Localization of low-energy states for semiclassical Toeplitz Operators

Alix Deleporte Advisor : Nalini Anantharaman

Institut de Recherche Mathématique Avancée Université de Strasbourg

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Classical mechanics	Quantum mechanics
Symplectic manifold M	Hilbert Space H
Function $\mathfrak{a} \in C^{\infty}(M,\mathbb{R})$	Self-adjoint operator $A \in L(H)$
Hamiltonien flow of a	Flow of e ^{itA/ħ}
Poisson Bracket	Lie Bracket

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- Quantization: for a given classical model, how to construct an associated quantum model?
- Semiclassics : the quantum model is \hbar -dependent. What can be said in the $\hbar \to 0$ limit ?



• Quantum spins: triplet of self-adjoint matrices $S_x, S_u, S_z \in M_{2S+1}(\mathbb{C})$, with

$$[S_a, S_b] = \frac{i}{S} \epsilon_{abc} S_c.$$

■ For $S = \frac{1}{2}$, one finds the Pauli matrices

$$S_x = \frac{1}{2} \left(\begin{smallmatrix} 0 & & 1 \\ 1 & & 0 \end{smallmatrix} \right) \ S_y = \frac{1}{2} \left(\begin{smallmatrix} 0 & & i \\ -i & & 0 \end{smallmatrix} \right) \ S_z = \frac{1}{2} \left(\begin{smallmatrix} 1 & & 0 \\ 0 & & -1 \end{smallmatrix} \right).$$

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■ For any finite graph G, we wish to study the following operator acting on $(\mathbb{C}^{2S+1})^{\otimes |G|}$:

$$H = \sum_{e \sim f} S_{x,e} S_{x,f} + S_{y,e} S_{y,f} + S_{z,e} S_{z,f}$$

as
$$S \to +\infty$$
.



Plan

- 1 Toeplitz operators on Bargmann spaces
 - Bargmann spaces
 - Definition
 - Semiclassical properties
- 2 Generalization to Kähler manifolds
 - Hardy spaces
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 - Wells
 - Miniwells
 - Conjectures



Bargmann spaces

The Bargmann spaces are L^2 spaces of holomorphic functions on \mathbb{C}^n , with a weight.

$$B_{\mathbf{N}}(\mathbb{C}^n) = \left\{z \mapsto \exp\left(-\frac{N}{2}|z|^2\right) f(z) \text{, f holomorphic}\right\} \cap L^2(\mathbb{C}^n)$$

Those are closed subspaces of $L^2(\mathbb{C}^n)$.

The Szegő projector

Let Π_N be the orthogonal projector from $L^2(\mathbb{C}^n)$ onto $B_N(\mathbb{C}^n)$. Π_N . It admits a Schwartz kernel:

$$\Pi_{\mathbf{N}}(z,w) = \frac{\mathbf{N}^{\mathbf{n}}}{\pi^{\mathbf{n}}} \exp \left[\mathbf{N}(-\frac{1}{2}|z-w|^2 + \mathrm{i}\Im(z\cdot\overline{w})) \right].$$

As $N\to +\infty$, the kernel is exponentially decreasing far from the diagonal. The typical interaction scale is $N^{-1/2}$.

Toeplitz operators

Definition

Let $h \in C^{\infty}(\mathbb{C}^n)$ a smooth bounded function, and $N \in \mathbb{N}$. We denote by $T_N(h)$ the Toeplitz operator associated to h:

$$\begin{split} T_N(h) : B_N(\mathbb{C}^n) & \mapsto & B_N(\mathbb{C}^n) \\ u & \mapsto & \Pi_N(hu). \end{split}$$

If h is not bounded, we can construct $T_N(h)$ as an unbounded operator on $B_N(\mathbb{C}^n)$.

The mapping $h \mapsto T_N(h)$ is linear and adjoint-preserving. If h is real-valued, then $T_N(h)$ is formally self-adjoint.



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- $\text{ If } h: z \mapsto \overline{z}^\alpha z^\beta \text{, then } \mathsf{T}_N(h) = \mathsf{N}^{-|\alpha|} \vartheta^\alpha z^\beta.$
- If q is a definite quadratic form on \mathbb{R}^{2n} , then $T_N(q)$ has a compact resolvent. The first eigenvalue $\mu_N(q) = N^{-1}\mu_1(q)$ is positive.

$$\mu_1(q) = \mathsf{min}\,\mathsf{Sp}(\mathsf{Op}_1^{W}(q)) + \frac{1}{2}\,\mathsf{tr}(q)$$

Composition and bracket

Proposition

Leta and b two smooth bounded functions on \mathbb{C}^n . Then there is a sequence $(c_i)_{i\in\mathbb{N}}$ of smooth bounded functions on \mathbb{C}^n , with $c_0=\mathfrak{ab}$ so that, as $N\to +\infty$, there holds:

$$T_N(a)T_N(b) = T_N(c_0) + N^{-1}T_N(c_1) + N^{-2}T_N(c_2) + \dots$$

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In particular,

$$[T_N(\mathfrak{a}),T_N(\mathfrak{b})]=\frac{i}{N}T_N(\{\mathfrak{a},\mathfrak{b}\})+O(N^{-2}).$$

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- We wish to generalize Bargmann spaces to other complex manifolds.
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- Instead of considering weighted spaces, we will consider spaces of holomorphic sections.

Notations

- M is a compact Kähler manifold, with symplectic form ω .
- L is a complex line bundle on M, endowed with a hermitian structure h, so that the curvature of the Chern connexion is ω .
- $N \geqslant 1$ is an integer.

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Then if s is a (continuous) section of $L^{\otimes N}$, one can compute

$$\|s\|_{L^2} := \int_M h_N(s(m)) \frac{\omega^{\wedge n}}{n!}.$$

By completion, one defines the Hilbert space of square-integrable sections of $L^{\otimes N}$.



Hardy spaces

Definition

The N-equivariant Hardy space is the space $H_N(M,L)$ of L^2 and holomorphic sections of $L^{\otimes N}$.

This space is finite-dimensional, the dimension is polynomial in N (Riemann-Roch).

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The Szegő projector S_N is the orthogonal projector from $L^2(M,L^{\otimes N})$ onto $H_N(M,L)$.

It always admits a Schwartz kernel (as a section of $L^{\otimes N}\boxtimes L^{\otimes -N}$) because $H_N(M,L)$ is finite-dimensional.

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 $T_N(h)$ acts on a finite-dimensional space, and it is symmetric when h is real-valued.

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 $T_N(h)$ acts on a finite-dimensional space, and it is symmetric when h is real-valued.

Observe that, for any $u, v \in H_N(M, L)$, there holds

$$\langle \mathbf{u}, \mathsf{T}_{\mathsf{N}}(\mathsf{h}) \mathbf{v} \rangle = \langle \mathsf{u}, \mathsf{h} \mathsf{v} \rangle.$$



Asymptotics for the Szegő projector

Proposition (Boutet-Sjostrand 74)

For every $\varepsilon > 0$ and $k \in \mathbb{N}$ there exists C such that, for every $N \in \mathbb{N}$:

$$d(x,y) > \epsilon \Rightarrow |S_N(x,y)| \leqslant CN^{-k}$$

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Proposition (Charles 00, Zelditch 02, Ma 06)

In a convenient system of local coordinates, near any point of the diagonal, there holds:

$$S_{N}(z,w) \simeq \Pi_{N}(z,w) \left[1 + \sum_{k=1}^{K} N^{-k/2} b_{k}(\sqrt{N}z, \sqrt{N}w) \right]$$

Composition and bracket

Proposition (Charles 00, Schlichenmaier 02)

Let α and b two smooth functions on M. Then there is a sequence $(c_i)_{i\in\mathbb{N}}$ of smooth functions on M, with $c_0=\alpha b$, such that, as $N\to +\infty$, there holds:

$$T_N(a)T_N(b) = T_N(c_0) + N^{-1}T_N(c_1) + N^{-2}T_N(c_2) + \dots$$

In particular,

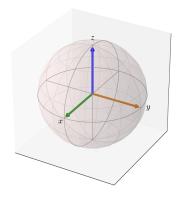
$$[T_N(\mathfrak{a}),T_N(\mathfrak{b})]=\frac{\mathfrak{i}}{N}T_N(\{\mathfrak{a},\mathfrak{b}\})+O(N^{-2}).$$

Hardy spaces on the sphere

- $H_N(\mathbb{CP}^1, L)$ corresponds to the the set of meromorphic functions on the sphere, with one pole of order at most N.
- Hence $H_N(\mathbb{CP}^1, L) \simeq \mathbb{C}_N[X]$, with dimension N+1.
- One Hilbert base is:

$$e_{k,N}(X) = \frac{\binom{k}{N}^{1/2}}{N} X^k.$$

Coordinate functions



- There are three coordinate functions on the sphere: x, y and z.
- The Toeplitz quantizations of these three functions are the spin operators, with $S = \frac{N}{2}$.

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A priori localization

- In the classical model, in order to minimize the energy, one picks any point where the energy is minimal.
- What happens for an eigenvector associated with the smallest eigenvalue of $T_N(h)$, as $N \to +\infty$?

Proposition (Charles 00)

An eigenvector with minimal eigenvalue is uniformly $O(N^{-\infty})$ outside any fixed neighbourhood of $\{h = min(h)\}$.

Can we get a more precise result?



A priori localization

- In the classical model, in order to minimize the energy, one picks any point where the energy is minimal.
- What happens for an eigenvector associated with the smallest eigenvalue of $T_N(h)$, as $N \to +\infty$?

Proposition (D. 16)

If the minimal set is non-degenerate, then for every $\delta \in [0,1/2),$ an eigenvector with minimal eigenvalue is uniformly $O(N^{-\infty})$ outside a neighbourhood of $\{h=min(h)\}$ with size $N^{-\delta}.$

Proof for localization speed

Let (\mathfrak{u}_N) be a sequence of unit eigenfunctions with minimal energy (λ_N) . Assume $\min(\mathfrak{h})=0$. We prove by induction on k that

$$\langle \mathfrak{u}_n, \mathfrak{h}^k \mathfrak{u}_n \rangle = O(N^{-k}).$$

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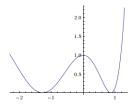
- Hard part: k = 1 (test $T_N(h)$ against a coherent state centred on a minimal point).
- Easy part: induction.

$$\langle u_n, h \star h u_n \rangle = \lambda_N^2 + O(N^{-\infty}),$$

where
$$h\star h=h^2+N^{-1}c_1(h,h)+O(N^{-2}).$$
 Now $c_1(h,h)\leqslant Ch.$

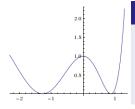
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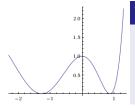


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What is minimized? The μ_1 of the hessian at this point.

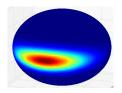
Case of wells: idea of proof

- By making more precise the previous argument, we have a lower bound for the first eigenvalue.
- The upper bound and a spectral gap are obtained by N^{-K} -quasimode for fixed K.

We remark that the quasimodes are exponentially localized, but this does not imply that the true eigenfunction is also localized.

Case of submanifold wells

- What can be said if h is minimal on a submanifold, with non-degenerate transverse hessian?
- \blacksquare \Rightarrow Same conclusion. (D.)



As N grows, the state concentrates on the miniwell and is more and more squeezed.

Miniwells in physics

It really happens in physics! For instance, with antiferromagnetic spins on a triangle graph.



It is conjectured that the minimal configurations are planar, in some cases.

Conjectures

Exponential Localization For now we only have $O(N^{-\infty})$ estimates for localisation. Can we hope for $O(\exp(-cN))$ estimates ?

Thermodynamical limit Instead of considering a fixed manifold M, we look at a particular symbol on M^n , and we let $n \to +\infty$. What is the behaviour vis-à-vis the semiclassical limit?

These two questions should be linked with each other.