# Localization of low-energy states for semiclassical Toeplitz Operators

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Classical mechanics	Quantum mechanics
Symplectic manifold M	Hilbert Space H
Function $\mathfrak{a}\in C^\infty(M,\mathbb{R})$	$ Self-adjoint\ operator\ A\in L(H) $
Hamiltonien flow of a	Flow of e <sup>itA/ħ</sup>
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- Quantization: for a given classical model, how to construct an associated quantum model?
- Semiclassics : the quantum model is  $\hbar$ -dependent. What can be said in the  $\hbar \to 0$  limit ?



• Quantum spins: triplet of self-adjoint matrices  $S_x, S_y, S_z \in M_{2S+1}(\mathbb{C})$ , with

$$[S_a, S_b] = \frac{i}{S} \epsilon_{abc} S_c.$$

■ For  $S = \frac{1}{2}$ , one finds the Pauli matrices

$$S_x = \frac{1}{2} \left( \begin{smallmatrix} 0 & & 1 \\ 1 & & 0 \end{smallmatrix} \right) \ S_y = \frac{1}{2} \left( \begin{smallmatrix} 0 & & i \\ -i & & 0 \end{smallmatrix} \right) \ S_z = \frac{1}{2} \left( \begin{smallmatrix} 1 & & 0 \\ 0 & & -1 \end{smallmatrix} \right).$$

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ight).$$

■ For any finite graph G, we wish to study the following operator acting on  $(\mathbb{C}^{2S+1})^{\otimes |G|}$ :

$$H = \sum_{e \sim f} S_{x,e} S_{x,f} + S_{y,e} S_{y,f} + S_{z,e} S_{z,f}$$

as 
$$S \to +\infty$$
.



### Plan

- 1 Toeplitz operators on Bargmann spaces
  - Bargmann spaces
  - Definition
  - Semiclassical properties
- 2 Generalization to Kähler manifolds
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## Bargmann spaces

The Bargmann spaces are  $L^2$  spaces of holomorphic functions on  $\mathbb{C}^n$ , with a weight.

$$B_{\mathbf{N}}(\mathbb{C}^n) = \left\{z \mapsto \exp\left(-\frac{N}{2}|z|^2\right) f(z) \text{, f holomorphic}\right\} \cap L^2(\mathbb{C}^n)$$

Those are closed subspaces of  $L^2(\mathbb{C}^n)$ .

# The Szegő projector

Let  $\Pi_N$  be the orthogonal projector from  $L^2(\mathbb{C}^n)$  onto  $B_N(\mathbb{C}^n)$ .  $\Pi_N$ . It admits a Schwartz kernel:

$$\Pi_{\mathbf{N}}(z,w) = \frac{\mathbf{N}^{\mathbf{n}}}{\pi^{\mathbf{n}}} \exp \left[ \mathbf{N}(-\frac{1}{2}|z-w|^2 + \mathrm{i}\Im(z\cdot\overline{w})) \right].$$

As  $N\to +\infty$ , the kernel is exponentially decreasing far from the diagonal. The typical interaction scale is  $N^{-1/2}$ .

### Toeplitz operators

#### Definition

Let  $h \in C^{\infty}(\mathbb{C}^n)$  a smooth bounded function, and  $N \in \mathbb{N}$ . We denote by  $T_N(h)$  the Toeplitz operator associated to h:

$$\begin{split} T_N(h) : B_N(\mathbb{C}^n) & \mapsto & B_N(\mathbb{C}^n) \\ u & \mapsto & \Pi_N(hu). \end{split}$$

If h is not bounded, we can construct  $T_N(h)$  as an unbounded operator on  $B_N(\mathbb{C}^n)$ .

The mapping  $h \mapsto T_N(h)$  is linear and adjoint-preserving. If h is real-valued, then  $T_N(h)$  is formally self-adjoint.



■ There holds  $T_N(1) = 1$ , and if h is holomorphic, then  $T_N(h) = h$ ; for instance  $T_N(z_i) = z_i$ .

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- $\text{ If } h: z \mapsto \overline{z}^\alpha z^\beta \text{, then } \mathsf{T}_N(h) = \mathsf{N}^{-|\alpha|} \vartheta^\alpha z^\beta.$
- If q is a definite quadratic form on  $\mathbb{R}^{2n}$ , then  $T_N(q)$  has a compact resolvent. The first eigenvalue  $\mu_N(q) = N^{-1}\mu_1(q)$  is positive.

$$\mu_1(q) = \mathsf{min}\,\mathsf{Sp}(\mathsf{Op}_1^W(q)) + \frac{1}{2}\,\mathsf{tr}(q)$$

## Composition and bracket

#### Proposition

Leta and b two smooth bounded functions on  $\mathbb{C}^n$ . Then there is a sequence  $(c_i)_{i\in\mathbb{N}}$  of smooth bounded functions on  $\mathbb{C}^n$ , with  $c_0=\mathfrak{a}\mathfrak{b}$  so that, as  $N\to +\infty$ , there holds:

$$T_N(a)T_N(b) = T_N(c_0) + N^{-1}T_N(c_1) + N^{-2}T_N(c_2) + \dots$$

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In particular,

$$[T_N(\mathfrak{a}),T_N(\mathfrak{b})]=\frac{i}{N}T_N(\{\mathfrak{a},\mathfrak{b}\})+O(N^{-2}).$$

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- We wish to generalize Bargmann spaces to other complex manifolds.
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- Instead of considering weighted spaces, we will consider spaces of holomorphic sections.

#### **Notations**

- M is a compact Kähler manifold, with symplectic form  $\omega$ .
- L is a complex line bundle on M, endowed with a hermitian structure h, so that the curvature of the Chern connexion is  $\omega$ .
- $N \geqslant 1$  is an integer.

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Then if s is a (continuous) section of  $L^{\otimes N}$ , one can compute

$$\|s\|_{L^2} := \int_M h_N(s(m)) \frac{\omega^{\wedge n}}{n!}.$$

By completion, one defines the Hilbert space of square-integrable sections of  $L^{\otimes N}$ .



## Hardy spaces

#### Definition

The N-equivariant Hardy space is the space  $H_N(M,L)$  of  $L^2$  and holomorphic sections of  $L^{\otimes N}$ .

This space is finite-dimensional, the dimension is polynomial in N (Riemann-Roch).

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#### **Definition**

The Szegő projector  $S_N$  is the orthogonal projector from  $L^2(M,L^{\otimes N})$  onto  $H_N(M,L)$ .

It always admits a Schwartz kernel ( as a section of  $L^{\otimes N}\boxtimes L^{\otimes -N}$ ) because  $H_N(M,L)$  is finite-dimensional.



### Toeplitz operators

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$$T_N(h): H_N(M, L) \mapsto H_N(M, L)$$
  
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 $T_N(h)$  acts on a finite-dimensional space, and it is symmetric when h is real-valued.

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Observe that, for any  $u, v \in H_N(M, L)$ , there holds

$$\langle \mathbf{u}, \mathsf{T}_{\mathsf{N}}(\mathsf{h}) \mathbf{v} \rangle = \langle \mathsf{u}, \mathsf{h} \mathsf{v} \rangle.$$



## Asymptotics for the Szegő projector

### Proposition (Boutet-Sjostrand 74)

For every  $\varepsilon > 0$  and  $k \in \mathbb{N}$  there exists C such that, for every  $N \in \mathbb{N}$ :

$$d(x,y) > \epsilon \Rightarrow |S_N(x,y)| \leqslant CN^{-k}$$

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#### Proposition (Charles 00, Zelditch 02, Ma 06)

In a convenient system of local coordinates, near any point of the diagonal, there holds:

$$S_{N}(z,w) \simeq \Pi_{N}(z,w) \left[ 1 + \sum_{k=1}^{K} N^{-k/2} b_{k}(\sqrt{N}z, \sqrt{N}w) \right]$$

# Composition and bracket

### Proposition (Charles 00, Schlichenmaier 02)

Let  $\alpha$  and b two smooth functions on M. Then there is a sequence  $(c_i)_{i\in\mathbb{N}}$  of smooth functions on M, with  $c_0=\alpha b$ , such that, as  $N\to +\infty$ , there holds:

$$T_N(a)T_N(b) = T_N(c_0) + N^{-1}T_N(c_1) + N^{-2}T_N(c_2) + \dots$$

In particular,

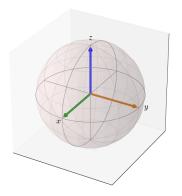
$$[T_N(\mathfrak{a}),T_N(\mathfrak{b})]=\frac{i}{N}T_N(\{\mathfrak{a},\mathfrak{b}\})+O(N^{-2}).$$

## Hardy spaces on the sphere

- $H_N(\mathbb{CP}^1, L)$  corresponds to the the set of meromorphic functions on the sphere, with one pole of order at most N.
- Hence  $H_N(\mathbb{CP}^1, L) \simeq \mathbb{C}_N[X]$ , with dimension N+1.
- One Hilbert base is:

$$e_{k,N}(X) = \frac{\binom{k}{N}^{1/2}}{N} X^k.$$

### Coordinate functions



- There are three coordinate functions on the sphere: x, y and z.
- The Toeplitz quantizations of these three functions are the spin operators, with  $S = \frac{N}{2}$ .

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### A priori localization

- In the classical model, in order to minimize the energy, one picks any point where the energy is minimal.
- What happens for an eigenvector associated with the smallest eigenvalue of  $T_N(h)$ , as  $N \to +\infty$ ?

#### Proposition (Charles 00)

An eigenvector with minimal eigenvalue is uniformly  $O(N^{-\infty})$  outside any fixed neighbourhood of  $\{h = \min(h)\}$ .

Can we get a more precise result?



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- In the classical model, in order to minimize the energy, one picks any point where the energy is minimal.
- What happens for an eigenvector associated with the smallest eigenvalue of  $T_N(h)$ , as  $N \to +\infty$ ?

#### Proposition (D. 16)

If the minimal set is non-degenerate, then for every  $\delta \in [0,1/2),$  an eigenvector with minimal eigenvalue is uniformly  $O(N^{-\infty})$  outside a neighbourhood of  $\{h=min(h)\}$  with size  $N^{-\delta}.$ 

## Proof for localization speed

Let  $(\mathfrak{u}_N)$  be a sequence of unit eigenfunctions with minimal energy  $(\lambda_N)$ . Assume  $\min(\mathfrak{h})=0$ . We prove by induction on k that

$$\langle \mathfrak{u}_n, \mathfrak{h}^k \mathfrak{u}_n \rangle = O(N^{-k}).$$

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- Hard part: k = 1 (test  $T_N(h)$  against a coherent state centred on a minimal point).
- Easy part: induction.

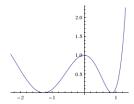
$$\langle u_n, h \star h u_n \rangle = \lambda_N^2 + O(N^{-\infty}),$$

where 
$$h\star h=h^2+N^{-1}c_1(h,h)+O(N^{-2}).$$
 Now  $c_1(h,h)\leqslant Ch.$ 



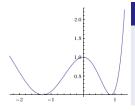
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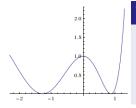


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What is minimized? The  $\mu_1$  of the hessian at this point.

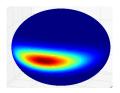
## Case of wells: idea of proof

- By making more precise the previous argument, we have a lower bound for the first eigenvalue.
- The upper bound and a spectral gap are obtained by  $N^{-K}$ -quasimode for fixed K.

We remark that the quasimodes are exponentially localized, but this does not imply that the true eigenfunction is also localized.

### Case of submanifold wells

- What can be said if h is minimal on a submanifold, with non-degenerate transverse hessian?
- $\blacksquare$   $\Rightarrow$  Same conclusion. (D.)



As N grows, the state concentrates on the miniwell and is more and more squeezed.

## Miniwells in physics

It really happens in physics! For instance, with antiferromagnetic spins on a triangle graph.



It is conjectured that the minimal configurations are planar, in some cases.

### Conjectures

Exponential Localization For now we only have  $O(N^{-\infty})$  estimates for localisation. Can we hope for  $O(\exp(-cN))$  estimates ?

Thermodynamical limit Instead of considering a fixed manifold M, we look at a particular symbol on  $M^n$ , and we let  $n \to +\infty$ . What is the behaviour vis-à-vis the semiclassical limit?

These two questions should be linked with each other.