

Lecture 22: Infinite Dimension Hilbert Spaces, Continued

Wed Oct 9 2019

Recall the position basis $\{|x\rangle\}$ and the position operator $X|x\rangle = x|x\rangle$. We are going to use these to represent an arbitrary function on the real line:

$$|\varphi\rangle = \int dx \varphi(x) |x\rangle \quad (1)$$

or

$$\varphi(x) = \langle x | \varphi \rangle \quad (2)$$

Let us now introduce a translation operator U . There could be many different transformations, so we will label ours $U(\alpha)$ such that $U(\alpha)|x\rangle = |x + \alpha\rangle$. We have chosen to symbolize this with a capital “U” because we suspect it’s unitary. Unitary transformations transform orthonormal bases to orthonormal bases. $\{|x\rangle\}$ is orthonormal, and the operator maps to states $\{|x + \alpha\rangle\} = \{|x\rangle\}$.

Let us translate the state $|\varphi\rangle$ from above. We want to evaluate some $\varphi(x)$ so we put the bra for the x states on each side. This shows that $U: \varphi(x) \rightarrow \varphi(x - \alpha)$:

$$\langle x | U(\alpha) | \varphi \rangle = \langle x | \int dx' \varphi(x') |x' + \alpha\rangle \quad (3)$$

$$= \int dx' \varphi(x') \delta(x - (x' + \alpha)) \quad (4)$$

$$= \varphi(x - \alpha) \quad (5)$$

What if we want to translate an operator? We say that $A: |\varphi\rangle \rightarrow |\chi\rangle = A|\varphi\rangle$. We want the following to happen

$$A': U|\varphi\rangle \rightarrow U|\chi\rangle = UA|\varphi\rangle \quad (6)$$

We know for a fact (from one line above) that $A': U|\varphi\rangle \rightarrow A'U|\varphi\rangle$. These must be equal, so

$$A' = UAU^\dagger \quad (7)$$

Now let us consider “infinitesimal” transformation. We consider a small δ such that $U(\delta)$ can be written as some Taylor series:

$$U(\delta) \approx U(0) + \delta \left. \frac{\partial U(\alpha)}{\partial \alpha} \right|_{\alpha=0} \quad (8)$$

$$= I - i\delta T \quad (9)$$

where

$$T \equiv i \left. \frac{\partial U}{\partial \alpha} \right|_{\alpha=0} \quad (10)$$

Additionally,

$$U^\dagger = I + i\delta T^\dagger + \dots \quad (11)$$

Together

$$UU^\dagger = I = I + i\delta(T^\dagger - T) + \mathcal{O}(\delta^2) \quad (12)$$

This order of δ must vanish, so $T = T^\dagger$, or T is hermitian.

Let us combine a finite and an infinitesimal transformation:

$$U(\alpha + \delta) = U(\delta)U(\alpha) = (I - i\delta T)U(\alpha) \quad (13)$$

Therefore,

$$\frac{\partial U}{\partial \alpha} = -iT U(\alpha) \quad (14)$$

We can solve this:

$$U(\alpha) = e^{-i\alpha T} \quad (15)$$

Now consider the infinitesimal operator acting on an arbitrary function

$$U(\delta)\varphi(x) = \varphi(x - \delta) \approx \varphi(x) - \delta\varphi'(x) \quad (16)$$

so

$$T = -i\frac{d}{dx} \quad (17)$$

so in general,

$$U(\alpha) = e^{-i\alpha\frac{d}{dx}} \quad (18)$$

We say that the derivative is the “generator” for the transformation group. This can be brought into three dimensions. Define $\vec{P} = -i\hbar\nabla$ such that $E^{-i\vec{\alpha}\cdot\vec{P}/\hbar}$.

$$X \xrightarrow{U} X' \Rightarrow X' = (I - i\delta P/\hbar)X(I + i\delta P/\hbar) \quad (19)$$

$$= X + (i\delta/\hbar)(XP - PX) + \mathcal{O}(\delta^2) \quad (20)$$

$$= X - \delta I \Rightarrow [X, P] \equiv XP - PX = i\hbar I \quad (21)$$

or by components,

$$[\vec{R}_j, \vec{P}_k] = i\hbar I \delta_{jk} \quad (22)$$

Theorem 1. Ehrenfest Theorem: Property $A = A(t)$ in $|\varphi(t)\rangle$. We want to look at

$$\langle A \rangle_{\varphi}(t) = \langle \varphi(t) | A(t) | \varphi(t) \rangle \quad (23)$$

$$\frac{d}{dt} \langle \varphi | A | \varphi \rangle = \left(\frac{d}{dt} \langle \varphi | \right) A | \varphi \rangle + \langle \varphi | \frac{d}{dt} A | \varphi \rangle + \langle \varphi | A \frac{d}{dt} | \varphi \rangle \quad (24)$$

This is equivalent to

$$\frac{1}{i\hbar} \langle \varphi | [A, H] | \varphi \rangle + \langle \varphi | \frac{d}{dt} A | \varphi \rangle \quad (25)$$