

First Principle 2017-Fall midterm Solution

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1. Al band structure using GGA calculation and free-electron band structure.

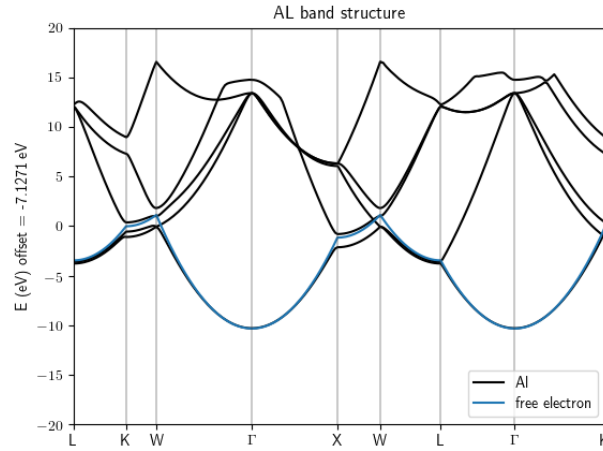


FIG. 1. Al band structure and free-electron band

In match the free electron band, we have to consider also the work function W (or equivalent the vacuum potential) of Al, in which the calculation gives : $W_{Al} = 3.207 eV$

2. Consider 1D case of KS ansatz. To derive the KS potential from the known density, we start with KS equation:

$$\left[-\frac{1}{2} \frac{\partial^2}{\partial x^2} + v_s(x) \right] \phi_i(x) = \epsilon_i \phi_i(x) \quad (1)$$

In here we set \hbar and mass m as 1.

Since there is a freedom to choose the basis, we can choose a basis that is real such that $\phi_i^* = \phi_i$. In

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which we can express the density and the kinetic density τ_L and τ as :

$$\rho(x) = \sum_i \phi_i(x)^2 \quad (2)$$

$$\tau_L = -\frac{1}{2} \sum_i \phi_i(x) \frac{\partial^2}{\partial x^2} \phi_i(x) \quad (3)$$

$$\tau = \frac{1}{2} \sum_i \left| \frac{\partial}{\partial x} \phi_i(x) \right|^2 \quad (4)$$

$$\tau_L = \tau - \frac{1}{4} \frac{\partial^2}{\partial x^2} \rho(x) \quad (5)$$

First, we multiply KS equation with ϕ_i and sum over i :

$$\tau_L + v_s(x)\rho(x) = \sum_i \epsilon_i \phi_i(x)^2 \quad (6)$$

then take derivation:

$$\frac{\partial}{\partial x} \tau_L + v_s(x) \frac{\partial}{\partial x} \rho(x) + \rho(x) \frac{\partial}{\partial x} v_s(x) = 2 \sum_i \epsilon_i \phi_i(x) \frac{\partial}{\partial x} \phi_i(x) \quad (7)$$

Now return to equation (1) and multiply by $2 \frac{\partial}{\partial x} \phi_i(x)$ and sum over i :

$$- \sum_i \frac{\partial^2}{\partial x^2} \phi_i(x) \frac{\partial}{\partial x} \phi_i(x) + v_s(x) \frac{\partial}{\partial x} \rho(x) = 2 \sum_i \epsilon_i \phi_i(x) \frac{\partial}{\partial x} \phi_i(x) \quad (8)$$

combine with equation (7) and (8), we have :

$$\frac{\partial}{\partial x} \tau_L + \sum_i \frac{\partial^2}{\partial x^2} \phi_i(x) \frac{\partial}{\partial x} \phi_i(x) + \rho(x) \frac{\partial}{\partial x} v_s(x) \quad (9)$$

$$= \frac{\partial}{\partial x} \tau_L + \frac{\partial}{\partial x} \tau + \rho(x) \frac{\partial}{\partial x} v_s(x) = 0 \quad (10)$$

With given density $\rho(x)$:

$$\rho(x) = A e^{-\alpha x^2} \quad (11)$$

the kinetic energy densities are obtained:

$$\tau = \frac{1}{2} A \alpha^2 x^2 e^{-\alpha x^2} \quad (12)$$

$$\tau_L = \frac{A \alpha}{2} e^{-\alpha x^2} [1 - \alpha x^2] \quad (13)$$

insert (12),(13) into (10) we have :

$$\rho(x) \frac{\partial}{\partial x} v_s(x) = \alpha^2 x \rho(x) \quad (14)$$

Thus we have derived the potential with constant c :

$$v_s(x) = \frac{\alpha^2}{2}x^2 + c \quad (15)$$

Finally, the normalization constant A can be derived with :

$$\int_{-\infty}^{\infty} \rho(x) = 1 \quad (16)$$

$$A \int_{-\infty}^{\infty} e^{-\alpha x^2} = A \sqrt{\frac{\pi}{\alpha}} = 1 \quad (17)$$

$$A = \sqrt{\frac{\alpha}{\pi}} \quad (18)$$

3. Car-Parrinello EOM

The Car-Parrinello lagrangian:

$$L = \sum_i \frac{1}{2} \mu \int |\dot{\phi}_i(r)|^2 dr + \sum_I \frac{1}{2} M_I \dot{R}_I^2 - E[\phi_i, R_i] - \sum_{ij} \Lambda_{ij} \left(\int \phi_i(r) \phi_j(r) dr - \delta_{ij} \right) \quad (19)$$

To impose the orthonormality, we add a term with lagrangian multiplier Λ_{ij} . With the Euler-Lagrangian equation:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} = \frac{\partial L}{\partial q} \quad (20)$$

The equation of motion of electronic and ionic DOF can be derived as:

$$\mu \ddot{\phi}_i(r) = - \sum_j \Lambda_{ij} \phi_j(r) - \frac{\partial E[\phi_i, R_I]}{\partial \phi_i} \quad (21)$$

$$M_I \ddot{R}_I = \frac{-\partial E[\phi_i, R_I]}{\partial R_I} \quad (22)$$

4. GGA, GGA+U of MnO in AF-II The crystal and magnetic structure of MnO in AF-II:

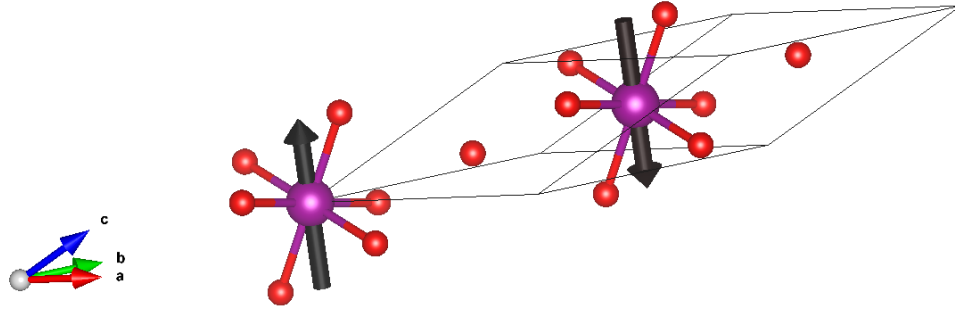


FIG. 2. MnO in AF-II, (**purple**: *Mn* atoms carry magnetic moment; **red**: non-magnetic *O* atoms)

• GGA

- (1) the band structure and total DOS, the energies are shifted by fermi energy $E_f = 5.2304\text{eV}$ to zero.

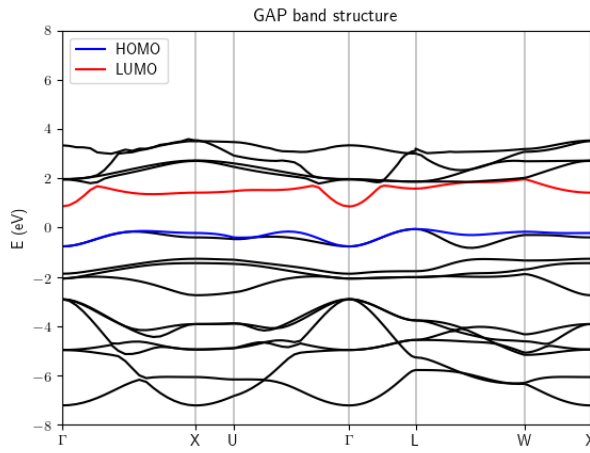


FIG. 3. MnO-AFII band-structure

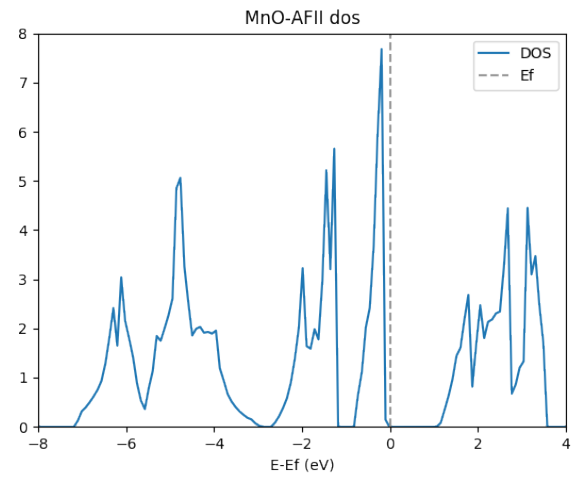


FIG. 4. MnO-AFII density of state

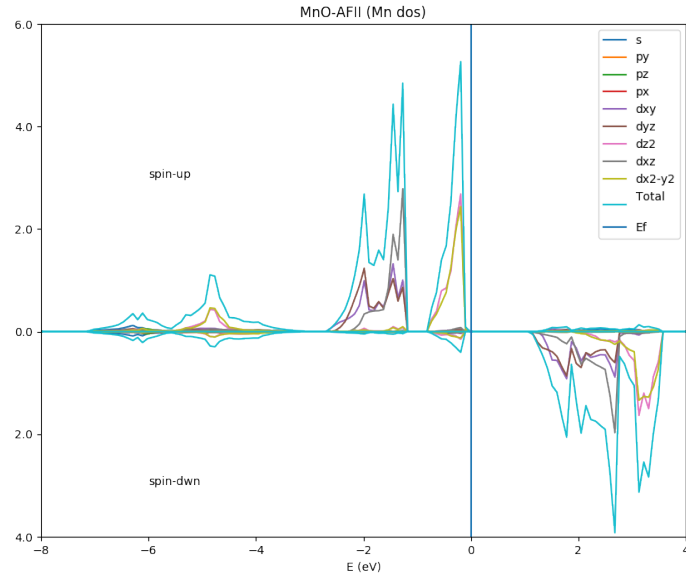


FIG. 5. Mn in MnO density of state

Energy gap E_g , total magnetic moment m and Mn moment m_{Mn} :

$$E_g = 0.918700 eV$$

$$m = 0.0000 \mu_B$$

$$m_{Mn} = 4.199 \mu_B$$

- GGA+U

(1) the band structure and total DOS, the energies are shifted by fermi energy $E_f = 4.0271/eV$ to zero.

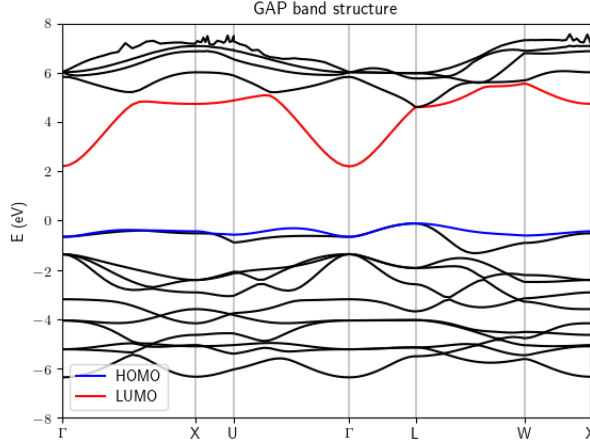


FIG. 6. MnO-AFII band-structure GGA+U

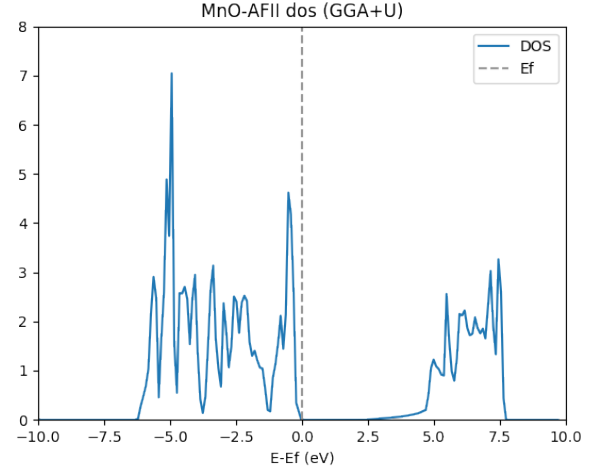


FIG. 7. MnO-AFII density of state GGA+U

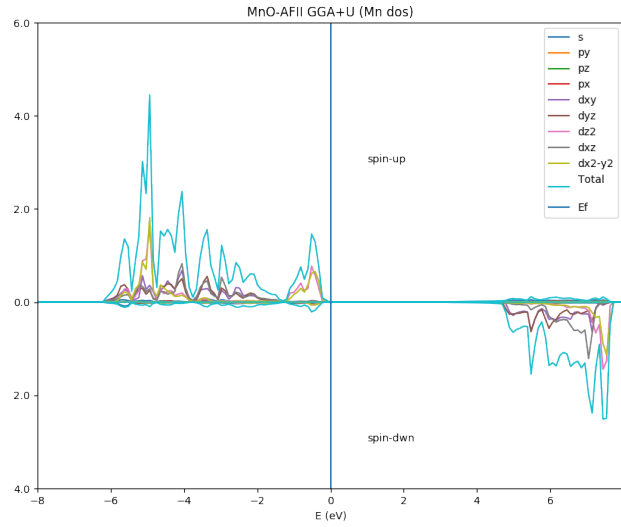


FIG. 8. Mn in MnO density of state GGA+U

Energy gap E_g , total magnetic moment m and Mn moment m_{Mn} :

$$E_g = 2.318900eV$$

$$m = 0.0000\mu_B$$

$$m_{Mn} = 4.686\mu_B$$

By apply U_{eff} on Mn, we decouple the Mn and O energy part, as a result, the AF magnetic property contributed from Mn can be calculate more accurate.

5. Finite difference algorithms.

In the following, we evaluate the harmonic oscillator with "Euler", "Predictor-Corrector" and "Velocity-verlet" method.

(a) Euler method with different dt (in unit of π) for initial condition $x(0) = 1, v(0) = 0$:

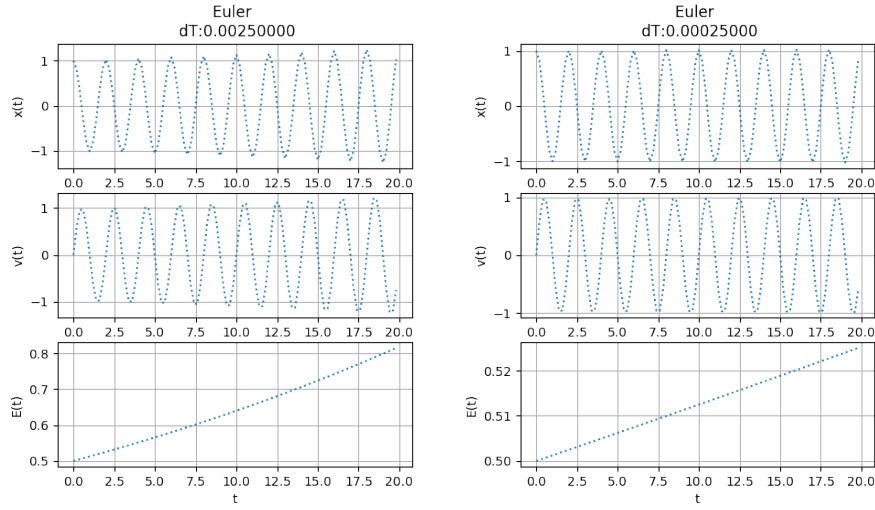


FIG. 9. Euler method, $dt = 2.5 \times 10^{-3}$ (left) and $dt = 2.5 \times 10^{-4}$ (right)

compare the result with (a), we can see that reducing the update time interval, the energy still not conserved, but the error is decreased.

- (b) Euler method compare with Predictor-Corrector method with $dt = 2.5 \times 10^{-3}$ (in unit of π), $x(0) = 1, v(0) = 0$

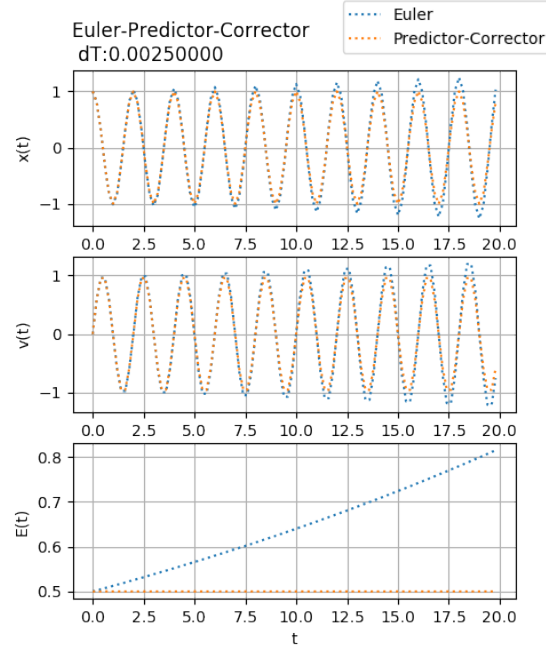


FIG. 10. Euler v.s. Predictor-Corrector method , $dt = 2.5 \times 10^{-3}$

compare the result with (a), we can see that with the same update time interval (dt), we can see that the energy increment error is significantly reduced.

- (c) Euler method compare with Velocity-Verlet method with $dt = 2.5 \times 10^{-3}$ (in unit of π), $x(0) = 1, v(0) = 0$

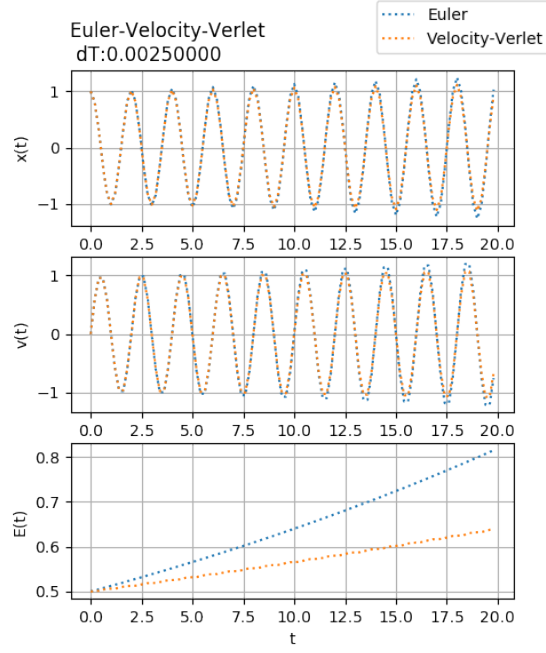


FIG. 11. Euler v.s. Velocity-Verlet method , $dt = 2.5 \times 10^{-3}$

compare the result with (a), we can see that with the same update time interval (dt), we can see that the energy increment error is reduced.