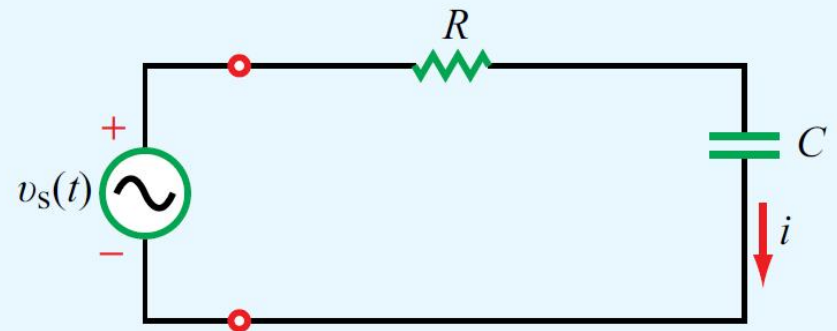


Figure 1-20 RC circuit connected to a voltage source $v_s(t)$.



Solution

Step 1: Adopt a cosine reference

$$\begin{aligned}v_s(t) &= V_0 \sin(\omega t + \phi_0) \\&= V_0 \cos\left(\frac{\pi}{2} - \omega t - \phi_0\right) \\&= V_0 \cos\left(\omega t + \phi_0 - \frac{\pi}{2}\right)\end{aligned}$$

Step 2: Express time-dependent variables as phasors

$$z(t) = \operatorname{Re}\left[\tilde{Z}e^{j\omega t}\right]$$

$$\begin{aligned}v_s(t) &= \operatorname{Re}\left[V_0 e^{j(\omega t + \phi_0 - \frac{\pi}{2})}\right] \\&= \operatorname{Re}\left[V_0 e^{j(\phi_0 - \frac{\pi}{2})} e^{j\omega t}\right] = \operatorname{Re}\left[\tilde{V}_s e^{j\omega t}\right]\end{aligned}$$

$$\tilde{V}_s = V_0 e^{j(\phi_0 - \frac{\pi}{2})}$$

$$i(t) = \operatorname{Re}(\tilde{I}e^{j\omega t})$$

$$\begin{aligned}\frac{di}{dt} &= \frac{d}{dt}\left[\operatorname{Re}(\tilde{I}e^{j\omega t})\right] \\&= \operatorname{Re}\left[\frac{d}{dt}(\tilde{I}e^{j\omega t})\right] = \operatorname{Re}\left[j\omega \tilde{I}e^{j\omega t}\right]\end{aligned}$$

$$\begin{aligned}\int i dt &= \int \operatorname{Re}(\tilde{I}e^{j\omega t}) dt \\&= \operatorname{Re}\left(\int \tilde{I}e^{j\omega t} dt\right) = \operatorname{Re}\left(\frac{\tilde{I}}{j\omega} e^{j\omega t}\right)\end{aligned}$$



Solution

Step 3: Recast the differential / integral equation in phasor form

$$R \operatorname{Re}(\tilde{I} e^{j\omega t}) + \frac{1}{C} \operatorname{Re}\left(\frac{\tilde{I}}{j\omega} e^{j\omega t}\right) = \operatorname{Re}\left[\tilde{V}_s e^{j\omega t}\right]$$

$$\operatorname{Re}\left\{\left[\left(R + \frac{1}{j\omega C}\right)\tilde{I} - \tilde{V}_s\right] e^{j\omega t}\right\} = 0$$

$$\operatorname{Im}\left\{\left[\left(R + \frac{1}{j\omega C}\right)\tilde{I} - \tilde{V}_s\right] e^{j\omega t}\right\} = 0$$

$$\tilde{I}\left(R + \frac{1}{j\omega C}\right) = \tilde{V}_s \quad \text{(phasor domain).}$$

Step 4: Solve the phasor-domain equation

$$\tilde{I} = \frac{\tilde{V}_s}{R + 1/(j\omega C)}$$

$$\begin{aligned}\tilde{I} &= V_0 e^{j(\phi_0 - \frac{\pi}{2})} \left[\frac{j\omega C}{1 + j\omega RC} \right] \\ &= V_0 e^{j(\phi_0 - \frac{\pi}{2})} \left[\frac{\omega C e^{j\pi/2}}{\sqrt{1 + \omega^2 R^2 C^2} e^{j\phi_1}} \right] \\ &= \frac{V_0 \omega C}{\sqrt{1 + \omega^2 R^2 C^2}} e^{j(\phi_0 - \phi_1)}\end{aligned}$$



Solution

Step 5: Find the instantaneous value

$$\begin{aligned} i(t) &= \operatorname{Re} \left[\tilde{I} e^{j\omega t} \right] \\ &= \operatorname{Re} \left[\frac{V_0 \omega C}{\sqrt{1 + \omega^2 R^2 C^2}} e^{j(\phi_0 - \phi_1)} e^{j\omega t} \right] \\ &= \frac{V_0 \omega C}{\sqrt{1 + \omega^2 R^2 C^2}} \cos(\omega t + \phi_0 - \phi_1) \end{aligned}$$

In summary, we converted all time-varying quantities into the phasor domain, solved for the phasor \tilde{I} of the desired instantaneous current $i(t)$, and then converted back to the time domain to obtain an expression for $i(t)$.



Table 1-5 Time-domain sinusoidal functions $z(t)$ and their cosine-reference phasor-domain counterparts \tilde{Z} , where $z(t) = \Re e [\tilde{Z}e^{j\omega t}]$.

$z(t)$		\tilde{Z}
$A \cos \omega t$	\longleftrightarrow	A
$A \cos(\omega t + \phi_0)$	\longleftrightarrow	$Ae^{j\phi_0}$
$A \cos(\omega t + \beta x + \phi_0)$	\longleftrightarrow	$Ae^{j(\beta x + \phi_0)}$
$Ae^{-\alpha x} \cos(\omega t + \beta x + \phi_0)$	\longleftrightarrow	$Ae^{-\alpha x} e^{j(\beta x + \phi_0)}$
$A \sin \omega t$	\longleftrightarrow	$Ae^{-j\pi/2}$
$A \sin(\omega t + \phi_0)$	\longleftrightarrow	$Ae^{j(\phi_0 - \pi/2)}$
$\frac{d}{dt}(z(t))$	\longleftrightarrow	$j\omega\tilde{Z}$
$\frac{d}{dt}[A \cos(\omega t + \phi_0)]$	\longleftrightarrow	$j\omega Ae^{j\phi_0}$
$\int z(t) dt$	\longleftrightarrow	$\frac{1}{j\omega}\tilde{Z}$
$\int A \sin(\omega t + \phi_0) dt$	\longleftrightarrow	$\frac{1}{j\omega} Ae^{j(\phi_0 - \pi/2)}$



Concepts

- Electromagnetics is the study of electric and magnetic phenomena and their engineering applications.
- The International System of Units consists of the six fundamental dimensions listed in **Table 1-1**. The units of all other physical quantities can be expressed in terms of the six fundamental units.
- The four fundamental forces of nature are the nuclear, weak-interaction, electromagnetic, and gravitational forces.
- The source of the electric field quantities **E** and **D** is the electric charge q . In a material, **E** and **D** are related by $\mathbf{D} = \epsilon \mathbf{E}$, where ϵ is the electrical permittivity of the material. In free space, $\epsilon = \epsilon_0 \approx (1/36\pi) \times 10^{-9}$ (F/m).
- The source of the magnetic field quantities **B** and **H** is the electric current I . In a material, **B** and **H** are related

by $\mathbf{B} = \mu \mathbf{H}$, where μ is the magnetic permeability of the medium. In free space, $\mu = \mu_0 = 4\pi \times 10^{-7}$ (H/m).

- Electromagnetics consists of three branches: (1) electrostatics, which pertains to stationary charges, (2) magnetostatics, which pertains to dc currents, and (3) electrodynamics, which pertains to time-varying currents.
- A traveling wave is characterized by a spatial wavelength λ , a time period T , and a phase velocity $u_p = \lambda/T$.
- An electromagnetic (EM) wave consists of oscillating electric and magnetic field intensities and travels in free space at the velocity of light $c = 1/\sqrt{\epsilon_0\mu_0}$. The EM spectrum encompasses gamma rays, X-rays, visible light, infrared waves, and radio waves.
- Phasor analysis is a useful mathematical tool for solving problems involving time-periodic sources.