

Two State

$$\langle \epsilon \rangle = 0 \cdot P_0 + \Delta P_1$$

$$P_n = \frac{e^{-\beta \epsilon_n}}{\sum_n e^{-\beta \epsilon_n}} = \frac{e^{-\beta \epsilon_n}}{1 + e^{-\beta \Delta}}$$

★ $\boxed{\langle \epsilon \rangle = \frac{\Delta e^{-\beta \Delta}}{1 + e^{-\beta \Delta}}}$

$$\langle \epsilon^2 \rangle = 0^2 P_0 + \Delta^2 \frac{e^{-\beta \Delta}}{(1 + e^{-\beta \Delta})}$$

So

$$\langle \delta \epsilon^2 \rangle = \langle \epsilon^2 \rangle - \langle \epsilon \rangle^2$$

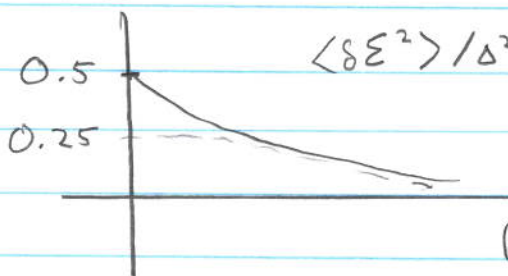
$$= \Delta^2 \frac{e^{-\beta \Delta}}{(1 + e^{-\beta \Delta})} - \frac{e^{-2\beta \Delta}}{(1 + e^{-\beta \Delta})^2}$$

$$= \Delta^2 \left[\frac{e^{-\beta \Delta} (1 + e^{-\beta \Delta}) - e^{-2\beta \Delta}}{(1 + e^{-\beta \Delta})^2} \right]$$

★ $\boxed{\delta \epsilon^2 = \Delta^2 \left(\frac{e^{-\beta \Delta}}{(1 + e^{-\beta \Delta})^2} \right)}$

$\langle \epsilon \rangle / \Delta$

So



$\langle \delta \epsilon^2 \rangle / \Delta^2$ is dashed line. When $\beta \Delta \approx 0$, then the atoms are equally likely to

be in either state $\langle \epsilon \rangle = 0 \cdot \frac{1}{2} + \Delta \cdot \frac{1}{2}$

Speed Distribution

So

$$a) \quad P(v) dv = \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-mv^2/2kT} 4\pi v^2 dv$$

• Maximizing $P(v)$ we have

$$P(v) = C e^{-v^2/2\sigma^2} v^2 \quad \text{with} \quad \sigma \equiv \left(\frac{kT}{m} \right)^{1/2}$$

$$P' = C e^{-v^2/2\sigma^2} \left(\frac{v^3}{\sigma^2} + 2v \right)$$

• So P is maximized when $P'(v) = 0$ or

$$v^2 = 2\sigma^2 \Rightarrow \boxed{v_* = \sqrt{\frac{2kT}{m}}}$$

$$b) \quad \underline{P} = \int_{v_*}^{2v_*} \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-mv^2/2kT} 4\pi v^2 dv$$

• Substituting

$$u = \frac{v}{(kT/m)^{1/2}} \quad \text{this becomes}$$

$$\rho = \int_{\sqrt{2}}^{2\sqrt{2}} \frac{1}{(2\pi)^{3/2}} 4\pi e^{-u^2/2} u^2 du$$

• So

$$\rho = \sqrt{\frac{2}{\pi}} \int_{\sqrt{2}}^{2\sqrt{2}} e^{-u^2/2} u^2 du$$

$$\approx 0.53$$

See program

```
from math import *

xmin = sqrt(2.)
xmax = sqrt(2.)*2.

n = 1000
dx = (xmax - xmin)/n

s = 0.
for i in range(0, n):
    x = i * dx + xmin
    s = s + dx * sqrt(2./pi) * exp(-x*x/2.) * x * x
print(s)
```

Energies

- $$P(v) dv = \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-mv^2/2kT} 4\pi v^2 dv$$

$$v = \left(\frac{2E}{m} \right)^{1/2} \quad dv = \frac{1}{2} \frac{2E}{\left(\frac{2E}{m} \right)^{1/2} m} = \frac{dE}{(2mE)^{1/2}}$$

• So

$$P(E) dE = 4\pi \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-E/kT} \frac{2E}{m} \frac{dE}{\sqrt{2mE}^{1/2}}$$

$$= 2\pi \frac{2^{3/2}}{2^{3/2}} \frac{m^{3/2}}{m^{3/2}} \frac{E^{1/2} dE}{(k_B T)^{3/2}} \frac{1}{\pi^{3/2}} e^{-E/kT}$$

$$P(E) dE = \frac{2}{\sqrt{\pi}} \beta^{3/2} e^{-\beta E} E^{1/2} dE$$

So

- $$\langle E \rangle = \int_0^{\infty} E P(E) dE$$

$$= \int_0^{\infty} \frac{2}{\sqrt{\pi}} \beta^{3/2} e^{-\beta E} E^{1/2} dE \times E$$

• Change vars $u = \beta \epsilon$

$$\langle \epsilon \rangle = \frac{1}{\beta} \int_0^{\infty} \frac{2}{\sqrt{\pi}} e^{-u} u^{3/2} du$$

$$\boxed{\langle \epsilon \rangle} = \frac{1}{\beta} \frac{2}{\sqrt{\pi}} \Gamma(5/2) = \frac{1}{\beta} \frac{2}{\sqrt{\pi}} \frac{3}{2} \Gamma\left(\frac{3}{2}\right)$$

Similarly

$$= \frac{1}{\beta} \cdot 2 \cdot \frac{3}{2} \cdot \frac{1}{2} = \boxed{\frac{3}{2} k_B T}$$

$$\langle \epsilon^2 \rangle = \frac{1}{\beta^2} \int_0^{\infty} \frac{2}{\sqrt{\pi}} e^{-u} u^{5/2} du$$

$$\langle \epsilon^2 \rangle = \frac{1}{\beta^2} \frac{2}{\sqrt{\pi}} \Gamma(7/2)$$

So

$$\boxed{\langle \epsilon^2 \rangle} = \frac{1}{\beta^2} \frac{2}{\sqrt{\pi}} \frac{5}{2} \cdot \frac{3}{2} \cdot \frac{1}{2} \Gamma(1/2)$$

$$= \frac{1}{\beta^2} \frac{2 \cdot 5}{2} \cdot \frac{3}{2} \cdot \frac{1}{2} = \boxed{\frac{1}{\beta^2} \frac{15}{4}}$$

So

$$\boxed{\langle \epsilon^2 \rangle - \langle \epsilon \rangle^2} = \frac{1}{\beta^2} \left(\frac{15}{4} - \frac{9}{4} \right) = \frac{3}{2} (k_B T)^2$$

Distribution on Sphere

a) The probability per area is constant

$$d\mathcal{P} \propto dA \propto R^2 d\Omega$$

• The normalization constant is $1/\text{total } A = 1/4\pi R^2$
So,

$$d\mathcal{P}_{\theta, \phi} = \frac{R^2 d\Omega}{4\pi R^2} = \frac{d\Omega}{4\pi} = \sin\theta \frac{d\theta d\phi}{4\pi}$$

• If we don't care about ϕ , we can integrate over ϕ

$$d\mathcal{P}_{\theta} = \int_0^{2\pi} d\phi d\mathcal{P}_{\theta, \phi} = \int_0^{2\pi} d\phi \frac{\sin\theta d\theta}{4\pi}$$

$$d\mathcal{P}_{\theta} = \frac{1}{2} \sin\theta d\theta$$

← This is how the book writes

$$d\mathcal{P}_{\theta} = \frac{1}{2} \sin\theta d\theta$$

So

$$\langle \cos^2\theta \rangle = \int_0^{\pi} \cos^2\theta \frac{1}{2} \sin\theta d\theta = \left. -\frac{1}{6} \cos^3\theta \right|_0^{\pi} = \frac{1}{3}$$

b) So

• $\int d\mathcal{P} = 1 \quad \leftarrow \text{normalization condition}$

$$\int_0^\pi C (1 + \cos^2 \theta) \frac{1}{2} \sin \theta d\theta = 1$$

$$C \left(-\frac{1}{2} \cos \theta - \frac{1}{6} \cos^3 \theta \right) \Big|_0^\pi = 1$$

$$\frac{1}{2} C \left(1 + \frac{1}{3} \right) = 1 \quad \text{or} \quad \boxed{C = 3/4}$$

• So

$$d\mathcal{P} = \frac{3}{4} (1 + \cos^2 \theta) \frac{\sin \theta d\theta d\phi}{4\pi}$$

Integrating over ϕ as in part (a), $\frac{d\Omega}{4\pi} \rightarrow \frac{1}{2} \sin \theta d\theta$
So the probability is

$$\boxed{d\mathcal{P}_\theta = \frac{3}{4} (1 + \cos^2 \theta) \frac{1}{2} \sin \theta d\theta}$$

Change vars, $u = \cos \theta$, $\frac{du}{d\theta} = -\sin \theta$.

$u = [-1, 1]$ since $\theta \in [0, \pi]$

Note $d\theta = du / |du/d\theta|$ when we use unoriented integrals.

$$d\rho_u = \frac{3}{8} (1 + \cos^2\theta) \frac{\sin\theta \, du}{|du/d\theta|}$$

$$d\rho_u = \frac{3}{8} (1 + u^2) \, du \quad \text{with } u = [-1, \dots, 1]$$

So

$$\langle \cos^2\theta \rangle_{\text{sphere}} = \langle u^2 \rangle_u = \int_{-1}^1 \frac{3}{8} (1 + u^2) u^2 \, du$$

$$= \frac{2}{5}$$