

Transforming Twists and Wrenches with the Adjoint Map in 3D

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March 11, 2020

1 3D Rigid Transformations

The Lie group $SE(3)$ is a subgroup of the general linear group $GL(4)$ of 4×4 invertible matrices of the form

$$T \triangleq \begin{bmatrix} R & t \\ 0 & 1 \end{bmatrix}$$

where $R \in SO(3)$ is a rotation matrix and $t \in \mathbb{R}^3$ is a translation vector.

Its Lie algebra $\mathfrak{se}(3)$ is the vector space of 4×4 twists $\hat{\mathcal{V}}$ parameterized by the *twist coordinates* $\mathcal{V} \in \mathbb{R}^6$, with the mapping [1]

$$\mathcal{V} \triangleq \begin{bmatrix} \Omega \\ v \end{bmatrix} \rightarrow \hat{\mathcal{V}} \in \mathfrak{se}(3) = \begin{bmatrix} [\Omega]_{\times} & v \\ 0 & 0 \end{bmatrix}$$

where $[\cdot]_{\times}$ denotes the skew-symmetric matrix of a 3-vector, i.e., the $SO(3)$ operator. Note we follow Lynch & Park's convention and reserve the first three components for rotation, and the last three for translation.

2 A Body moving in Space

For a body undergoing a constant twist \mathcal{V}^s , all the points p^s on that body are moving according to, in homogeneous coordinates:

$$v^s = \widehat{\mathcal{V}^s} p^s$$

hence in non-homogeneous coordinates:

$$v^s = \Omega \times p^s + v$$

Which coordinates we use should be obvious from context.

3 The Adjoint Map

Suppose we want to express the twist in the body frame B ? We start from

$$v^s = \widehat{\mathcal{V}^s} p^s$$

but we know that at any given moment $v^s = T_b^s v^b$, and $p^s = T_b^s p^b$. Substituting in, we have

$$T_b^s v^b = \widehat{\mathcal{V}^s} T_b^s p^b$$

$$v^b = [T_b^s]^{-1} \widehat{\mathcal{V}}^s T_b^s p^b$$

or, in other words

$$\widehat{\mathcal{V}}^b = [T_b^s]^{-1} \widehat{\mathcal{V}}^s T_b^s = T_s^b \widehat{\mathcal{V}}^s T_b^s = T_s^b \widehat{\mathcal{V}}^s [T_s^b]^{-1}$$

The last form above we call **Adjoint Map**, with symbol and dropping the indices for now:

$$\begin{aligned} Ad_T \widehat{\mathcal{V}} &= T \widehat{\mathcal{V}} T^{-1} \\ &= \begin{bmatrix} R & t \\ 0 & 1 \end{bmatrix} \begin{bmatrix} [\Omega]_{\times} & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R^T & -R^T t \\ 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} [R\Omega]_{\times} & -[R\Omega]_{\times} t + Rv \\ 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} [R\Omega]_{\times} & t \times R\Omega + Rv \\ 0 & 0 \end{bmatrix} \end{aligned}$$

We recognize the last line as the hat of another twist, with $\Omega' = R\Omega$ and $v' = t \times R\Omega + Rv$. Hence, we can express the body twist \mathcal{V}^b in terms of the spatial twist \mathcal{V}^s by the Adjoint map as follows:

$$\mathcal{V}^b = \begin{bmatrix} \Omega^b \\ v^b \end{bmatrix} = \begin{bmatrix} R_s^b & 0 \\ [t_s^b]_{\times} R_s^b & R_s^b \end{bmatrix} \begin{bmatrix} \Omega^s \\ v^s \end{bmatrix}$$

We call this 6×6 matrix the Adjoint map matrix and write

$$\mathcal{V}^b = [Ad_{T_s^b}] \mathcal{V}^s$$

and of course vice versa:

$$\mathcal{V}^s = [Ad_{T_b^s}] \mathcal{V}^b$$

Note that there is a subtle but important distinction: Ad_T operates on 3×3 Lie algebra elements $\widehat{\mathcal{V}}$, and $[Ad_T]$ multiplies 6×1 twist coordinates \mathcal{V} .

4 Wrenches

Wrenches \mathcal{F} are dual to twists, and so they transform inversely. The dot product of a twist and a wrench is power, and of course the power has to be the same in both spatial and body frame. Hence we have

$$\mathcal{F}^{sT} \mathcal{V}^s = \mathcal{F}^{bT} \mathcal{V}^b$$

and if we substitute in $\mathcal{V}^b = \left[Ad_{T_s^b} \right] \mathcal{V}^s$, we obtain

$$\mathcal{F}^{sT} \mathcal{V}^s = \mathcal{F}^{bT} \left[Ad_{T_s^b} \right] \mathcal{V}^s$$

Since this must hold for all possible twists \mathcal{V}^s , we have

$$\mathcal{F}^{sT} = \mathcal{F}^{bT} \left[Ad_{T_s^b} \right]$$

or (transposing everything)

$$\mathcal{F}^s = \left[Ad_{T_s^b} \right]^T \mathcal{F}^b$$

Switching the indices yields the other direction:

$$\mathcal{F}^b = \left[Ad_{T_b^s} \right]^T \mathcal{F}^s$$

References

- [1] R.M. Murray, Z. Li, and S. Sastry. *A Mathematical Introduction to Robotic Manipulation*. CRC Press, 1994.