

Review Lecture 3.

Classical Correspondences, Superposition, Particles as wave packets. —

Planck & Einstein found that light previously thought as a wave, can behave on the subatomic scale like a collection of particles, each with an energy associated with  $h\nu$ .

de Broglie argued that a nonrelativistic electron a particle, acts like a wave with a characteristic wavelength  $\lambda = \frac{h}{m_e v}$ .

In the "old" quantum mechanics of Bohr — he raised some conceptual difficulties:

- The angular momentum of an orbiting electron is quantized in units of  $\hbar$ , for the orbits to be "stationary"

Bohr's condition for motion to be circular can be written

$$\oint p \, dq = n\hbar \implies m_e v (2\pi r) = n\hbar$$

$$m_e v r = \frac{\hbar}{2\pi} n = n\hbar.$$

de Broglie gave us another way of thinking about Bohr's circular orbits

$$\oint dy = n\lambda \implies 2\pi r = n \left( \frac{h}{m_e v} \right)$$

$m_e v r = n\hbar.$

The stationary orbits correspond to an integer # of wavelengths.

The stationary object is therefore accounted for by de Broglie by stipulating that  $e$  behave as waves with a given wavelength  $\lambda_e$ .

$$\lambda_e = \frac{h}{m_e v}$$

Correspondence Principle —

Classical Mechanics must emerge as a limit for quantum mechanics at large  $n$ 's.

Classical Mechanics also emerges from quantum mechanics as  $h \rightarrow 0$ .

\* QM becomes Classical as  $n \rightarrow \infty$ ,  $h \rightarrow 0$

The principle of superposition & wave packets —

The accumulation of physics in the early 20<sup>th</sup> century indicating that light: particles  $\leftrightarrow$  waves led us to a point where particles w/ definite positions & momentum are way to a description in terms of waves & wave packets. This seemed surprising & some of the properties achievable via waves.

Familiar waves, like those for sound or water, in 1D take the form

$$\frac{\partial^2 \psi}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2}$$

We look for solutions in the form of an oscillation wave.

$$\psi(x, t) = \phi(x) e^{i\omega t} \Rightarrow \frac{\partial^2 \phi(x)}{\partial x^2} = -\frac{\omega^2}{c^2} \phi(x)$$



"fine"

$\phi(x) = e^{ikx}$  where  $\hbar k^2 = \omega^2$  so that  $\psi(x,t) = e^{i(kx - \omega t)}$  w/  $\hbar\omega = \pm \frac{\hbar^2 k^2}{2m}$

A property of this equation is that it is linear in  $\psi$ , which yields to the principle of superposition:

if  $\psi_1(x,t)$  &  $\psi_2(x,t)$  satisfy the wave equation, so does

$$\psi_1(x,t) + \psi_2(x,t)$$

This property allows us to build wave packets as one would get by throwing a rock into the middle of a pond. It would give us difficulty trying to envision quantum mechanics w/o this property.

The correspondence principle requires us to be able to create localized particles in quantum mechanics, so we know how to build localized wave packets from waves through Fourier analysis:

$$\psi(x,t=0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk \quad \phi(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \psi(x,t=0) e^{-ikx} dx$$

for example,

$$\psi(x,t=0) = e^{-x^2/2a^2} \implies \psi(x,t=0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-a^2 k^2/4} e^{ikx} dk$$

However, this infinite sum of plane waves will not be a solution of our quantum mechanics wave equation unless our superposition principle holds.

## Wave Packets & Uncertainty Relationships

There is a size scale associated w/ our coordinate space wave packet,  $\Delta x \sim a$ .  
The smaller  $a$  — the more localized in  $x$  — the broader range of contributing momenta — space waves.

$$e^{-a^2 k^2 / 4} = e^{-k^2 / (2/a)^2} \Rightarrow \Delta k = \frac{2}{a} \text{ consequently } \Delta x \Delta k \sim 1.$$

The more localized wave packets are in coordinate space, the more they disperse in momentum space.

As the de Broglie relation gives us  $p = \frac{h}{\lambda} = \frac{h}{2\pi} k$  so that  $\Delta k = \frac{\Delta p}{\hbar}$ ,  
Quantum Mechanics has an uncertainty principle relating the product of coordinate & momentum uncertainties to  $\hbar$ .

## Heisenberg's Uncertainty Principle

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

## Schrodinger's Equation

These kinds of conservation helped Schrodinger create a wave equation that might account for the apparent dual nature of the  $e^-$  & other elementary particles. In 1D,

$$\left[ \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \psi(x,t) = i\hbar \frac{\partial}{\partial t} \psi(x,t)$$



So what does this equation mean? Standing at the LHS, the potential is here while the derivative term can be rewritten for a very different units.

$$\frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} = \frac{-\hbar^2 c^2}{2mc^2} \frac{\partial^2}{\partial x^2}$$

As  $(\hbar c)^2$  has units of (Energy-distance)<sup>2</sup>,  $mc^2$  is an energy, and  $\frac{\partial^2}{\partial x^2}$  has units of (distance)<sup>-2</sup>. The first term on the LHS must be energy. More specifically, Kinetic Energy.

Classically KE is  $\frac{P^2}{2m}$  and since  $\frac{P^2}{2m} + V(x) = E_{\text{total}}$  the RHS must be energy  $E$ . This requirement that Schrödinger's Equation  $\Rightarrow$  its differential operators correspond w/ classical mechanics allows us to define the Equation's differential operators

Schrödinger's Equation Operators

$$\hat{P} \longleftrightarrow \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\hat{E} \longleftrightarrow i\hbar \frac{\partial}{\partial t}$$

Schrödinger's Equation can be rewritten

$$\left[ \frac{\hat{P}^2}{2m} + V(x) \right] \psi(x,t) = \hat{E} \psi(x,t)$$

$\hat{P} \hat{E}$  are differential operators that act on the wave function

Consider a free particle, so  $V \rightarrow 0$ , w/ a wavefunction

$$\psi(x,t) = e^{i(\mathcal{E}t - px)/\hbar}$$

where  $\mathcal{E}$  &  $p$  are numbers. We note that

$$\hat{p}\psi(x,t) = p\psi(x,t) \quad \hat{\mathcal{E}}\psi(x,t) = \mathcal{E}\psi(x,t)$$

while plugging into Schrodinger's Equation yields the constraint

$$\frac{\mathcal{E}}{2m} = \mathcal{E}$$

Therefore, there are an infinite set of plane-wave functions/solutions that are labeled by the number  $p$ , with  $\mathcal{E}(p) = \frac{p^2}{2m}$

The plane-wave solution of the Schrodinger Equation has a characteristic wavelength. The length of a wave corresponds to the distance required to change the phase by  $2\pi$  at a fixed time  $t$ . That is,

$$\lambda_{\pi} = \frac{p \lambda}{\hbar} \equiv \frac{p \lambda}{\hbar} \Rightarrow \lambda = \frac{2\pi \hbar}{p} = \frac{h}{p}$$

de Broglie wavelength.