

## Phase Space

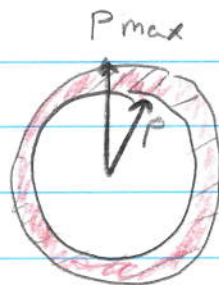
$$\begin{aligned} \text{a)} \quad V_{ps} &= \int_{[E, E+\delta E]} d^3r d^3p = V \int d^3p \\ & \quad \text{shell with energy in range } [E, E+\delta E] \\ &= V 4\pi p^2 \delta p \end{aligned}$$

Then

$$p = \sqrt{2mE}$$

See lecture!

$$p_{\max} = \sqrt{2m(E+\delta E)}$$



$$\begin{aligned} \delta p &= p_{\max} - p = \sqrt{2mE} (1 + \delta E/E)^{1/2} - \sqrt{2mE} \\ &\approx \sqrt{2mE} \left( 1 + \frac{1}{2} \frac{\delta E}{E} \right) - \sqrt{2mE} \end{aligned}$$

$$\delta p = p \frac{\delta E}{2E} \quad \text{i.e.} \quad \frac{\delta p}{p} = \frac{\delta E}{2E}$$

So

$$V_{ps} = V 4\pi p^3 \frac{\delta E}{2E} = V [4\pi (2mE)^{3/2} \frac{\delta E}{2E}]$$

b) Then

$$\begin{aligned}\Omega(E, V) &= \frac{1}{2!} \int \frac{d^3 r_1 d^3 p_1}{h^3} \frac{d^3 r_2 d^3 p_2}{h^3} \\ &= \frac{V^2}{2!} \int \frac{d^3 p_1 d^3 p_2}{h^6} \\ &\quad \underbrace{[E, E+\delta E]}_{\text{shell volume}/h^6}\end{aligned}$$

The momentum integral is a spherical shell in 6 dimensions  
The momentum space volume is :

$$\frac{\text{shell volume}}{h^6} = \frac{2 \pi^{6/2}}{\Gamma(6/2)} p^5 \frac{\delta p}{h^6}$$

As before

$$\delta p = p \frac{\delta E}{2E}$$

So

$$p = \sqrt{2mE}$$

$$\Omega(E, V) = \frac{V^2}{2!} \frac{2 \pi^3}{2!} \left( \frac{p}{h} \right)^6 \frac{\delta E}{2E}$$

$$\Omega(E, V) = V^2 \pi^3 \left( \frac{2mE}{h^2} \right)^3 \frac{\delta E}{4E}$$

c) Now the general case

$$\Omega(E, V) = \frac{1}{N!} \int \frac{d^3 r_1}{h^3} \dots \frac{d^3 r_N d^3 p_N}{h^3}$$

$$= \left( \frac{V^N}{N!} \right) \int_{[E, E+\delta E]} (d^3 p_1 \dots d^3 p_N) / h^{3N}$$

↑ coordinate space      ← momentum space  
shell volume /  $h^{3N}$

So this is again a shell of dimension  $3N$

$$\frac{\text{Shell volume}}{h^{3N}} = \frac{2 \pi^{3N/2}}{\Gamma(3N/2)} \left( \frac{p}{h} \right)^{3N} \frac{\delta E}{\bar{E}}$$

Using  $\frac{p}{h} = \left( \frac{2mE}{h^2} \right)^{1/2}$  we get:

$$\Omega(E, V) = \left( \frac{V^N}{N!} \right) \left( \frac{1}{\Gamma(3N/2)} \left( \frac{2\pi m E}{h^2} \right)^{3N/2} \right) \frac{\delta E}{\bar{E}}$$

↑ coordinate space      ← momentum space

d) Writing

$$N! = \left( \frac{N}{e} \right)^N, \quad \Gamma\left(\frac{3N}{2}\right) = \left( \frac{3N}{2e} \right)^{3N/2}$$

We have

$$\Omega(E, V) = \left[ e^N \left( \frac{V}{N} \right)^N \right] \left[ e^{3N/2} \left( \frac{4\pi m E}{3N h^2} \right)^{3N/2} \right]$$

this can  
be set to  
one  
✓  
see  
problem  
+  
lecture

Taking the log of this expression, and dropping the  $\delta E/E$  term, we have

$$S = k \ln \Omega$$

$$S = k \left[ \frac{5N}{2} + N \ln \frac{V}{N} + N \ln \left( \frac{4\pi m E}{3N h^2} \right)^{3/2} \right]$$

The term:  $\left( \frac{4\pi m E}{N h^2} \right)^{3/2} = \left( \frac{2\pi m kT}{h^2} \right)^{3/2} = \left( \frac{1}{\lambda_{th}^2} \right)^{3/2} = \frac{1}{\lambda_{th}^3}$

where we used  $\frac{E}{N} = \frac{3}{2} kT$

So finally

$$S = Nk \left[ \ln \left( \frac{V}{N \lambda_{th}^3} \right) + \frac{5}{2} \right]$$

↑ we will derive this in other ways as the course progresses. Here we have derived it from the microcanonical ensemble. Later we will derive it from the canonical ensemble.



## Entropy Changes of Ideal Gas

✓ This is the entropy of a MAIG, which we derived. We will use this in this problem

$$(*) \quad S = Nk \ln V + \frac{3Nk}{2} \ln E + \text{const}$$

$$a) \quad PV = NkT, \text{ so for } P \text{ const, } \frac{V_1}{V_2} = \frac{T_1}{T_2} = \frac{E_1}{E_2}$$

$$\text{Since } E = \frac{3}{2} NkT \propto T,$$

$S_0$

$$\Delta S = S_2 - S_1$$

$$= Nk \ln \frac{V_2}{V_1} + \frac{3Nk}{2} \ln \frac{E_2}{E_1}$$

$$= Nk \ln \frac{T_2}{T_1} + \frac{3}{2} Nk \ln \frac{T_2}{T_1}$$

$$\Delta S = \frac{5Nk}{2} \ln \frac{T_2}{T_1}$$

$$b) \quad \Delta S = S_2 - S_1$$

$$= Nk \ln \frac{V_2}{V_1} + \frac{3Nk}{2} \ln \frac{E_2}{E_1} \quad \text{now } PV = NkT$$

$$P \propto T \propto E$$

constant volume

$S_0$

$$\Delta S = \frac{3}{2} Nk \ln \left( \frac{P_2}{P_1} \right)$$

(c) Now for an adiabatic expansion

$$TV^{\gamma-1} = \text{const}$$

this is for a  
MAIG

$$\frac{T_2}{T_1} = \left( \frac{V_1}{V_2} \right)^{\gamma-1} \quad \text{where } \gamma = \frac{5}{3}$$

or

$$\frac{T_2}{T_1} = \left( \frac{V_1}{V_2} \right)^{2/3}$$

$$\frac{E_2}{E_1} = \frac{T_2}{T_1} = \left( \frac{V_2}{V_1} \right)^{-2/3}$$

$S_0$

$$\Delta S = Nk \ln \frac{V_2}{V_1} + \frac{3}{2} Nk \ln \frac{T_2}{T_1}$$

$$(i) = Nk \ln \frac{V_2}{V_1} + \frac{3}{2} Nk \ln \left( \frac{V_2}{V_1} \right)^{-2/3} = 0$$

(ii) The coordinate space volume is getting bigger, but the momentum space volume is getting smaller as the temperature drops. We have  $EV^{3/2} = \text{const}$  during the adiab expansion of a MAIG.

## Ball and Lake

$$a) \quad Q = C(T - T_B^o)$$

$$b) \quad \Delta S = \int_{T_B^o}^T \frac{C dT_B}{T_B}$$

$$c) \quad \Delta S = C \ln \frac{T}{T_B^o}$$

(ii) Now

$$\Delta S = C \ln \frac{T}{T_B^o} = -C \ln \frac{T_B^o}{T} = C \ln \left( 1 - \frac{T - T_B^o}{T} \right)$$

$$\frac{Q}{T} = C \frac{(T - T_B^o)}{T}$$

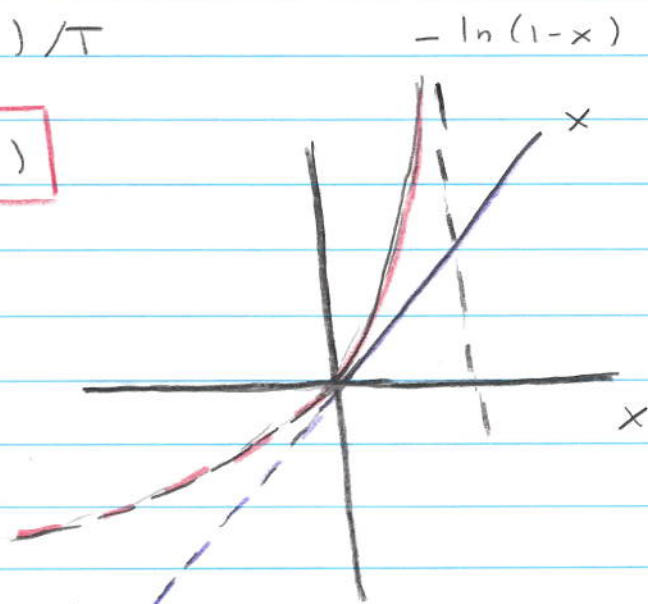
So define  $x \equiv (T - T_B^o) / T$

$$\Delta S = -C \ln(1 - x)$$

$$\frac{Q}{T} = Cx$$

So  $\Delta S > \frac{Q}{T}$

as is clear from graph



iii) For small  $x$

$$-\ln(1-x) \approx -\left(-x - \frac{(-x)^2}{2}\right)$$

$$\approx x + \frac{x^2}{2}$$

So

$$\Delta S \approx C \frac{\Delta T}{T} + \frac{C}{2} \left(\frac{\Delta T}{T}\right)^2 = \frac{Q}{T} + \frac{C \Delta T^2}{2T^2} > \frac{Q}{T}$$

$$\frac{Q}{T} = C \frac{\Delta T}{T}$$



## Heating Water

a) We have one steps,  $T_0 = 0^\circ\text{C} = 273\text{K}$   
and  $T_1 = 100^\circ\text{C} = 373\text{K}$ .

The specific heat of one mole of  $\text{H}_2\text{O}$ , is  
the mass of 1 mol  $M_m$  times  $C_m = 4180 \frac{\text{J}}{\text{kg}^\circ\text{K}}$

The molar mass is

$$M_m = 18 \text{ NA } m_p = 18 \text{ g}$$

↑  
two hydrogen 1+1 + 16 ← oxygen

So  $C = 18 \text{ g } C_m = 75.2 \text{ J}/^\circ\text{K}$

Then from the previous problem

$$\begin{aligned}\Delta S_{\text{univ}} &= \Delta S_{\text{water}} + \Delta S_{\text{res}} = C \ln \frac{T_1}{T_0} - C \frac{(T_1 - T_0)}{T_0} \\ &= 3.3 \text{ J}/^\circ\text{K}\end{aligned}$$

b) In two steps  $T_0 = 0^\circ\text{C}$ ,  $T_1 = 50^\circ\text{C}$ ,  $T_2 = 100^\circ\text{C}$

$$\Delta S = C \left[ \underbrace{\ln \frac{T_1}{T_0} - \frac{C(T_1 - T_0)}{T_1}}_{\text{Step 1}} + \underbrace{\ln \frac{T_2}{T_1} - \frac{C(T_2 - T_1)}{T_2}}_{\text{step 2}} \right]$$

combine

So

$$\Delta S_{\text{uni}} = C \left[ \ln \frac{T_2}{T_0} - \sum_{i=1}^2 \frac{\Delta T_i}{T_i} \right] = 1.75 \text{ J/K}$$

with  $\Delta T_i = T_i - T_{i-1}$ . We note this can be written:

$$\Delta S_{\text{uni}} = \Delta S_{\text{sys}} - \sum_i \frac{Q_i}{T_i}$$

c) For four steps we have  $T_0 = 0^\circ\text{C}$ ,  $T_1 = 25^\circ\text{C}$ ,  $T_2 = 50^\circ\text{C}$ ,  $T_3 = 75^\circ\text{C}$ ,  $T_4 = 100^\circ\text{C}$

$$\Delta S_{\text{univ}} = \Delta S_{\text{sys}} - \sum_{i=1}^4 Q_i / T_i$$

$$= C \left[ \ln \left( \frac{T_4}{T_0} \right) - \sum_{i=1}^4 \frac{\Delta T_i}{T_i} \right] = 0.9 \text{ J/K}$$

d) To make a reversible process we must take the number of steps very large

$$\Delta S_{\text{univ}} = C \left[ \ln \left( \frac{T_N}{T_0} \right) - \sum_{i=1}^N \frac{\Delta T_i}{T_i} \right]$$

In the limit, the sum is replaced by an integral

$$\sum_i \frac{\Delta T_i}{T_i} \approx \int_{T_0}^{T_N} \frac{dT}{T} = \ln \left( \frac{T_N}{T_0} \right)$$

And then

$$\Delta S_{\text{univ}} = C \left[ \ln \left( \frac{T_N}{T_0} \right) - \ln \left( \frac{T_N}{T_0} \right) \right] = 0$$