Lecture 22: Ring of Current in Cylindrical Coordinates

Mon Oct 7 2019

0.1 Ring of Current in Cylindrical Coordinates

From last time

$$\frac{1}{|\vec{x} - \vec{x}'|} = \frac{4}{\pi} \int_0^\infty dk \cos(k(z - z')) \left[\frac{1}{2} I_0(k\rho_<) K_0(k\rho_>) + \sum_{m=1}^\infty I_m(k\rho_<) K_m(k\rho_>) \cos(m(\varphi - \varphi')) \right]$$
(1)

$$\vec{A}(\vec{x}) = \frac{\mu_0}{4\pi} \int \vec{J}(\vec{x}') \frac{1}{|\vec{x} - \vec{x}'|} d^3 x'$$
 (2)

$$= \frac{\mu_0}{4\pi} \int I_0 \delta(\phi' - a) \delta(z') [-\sin(\varphi')\hat{i} + \cos(\varphi')\hat{j}]$$
 (3)

$$\times \frac{1}{|\vec{x} - \vec{x}'|} \rho' \, \mathrm{d}\rho' \, \mathrm{d}z' \, \mathrm{d}\varphi' \tag{4}$$

where $\hat{\varphi} = [-\sin(\varphi)\hat{i} + \cos(\varphi)\hat{j}]$. We can choose $\varphi = 0$ since we believe the system is symmetric about φ . By doing this, we can reduce the equation to

$$\vec{A}(\vec{x}) = \frac{\mu_0 I_0 a}{\pi} \int_0^\infty \mathrm{d}k \cos(kz) I_1(k\rho_<) K_1(k\rho_>) \hat{j} \tag{5}$$

We can then use the previous formulation to write down the elements of the \vec{B} field using the curl:

$$B_{\rho} = \frac{1}{\rho} \partial_z A_{\varphi} \tag{6}$$

and

$$B_z = \frac{1}{\rho} \partial_\rho (\rho A_\varphi) \tag{7}$$

or if we rewrite the potential with $\hat{j} = \hat{\varphi}$

$$B_{\rho} = \frac{1}{\rho} \frac{\mu_0 I_0 a}{\pi} \int_0^{\infty} dk \, (-k) \sin(kz) I_1(k\rho_{<}) K_1(k\rho_{>}) \tag{8}$$

and

$$B_z = \frac{\mu_0 I_0 a}{\pi} \int_0^\infty dk \cos(kz) \begin{cases} \frac{1}{\rho} \frac{\partial}{\partial \rho} \left[\rho I_1(ka) K_1(k\rho) \right] & \rho > a \\ \frac{1}{\rho} \frac{\partial}{\partial \rho} \left[\rho I_1(k\rho) K_1(ka) \right] & \rho < a \end{cases}$$
(9)