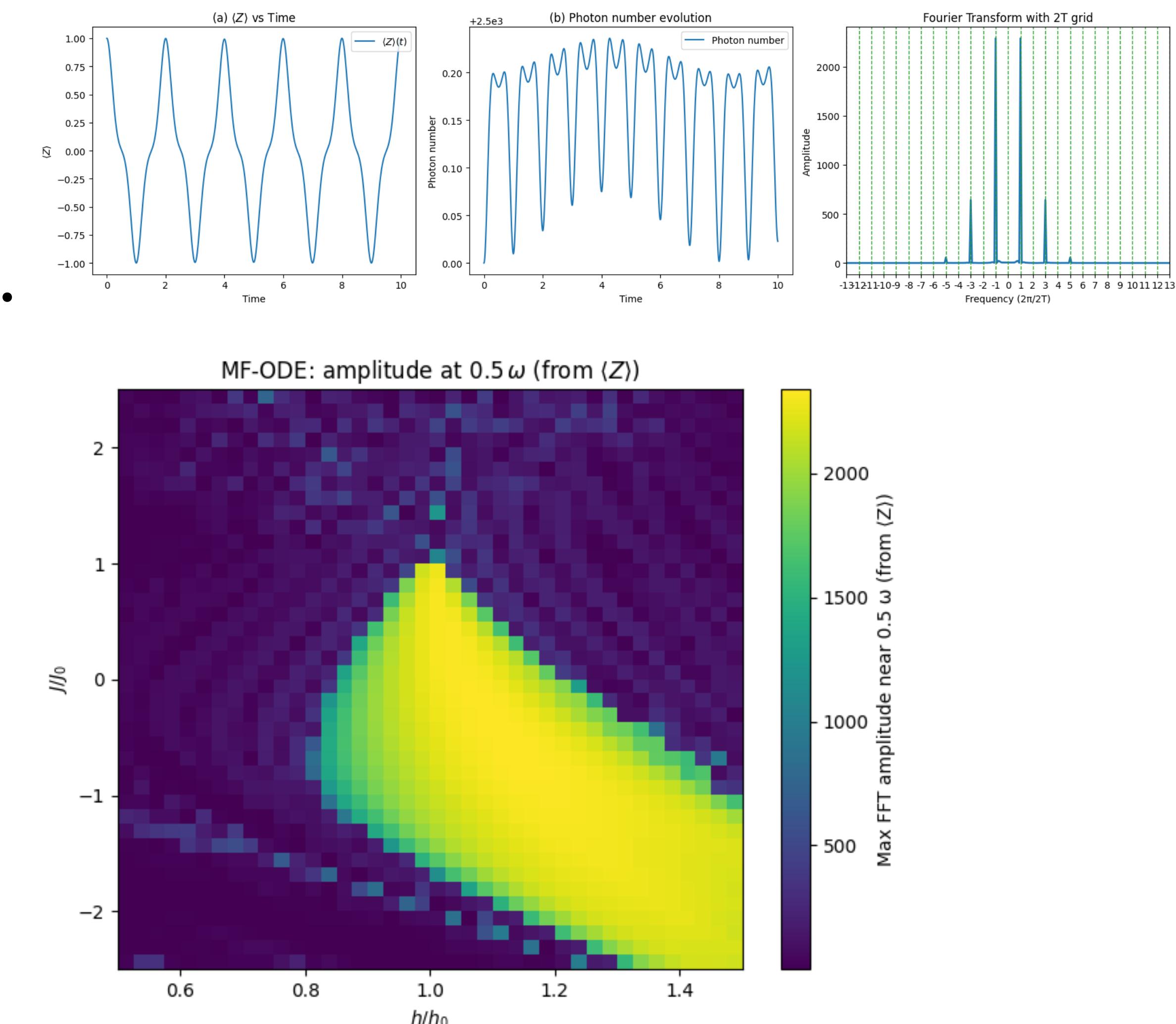


My results

Time Crystals Induced by Cavity (not published yet)

Subharmonic response in a driven all-to-all Ising model. Top: $\langle Z \rangle(t)$, photon number, and 2T-stroboscopic FFT display a pronounced $\omega/2$ peak (period-doubling) for the Hamiltonian in Eq. (1). Bottom: heatmap of the $\omega/2$ FFT amplitude from $\langle Z \rangle$ versus h/h_0 and J/J_0 , delineating a broad lobe of robust time-crystal behavior.

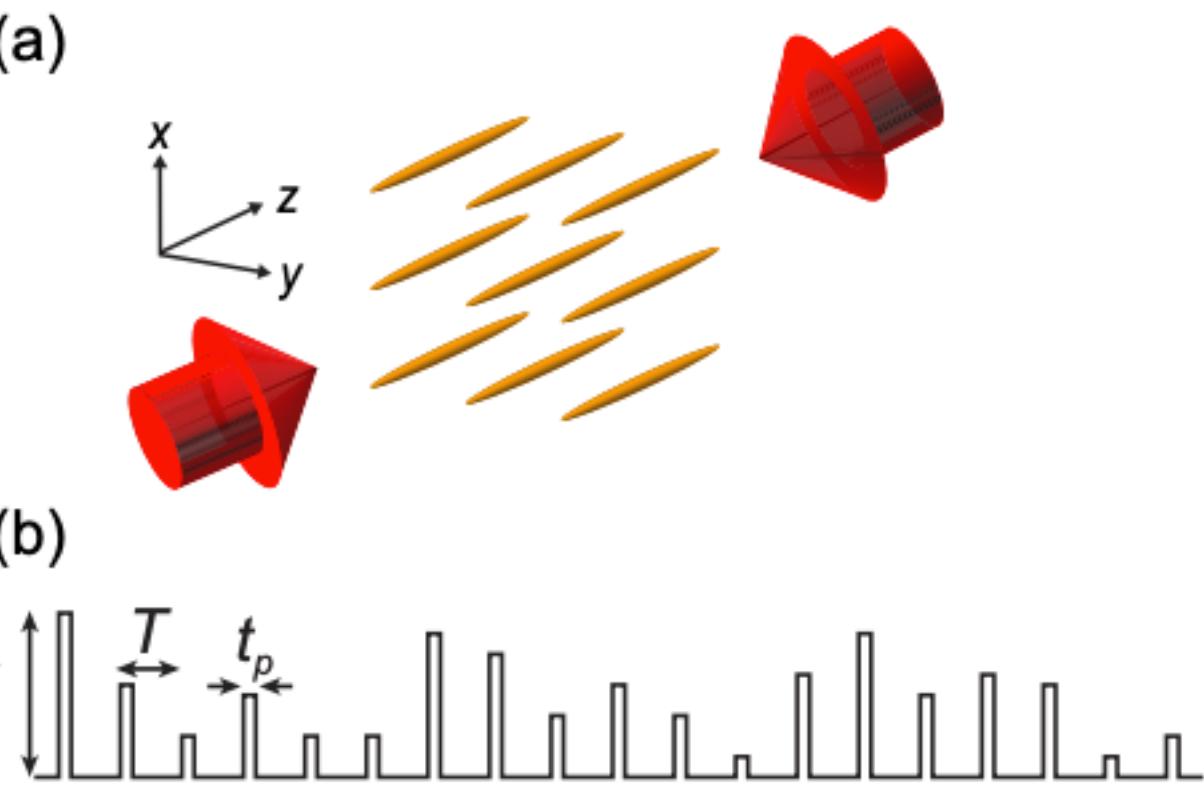
With Prof. Michael Kolodrubetz, I investigated how cavity-mediated long-range interactions can generate persistent, period-doubled correlations without external driving. By mapping between Majorana formulations and an Ising-spin description, we identified observables---stable long-time oscillations and correlations---that survive photon loss and are accessible in realistic measurements. The resulting time-independent Hamiltonian realizes the phenomenology of discrete time crystals via self-organized feedback, connecting Floquet physics to cavity QED and highlighting collective interactions as a resource for nonequilibrium order.

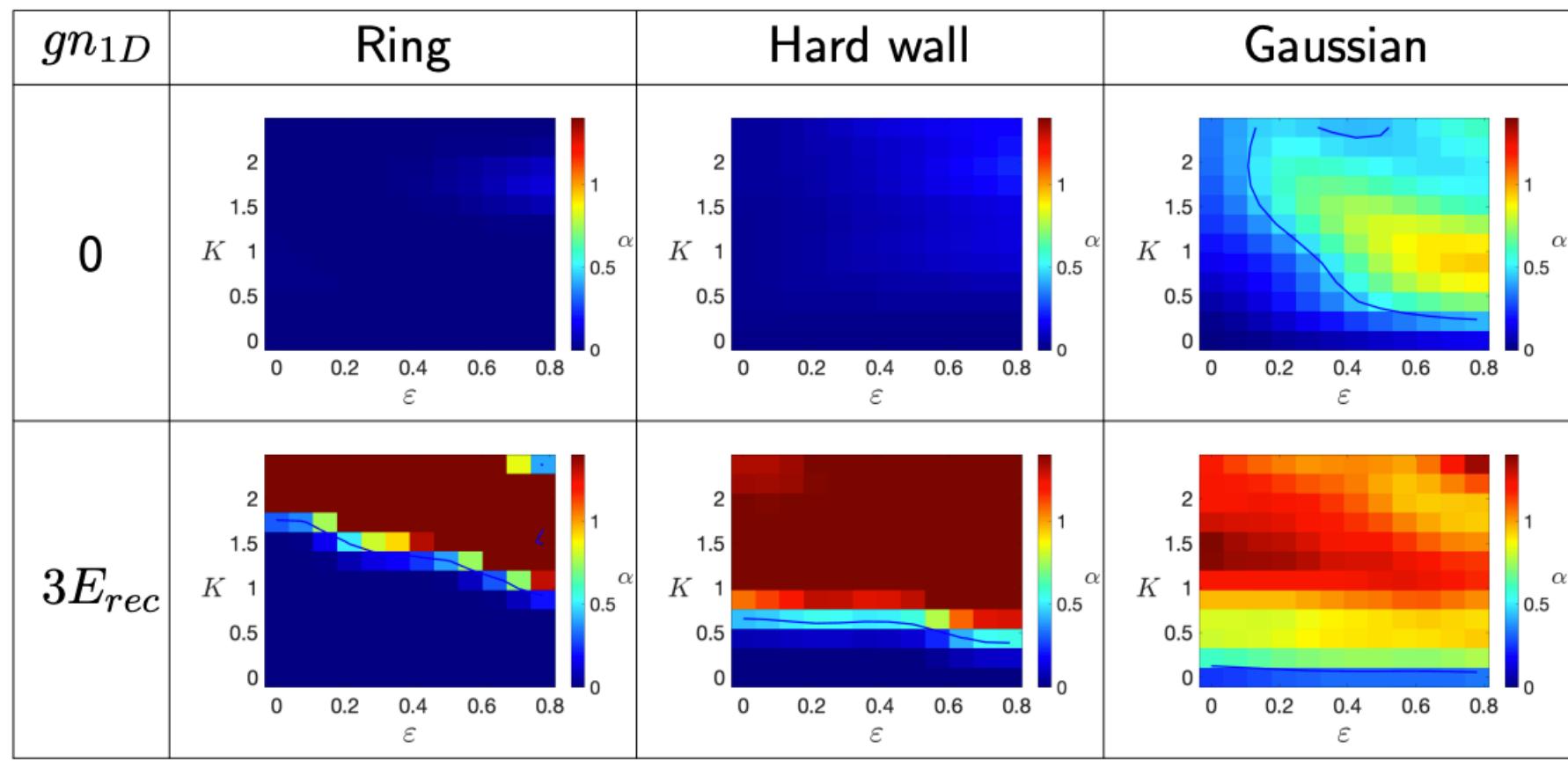


(Effect of Trap in) Quantum Kicked Rotor with Interaction

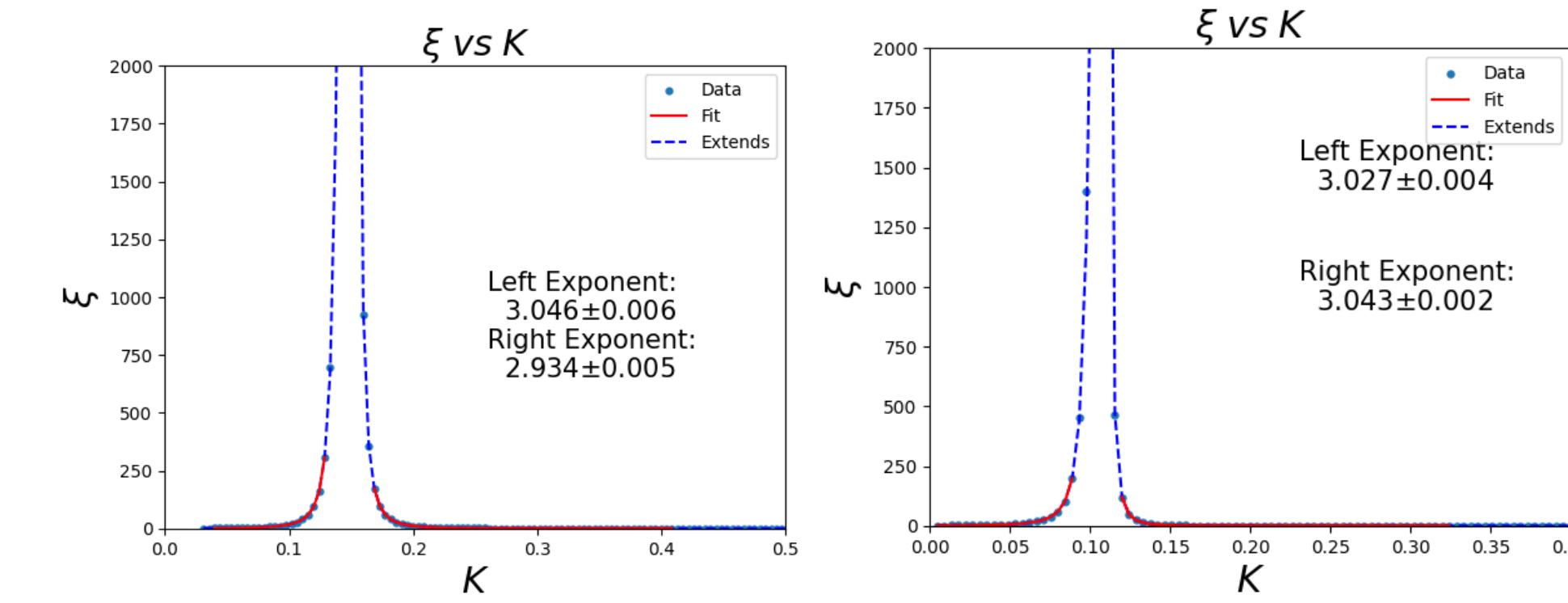
Working with Prof. Chuanwei Zhang and Prof. Subhadeep Gupta's group (UW), I studied ultracold atomic gases under quasi-periodic kicks as a controlled setting for Anderson metal--insulator physics. Using a Bose--Einstein condensate in an optical lattice, we quantified how repulsive interactions and realistic trapping potentials reshape transport. Mean-field simulations matched experiment and revealed a distinct subdiffusive regime that departs from single-particle predictions. From data--theory comparisons I extracted critical exponents indicative of an interaction-driven transition. This clarifies how many-body effects govern dynamical phases in driven disordered systems and provides a blueprint for using kicks as tunable probes of transport universality.

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- Experimental realization of the quantum kicked rotor. (a) Ultracold atoms interact with a pulsed, far-detuned standing wave formed by two counter-propagating laser beams (axes x-y-z). The optical lattice imparts the “kicks.” (b) Pulse train with period (T), duration (t_p), and kick strength (K) setting the effective stochasticity.
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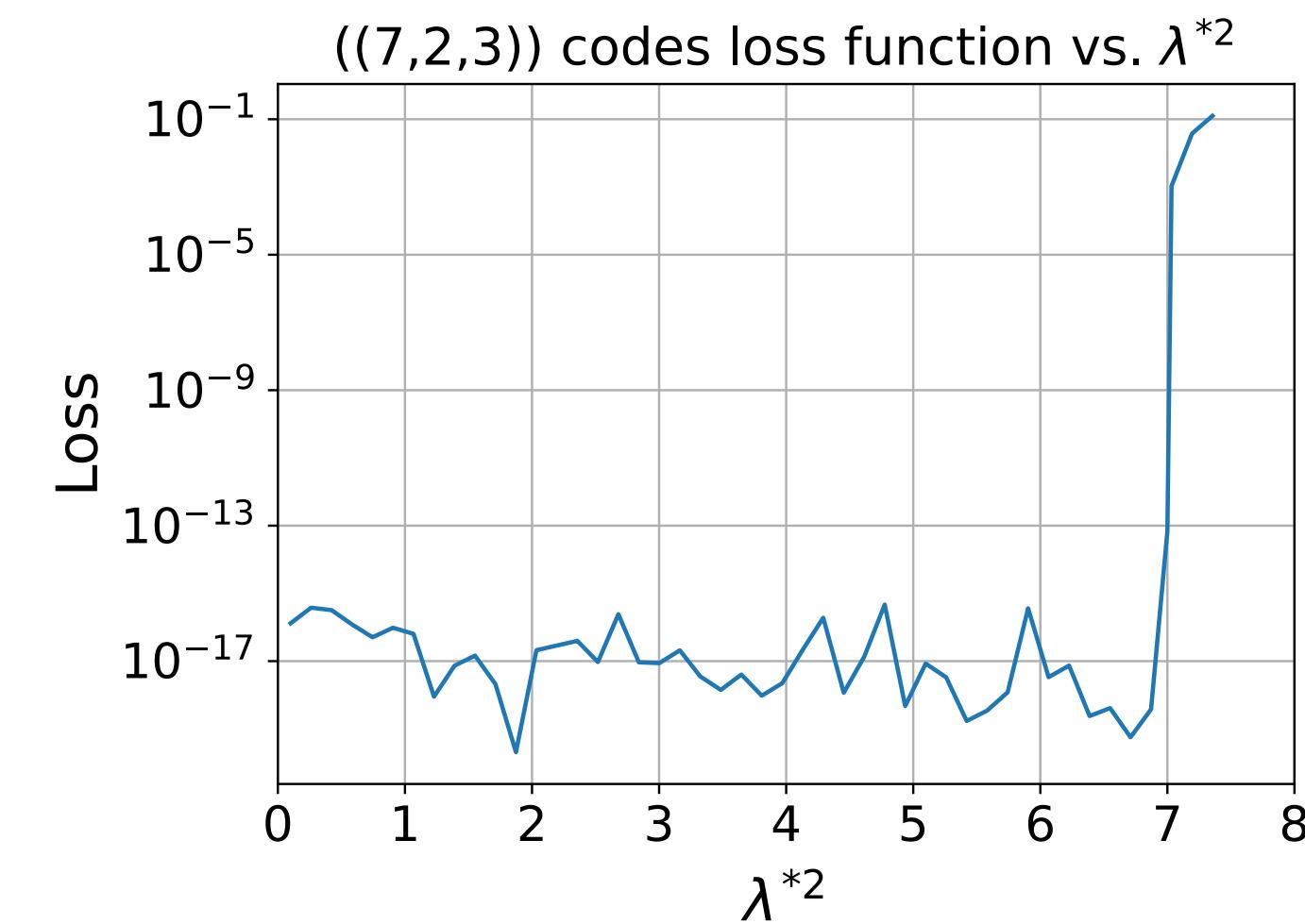
- Comparison across setups. This shows heatmaps of diffusion exponent (α) versus kick strength (K) and detuning (ε) for three geometries (Ring, Hard wall, Gaussian) and two interaction cases ($gn_{1D}=0$) and ($3E_{rec}$)).
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Localization length ($\xi(K)$) for Ring and Hard-wall cases; fits near the delocalization point yield critical exponents ($\nu \approx 3$) on each side, indicating robust, setup-insensitive behavior.

Finding Quantum Error Correction code

In collaboration with Prof.\ Bei Zeng, Prof.\ Yiu-Tung Poon, and Dr.\ Chao Zhang, I developed a framework that analyzes quantum codes through the off-diagonal Knill--Laflamme conditions. We introduced a \emph{signature vector} with locally unitary--invariant norm $\|\lambda^*\|$, which diagnoses error-correcting capability and guides code discovery. Using this tool, we uncovered new non-additive code families, including $((6,2,3))$ and $((7,2,3))$, and charted their relationships on a Stiefel-manifold picture of code spaces. The method supplies analytic traction in high-dimensional landscapes and suggests routes to robust, low-overhead protocols.



Loss landscape for $((7,2,3))$ nonadditive codes versus signature-vector magnitude ($\|\lambda^*\|^2$) (log scale). The KL-violation loss remains at machine precision across most ($\|\lambda^*\|^2$), then rises sharply near ($\|\lambda^*\|^2=7$), marking the boundary where no code simultaneously satisfies the constraints for this target ($\|\lambda^*\|$).

$$c_0 = \frac{\sqrt{\sqrt{7}\lambda^* + 8}}{8},$$

$$c_1 = -\frac{\sqrt{\sqrt{7}\lambda^*}}{8},$$

$$c_4 = -\sqrt{3}c_1,$$

$$c_3 = \frac{2}{5} \left(\sqrt{7}c_0 \pm \sqrt{7c_0^2 - \frac{15\sqrt{7}\lambda^*}{64}} \right),$$

$$c_2 = -2c_3 + \sqrt{7}c_0$$