

Then \rightarrow RNC at $(0,0)$

①

$$ds^2 = -dt^2 + a(t)^2(dx^2 + dy^2)$$

Non-zero $\Gamma_{\rho\sigma}^{\mu}$:

$$\Gamma_{ij}^0 = -\frac{1}{2}(-g_{ij,0}) = \dot{a} a \delta_{ij}$$

$$\begin{aligned}\Gamma_{0j}^i &= \Gamma_{j0}^i = \frac{1}{2} g^{ik} (g_{kj,0}) = \frac{1}{2} a^{-2} \delta^{ik} \delta_{kj} 2\dot{a}a \\ &= \delta_j^i \frac{\dot{a}}{a} = \delta_j^i H.\end{aligned}$$

New coords $x^{\bar{\mu}} = A^{\bar{\mu}}_{\nu} x^{\nu} + \frac{1}{2} B^{\bar{\mu}}_{\rho\sigma} x^{\rho} x^{\sigma}$

$$\text{s.t. } g_{\bar{\mu}\bar{\nu}} = g_{\mu\nu} \Big|_p = \frac{\partial x^{\bar{\mu}}}{\partial x^{\mu}} \frac{\partial x^{\bar{\nu}}}{\partial x^{\nu}} g_{\bar{\mu}\bar{\nu}} \Big|_p \text{ at } (x^{\mu}) = (0,0)$$

$$\Rightarrow A^{\bar{\mu}}_{\nu} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & a(0) & 0 \\ 0 & 0 & a(0) \end{pmatrix}$$

$$B^{\bar{\mu}}_{\rho\sigma} = A^{\bar{\mu}}_{\rho} \Gamma^{\rho}_{\sigma} \Big|_p$$

Non-zero components of $\Gamma^{\mu}_{\rho\sigma}$ give

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$$B_{ij}^0 = \Gamma_{ij}^0|_p = \dot{a} a \delta_{ij}$$

$$B_{oj}^{\bar{t}} = A_{\bar{t}i}^{\bar{t}} B_{oj}^i|_p = a(t) \delta_{\bar{t}i}^{\bar{t}} \Gamma_{oj}^i$$

$$= a(t) \delta_{\bar{t}i}^{\bar{t}} \delta_j^i H = \delta_j^{\bar{t}} \dot{a}(t)$$

$$\Rightarrow \tilde{t} = t + \frac{1}{2} \dot{a}(t) a(t) (x^2 + y^2)$$

$$\tilde{x} = a(t) x + \frac{1}{2} \dot{a}(t) t x$$

$$\tilde{y} = a(t) y + \frac{1}{2} \dot{a}(t) t y$$

(+)

$$g^{\tilde{\mu}\tilde{\nu}} = \frac{\partial x^{\tilde{\mu}}}{\partial x^{\mu}} \frac{\partial x^{\tilde{\nu}}}{\partial x^{\nu}} g^{\mu\nu}$$

$$g^{\tilde{0}\tilde{0}} = -1 + \frac{\dot{a}(t)^2 a(t)^2}{a(t)^2} (x^2 + y^2)$$

$$g^{\tilde{1}\tilde{1}} = -\dot{a}(t)^2 x^2 + \frac{1}{a(t)^2} (a(t) + \frac{1}{2} \dot{a}(t) t)^2$$

$$= \frac{a(t)^2}{a(t)^2} (1 + \frac{\dot{a}(t)}{a(t)} t)^2 - \frac{\dot{a}(t)^2}{4} x^2$$

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$$g^{\tilde{t}\tilde{t}} = \frac{a_0^2}{a_c^2} (1 + H_0 t)^2 - \dot{a}_0^2 x^2 \quad \text{where we've defined } a_0 \equiv a(0) \\ a_c \equiv a(t) \text{ etc.}$$

Check that this has zero first-order contributions in t, x, y :

$$\begin{aligned} g^{\tilde{t}\tilde{t}} &= \frac{a_0^2}{a_0^2 (1 + H_0 t)^2} (1 + H_0 t)^2 - \dot{a}_0^2 x^2 + \text{h.o.} \\ &= \frac{a_0^2}{a_0^2 (1 + H_0 t + \frac{1}{2} \frac{\ddot{a}_0}{a_0} t^2)^2} (1 + H_0 t)^2 - \dot{a}_0^2 x^2 + \text{h.o.} \\ &= (1 - 2H_0 t - \frac{1}{2} \frac{\ddot{a}_0}{a_0} t^2) (1 + 2H_0 t + H_0^2 t^2) - \dot{a}_0^2 x^2 + \text{h.o.} \\ &= (1 - \frac{\ddot{a}_0}{a_0} t^2 + H_0^2 t^2) - \dot{a}_0^2 x^2 + \text{h.o.} \\ &= 1 + (H_0^2 - \frac{\ddot{a}_0}{a_0}) t^2 - \dot{a}_0^2 x^2 \quad \checkmark \end{aligned}$$

$$g^{\tilde{x}\tilde{x}} = \frac{a_0^2}{a_c^2} (1 + H_0 t)^2 - \dot{a}_c^2 y^2$$

$$\begin{aligned} g^{\tilde{0}\tilde{1}} &= -\frac{\partial \hat{t}}{\partial t} \frac{\partial \hat{x}}{\partial t} + \frac{1}{a_c^2} \frac{\partial \hat{t}}{\partial x} \frac{\partial \hat{x}}{\partial x} + \frac{1}{a_c^2} \frac{\partial \hat{t}}{\partial y} \frac{\partial \hat{x}}{\partial y} \\ &= -\dot{a}_0 x + \frac{1}{a_c^2} \dot{a}_0 a_0 x a_0 (1 + H_0 t) + \frac{1}{a_c^2} \dot{a}_0 a_0 y \times 0 \\ &= -\dot{a}_0 x + \frac{a_0^2}{a_c^2} (\dot{a}_0 x (1 + H_0 t)) \dot{a}_0 x \\ &= \dot{a}_0 \left(\frac{a_0^2}{a_c^2} - 1 \right) x + \frac{\dot{a}_0 H_0 a_0^2}{a_c^2} t x \end{aligned}$$

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$$g^{\hat{0}\hat{2}} = \dot{a}_0 \left(\frac{a_0^2}{a_e^2} - 1 \right) y + \frac{a_0^2 \dot{a}_0 H_0}{a_e^2} t y$$

$$g^{\hat{1}\hat{2}} = - \frac{\partial \hat{x}}{\partial t} \frac{\partial y}{\partial t} + \frac{\partial \hat{t}}{\partial x} 0 + 0$$

$$= - (1 + H_0 t) (1 + H_0 t) - \dot{a}_0 x \dot{a}_0 y$$

$$= - \dot{a}_0^2 x y$$

let's compute the Riemann tensor in (t, x, y) coordinates.

$$R^{\mu}_{\nu\rho\sigma} = \Gamma^{\mu}_{\nu\sigma, \rho} - \Gamma^{\mu}_{\nu\rho, \sigma} + \Gamma^{\mu}_{\nu\sigma} \Gamma^{\rho}_{\rho\sigma} - \Gamma^{\mu}_{\nu\rho} \Gamma^{\rho}_{\rho\sigma}$$

Non-zero derivative of Γ :

$$\Gamma^0_{ij,0} = \delta_{ij} (\ddot{a}a + \dot{a}^2)$$

$$\Gamma^i_{0j,0} = \delta^i_j \left(\frac{\ddot{a}}{a} - \frac{\dot{a}^2}{a^2} \right) = \delta^i_j \left(\frac{\ddot{a}}{a} - H^2 \right)$$

Then

$$\begin{aligned} R^0_{\nu\rho\sigma} &= \Gamma^0_{\nu\rho, \sigma} - \Gamma^0_{\nu\sigma, \rho} + \Gamma^{\mu}_{\nu\sigma} \Gamma^{\rho}_{\rho\sigma} - \Gamma^{\mu}_{\nu\rho} \Gamma^{\rho}_{\rho\sigma} \\ &= \cancel{\delta^i_v \delta^j_\sigma \delta^0_\rho (\delta_{ij} (\ddot{a}a + \dot{a}^2))} \end{aligned}$$

Check the expansion:

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$$\underline{\hat{t} = \hat{x}^0} \quad g^{\hat{0}\hat{0}} = 1^{00} + \alpha R^0{}_\nu{}^\sigma x^\nu x^\sigma$$

$$R^0{}_\nu{}^\sigma = g^{0\rho} R^0{}_{\nu\rho\sigma}$$

$$= -R^0{}_{\nu\sigma\sigma}$$

$$= - \left[\Gamma_{\nu\sigma,0}^0 - \Gamma_{\nu 0,\sigma}^0 + \cancel{\Gamma_{\nu\sigma}^k \Gamma_{k0}^0} - \cancel{\Gamma_{\nu 0}^k \Gamma_{k\sigma}^0} \right]$$

$$\Rightarrow R^0{}_i{}^0{}_j = -R^0{}_{i0j} = - \left[\Gamma_{ij,0}^0 + \cancel{\Gamma_{ij}^0 \Gamma_{00}^0} - \cancel{\Gamma_{i0}^k \Gamma_{kj}^0} \right]$$

$$= - \left[\delta_{ij} (\ddot{a}a + \dot{a}^2) - \delta_i^k H \delta_{kj} \dot{a}a \right]$$

$$= - \left(\delta_{ij} (\ddot{a}a + \dot{a}^2) - \delta_{ij} \dot{a}^2 \right)$$

$$= - \underline{\delta_{ij} \ddot{a}a}$$

$$R^0{}_0{}^0{}_0 = 0, \quad R^0{}_i{}^0{}_0 = 0, \quad R^0{}_0{}^0{}_i = 0$$

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$R'_{\nu\sigma}$:

$$R'_{00} = \frac{1}{a_e^2} R'_{010}$$

$$= \frac{1}{a_e^2} [0 - \Gamma'_{01,0} + 0 - \Gamma_{01}^k \Gamma'_{0k}]$$

$$= \frac{1}{a_e^2} [-\left(\frac{\ddot{a}}{a} - H^2\right) - \delta^k_1 H \delta'_k H]$$

$$= \frac{1}{a_e^2} \left[-\frac{\ddot{a}_e}{a_e}\right] = -\frac{\ddot{a}_e}{a_e^3}$$

$$R'_{0i} = \frac{1}{a_e^2} R'_{0ii}$$

$$= \frac{1}{a_e^2} \left[\cancel{\Gamma'_{0i,1}}^{\rightarrow 0} - 0 + \cancel{\Gamma_{0i}^k \Gamma'_{k1}}^{\rightarrow 0} - \cancel{\Gamma_{01}^k \Gamma'_{ki}}^{\rightarrow 0} \right]$$

$$= 0$$

$$R'_{ij} = \frac{1}{a_e^2} R'_{iij} = \frac{1}{a_e^2} [\Gamma_{ij}^k \Gamma'_{k1} - \Gamma_{i1}^k \Gamma'_{kj}]$$

$$= \frac{1}{a_e^2} [\Gamma_{ij}^0 \Gamma'_{01} - \Gamma_{i1}^0 \Gamma'_{0j}]$$

$$= \frac{1}{a_e^2} [\dot{a}_e a_e \delta_{ij} H_e - \dot{a}_e a_e \delta_{1i} \delta'_j \dot{a}_e H]$$

$$R'_{11} = \frac{1}{a_e^2} [\dot{a}_e^2 - \dot{a}_e^2] = 0$$

$$R'_{22} = \frac{1}{a_e^2} [\dot{a}_e^2] = \frac{\dot{a}_e^2}{a_e^2} = H_e^2 \quad R'_{12} = R'_{21} = 0.$$

New



$$g^{\mu\nu} = 1 + \left(H_0^2 - \frac{\ddot{a}_0}{a_0} \right) t^2 - \dot{a}_0^2 x^2 + \text{h.o.}$$

whereas

$$\eta'' + \alpha R'_{\mu\nu} x^\mu x^\nu = 1 + \alpha \left(H_t^2 g^2 - \frac{\ddot{a}_t}{a_t^3} t^2 \right) \\ = 1 + \alpha ($$

Don't agree!

To second order, let's invert (1) (8)
for τ (2).

$$t = \hat{t} - \frac{1}{2} H(0) \frac{\hat{x}^2 + \hat{y}^2}{(1 + H(0)t)^2}$$

~~$$\hat{t} = \left(t + \frac{1}{2} H(0) (\hat{x}^2 + \hat{y}^2) \right) (1 - 2H(0)t + 3H(0)^2 t^2) + O(t^3)$$~~

~~$$\Rightarrow \frac{3}{2} H_0^3 (\hat{x}^2 + \hat{y}^2) t^2 + (1 - H_0^2 (\hat{x}^2 + \hat{y}^2)) t + \frac{1}{2} H_0 (\hat{x}^2 + \hat{y}^2) - \hat{t} = 0$$~~

$$t(1 + H(0)t)^2 - \hat{t}(1 + H(0)t)^2 = -\frac{1}{2} H(0)(\hat{x}^2 + \hat{y}^2)$$

• ~~Figure out~~

GNC: $ds^2 = -d\hat{t}^2 + h_{ij}(t, \underline{\hat{x}}) d\hat{x}^i d\hat{x}^j$

↓

RNC: ~~need $g_{\mu\nu}$ at $(0, \underline{\hat{x}}_0)$.~~

$$ds^2 = -dt^2$$

Do inversion systematically.

We have $x^{\bar{\mu}} = A^{\bar{\mu}}_{\nu} x^{\nu} - \frac{1}{2} A^{\bar{\mu}}_{\rho} \Gamma^{\rho}_{\alpha\beta} x^{\alpha} x^{\beta}$.

~~$x^P = A^P_v x^v$~~ + RNC inverse relations: (9)

$$\# y^P = A^P_v x^v + B^P_{\rho\sigma} x^P x^\sigma. \text{ Expanded to 1}^{st} \text{ order.}$$

$$x^P_\sigma = (A^{-1})^v_\rho y^v + \# B^P_{\rho\sigma} y^P y^\sigma$$

$$\begin{aligned} y^P &= A^P_v (A^{-1})^v_\rho y^P + A^P_v C^v_{\rho\sigma} y^P y^\sigma \\ &\quad + B^P_{\rho\sigma} (A^{-1})^\rho_\alpha y^\alpha + C^\rho_{\alpha\beta} y^\alpha y^\beta \\ &\quad \# ((A^{-1})^\sigma_\alpha y^\alpha + \mathcal{O}(y^2)) \end{aligned}$$

$$\begin{aligned} y^P &= y^P + A^P_v C^v_{\rho\sigma} y^P y^\sigma \\ &\quad + B^P_{\rho\sigma} (A^{-1})^\rho_\alpha (A^{-1})^\sigma_\beta y^\alpha y^\beta \end{aligned}$$

$$\Rightarrow A^P_v C^v_{\rho\sigma} + B^P_{\alpha\beta} (A^{-1})^\alpha_\rho (A^{-1})^\beta_\sigma = 0$$

$$\Rightarrow C^v_{\rho\sigma} = (A^{-1})^P_v B^v_{\alpha\beta} (A^{-1})^\alpha_\rho (A^{-1})^\beta_\sigma$$

Now $\mathcal{B}_{\rho\sigma}^{\mu} = -\frac{1}{2} A_{\nu}^{\mu} \Gamma_{\rho\sigma}^{\nu}$ so

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$$\begin{aligned} C_{\rho\sigma}^{\mu} &= -\frac{1}{2} (A^{-1})_{\nu}^{\mu} A_{\gamma}^{\nu} \Gamma_{\alpha\beta}^{\gamma} (A^{-1})_{\rho}^{\alpha} (A^{-1})_{\sigma}^{\beta} \\ &= -\frac{1}{2} \Gamma_{\alpha\beta}^{\mu} (A^{-1})_{\rho}^{\alpha} (A^{-1})_{\sigma}^{\beta} \end{aligned}$$

So to second order we have

$$x^{\mu} = (A^{-1})_{\nu}^{\mu} y^{\nu} - \frac{1}{2} \Gamma_{\alpha\beta}^{\mu} (A^{-1})_{\rho}^{\alpha} (A^{-1})_{\sigma}^{\beta} y^{\rho} y^{\sigma}$$

If we denote y^{μ} by x^{μ} and " $(A^{-1})_{\nu}^{\mu} = A_{\nu}^{\mu}$ ",
i.e. $A_{\nu}^{\mu} A_{\rho}^{\nu} = \delta_{\rho}^{\mu}$ we have

~~$$x^{\mu} = A_{\nu}^{\mu} x^{\nu} - \frac{1}{2} \Gamma_{\alpha\beta}^{\mu} A_{\rho}^{\alpha} A_{\sigma}^{\beta} x^{\rho} x^{\sigma}$$~~

$$x^{\mu} = A_{\nu}^{\mu} x^{\nu} - \frac{1}{2} \Gamma_{\alpha\beta}^{\mu} A_{\rho}^{\alpha} A_{\sigma}^{\beta} x^{\rho} x^{\sigma}$$

Where A_{ν}^{μ} and $\Gamma_{\alpha\beta}^{\mu}$ denote the constant
matrices that satisfy:

$$g_{\mu\nu} = A_{\mu}^{\rho} A_{\nu}^{\sigma} g_{\rho\sigma}(0, \underline{x}_0)$$

$$\Gamma_{\alpha\beta}^{\mu} = \Gamma_{\alpha\beta}^{\mu}(0, \underline{x}_0).$$

What does the $t=0$ surface look like in RNCs? Also since $g_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 \\ 0 & h \end{pmatrix}$

~~$X^\mu = A^\mu_{\bar{\mu}}$~~ we have ~~$\eta_{\bar{\mu}\bar{\nu}} = A^\mu_{\bar{\mu}} A^\nu_{\bar{\nu}} g_{\mu\nu}$~~

$A = \begin{pmatrix} 1 & 0 & 0 \\ 0 & a \end{pmatrix}$ and a satisfies

~~$1 = \eta_{\bar{\mu}\bar{\nu}} a^{\bar{\mu}} a^{\bar{\nu}}$~~ where ~~$g_{\mu\nu}$~~

$1 = \underline{a}^{-1} \underline{h} \underline{a}^{-1}$ where $h_{ij} \equiv h_{ij}(a, x_0)$.

$\underline{a}^{-1} = \underline{h}^{-1}$

$1 = \underline{h} \underline{a}^{-2} \Rightarrow \underline{a}^2 = \underline{h} \Rightarrow \underline{a} = \sqrt{\underline{h}}$

• Since \underline{h} is positive-definite it has a unique positive-definite square root.

• Then for ~~$\bar{t}=0$~~

~~$X^{\bar{\mu}} =$~~

• Recall that A represents the matrix $A^{\bar{\mu}}_{\mu}$ in

$X^{\bar{\mu}} = A^{\bar{\mu}}_{\nu} X^{\nu} \neq -\frac{1}{2} A^{\bar{\mu}}_{\rho} \Gamma^{\rho}_{\mu\nu} X^{\mu} X^{\nu}$

where $X^{\bar{\mu}} = \text{RNC}$, $X^{\mu} = \text{GNC}$.

Clarify notation: as at the beginning \tilde{x}^μ denotes RNCS, x^μ denotes GNCS. 1

• We want to find the equation for the $t \rightarrow \infty$ surface in RNCS. For that we need the inverse relations derived on the last pages. We have:

$$0 = t = x^0 = A^0_{\bar{\nu}} x^{\bar{\nu}} - \frac{1}{2} \Gamma^0_{\alpha\beta} A^\alpha_{\bar{F}} A^\beta_{\bar{J}} x^{\bar{F}} x^{\bar{J}} + O(x^3)$$

where $A^\alpha_{\bar{F}}$ are the components of A^{-1} .

Now $A^0_{\bar{t}} = 0$ and

$$\Gamma^0_{\alpha\beta} = \frac{1}{2} g^{00} (g_{\alpha\beta} + g_{\beta\alpha} - g_{\alpha 0} g_{\beta 0})$$
$$= + \frac{1}{2} g_{\alpha\beta 0} \quad (= 0 \text{ unless } \alpha=i, \beta=j)$$

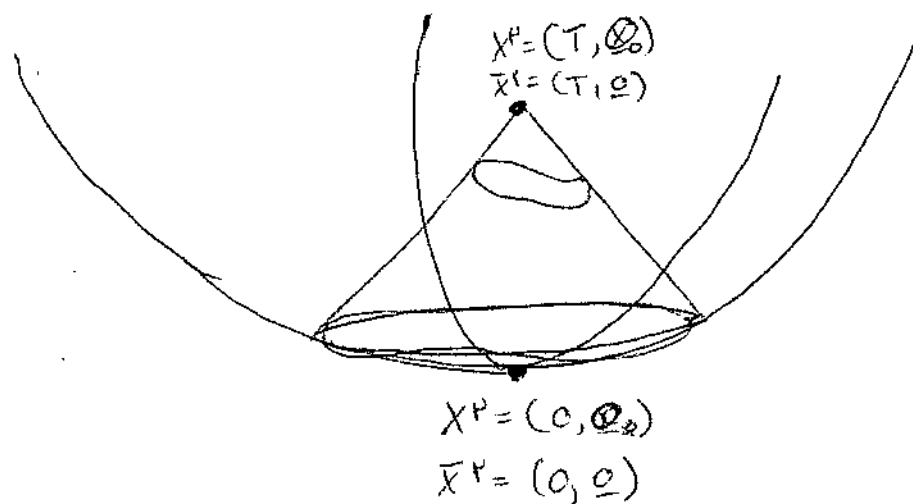
\Rightarrow only non-zero components in $\Gamma^0_{\alpha\beta}$ are Γ^0_{ij} .

So $0 = t = \frac{1}{2} \Gamma^0_{ij} A^i_{\bar{F}} A^j_{\bar{J}} x^{\bar{F}} x^{\bar{J}} + O(x^3)$

$\Rightarrow \boxed{t = \frac{1}{2} \Gamma^0_{ij} A^i_{\bar{F}} A^j_{\bar{J}} x^{\bar{F}} x^{\bar{J}}}$

So the $t=0$ surface is a
quadric surface in RNCs.

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Note that we've rescaled the GNCs ~~etc~~
to the particular point $(0, x_0)$ by defining
 $\underline{x} \rightarrow \underline{x}' = \underline{x} - \underline{x}_0$ so that $(0, x_0) \rightarrow (0, 0)$.
Note Then $ds^2 = -dt^2 + h_{ij}(t, \underline{x}' + \underline{x}_0) dx'^i dx'^j$
so the h_{ij} that enters in A & Γ is
that evaluated at $(0, x_0)$ in the original GNC.

Also note that $X^\mu = (T, 0)$ corresponds to $\bar{X}^\mu(T, 0)$

hence $\bar{X}^\mu = A^\mu_\nu X^\nu - \frac{1}{2} A^\mu_\nu \Gamma^\nu_{\rho\sigma} X^\rho X^\sigma$

$$X^0 = X^0 - \frac{1}{2} \Gamma^0_{\rho\sigma} X^\rho X^\sigma$$

$$= X^0 - \frac{1}{2} \Gamma^0_{ij} X^i X^j = 0.$$

$$X^i = 0 - A^\mu_i \Gamma^\mu_{\rho\sigma} X^\rho X^\sigma \text{ but } \Gamma^i_{00} = 0 \text{ so}$$

$$= 0 \text{ A.O.}$$

- We need to integrate from the bottom geodesic to the topmost lightcone emanating at $(T, 0)$. What's the equation for this lightcone?
- Sumati's analysis \Rightarrow only need to find lightcone in target space at $(T, 0)$; then treat light-rays as straight as any corrections due to geodesic acceleration occur at soft orders higher than T^2 .
- To find the light-cone note that at $(T, 0)$

$$g_{\bar{\mu}\bar{\nu}}(T, 0) = \sqrt{-\det \bar{g}} \frac{1}{3} T^2 R_{\bar{\mu}0\bar{\nu}0} + \mathcal{O}(T^3).$$

so the null geodesic vector $\bar{z}^{\bar{\mu}}$ satisfies

$$-(\bar{z}^0)^2 + \sum_{\bar{i}} \bar{z}^{\bar{i}2} = \frac{1}{3} R_{\bar{\mu}0\bar{\nu}0} \bar{z}^{\bar{\mu}} \bar{z}^{\bar{\nu}}.$$

Recall $R^{\bar{\mu}}_{\bar{\nu}\bar{\sigma}\bar{\omega}} = \Gamma^{\bar{\mu}}_{\bar{\sigma}\bar{\nu},\bar{\omega}} - \Gamma^{\bar{\mu}}_{\bar{\nu}\bar{\sigma},\bar{\omega}} + \Gamma^{\bar{\mu}}_{\bar{\nu}\bar{\omega}} \Gamma^{\bar{\omega}}_{\bar{\sigma}} - \Gamma^{\bar{\mu}}_{\bar{\sigma}\bar{\omega}} \Gamma^{\bar{\omega}}_{\bar{\nu}}$ so

$$R^{\bar{\mu}}_{0\bar{\omega}0} = \Gamma^{\bar{\mu}}_{00,\bar{\omega}} - \Gamma^{\bar{\mu}}_{0\bar{\omega},0} + \Gamma^{\bar{\mu}}_{\bar{\omega}0} \Gamma^0_{00} - \Gamma^{\bar{\mu}}_{0\bar{\omega}} \Gamma^0_{00}$$