

Manipulating Taylor Series

$$a) \quad \frac{1}{1+x} = 1 - x + x^2 - x^3 + \dots$$

This follows from the geometric series
 $1/(1-u) = 1 + u + u^2 + \dots$ with $u = -x$

Integrating

$$\int_0^x \frac{dx'}{1+x'} = \log(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} - \frac{x^4}{4} + \dots$$

b)

c) For x large we have $e^{-x} \ll 1$

$$\frac{1}{e^x - 1} = \frac{e^{-x}}{1 - e^{-x}}$$

call $u = e^{-x} \ll 1$

$$\frac{1}{e^x - 1} = \frac{u}{(1-u)} = u (1 + u + u^2 + \dots)$$

$$\boxed{\frac{1}{e^x - 1} = e^{-x} (1 + e^{-x} + e^{-2x} + O(e^{-3x}))}$$

d) $\frac{1}{e^x - 1}$

we expand $e^x = 1 + x + \frac{x^2}{2} + \frac{x^3}{6}$

$$\frac{1}{e^x - 1} \approx \frac{1}{x + x^2/2 + x^3/6} = \frac{1}{x} \frac{1}{(1 + x/2 + x^2/6 + O(x^3))}$$

calling, $u = x/2 + x^2/6$, we have

$$\frac{1}{e^x - 1} \approx \frac{1}{x} \left(\frac{1}{1+u} + O(x^3) \right)$$

$$\approx \frac{1}{x} (1 - u + u^2 + O(x^3))$$

$$\approx \frac{1}{x} \left(1 - \left(\frac{x}{2} + \frac{x^2}{6} \right) + \left(\frac{x^2}{4} \right) + O(x^3) \right)$$

$$\frac{1}{e^x - 1} \approx \frac{1}{x} \left(1 - \frac{x}{2} + \frac{x^2}{12} \right) + O(x^2)$$

e) $\frac{1}{e^{-x} + 1} \approx 1 - e^{-x} + e^{-2x}$ set $u = e^{-x}$

f) $\log(1 - e^{-x}) \approx \log(1 - (1 - x + \frac{x^2}{2} - \frac{x^3}{6}))$

$$\approx \log x \left(1 - \frac{x}{2} + \frac{x^2}{6} \right) = \log x + \log \left(1 - \frac{x}{2} + \frac{x^2}{6} \right)$$

So

$$\log(1 - e^{-x}) = \log(x) + \log\left(1 - \overbrace{\frac{x}{2} + \frac{x^2}{6}}^{\text{call it } u}\right) + O(x^3)$$

Setting $u = -\frac{x}{2} + \frac{x^2}{6}$ we have

$$\log(1 + u) = u - \frac{u^2}{2} + O(u^3) \quad \text{with } x \text{ of order } u$$

So

$$\log(1 - e^{-x}) \approx \log(x) + \left(-\frac{x}{2} + \frac{x^2}{6}\right) - \frac{1}{2}\left(-\frac{x}{2}\right)^2 + O(x^3)$$

$$\log(1 - e^{-x}) \approx \log x - \frac{x}{2} + \frac{x^2}{24} + O(x^3)$$

Energy of SHO

a) We have

$$Z = \frac{1}{1 - e^{-\beta \hbar \omega_0}}$$

Then

$$\langle E \rangle = - \frac{\partial}{\partial \beta} \log Z = + \frac{\partial}{\partial \beta} \log (1 - e^{-\beta \hbar \omega_0})$$

$$= \frac{1}{1 - e^{-\beta \hbar \omega_0}} e^{-\beta \hbar \omega_0} \hbar \omega_0$$

$$\boxed{\langle E \rangle = \frac{\hbar \omega_0}{e^{\beta \hbar \omega_0} - 1}}$$

b) Then

$$\frac{\langle E \rangle}{\hbar \omega_0} = \frac{1}{e^{\hbar \omega_0 / k_B T} - 1} = \langle n \rangle$$

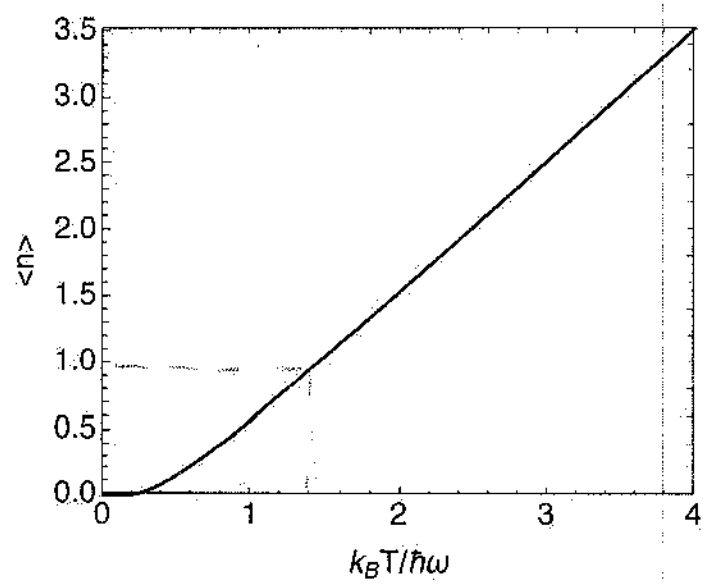
Then



Then from graph

$$\langle n \rangle = 1 \quad \text{when}$$

$$k_B T / \hbar \omega_0 = 1.45$$



or

$$k_B T = 1.45 \hbar \omega_0$$

(c) Then a nice plot of $\langle E \rangle$ is given in the problem statement.

(d) Using the series of problem 1

with $x \equiv \hbar \omega_0 / k_B T$

at low temperature $k_B T \ll \hbar \omega_0$ then $x \gg 1$,
and

$$\frac{1}{e^x - 1} = \frac{e^{-x}}{1 - e^{-x}} \approx e^{-x} (1 + e^{-x} + \dots)$$

And

$$\langle E \rangle = \hbar \omega_0 e^{-\beta \hbar \omega_0} (1 + e^{-\beta \hbar \omega_0} + \dots)$$

At high temperature $x \ll 1$

$$\frac{1}{e^x - 1} \approx \frac{1}{x} - \frac{1}{2}$$

$$\langle E \rangle = \hbar \omega_0 \left(\frac{k_B T}{\hbar \omega_0} - \frac{1}{2} \right) \approx k_B T \left(1 - \frac{\hbar \omega_0}{2 k_B T} \right)$$

e) At high temperature the number of quanta $\langle n \rangle$ is very large. In this regime $\langle n \rangle \gg 1$, quantum mechanics becomes continuous, $\frac{\Delta E}{E} \ll 1$, and it approaches classical mechanics.

This is the Bohr correspondence principle

f) We have

$$i) \quad U = N \left[\frac{5}{2} kT + \frac{\hbar \omega_0}{e^{\beta \hbar \omega_0} - 1} \right]$$

this is $f_0(T)$

Then

$$ii) \quad C_V = \left(\frac{dU}{dT} \right)_V = N \left[\frac{5}{2} k + \frac{-\hbar \omega_0 e^{\beta \hbar \omega_0}}{(e^{\beta \hbar \omega_0} - 1)^2} \hbar \omega_0 \frac{2}{dT} \frac{1}{kT} \right]$$

$$= N \left[\frac{5}{2} k + \frac{(\beta \hbar \omega_0)^2 e^{\beta \hbar \omega_0}}{(e^{\beta \hbar \omega_0} - 1)} k \right]$$

$$C_V = Nk \left[\frac{5}{2} + \frac{(\beta \hbar \omega_0)^2 e^{\beta \hbar \omega_0}}{(e^{\beta \hbar \omega_0} - 1)^2} \right]$$

So

$$C_p = C_v + Nk_B$$

$$C_p = Nk_B \left[\frac{7}{2} + \frac{(\beta \hbar \omega_0)^2}{(e^{\beta \hbar \omega_0} - 1)^2} \right]$$

iii) So we see that the model nicely captures the transition from $C_p = \frac{7}{2} = 3.5$ to $\frac{9}{2} = 4.5$

but misses the transition to $\frac{5}{2}$ at low temperatures