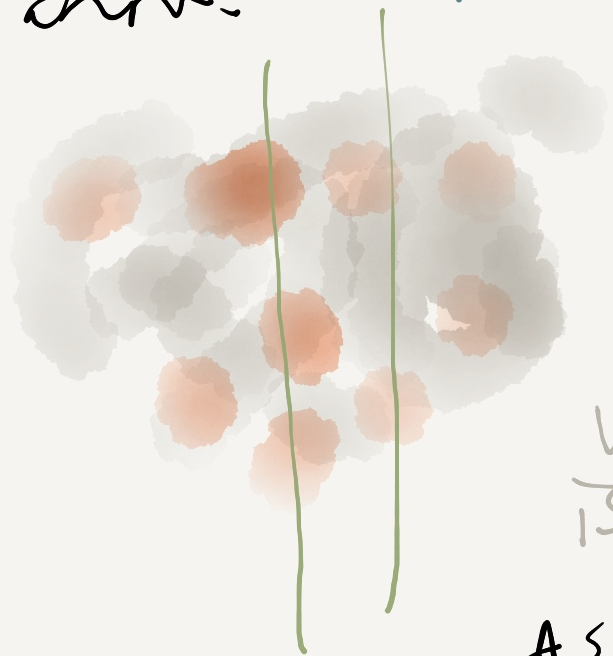


TS4

How does the spatial scale of the T fluctuations affect the plane of-sky variation?

Turbulence during the approach to IAH



If the line of sight crosses many regions of different T , then the variation in the average T is reduced.

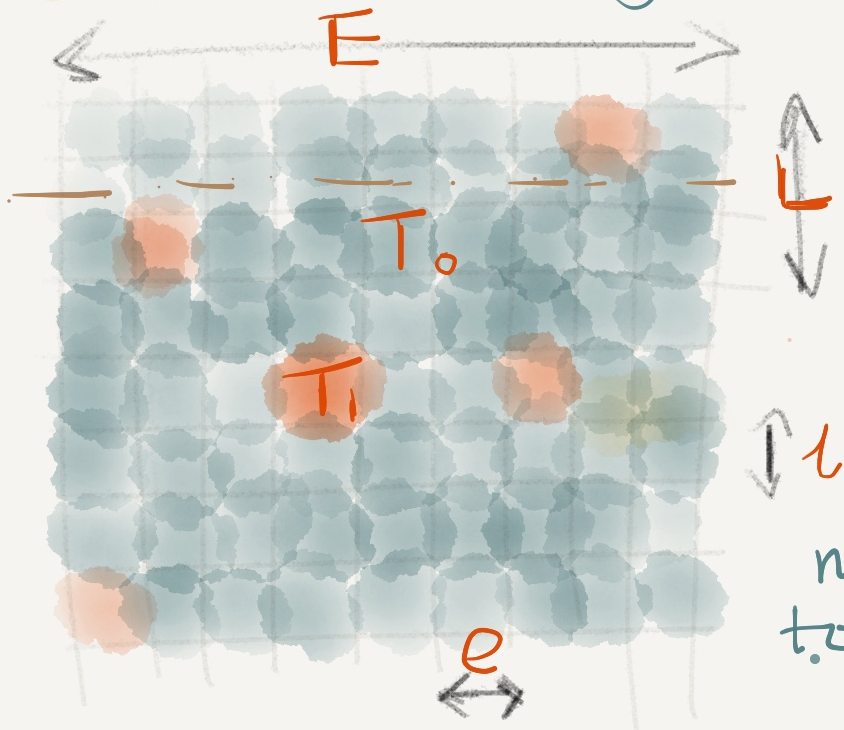
Assume that a fraction f

of the emissivity comes

from gas with $T = T_0(1 + \delta)$ where as

the remaining fraction $(1-f)$ comes from $T = T_0$.

Better yet: assume that we have TSS
a large number of emission blocks,
each with exactly the same emiss measure e
Each block has a probability p of being
"anomalous" $T = T_1 = T_0(1 + \delta)$ and a probability
 $1 - p$ of being "normal": $T = T_0$.



We observe on an
angular scale L
whereas the angular scale
of the blocks is 1 . The
total emiss measure
is E . Therefore the
number of blocks that contribute
to each observation is

$$N = \max(1, EL^2/e)$$

Mean temperature: $\bar{T} = [1 - p + p(1 + \delta)] T_0$ T56

$$\bar{T} = (1 + p\delta) T_0$$

Fluctuations: $t^2 = \frac{(T_0 - \bar{T})^2 (1 - p) + (T_1 - \bar{T})^2 p}{\bar{T}^2}$

$$(T_0 - \bar{T})^2 = p^2 \delta^2 T_0^2$$

$$(T_1 - \bar{T})^2 = (1 - p)^2 \delta^2 T_0^2 \Rightarrow \text{numerator is}$$

$$\Rightarrow t^2 = \frac{p(1 - p) \delta^2}{(1 + p\delta)^2}$$

If $p \ll 1$ and $p\delta \ll 1$, then

$$t^2 \simeq p\delta^2$$

Note that we

do not necessarily have δ being small — it can be large so long as $p \ll 1$ and t^2 will still be small. Also, we can have $\delta < 0$, corresponding to cool spots.

So there are two different ways of getting a small t^2 value: $t^2 \ll 1$ T57

A. A checkerboard ($p \sim 0.5$) of small fluctuations ($\delta^2 \ll 1$) $\Rightarrow t^2 \simeq 0.25 \delta^2$

B. A sprinkling ($p \ll 1$) of large fluctuations ($\delta \sim 1$) $\Rightarrow t^2 \sim p$ Note that it is not reasonable to have

~~δ much larger than unity since classical minimization would change the minimization state and the true spectrum.~~

When we measure the T of a region that contains N blocks (see T55), then on average there will be $\bar{n} = pN$ excitations blocks, but there will really be a spread of values $n = \bar{n} \pm \sqrt{\bar{n}}$ (if \bar{n} is larger than a few).

So the mean T of a ^{single} patch of scale L will be T_{S8} .

$T(L) = (1 + f\delta)T_0$ where $f = \frac{n}{N}$ is the actual fraction of anomalous blocks in the patch and where n is drawn from a Poisson distribution of mean $\bar{n} = pN$. $P(n, \bar{n})$

$$(T(L) - \bar{T})^2 = (f - p)^2 \delta^2 T_0^2$$

The plane-disk temp fluctuations are therefore

$$\hat{t}^2(L) = \frac{1}{T_0^2} \langle (T(L) - \bar{T})^2 \rangle$$

$$= \int_0^\infty dn P(n, \bar{n}) \left(\frac{n}{N} - p \right)^2 \delta^2$$

$$= \frac{\delta^2}{N^2} \text{Var}(n) = p \frac{\delta^2}{N} = t^2/N$$

For the two phase model we had $t^2 = \frac{\phi(1-\phi)\delta^2}{(1+\phi\delta)^2}$ TS9

Whereas the extra heating or cooling can be found as

$$\approx \phi\delta^2$$

$$L = L_0(1-\phi) + \phi L_0(1+\delta)^a \quad a \equiv \frac{d \ln L}{d \ln T}$$

In the case that $\delta \ll 1$ then we can

simplify this $L = L_0(1 + a\phi\delta)$ where L_0 is the equilibrium cooling @ $T = T_0$

Using fact that $\delta \approx (t^2/\phi)^{1/2}$

$\Rightarrow L \approx L_0(1 + a\phi^{1/2}(t^2)^{1/2})$ so assuming thermal equilibrium $H = L$ we have the fractional extra heating required as $\frac{H-H_0}{H_0} = \frac{L-L_0}{L_0} = a\phi^{1/2}(t^2)^{1/2} \equiv h$

TS10

If $\delta \approx 1$ then $t^2 \approx p$ and we would have $h \approx ap \approx at^2$ in the above approx, (of course we don't satisfy $\delta \ll 1$ so we need to do it properly).

On the other hand, if $p \approx 0.5$ then $t^2 \approx \frac{1}{4}\delta^2$ and $h \approx \frac{1}{2}a\delta \approx a(t^2)^{1/2}$

For the Ring Nebula, we have $t^2 \approx 0.04$ from the ADF. Or $t_{\uparrow}^2 \approx 0.005$ for the part at small/medium scales. The cooling sensitivity index $a \approx 2$ I think, which would imply $h \approx 0.4$ to reproduce the ADF t^2 , if we had high filling factor, small amplitude fluctuations. This is a lot of extra heating!

TS11

Going back to the rare, large amplitude case, we had $h = p[(1+s)^a - 1]$

If we take $s \simeq 1$ and $a \simeq 2$ this is $h \simeq 3p \simeq 3t^2$. This is much easier to satisfy since $t^2 = 0.04 \Rightarrow h \simeq 0.12$

Although this is still rather high.

If we only needed to explain the observed $t_A^2 \simeq 0.005$ then we would only require $h \simeq 1.5\%$, which is much more reasonable.

What fraction of the total radiative energy absorbed by an HII region is retained as thermal energy?