MHPC Molecular dynamics

Sabine Reißer

Thermostat

```
// Read initial positions
read positions
// Set random velocities
randomize velocities
// Compute initial neighbor list
compute list
// compute initial forces
compute forces
for(istep=0;istep<nstep;istep++){</pre>
    thermostat (dt/2)
    velocity += (force*dt/2)/mass
    position += velocity*dt
    // Check whether the neighbour list has to be recomputed
    check list
    if(recompute list) {
      compute list
    compute forces
    velocity += (force*dt/2)/mass
    thermostat (dt/2)
```

Thermostat

Energy constant

Number of particles constant

Volume constant

NVE

(microcanonical ensemble)

Temperature constant

Number of particles constant

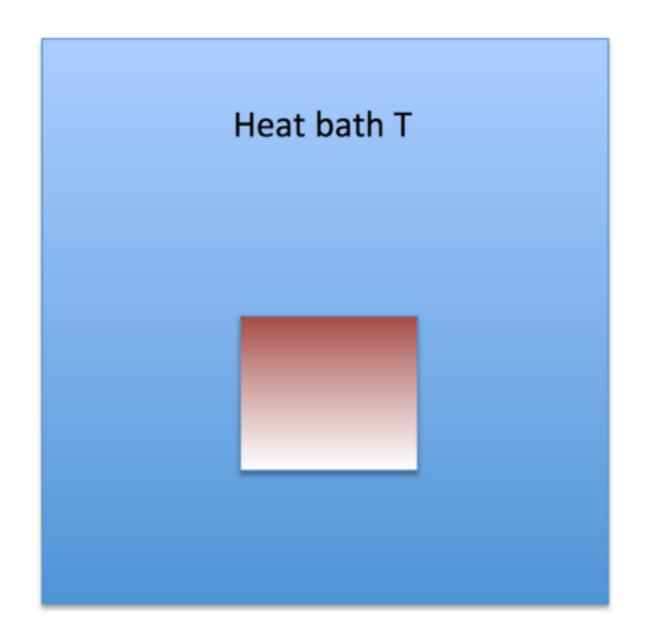
Volume constant

NVT

(Canonical ensemble)

Thermostat: algorithm to control temperature in MD simulation

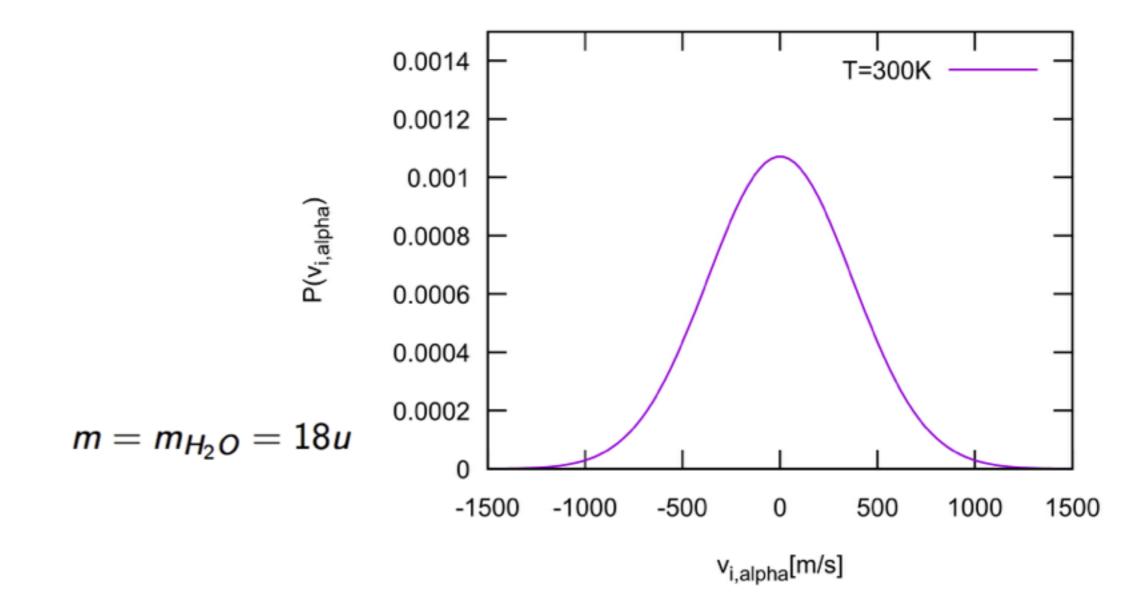
Canonical ensemble



System at thermal equilibrium with heat bath at temperature T

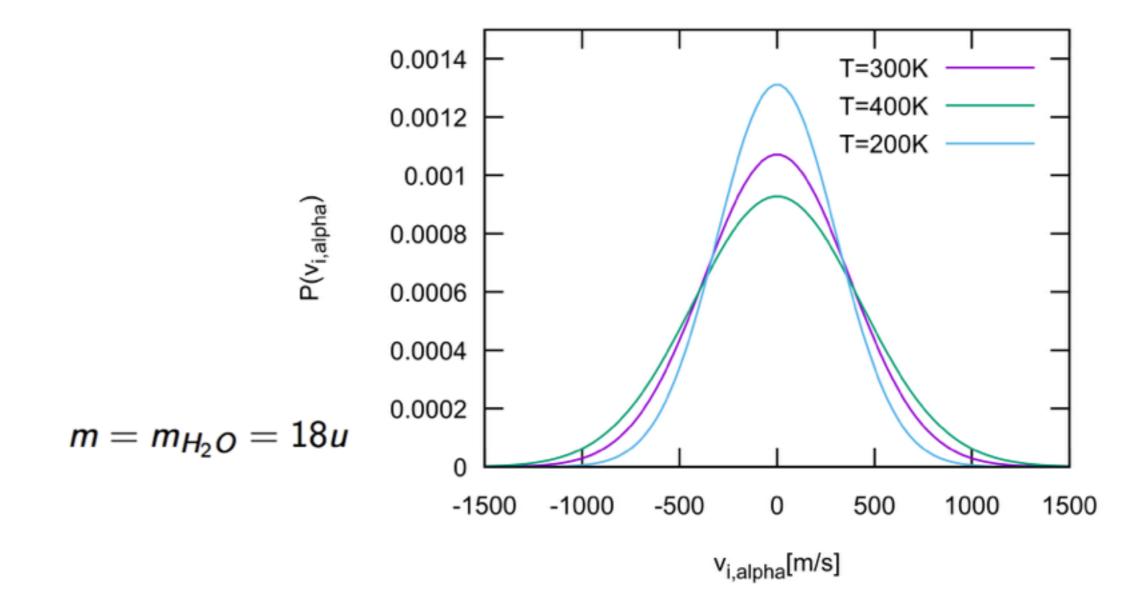
Velocities according to Maxwell-Boltzmann distribution

$$P(v_{i,\alpha}) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{m v_{i,\alpha}}{2k_B T}\right)$$



Velocities according to Maxwell-Boltzmann distribution

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Velocities according to Maxwell-Boltzmann distribution

$$P(v_{i,\alpha}) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{mv_{i,\alpha}}{2k_B T}\right)$$

$$0.002$$

$$P(v_i) = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} v_i^2 \exp\left(-\frac{mv_i^2}{2k_B T}\right)$$

$$0.0015$$

$$\frac{2}{\Delta}$$

$$0.0005$$

$$0.0005$$

$$0.0005$$

$$0.0005$$

$$0.0005$$

$$0.0000$$

$$0.0000$$

$$0.0000$$

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$$0.0015$$

$$\frac{2}{\alpha}$$

$$0.0005$$

$$m = m_{H_2O} = 18u$$

$$0$$

$$0.0005$$

$$0.0005$$

$$0.0005$$

$$0.0005$$

 $v_i[m/s]$

Typical speed ≈ speed of sound

Velocities according to Maxwell-Boltzmann distribution

$$P(v_{i,\alpha}) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{m v_{i,\alpha}}{2k_B T}\right)$$

Equipartition theorem:

average kinetic energy per degree of freedom related to T

$$\frac{1}{2}m_i\left\langle v_{i,\alpha}^2\right\rangle = \frac{1}{2}k_BT$$

Calculate instantaneous temperature:

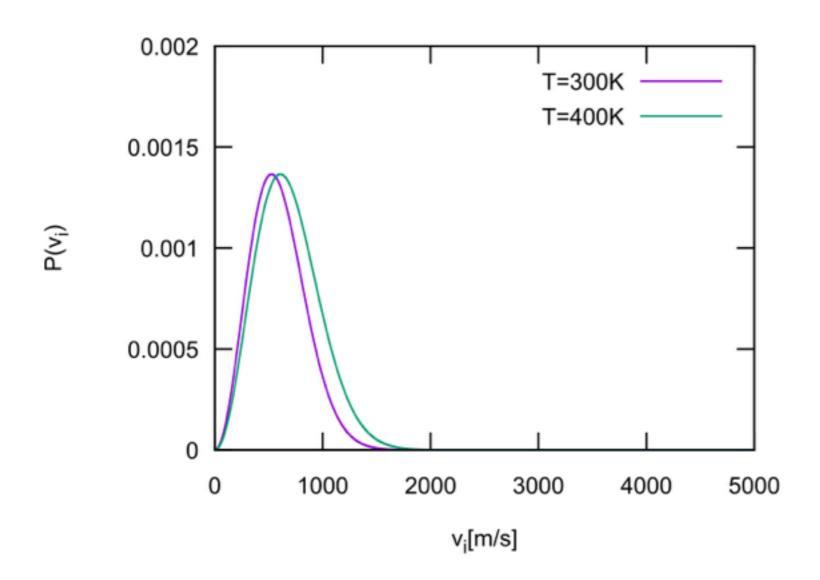
$$k_B T(t) = \frac{1}{N_f} \sum_{i,\alpha} m_i v_{i,\alpha}^2(t)$$

$$v_{i,\alpha}^{new} = \sqrt{\frac{T_{ref}}{T}}v_{i,\alpha}^{old}$$

Calculate instantaneous temperature:

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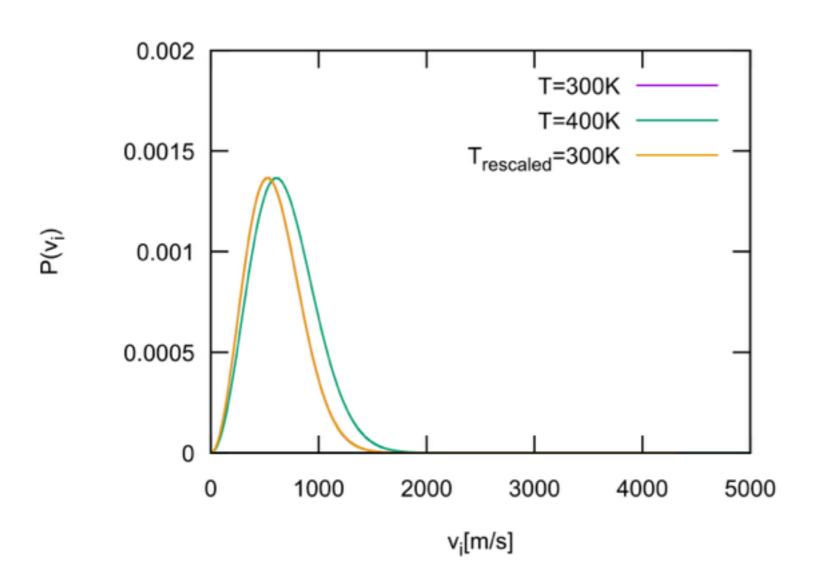
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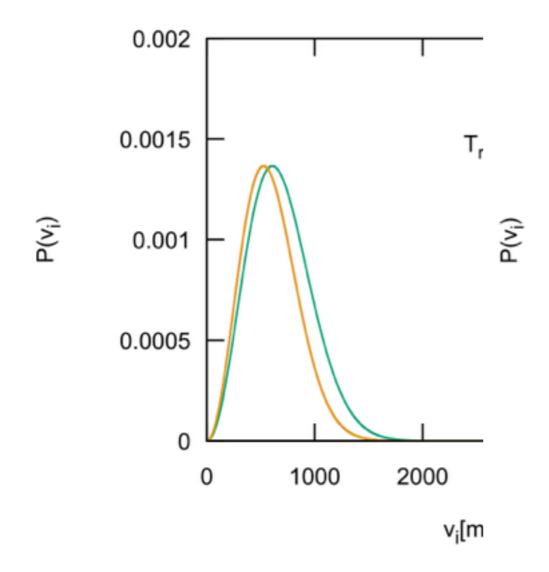
$$v_{i,\alpha}^{new} = \sqrt{\frac{T_{ref}}{T}}v_{i,\alpha}^{old}$$

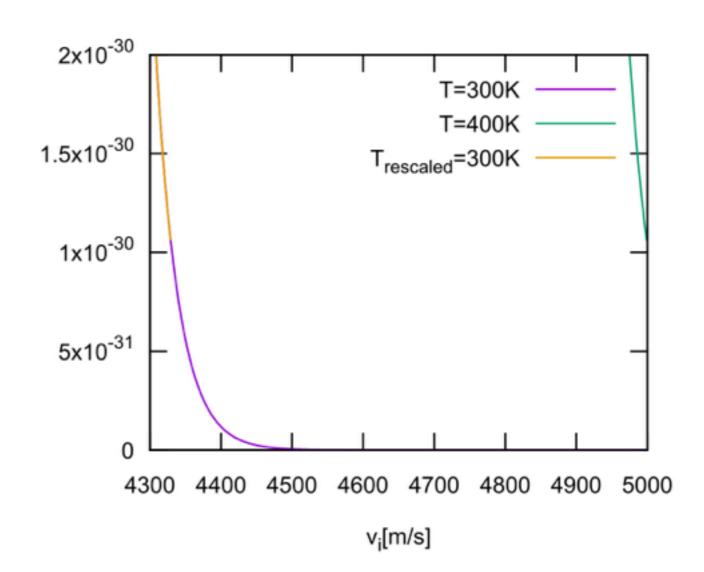


Calculate instantaneous temperature:

$$k_B T(t) = \frac{1}{N_f} \sum_{i,\alpha} m_i v_{i,\alpha}^2(t)$$

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Calculate instantaneous temperature:

$$k_B T(t) = \frac{1}{N_f} \sum_{i,\alpha} m_i v_{i,\alpha}^2(t)$$

Rescale velocities:

$$v_{i,\alpha}^{new} = \sqrt{\frac{T_{ref}}{T}}v_{i,\alpha}^{old}$$

Tail of velocity distribution is cut off!

The Flying Ice Cube: Velocity Rescaling in Molecular Dynamics Leads to Violation of Energy Equipartition

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Calculate instantaneous temperature:

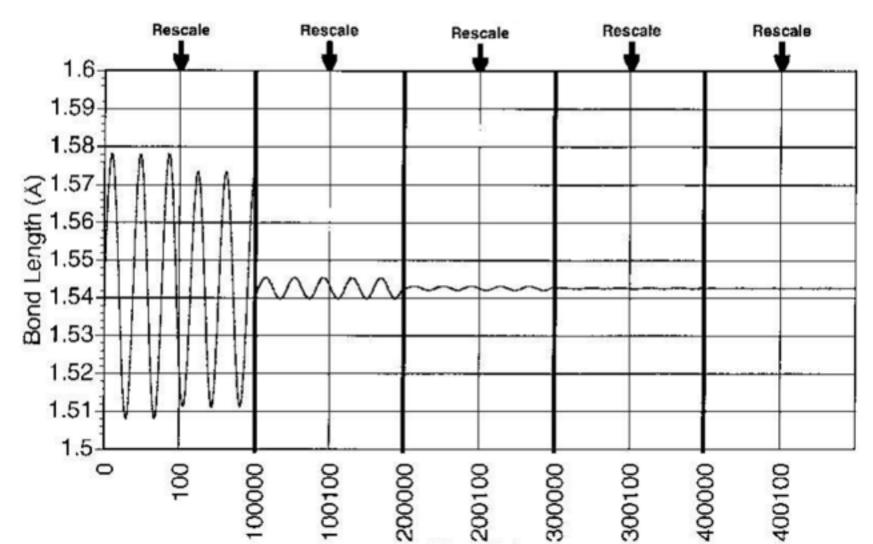
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Rescale velocities:

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Tail of velocity dis

The Flying Ice Cube Rescaling in Molecu Leads to Violation of Energy Equipartition



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Received 31 July 1997; accepted 4 December 1997

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Berendsen thermostat

Scale velocities by
$$c_t = \sqrt{1-rac{\Delta t}{ au}}\left(1-rac{T_{ref}}{T}
ight)$$
 $au o imes \infty, \ c_t o 1$ $au o imes \Delta t, \ c_t o \sqrt{rac{T_{ref}}{T}}$

Does not yield canonical ensemble Controls temperature very reliably → is still used for equilibration (short)

Andersen thermostat

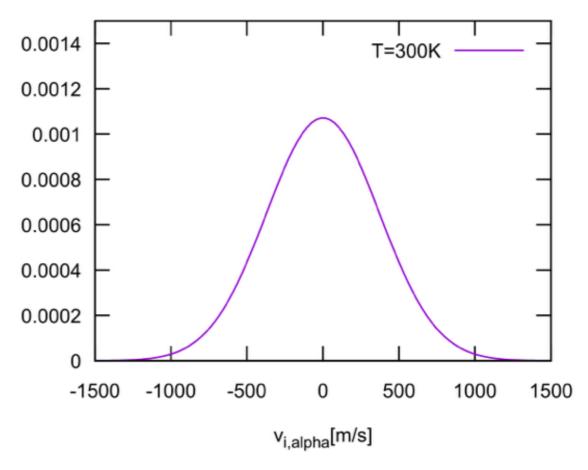
Perturbation with a stochastic term

- 1. Initialize system (positions and velocities)
- 2. Select n particles for collision with heat bath. Probability for a particle to be selected in Δt is $\gamma \Delta t$, γ = strength of coupling
- For selected particle draw new velocity from Maxwell-Boltzmann distribution

$$P(v_{i,\alpha}) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{mv_{i,\alpha}}{2k_B T}\right)$$

Large $\gamma \rightarrow$ strong coupl. \rightarrow dynamics are altered

Small $\gamma \rightarrow$ weak coupl. \rightarrow poor T control



Langevin thermostat

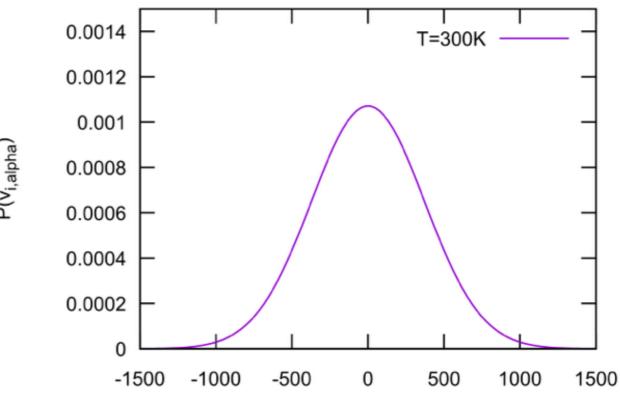
Perturb velocities with a Gaussian term

$$v_{new} = c_1 v_{old} + c_2 N(0, \sqrt{k_B T/m})$$
 $c_1^2 + c_2^2 = 1$
 $c_1 = exp(-\gamma \Delta t)$

$$\gamma \rightarrow 0, c_1 \rightarrow 1$$
 no thermostat

$$0 \to \infty$$
, $c_1 \to 0$ Andersen thermostat with strong coupling

$$P(v_{i,\alpha}) = \left(\frac{m}{2\pi k_B T}\right)^{1/2} \exp\left(-\frac{m v_{i,\alpha}}{2k_B T}\right)$$



v_{i,alpha}[m/s]

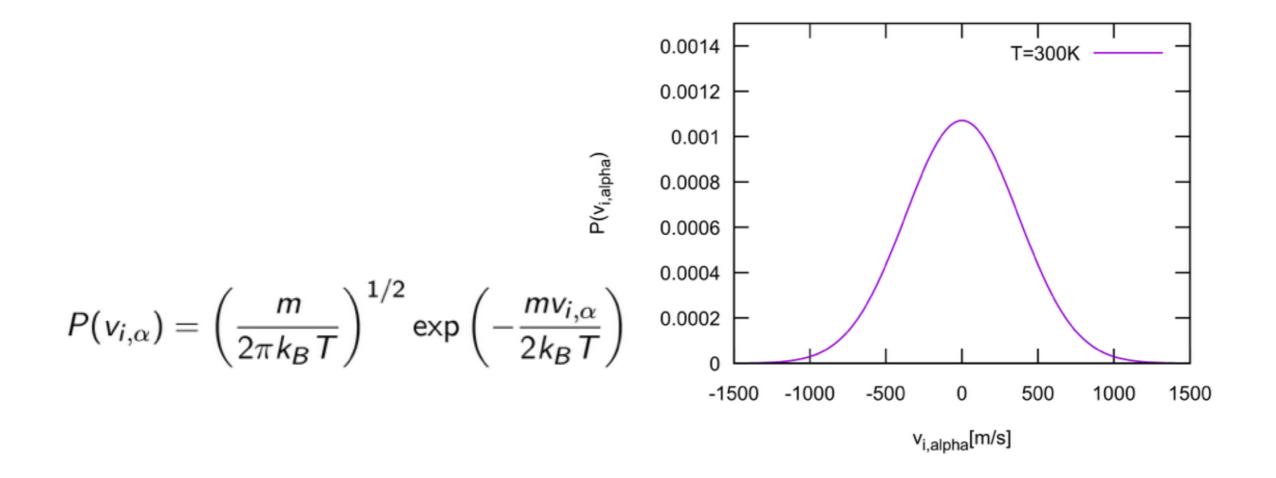
Yields correct ensemble!

Langevin thermostat in simplemd

```
// Read initial positions
read positions
// Set random velocities
randomize velocities
// Compute initial neighbor list
compute list
// compute initial forces
compute forces
for(istep=0;istep<nstep;istep++){</pre>
    thermostat (dt/2)
    velocity += (force*dt/2)/mass
    position += velocity*dt
    // Check whether the neighbour list has to be recomputed
    check list
    if(recompute list) {
      compute list
    compute forces
    velocity += (force*dt/2)/mass
    thermostat (dt/2)
```

Initial randomization of velocities

```
void randomize_velocities()
{
// randomize the velocities according to the temperature
  for(int iatom=0;iatom<natoms;iatom++) for(int i=0;i<3;i++)
     velocities[iatom][i]=sqrt(temperature/masses[iatom])*random.Gaussian();
}</pre>
```



Langevin thermostat

```
Thermostat() {
// Langevin thermostat, implemented as decribed in Bussi and Parrinello, Phys. Rev. E (2007)
// it is a linear combination of old velocities and new, randomly chosen, velocity,
// with proper coefficients
  double c1,c2;
  c1=exp(-friction*dt);
  for(int iatom=0;iatom<natoms;iatom++){</pre>
    c2=sqrt((1.0-c1*c1)*temperature/masses[iatom]);
    for (int i=0; i<3; i++) {
      engint+=0.5*masses[iatom]*velocities[iatom][i]*velocities[iatom][i];
      velocities[iatom][i]=c1*velocities[iatom][i]+c2*random.Gaussian();
      engint-=0.5*masses[iatom]*velocities[iatom][i]*velocities[iatom][i];
```

K+V+engint is conserved (useful for checking the timestep)

Time average and ensemble average

$$A = \langle A(r(t), p(t)) \rangle_{time} = \frac{1}{N_{steps}} \sum_{t=0}^{N_{steps}} A(r(t), p(t))$$

$$A = \langle A(r,p) \rangle_{ensemble} = \sum_{t=0}^{N_{steps}} A(r,p) \rho_{ensemble}(r,p)$$

 ρ = probability to find system in (r, p)

Ergodicity: $\langle A(r(t), p(t)) \rangle_{time} = \langle A(r, p) \rangle_{ensemble}$

In canonical ensemble

$$\rho(r,p) = \frac{1}{Z}e^{-E(r,p)/k_BT}$$

$$E(r,p) = K(p) + V(r)$$

$$(r,p) \rightarrow \text{microstate } i, Z = \sum_{i} e^{-E_i/k_BT}$$
 canonical partition function

if
$$A = A(r) \rightarrow \langle A \rangle = \sum_{i} \rho(r_i) A(r_i)$$

$$\rho(r_i) = \frac{1}{Z_V} e^{-V(r_i)/k_B T}$$

$$Z_V = \sum_{i}^{r} e^{-V(r_i)/k_BT}$$

How to generate correct ensemble if we don't know Z?

Metropolis

Use detailed balance

$$\rho(r_i)\pi(r_i \to r_j) = \rho(r_j)\pi(r_j \to r_i)$$

$$\pi(r_i \to r_j) = \alpha(r_i \to r_j) \operatorname{acc}(r_i \to r_j)$$

† trial probability (e.g. random)

$$\operatorname{acc}(r_i o r_j) = egin{cases}
ho(r_j)/
ho(r_i) & ext{if }
ho(r_j) <
ho(r_i) \ 1 & ext{if }
ho(r_j) \geq
ho(r_i) \end{cases}$$
 acceptance

$$\frac{\rho(r_j)}{\rho(r_i)} = \frac{\pi(r_i \to r_j)}{\pi(r_j \to r_i)} = e^{-(V(r_j) - V(r_i))/k_BT}$$

no partition function

Metropolis, Rosenbluth, Rosenbluth, Teller and Teller. Equation of State Calculations by Fast Computing Machines. 1953

Monte Carlo - Metropolis algorithm

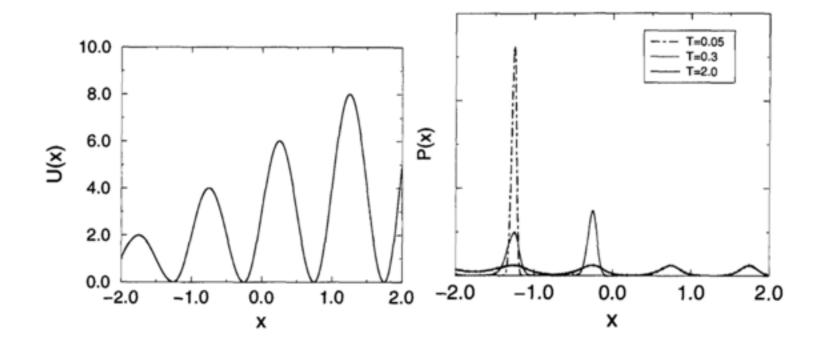
```
1. Start from initial condition r<sub>1</sub>
2. Calculate potential energy V(\mathbf{r}_1)
3. Propose a move \mathbf{r}_1 \rightarrow \mathbf{r}_2 (e.g. random)
4. Calculate V(\mathbf{r}_2)
5. Calculate acc(\mathbf{r}_1 \rightarrow \mathbf{r}_2) = exp(-(V(\mathbf{r}_2) - V(\mathbf{r}_1) / kBT)
6. if V(\mathbf{r}_2) \leq V(\mathbf{r}_1)
           accept the move
    else
           extract a random number n in [0,1]
            if n < acc(\mathbf{r}_1 \rightarrow \mathbf{r}_2)
                                                                     \exp(-\beta\delta\mathscr{V})
                 accept the move
           else
                  reject
                                                                                 Reject
                                                          Always
                                                          accept
                                                                            Accept
                                                                                           * ξ<sub>1</sub>
                                                                                          δ¥ ....
```

Monte Carlo - Metropolis algorithm

Possible to generate samples with desired distribution Allows evaluation of average

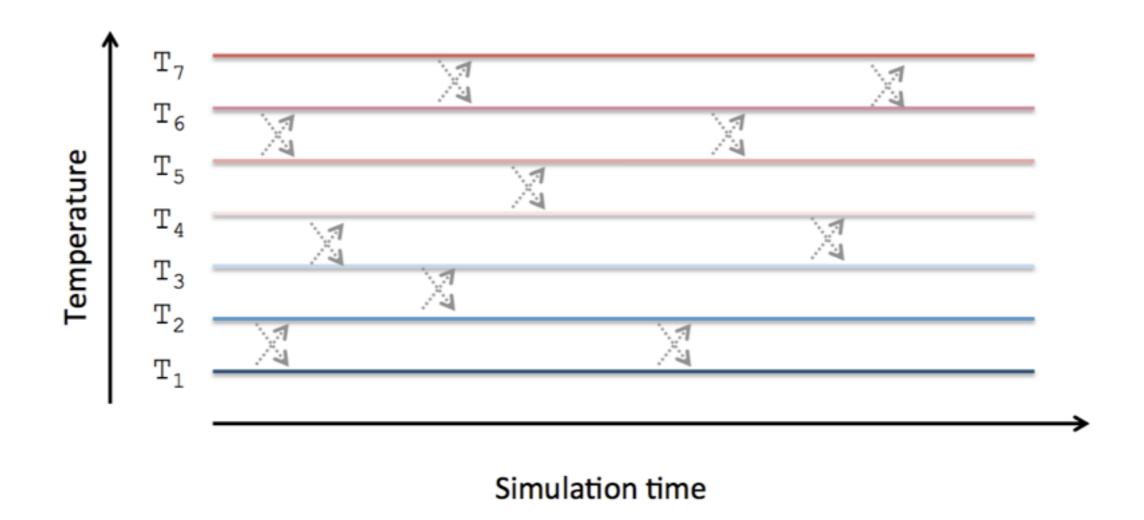
$$\langle A \rangle = \sum_{i} \rho_{NVT}(r_i) A(r_i)$$

Sampling problem - stuck in a minimum

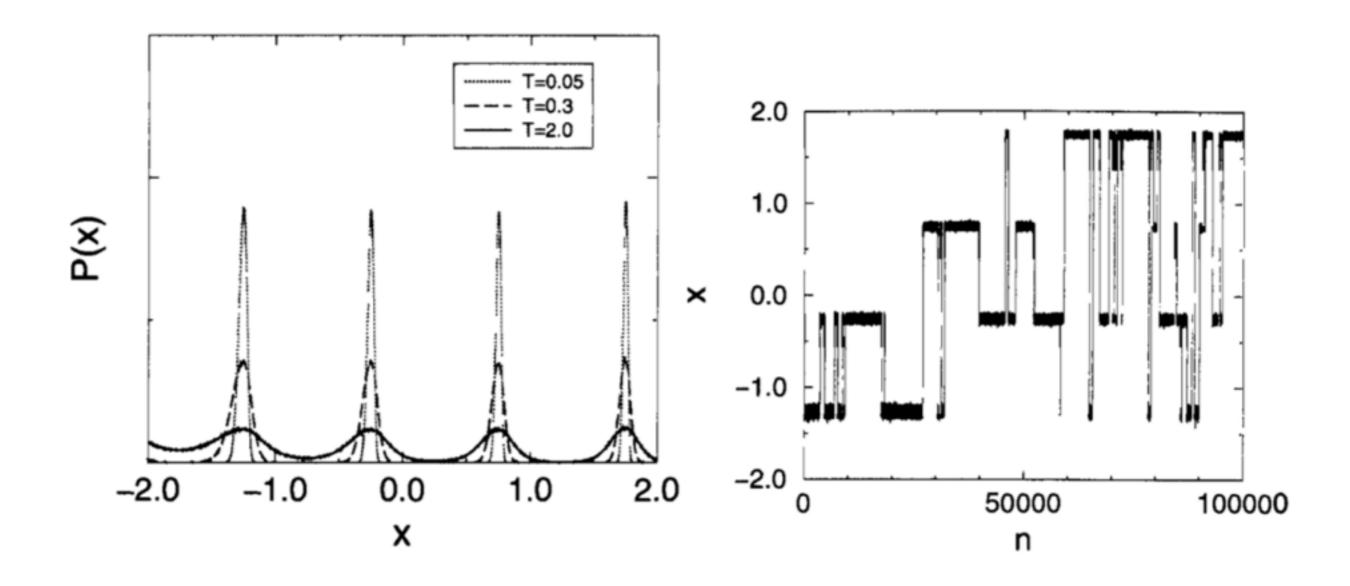


One-dimensional example potential

- If initially in one minimum, you will stay there
- at high temperatures you can overcome energy barriers



Idea: run parallel simulations at different temperatures. Use a "swap" move to exchange between replicas Accept/reject using the standard Metropolis criterion

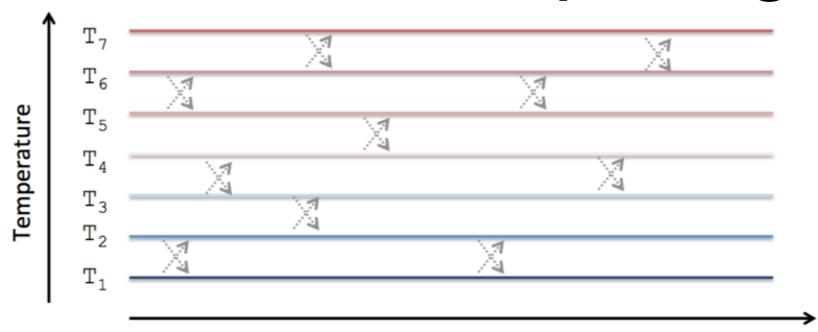


One-dimensional example potential

Now the whole space is sampled, even by the replica at low T

$$\operatorname{acc}(r_i o r_j) = egin{cases}
ho(r_i)/
ho(r_j) & ext{if }
ho(r_j) <
ho(r_i) \ & ext{acceptance} \end{cases}$$

$$acc = e^{-(1/k_BT_j-1/k_BT_i)(V(r_i)-V(r_j))}$$



Simulation time

```
1. Create different replicas running at different temperatures  
2. Every t steps attempt swap move between replica i and j calculate acc = \exp((1/kBT_i-1/kBT_j)(V(\mathbf{r}_i)-V(\mathbf{r}_j))) if acc > 1  
   swap the coordinates else  
   if random.U01()<acc swap coordinates
```

```
ATT! If swap is accepted, rescale velocities v_i(\text{new}) = \text{sqrt}(T_i/T_j); v_j(\text{new}) = \text{sqrt}(T_j/T_i)
```

Books on MD

Tamar Schlick: Molecular Modeling an Simulation: an Interdisciplinary guide (2010)

Allen & Tildesley: Computer Simulations of liquids (1991)