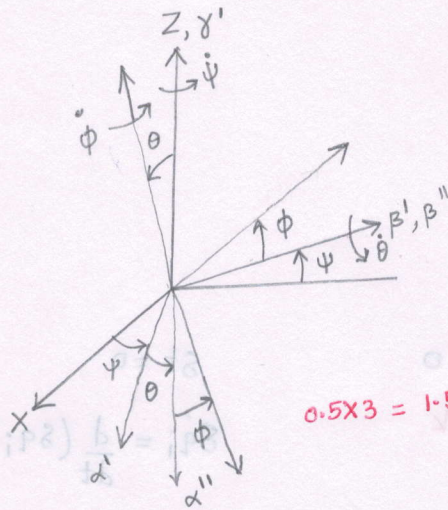


1.

$$\begin{bmatrix} \omega_{\alpha'} \\ \omega_{\beta'} \\ \omega_{\gamma'} \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \dot{\psi} \end{bmatrix}$$

$$\begin{bmatrix} \omega_{\alpha''} \\ \omega_{\beta''} \\ \omega_{\gamma''} \end{bmatrix} = \underbrace{\begin{bmatrix} c\theta & 0 & -s\theta \\ 0 & 1 & 0 \\ s\theta & 0 & c\theta \end{bmatrix}}_{Q_\theta^T} \begin{bmatrix} 0 \\ 0 \\ \dot{\psi} \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ \dot{\theta} \end{bmatrix}$$

$$= \begin{bmatrix} -\dot{\psi} s\theta \\ \dot{\theta} \\ \dot{\psi} c\theta \end{bmatrix} \quad 1 \text{ Mark}$$



0.5X3 = 1.5 Mark

$$\begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} = \begin{bmatrix} c\phi & s\phi & 0 \\ -s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} -\dot{\psi} s\theta \\ \dot{\theta} \\ \dot{\psi} c\theta \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ \dot{\phi} \end{bmatrix}$$

$$= \begin{bmatrix} -\dot{\psi} s\theta c\phi + \dot{\theta} s\phi \\ \dot{\psi} s\theta s\phi + \dot{\theta} c\phi \\ \dot{\psi} c\theta + \dot{\phi} \end{bmatrix} = \begin{bmatrix} -s\theta c\phi & s\phi & 0 \\ s\theta s\phi & c\phi & 0 \\ c\theta & 0 & 1 \end{bmatrix} \begin{bmatrix} \dot{\psi} \\ \dot{\theta} \\ \dot{\phi} \end{bmatrix}$$

1 Mark

2. The actual path in configuration space followed by a Holonomic dynamical system during the fixed time interval t_1 to t_2 is such that the integral

$$I = \int_{t_1}^{t_2} L(q, \dot{q}, t) dt$$

is stationary end points.

with respect to path variations which vanish at the 1.5 Mark

Mathematical Form

$$\delta I = 0, \text{ i.e., } \delta \int_{t_1}^{t_2} L dt = 0 \quad 0.5 \text{ Mark}$$

3. From Hamilton's principle, we have $\delta I = \int_{t_1}^{t_2} \delta L dt = 0$

$$\Rightarrow \int_{t_1}^{t_2} L(q_1, q_2, \dots, q_n; \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t) dt = 0$$

For holonomic system we have

$$\int_{t_1}^{t_2} \delta L(q_1, q_2, \dots, q_n; \dot{q}_1, \dot{q}_2, \dots, \dot{q}_n, t) dt = 0$$

$$\Rightarrow \int_{t_1}^{t_2} \left\{ \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i + \sum_{i=1}^n \frac{\partial L}{\partial \dot{q}_i} \delta \dot{q}_i + \frac{\partial L}{\partial t} \delta t \right\} dt = 0$$

1 Mark

$$\delta t = 0$$

$$\delta \dot{q}_i = \frac{d}{dt}(\delta q_i)$$

$$\Rightarrow \int_{t_1}^{t_2} \left\{ \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i + \sum_{i=1}^n \frac{\partial L}{\partial \dot{q}_i} \frac{d}{dt}(\delta q_i) \right\} dt = 0$$

$$\Rightarrow \int_{t_1}^{t_2} \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i dt + \sum_{i=1}^n \int_{t_1}^{t_2} \frac{\partial L}{\partial \dot{q}_i} \frac{d}{dt}(\delta q_i) dt = 0$$

$$\Rightarrow \int_{t_1}^{t_2} \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i dt + \sum_{i=1}^n \left. \frac{\partial L}{\partial \dot{q}_i} \delta q_i \right|_{t=t_1}^{t=t_2} - \sum_{i=1}^n \int_{t_1}^{t_2} \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) \delta q_i dt = 0$$

1 Mark

$$\Rightarrow \int_{t_1}^{t_2} \sum_{i=1}^n \frac{\partial L}{\partial q_i} \delta q_i dt - \int_{t_1}^{t_2} \sum_{i=1}^n \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) \delta q_i dt = 0$$

$$\Rightarrow \int_{t_1}^{t_2} \sum_{i=1}^n \left\{ \frac{\partial L}{\partial q_i} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) \right\} \delta q_i dt = 0$$

1 Mark

From Fundamental lemma, the integrand should vanish. Further, since δq_i are all independent, we get

1 Mark (Reasoning)

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0.$$

4. Constraint : $r\dot{\theta} - \dot{s} = 0$ —① 0.5 Mark

Kinetic Energy $T = \frac{1}{2} m \dot{s}^2 + \frac{1}{2} m r^2 \dot{\theta}^2$ } 0.5 Mark

Potential Energy $V = mg(L-s)\sin\phi + r\cos\phi$

where L is the length of the inclined plane

Lagrangian $L = T - V$
 $= \frac{1}{2} m \dot{s}^2 + \frac{1}{2} m r^2 \dot{\theta}^2 - mg(L-s)\sin\phi$ ~~0.5 Mark~~

The two Lagrange's equations of second-kind are

$m\ddot{s} - mg\sin\phi = -\mu$ —②

$mr^2\ddot{\theta} = \mu r$ —③

Equations ①, ②, ③ constitute three equations in unknowns θ, s, λ 2 Mark

From ① $r\dot{\theta} - \dot{s} = 0$

Differentiating, we get $r\ddot{\theta} - \ddot{s} = 0$ —④

From ③ and ④ we get $m\ddot{s} = \mu$ —⑤

From ⑤ and ②, ③ $\ddot{s} = \frac{g\sin\phi}{2}$

$\mu = \frac{mg\sin\phi}{2}$ 100 Mark

$\ddot{\theta} = \frac{g\sin\phi}{2r}$

Friction Force and Moment enforce the constraint $\dot{s} - r\dot{\theta} = 0$

Constraint Forces

$C_\theta = \frac{Mg\sin\phi}{2} r$ ~~0.5 Mark~~

$C_s = -\frac{Mg\sin\phi}{2}$ ~~1 Mark~~

$C_\theta = \mu r$

$C_s = -\mu$ 0.5 Mark

C_θ is frictional torque about c.m.
 C_s is frictional force at contact

5. $H = \frac{p^2}{2\alpha} - bqpe^{-\alpha t} + \frac{ba}{2} q^2 e^{-\alpha t} (\alpha + be^{-\alpha t}) + \frac{kq^2}{2}$

$\dot{q} = \frac{\partial H}{\partial p} \Rightarrow \dot{q} = \frac{p}{\alpha} - bq e^{-\alpha t}$ —① 0.5 Mark

From ① $p = \alpha \{ \dot{q} + bq e^{-\alpha t} \}$ —② 0.5 Mark

$H = p\dot{q} - L \Rightarrow L = p\dot{q} - H$ 0.5 Mark

$$\Rightarrow L = p\dot{q} - \left\{ \frac{p^2}{2\alpha} - bqpe^{-\alpha t} + \frac{ba}{2} \dot{q}^2 e^{-\alpha t} (\alpha + be^{-\alpha t}) + \frac{kq^2}{2} \right\} \quad 0.5 \text{ Mark}$$

Now substitute $p = \alpha \{ \dot{q} + bq e^{-\alpha t} \}$

$$\Rightarrow L = \alpha \dot{q} \{ \dot{q} + bq e^{-\alpha t} \} - \left\{ \frac{1}{2\alpha} \alpha^2 (\dot{q} + bq e^{-\alpha t})^2 - bq e^{-\alpha t} \alpha (\dot{q} + bq e^{-\alpha t}) + \frac{ba}{2} \dot{q}^2 e^{-\alpha t} (\alpha + be^{-\alpha t}) + \frac{kq^2}{2} \right\} \quad 0.5 \text{ Mark}$$