

Physics 306: Thermal Physics

Midterm Exam

Stony Brook University

Fall 2024

General Instructions:

You may use one page (front and back) of handwritten notes and a calculator. Graphing calculators are allowed. **No other materials may be used.**

1 Integrals

Gamma Function:

$$\Gamma(z) \equiv \int_0^{\infty} x^{z-1} e^{-x} dx \quad (1)$$

with specific results

$$\Gamma(z+1) = z\Gamma(z) \quad \Gamma(n) = (n-1)! \quad \Gamma(\tfrac{1}{2}) = \sqrt{\pi} \quad (2)$$

Gaussian Integrals:

$$I_n = \frac{1}{\sqrt{2\pi}\sigma} \int_{-\infty}^{\infty} dx e^{-x^2/2\sigma^2} x^n \quad (3)$$

with specific results

$$I_0 = 1 \quad I_2 = \sigma^2 \quad I_4 = 3\sigma^4 \quad I_6 = 15\sigma^6 \quad (4)$$

Other integrals:

$$\int e^u du = e^u + C$$

$$\int e^u u du = e^u(u-1) + C \quad (5)$$

$$\int e^u u^2 du = e^u(u^2 - 2u + 2) + C$$

Here α is any real number (positive or negative).

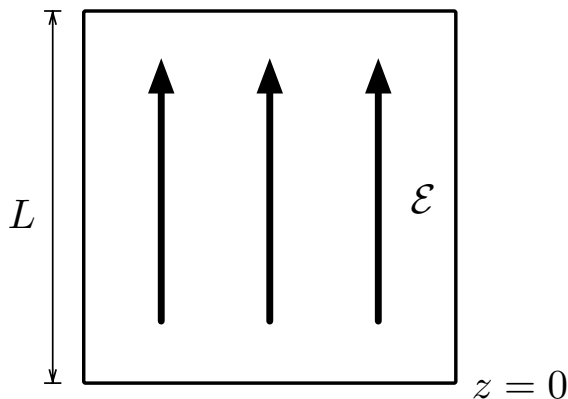
Problem 1. Plasma in an electric field

Consider a plasma made up of N electrons in a 3D box with volume $V = L^3$. Treat the plasma as an ideal gas maintained at temperature T , and ignore the Coulomb interaction between the electrons.

A uniform electric field of magnitude \mathcal{E} is applied in the z -direction. Recall that the potential energy $V(z)$ of a negatively charged particle in a constant electric field is given by¹

$$V(z) = q\mathcal{E}z \quad (6)$$

where q is the absolute value of the electron's charge.



Some potentially useful integrals are given on the first page.

- Determine the most probable speed for the electrons.
- Determine the charge density $\rho(z)$, i.e. the density of electrons times $-q$, as a function of z .
Hint: First determine the probability density of an electron having height z .
- What is the average energy of the plasma. Taylor expand your result in the limit that the electric field is weak to leading (non-zero) order in the field strength \mathcal{E} .

¹The potential energy is the charge times the voltage. The voltage is $\Phi(z) = -\mathcal{E}z$ for a constant field.

Solution

(a) The speed distribution is

$$\frac{d\mathcal{P}}{dv} = P(v) = C e^{-mv^2/2kT} v^2 \quad (7)$$

where C is an unimportant constant. Maximizing $P(v)$ by finding where the derivative is zero yield

$$P'(v) = C e^{-mv^2/2kT} \left(\frac{mv^3}{kT} - 2v \right) = 0 \quad (8)$$

so

$$v_{\max} = \sqrt{2kT/m} \quad (9)$$

(b) The probability distribution in z is

$$\frac{d\mathcal{P}}{dz} \equiv P(z) = C e^{-\beta V(z)} = C e^{-\beta q\mathcal{E}z} = C e^{-az} \quad (10)$$

where we defined, $a \equiv \beta q\mathcal{E}$. The normalization constant can be found by normalizing the probability

$$\int_0^L P(z) dz = \int_0^L C e^{-az} = \frac{C}{a} (1 - e^{-aL}) = 1 \quad (11)$$

So

$$C = \frac{a}{1 - e^{-aL}} \quad (12)$$

Now we can find the probability per dz

$$\frac{d\mathcal{P}}{dz} \equiv P(z) = \frac{a e^{-az}}{1 - e^{-aL}}, \quad (13)$$

Multiplying $d\mathcal{P}/dz$ by N would give the number of particles per dz . Dividing by the area L^2 and noting that $dV = L^2 dz$ we find the electron density

$$\boxed{n(z) = \frac{dN}{dV} = \frac{1}{L^2} \left(\frac{dN}{dz} \right) = \frac{Na}{L^2} \frac{e^{-az}}{1 - e^{-aL}}.} \quad (14)$$

(c) The energy of an electron is

$$\bar{\epsilon} = \langle KE \rangle + \langle PE \rangle = \left\langle \frac{p_x^2}{2m} + \frac{p_y^2}{2m} + \frac{p_z^2}{2m} \right\rangle + \langle PE \rangle = \frac{3}{2}kT + \langle PE \rangle \quad (15)$$

where we used the equipartition theorem to evaluate the kinetic energy. The average potential energy is

$$\langle PE \rangle = \int_0^L q\mathcal{E}z C e^{-az} dz = kTC \int_0^L az e^{-az} dz \quad (16)$$

Performing the integral using the table, and putting in the value of C

$$\langle PE \rangle = kTC \left[-\frac{1}{a} e^{-az} (1 + az) \right]_0^L \quad (17)$$

$$= kT \frac{C}{a} [1 - e^{-aL} (1 + aL)] \quad (18)$$

Putting in the value of C and multiplying by the total number of particles gives the total energy

$$\boxed{U = \frac{3}{2} NkT + NkT \left[\frac{1 - e^{-aL} (1 + aL)}{1 - e^{-aL}} \right]} \quad (19)$$

The series expansion works as follows. Writing

$$e^{-aL} = 1 - (aL) + \frac{1}{2}(aL)^2 \quad (20)$$

We find

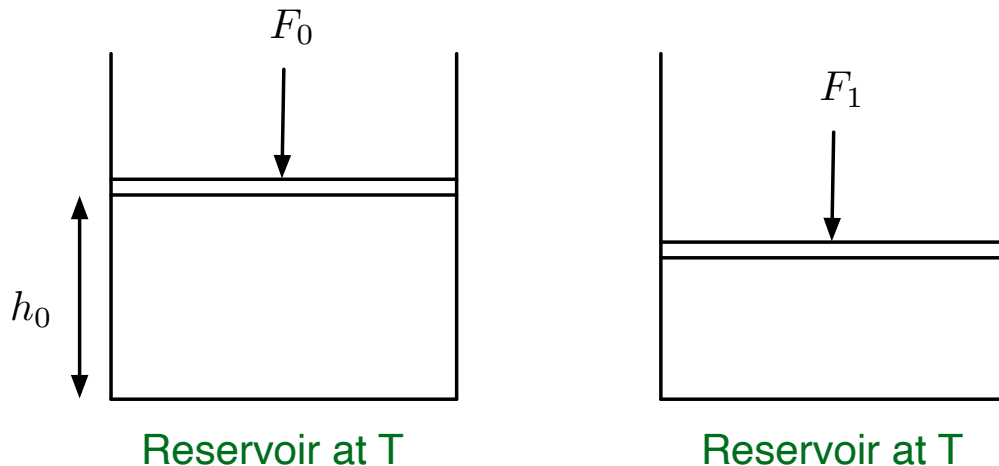
$$\langle PE \rangle = \frac{1}{\beta} \left[\frac{1 - (1 - (aL) + \frac{(aL)^2}{2}) (1 + (aL))}{aL} \right] \simeq kT \frac{aL}{2} \quad (21)$$

Then multiplying by N we find

$$U = \frac{3NkT}{2} + \frac{q\mathcal{E}L}{2} \quad (22)$$

Problem 2. A Sudden Force

A cylinder has cross-sectional area A and initial height h_0 . The cylinder contains an ideal gas and is surrounded by a reservoir at temperature T . Initially, a piston applies a force F_0 on the gas in the cylinder. At time $t = 0$, the applied force is suddenly increased to F_1 and held constant. The cylinder then contracts until it reaches a new equilibrium volume. Do not assume that the gas is in equilibrium during the contraction.



- What is the final volume? What is the work by the external force done during the process? What is the heat that flows out of the cylinder during the process?
- Determine the change in entropy of the gas in the cylinder.
- What is the change in entropy of the universe ΔS_{univ} ? Let $F_1 = F_0(1 + \delta)$. Make a series expansion of ΔS_{univ} for $\delta \ll 1$ to quadratic order in δ .

Solution

(a) Since $PV = NkT$ is a constant before and after, and since $pV = Fh$ we have

$$F_0 h_0 = F_1 h_1 = NkT, \quad (23)$$

or

$$h_1 = h_0 \frac{F_0}{F_1}. \quad (24)$$

The work done on the gas

$$W_{in} = F_1 \Delta x = F_1 (h_0 - h_1). \quad (25)$$

Using the first law

$$\Delta U = Q_{in} + W_{in} = 0, \quad (26)$$

Thus the heat that flows out is

$$Q_{in} = -W_{in} \quad Q_{out} = -Q_{in} = F_1 (h_0 - h_1). \quad (27)$$

(b) The entropy is

$$S = \text{const} + \frac{3}{2} Nk \ln U + Nk \ln V, \quad (28)$$

The temperature is fixed so $U_f = U_i$ and thus $\Delta S = S_f - S_i$ is

$$\Delta S = Nk [\ln(V_f) - \ln(V_i)] = Nk \ln(V_f/V_i) = Nk \ln(h_1/h_0) = Nk \ln(F_0/F_1). \quad (29)$$

We note that $\Delta S < 0$ since $F_0/F_1 < 1$.

(c) The change in entropy of the universe is

$$\Delta S_{univ} = \Delta S_{sys} + \Delta S_R = \Delta S_{sys} - \frac{Q_{in}}{T} = Nk \ln(F_0/F_1) + \frac{F_1 (h_0 - h_1)}{T}. \quad (30)$$

Now using the results from part (a) e.g. Eq. (24) and Eq. (23) we find

$$\Delta S_{univ} = Nk [\ln(F_0/F_1) + F_1/F_0 - 1]. \quad (31)$$

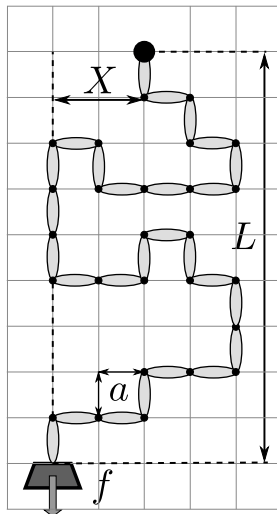
Writing $F_1 = F_0(1 + \delta)$ we expand

$$\Delta S_{univ} = Nk [-\ln(1 + \delta) + \delta] \simeq Nk \left[\left(-\delta + \frac{\delta^2}{2} \right) + \delta \right] \simeq Nk \frac{\delta^2}{2}. \quad (32)$$

We see that ΔS_{univ} is positive as it must be.

Problem 3. A Model For an Elastic Band

In a simple model of an elastic polymer chain at temperature T , the polymer consists of N links, each of length a with $N \gg 1$. The first link is attached to a fixed point (the black dot), while the last link is pulled down by a force f in the z -direction.



Each link in the chain is independent and can point in one of four directions: Right (R), Left (L), Up (U), or Down (D). The links do not interact with each other and can cross without tangling. The displacement $(\Delta x, \Delta y)$ of the Right links is $(a, 0)$. The displacement of the Down link (D) is $(0, -a)$.

(a) What are the displacements $(\Delta x, \Delta y)$ of the Left and Up links²?

Because of the external force, the downward directed links have lower energy and are more probable. The energies of the four link configurations are tabulated below:

State	Energy
R	0
L	0
U	fa
D	$-fa$

(b) Determine the probability for a link be R , L , U , or D .

(c) In each step of the chain, the y position of the next link is displaced by Δy . Determine the mean Δy per link.

²Not a trick question. Just a question to make the setup clear.

- (d) What is the mean length L in the y direction (see figure). What is the variance in L ?

Hint: relate the variance in L to the variance of Δy .

- (e) What approximately is the probability distribution for L ?

Solution

(a) For left $(\Delta x, \Delta y) = (-a, 0)$. For up $(\Delta x, \Delta y) = (0, a)$.

(b) The partition function is

$$Z = \sum_{s \in RLUD} e^{-\beta \epsilon_s} = 1 + 1 + e^{-\beta f a} + e^{\beta f a} = 2 + 2 \cosh(\beta f a). \quad (33)$$

The probabilities are

$$P_R = P_L = \frac{1}{2 + 2 \cosh(\beta f a)} \quad P_U = \frac{e^{-\beta f a}}{2 + 2 \cosh(\beta f a)} \quad P_D = \frac{e^{\beta f a}}{2 + 2 \cosh(\beta f a)}. \quad (34)$$

(c) The mean Δy is

$$\langle \Delta y \rangle = 0 \cdot P_R + 0 \cdot P_L + a P_U - a P_D = a(P_U - P_D) = -a \frac{\sinh(\beta f a)}{1 + \cosh(\beta f a)}. \quad (35)$$

(d) The mean length is

$$\bar{L} = N \langle \Delta y \rangle = -N a \frac{\sinh(\beta f a)}{1 + \cosh(\beta f a)}. \quad (36)$$

where the minus sign indicates a downward direction. The average of $(\Delta y)^2$ is

$$\langle (\Delta y)^2 \rangle = 0^2 P_R + 0^2 P_L + a^2 P_U + (-a)^2 P_D = \frac{a^2 \cosh(\beta f a)}{1 + \cosh(\beta f a)}. \quad (37)$$

Using that

$$\delta y^2 = \langle (\Delta y)^2 \rangle - \langle \Delta y \rangle^2, \quad (38)$$

leads ultimately to the variance

$$\langle \delta y^2 \rangle = a^2 \left[\frac{\cosh(\beta f a)}{1 + \cosh(\beta f a)} - \frac{\sinh^2(\beta f a)}{(1 + \cosh(\beta f a))^2} \right] = \frac{a^2 \cosh(\beta f a)}{(1 + \cosh(\beta f a))^2}, \quad (39)$$

where we used $\cosh^2(x) - \sinh^2(x) = 1$. The variance of the total length is just N times the variance for each link

$$\langle \delta L^2 \rangle = N \langle \delta y^2 \rangle. \quad (40)$$

We used that the variance of a sum is the sum of the variances for independently distribution quantities.

(e) The probability is Gaussian

$$d\mathcal{P} = \frac{1}{\sqrt{2\pi \langle \delta L^2 \rangle}} \exp(-(L - \bar{L})^2 / 2 \langle \delta L^2 \rangle). \quad (41)$$

This is the central limit theorem.