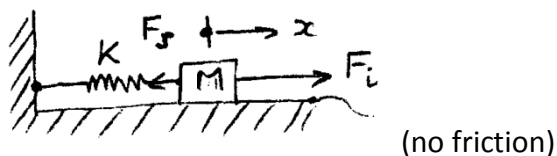


LECTURE 4 – THE DERIVATION OF SYSTEM EQUATIONS

To MODEL a physical system into a LINEAR system, we must describe each of the components of the system with its governing mathematical equation. Then we combine the equations into a complete system specifying the input and output.

Then we ANALYSE the LINEAR system by discovering its behaviour as it goes from one state to another.

An example with a simple system:



The mass m is initially at rest. Assume the spring is neither compressed nor extended. Put a disturbing force F_i on the mass but unspecified in $F_i(t)$ and allow the mass to displace by $x(t)$.

Newton's 2nd Law says $F_i - F_s = M\ddot{x}$ (Ignore (t) in $F(t)$ and $x(t)$)

However, the spring element has an equation $F_s = kx$

$\therefore M\ddot{x} + kx = F_i$, a second order differential equation where F_i is the input and x is the output.

Take LAPLACE TRANSFORM of the differential equation:

$M(s^2X(s) - sx(o) - \dot{x}(o)) + KX(s) = F_i(s)$ into the complex domain.

We know $x(o) = 0$ (not compressed), $\dot{x}(o) = 0$ (not moving)

$\therefore X(s)(Ms^2 + K) = F_i(s)$ or

$\frac{X(s)}{F_i(s)} = \frac{1}{Ms^2 + K} \equiv G(s)$ This is the TRANSFER FUNCTION of the system given.

Note the relationship is independent of magnitude of the input $F_i(s)$.

We are ready to analyse the system but must specify our input to see what it does to the output.

Imagine F_i is a unit impulse (you give M a quick flick)

$$\therefore F_i(s) = \mathcal{L}\{\delta(t)\} = 1 \quad \therefore X(s) = \frac{1}{Ms^2 + K} \cdot 1$$

$$\therefore x(t) = \mathcal{L}^{-1} \left\{ \frac{\frac{1}{M}}{s^2 + \frac{K}{M}} \right\} = \mathcal{L}^{-1} \left\{ \frac{\frac{1}{\sqrt{MK}} \cdot \sqrt{\frac{K}{M}}}{s^2 + \frac{K}{M}} \right\} = \frac{1}{\sqrt{MK}} \cdot \sin \left(\sqrt{\frac{K}{M}} \cdot t \right),$$

which is taken from the standard form:

$$\mathcal{L}^{-1} \left\{ \frac{\omega}{s^2 + \omega^2} \right\} = \sin \omega t$$

Notice we fed in boundary conditions at $t = 0$. They happened to be all zero but this does not always occur.

A closer look at the TRANSFER FUNCTION

There are basically two ways to arrange the transfer function terms:

1. Non-dimensional form, e.g. $\frac{1/K}{\frac{M}{K}s^2 + 1}$
2. Normalised form, e.g. $\frac{1/M}{s^2 + \frac{K}{M}}$

From the normalised form, the roots are obvious $s = \pm \sqrt{-\frac{K}{M}}$

From the non dimensional form we can deduce a few properties. Take $G(S) = \frac{A}{(s+a)(s+b)}$ as an example (normalised). Non-dimensionally $G(S) = \frac{A/ab}{\left(\frac{s}{a}+1\right)\left(\frac{s}{b}+1\right)}$. From this form, we can pick at the following:

A/ab is the DC gain (or steady state gain) of the system (the gain at zero frequency). $\frac{1}{a}, \frac{1}{b}$ are the TIME CONSTANTS of the system (the exponential coefficients that will appear in time domain answers).

To prove A/ab is the DC gain, let us use the FINAL VALUE THEOREM to find this gain, i.e. when $s \rightarrow 0$ or $t \rightarrow \infty$. We must specify what the input is first, i.e. input = $\frac{1}{s}$, a unit step.

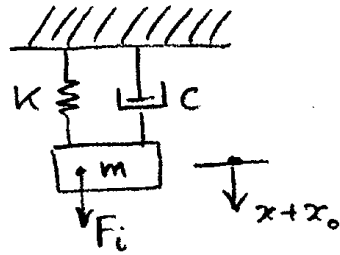
$$\therefore \lim_{t \rightarrow \infty} G(S) \cdot \frac{1}{s} = \lim_{s \rightarrow 0} \left(s \cdot \frac{A}{(s+a)(s+b)} \cdot \frac{1}{s} \right) = \frac{A}{ab} \text{ QED.}$$

HELPFUL HINTS TO GET TO A TRANSFER FUNCTION

1. Specify exactly what the start ($t = 0$) conditions are for the system; the boundary conditions. Make as many of them equate to zero as possible by moving datum.
2. Try to write equations that deal with a change of state rather than going from state 1 to state 2 e.g. look for a change of temperature on ambient, not T_1 to T_2 .

Some examples of system modelling to a transfer function

1. For the mass-spring-damper system, derive the differential equation relating displacement x to force F_i . Assume the equipment is in equilibrium before F_i is applied and x_o is the extension of the spring due to the dead weight of the mass m .



$$\therefore F_i + mg = K(x + x_o) + C(\dot{x} + \dot{x}_o) + m(\ddot{x} + \ddot{x}_o)$$

$$\therefore F_i = Kx + Kx_o - mg + m\ddot{x} + C\dot{x}$$

$$\text{But } Kx_o = mg \quad \therefore m\ddot{x} + C\dot{x} + Kx = F_i$$

Here we could have assumed x was the incremental extension.

Now take Laplace to get $X(s)(ms^2 + Cs + K) = F_i(s)$ where $x(0) = \dot{x}(0) = 0$.

$$\therefore \frac{X(s)}{F_i(s)} = \frac{1}{ms^2 + Cs + K} \equiv \frac{\frac{1}{m}}{s^2 + \frac{C}{m}s + \frac{K}{m}}$$

Hence $\frac{1}{K}$ is the DC gain, i.e. if F_i is not changing, $x = \frac{1}{K} \cdot F_i$ QED

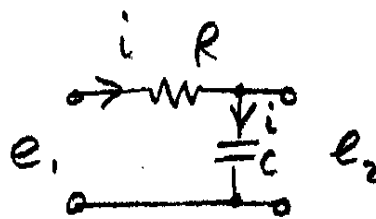
2. For the electrical circuit, derive the transfer function.

Assume zero boundary conditions

$$\therefore e_1 = iR + e_2, \quad i = C\dot{e}_2$$

$$\therefore RC\dot{e}_2 + e_2 = e_1$$

e_2 is output and e_1 is input.



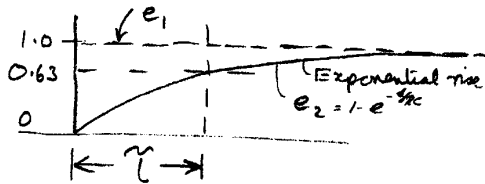
$$\therefore E_2(s)(RCs + 1) = E_1(s) \quad \therefore \frac{E_2(s)}{E_1(s)} = \frac{1}{RCs + 1}$$

We see RC is the time constant of this system and its DC gain is 1. For a unit step input:

$$E_2(s) = \frac{1}{s(RCs + 1)} = \frac{\frac{1}{RC}}{s(s + \frac{1}{RC})}$$

$$e_2 = 1 - e^{-\frac{t}{RC}}$$

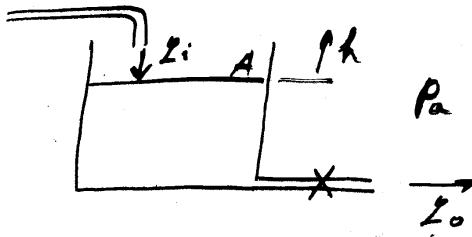
Graphically:



$\tau = RC$ is the time taken to get 63% towards its final value at $t \rightarrow \infty$.

3. Water tank system:

Create a model in the time domain. Derive the transfer function. Analyse the unit step response.



Flow through the valve:

$$\begin{aligned} q_0 &= K\Delta p \\ &= K[(p_a + \rho gh) - p_a] \\ &= K\rho gh \end{aligned}$$

Differential equations:

$$q_0 = \frac{h}{R} \quad (1)$$

$$q_i - q_0 = A\dot{h} \quad (2)$$

Combining (1) and (2) yields:

$$A\dot{h} + \frac{h}{R} = q_i$$

Laplace transform:

$$\left(As + \frac{1}{R}\right)H(s) = Q_i(s)$$

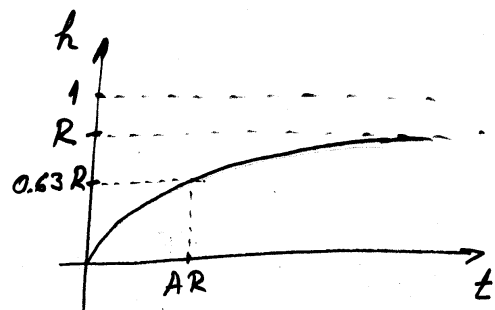
$$\frac{H(s)}{Q_i(s)} = \frac{1}{As + \frac{1}{R}} = \frac{R}{ARS + 1} = \frac{\frac{1}{A}}{s + \frac{1}{AR}}$$

DC gain: R ; time constant: AR

Assume:

$$Q_i(s) = \frac{1}{s}$$

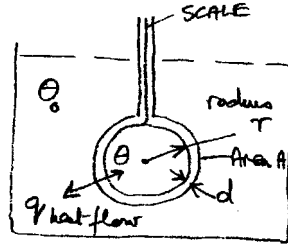
$$H(s) = \frac{\frac{1}{A}}{s(s + \frac{1}{AR})}$$



$$h(t) = \mathcal{L}^{-1}\{H(s)\}$$

$$h(t) = R \left(1 - e^{-\frac{t}{AR}} \right)$$

4. For a thermometer (thermal system) we will look at the dynamics of the temperature indicated on the scale. The bulb is glass of conductivity K containing mercury of characteristics $\rho c V$. The bulb is at θ_1 prior to being placed in the bath at θ_0 . Let us find out how the temperature indicated changes with time.



Conductivity Law:

$$q_c = -\frac{KA}{d} \cdot \frac{d\theta}{dx} \equiv \frac{KA}{d} (\theta_0 - \theta) \quad (1)$$

Internal Energy Law:

$$q_I = C_t \frac{d\theta}{dt} \equiv \rho c V \dot{\theta} \quad C_t \text{ is heat capacity } \rho c V \quad (2)$$

Balancing 1 against 2 gives

$$q_c - q_I = \frac{KA}{d} (\theta_0 - \theta) - \rho c V \dot{\theta} \equiv 0$$

$$\therefore \dot{\theta} - \frac{KA}{\rho c V d} (\theta_0 - \theta) = 0 \quad \text{and let } \frac{A}{V} \cong \frac{4\pi r^2}{\frac{4}{3}\pi r^3} = \frac{3}{r}$$

Rearranging to treat θ_0 as the input and θ as output

$$\dot{\theta} + \alpha \theta = \alpha \theta_0, \text{ where } \alpha = \frac{3K}{\rho c d r}$$

Take Laplace to give

$$s\theta(s) - \theta(0) + \alpha\theta(s) = \alpha\theta_0(s)$$

The trick is to treat the system as for incremental changes by making $\theta(0) = 0$, in which case $\theta(s)$ is the difference over ambient, not an absolute value.

$$\therefore \text{The transfer function is } \frac{\theta(s)}{\theta_0(s)} = \frac{\alpha}{s+\alpha}$$

Now specify the driving function $\theta_0(s)$ as a step function starting at $t = 0$ of magnitude θ_B (difference between ambient and fluid temperatures). This is like saying at $t=0$ we plunge the bulb into a water bath and swish it about all the time. $\therefore \theta(s) \equiv \frac{\theta_B}{s}$

$$\therefore \theta(s) = \frac{\alpha \theta_B}{s(s+\alpha)} \equiv \frac{\theta_B}{s(Ts+1)} \quad \text{where} \quad T = \frac{1}{\alpha} = \frac{\rho c d r}{3K}$$

With inverse Laplace, we get

$$\theta(t) = \theta_B(1 - e^{-\alpha t}), \text{ a normal exponential change.}$$

Finally, let us calculate α or T for a typical thermometer.

For mercury, $\rho = 1.35 \times 10^3 \text{ Kg m}^{-3}$, $c = 1.37 \times 10^2 \text{ J kg}^{-1} \text{ K}^{-1}$.

Let $d = 0.5 \text{ mm}$, $K = 1.1 \text{ Js}^{-1} \text{ m}^{-1} \text{ K}^{-1}$, $r = 2.0 \text{ mm}$.

$$T = \frac{\rho c d r}{3K} = \frac{1.35 \cdot 10^4 \cdot 1.37 \cdot 10^2 \cdot 5 \cdot 10^{-4} \cdot 2 \cdot 10^{-3}}{3 \cdot 1 \cdot 1} = 5.6 \cdot 10^{-1} \text{ s}$$

The time constant is $\sim \frac{1}{2}$ sec, so if we wait ~ 5 time constants, i.e. 2.5 seconds, the reading should be virtually at its final value. In other words, this thermometer stabilises very quickly.

