MA-207 Differential Equations II Lecture-8 Eigenvalue Problem and Fourier Expansion

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Consider the following Eigen Value Problems, where $\lambda \in \mathbb{R}$ and L > 0.

- Problem 1. $y'' + \lambda y = 0$, y(0) = 0, y(L) = 0.
- Problem 2. $y'' + \lambda y = 0$, y'(0) = 0, y'(L) = 0.
- **9** Problem 3. $y'' + \lambda y = 0$, y(0) = 0, y'(L) = 0.
- Problem 4. $y'' + \lambda y = 0$, y'(0) = 0, y(L) = 0.
- Problem 5. $y'' + \lambda y = 0$, y(-L) = y(L), y'(-L) = y'(L).

The boundary condition in problem 5 is called periodic. Obviously, $y \equiv 0$ is a (trivial) solution for all 5 problems for all λ .

For most values of λ , there are no other solutions.

Question. For what values of λ does the "EVP" have a non-trivial solutions and what are the solutions?

Any λ for which the EVP 1-5 has a non-trivial solution is called an eigenvalue of that problem and non-trivial solutions for an eigenvalue λ are called λ -eigenfunction, or eigenfunction associated with λ .

A non-zero constant multiple of a λ -eigenfunction is again a λ -eigenfunction.

Solving an eigenvalue problem means finding all its eigenvalues and associated eigenfunctions.

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THEOREM.

- ullet Problems 1-5 have no negative eigenvalues.
- $\lambda = 0$ is an eigenvalue of Problems 2 and 5 with associated eigenfunctions $y_0 = 1$.
- $\lambda = 0$ is not an eigenvalue of Problems 1, 3, 4.

Proof. Let us prove first two; third is left as an exercise. Suppose $\lambda < 0$. Write $\lambda = -a^2$, a > 0.

The general solution of ODE $y'' = a^2y$ is $y(x) = Ce^{ax} + De^{-ax}$.

In problem 1, y(0)=y(L)=0 gives C+D=0 and $Ce^{aL}+De^{-aL}=0$.

Check $\implies C = D = 0$.

In problem 2, y'(0)=y'(L)=0 gives aC-aD=0 and $aCe^{aL}-aDe^{-aL}=0$.

Since $a \neq 0$, this forces C = D = 0.

Similarly, do the other problems.

Now consider the second statement in the theorem.

If $\lambda = 0$, the solution of y'' = 0 is y(x) = ax + b.

In problem 2, the derivative vanishes, so a = 0.

Thus, y(x) = constant is the solution.

In problem 5, y(-L) = y(L) gives -aL + b = aL + b. Thus a = 0 and y(x) = constant is the solution.

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Eigenvalue Problem 1

Theorem

The eigenvalue problem

$$y'' + \lambda y = 0$$
, $y(0) = 0$, $y(L) = 0$

has infinitely many positive eigenvalues

$$\lambda_n = \frac{n^2 \pi^2}{L^2}$$

with associated eigenfunctions

$$y_n = \sin \frac{n\pi x}{L}, \quad n = 1, 2, \dots.$$

There are no other eigenvalues.

Proof. Any eigen value must be positive (previous theorem).

If y is a solution of $y'' + \lambda y = 0$ with $\lambda > 0$, then

$$y(x) = c_1 \cos \sqrt{\lambda} x + c_2 \sin \sqrt{\lambda} x$$

Now
$$y(0) = 0 \implies c_1 = 0$$

$$\implies y(x) = c_2 \sin \sqrt{\lambda} x$$
, $c_2 \neq 0$ as $y \neq 0$.

$$y(L) = 0 \implies \sin \sqrt{\lambda} L = 0$$

$$\implies \sqrt{\lambda}L = n\pi$$
. Therefore, for $n \ge 1$, $\lambda_n = \frac{n^2\pi^2}{L^2}$

is an eigenvalue with an associated eigenfunction

$$y_n = \sin \frac{n\pi x}{L}.$$

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Theorem

The eigenvalue problem 2

$$y'' + \lambda y = 0$$
, $y'(0) = 0$, $y'(L) = 0$

has an eigenvalue $\lambda_0 = 0$ with eigenfunction $y_0 = 1$ and infinitely many positive eigenvalues

and infinitely many positive eigenvalues
$$\lambda_n = \frac{n^2\pi^2}{L^2}, \quad n=1,2,\dots$$

with associated eigenfunctions

$$y_n = \cos \frac{n\pi x}{L}.$$

There are no other eigenvalues.

Proof. Similar to the proof of Problem 1, hence is left as an exercise.

Theorem

The eigenvalue problem 3

$$y'' + \lambda y = 0$$
, $y(0) = 0$, $y'(L) = 0$

has infinitely many positive eigenvalues

$$\lambda_n = \frac{(2n-1)^2 \pi^2}{4L^2}, \quad n = 1, 2, \dots$$

with associated eigenfunctions

$$y_n = \sin\frac{(2n-1)\pi x}{2L}.$$

There are no other eigenvalues.

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Proof. Any eigenvalue must be positive.

If y is a solution of $y'' + \lambda y = 0$ with $\lambda > 0$, then

$$y(x) = c_1 \cos \sqrt{\lambda}x + c_2 \sin \sqrt{\lambda}x$$

$$y(0) = 0 \implies c_1 = 0$$

So $y(x) = c_2 \sin \sqrt{\lambda}x$ with $c_2 \neq 0$, as $y \neq 0$.

$$y'(L) = 0 \implies \cos\sqrt{\lambda}L = 0$$

$$\implies \sqrt{\lambda}L = \frac{(2n-1)\pi}{2}.$$

Therefore, for $n \geq 1$, $\lambda_n = \frac{(2n-1)^2\pi^2}{4L^2}$ is an eigenvalue with an associated eigenfunction

$$y_n = \sin\frac{(2n-1)\pi x}{2L}$$

Theorem. The eigenvalue problem 4

$$y'' + \lambda y = 0$$
, $y'(0) = 0$, $y(L) = 0$

has infinitely many positive eigenvalues

$$\lambda_n = \frac{(2n-1)^2 \pi^2}{4L^2}, \quad n = 1, 2, \dots$$

with associated eigenfunctions

$$y_n = \cos\frac{(2n-1)\pi x}{2L}$$

There are no other eigenvalues.

Proof. Similar to proof of Problem 3, hence is left as an exercise.

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Theorem. The eigenvalue problem 5

$$y'' + \lambda y = 0$$
, $y(-L) = y(L)$, $y'(-L) = y'(L)$

has an eigenvalue $\lambda_0=0$ with associated eigenfunction $y_0=1$

and infinitely many positive eigenvalues

$$\lambda_n = \frac{n^2 \pi^2}{L^2}, \quad n = 1, 2, \dots$$

with associated eigenfunctions

$$y_{1n} = \cos \frac{n\pi x}{L}$$
 and $y_{2n} = \sin \frac{n\pi x}{L}$.

There are no other eigenvalues.

Proof. We know that $\lambda = 0$ is an eigenvalue with eigenfunction $y_0 = 1$

and any other eigenvalue is positive.

If y is a solution of $y'' + \lambda y = 0$ with $\lambda > 0$, then

$$y(x) = c_1 \cos \sqrt{\lambda} x + c_2 \sin \sqrt{\lambda} x$$

The boundary condition y(-L) = y(L) implies

$$c_1 \cos(-\sqrt{\lambda}L) + c_2 \sin(-\sqrt{\lambda}L) = c_1 \cos(\sqrt{\lambda}L) + c_2 \sin(\sqrt{\lambda}L)$$

Since $\cos(-\sqrt{\lambda}x) = \cos(\sqrt{\lambda}x)$, and $\sin(-\sqrt{\lambda}x) = -\sin(\sqrt{\lambda}x)$

We get
$$c_2 \sin \sqrt{\lambda} x = 0$$

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Similarly, using y'(-L) = y'(L) for

$$y'(x) = \sqrt{\lambda} \left(-c_1 \sin \sqrt{\lambda} x + c_2 \cos \sqrt{\lambda} x \right)$$

we get

$$c_1 \sin \sqrt{\lambda} L = 0.$$

For a non-trivial solution, $(c_1, c_2) \neq (0, 0)$, we get

$$\sin \sqrt{\lambda} L = 0 \implies \sqrt{\lambda} L = n\pi \text{ for } n \ge 1.$$

The eigenvalues are $\lambda_n = \frac{n^2 \pi^2}{L^2}$ $n = 1, 2, \dots$

and each eigenvalue has two linearly independent associated eigenfunctions

$$\cos \frac{n\pi x}{L}$$
 and $\sin \frac{n\pi x}{L}$.

Definition. Recall that two square integrable functions f and g in $L^2([a,b])$ are orthogonal if

$$\langle f, g \rangle := \int_a^b f(x)g(x) \, dx = 0$$

More generally, a set of functions $\{\phi_1, \phi_2, \dots, \phi_n, \dots\}$ (finite or infinitely many) are orthogonal on [a, b] if

$$\int_{a}^{b} \phi_{i}(x)\phi_{j}(x) dx = 0, \quad \text{whenever} \quad i \neq j$$

We have seen orthogonality of Legendre function. We will study Fourier series with respect to different orthogonal systems.

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Exercise. Consider the eigenfunctions of EVP's 1-5.

$$\bullet \sin \frac{\pi x}{L}, \sin \frac{2\pi x}{L}, \dots, \sin \frac{n\pi x}{L}, \dots$$

$$2 1, \cos \frac{\pi x}{L}, \cos \frac{2\pi x}{L}, \dots, \cos \frac{n\pi x}{L}, \dots$$

$$3\sin\frac{\pi x}{2L}$$
, $\sin\frac{3\pi x}{2L}$, ..., $\sin\frac{(2n-1)\pi x}{2L}$, ...

$$\bullet \cos \frac{\pi x}{2L}, \cos \frac{3\pi x}{2L}, \dots, \cos \frac{(2n-1)\pi x}{2L}, \dots$$

1,
$$\cos \frac{\pi x}{L}$$
, $\sin \frac{\pi x}{L}$, $\cos \frac{2\pi x}{L}$, $\sin \frac{2\pi x}{L}$, ..., $\cos \frac{n\pi x}{L}$, $\sin \frac{n\pi x}{L}$, ...

Show that eigenfunctions of (1-4) are orthogonal on [0, L] and of (5) is orthogonal on [-L, L].

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Theorem. Suppose the functions ϕ_1, ϕ_2, \ldots , are orthogonal on [a, b] and

$$\|\phi_n\|^2 := \int_a^b \phi_n^2(x) \, dx \neq 0 \quad n = 1, 2, \dots$$

Let c_1, c_2, \ldots , be constants and $0 < M < \infty$ s.t.

$$\left| \sum_{m=1}^{N} c_m \, \phi_m(x) \right| \le M, \quad \text{for } a \le x \le b \text{ and } N \ge 1.$$

Suppose also that the series $\sum_{m=1}^{\infty} c_m \phi_m(x)$

converges to f(x) and f(x) is integrable on [a,b]. Then

$$c_n = \frac{1}{\|\phi_n\|^2} \int_a^b f(x)\phi_n(x) dx, \quad n = 1, 2, \dots$$

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Proof.
$$f(x) = \sum_{m=1}^{\infty} c_m \phi_m(x)$$
 gives

$$\int_a^b f(x)\phi_n(x) dx = \int_a^b \left(\sum_{m=1}^\infty c_m \phi_m(x)\right) \phi_n(x) dx$$

Boundedness of partial sums and integrability of f allows us to interchange the operation of integration and summation, so

$$\int_{a}^{b} f(x)\phi_{n}(x) dx = \sum_{m=1}^{\infty} c_{m} \int_{a}^{b} \phi_{m}(x)\phi_{n}(x) = c_{n} \|\phi_{n}\|^{2}$$

by orthogonality of $\{\phi_1, \phi_2, \ldots\}$.

Definition. Suppose $\{\phi_1, \phi_2, \ldots\}$ is orthogonal on

$$[a,b] \text{ and } \|\phi_n\|^2 := \int_a^b \phi_n^2(x) \, dx \neq 0, \ n = 1, 2, \dots$$

For $f \in L^2([a,b])$, the series

$$\sum_{n=1}^{\infty} c_n \, \phi_n(x), \text{ where } c_n = \frac{1}{\|\phi_n\|^2} \int_a^b f(x) \phi_n(x) \, dx$$

is Fourier series of f and c_1, c_2, \ldots the Fourier coefficients of f w.r.t. orthogonal set $\{\phi_n\}_{n=1}^{\infty}$.

We write
$$f(x) \sim \sum_{n=1}^{\infty} c_n \, \phi_n(x), \quad a \leq x \leq b$$

Qn. What about the convergence of series to f(x)? The answer depends on the orthogonal system $\{\phi_1, \phi_2, \ldots\}$.

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Fourier Series.

The set of eigenfunctions B =

$$\left\{1, \cos\frac{\pi x}{L}, \sin\frac{\pi x}{L}, \cos\frac{2\pi x}{L}, \dots, \cos\frac{n\pi x}{L}, \sin\frac{n\pi x}{L}, \dots\right\}$$

of Eigenvalue Problem 5 is orthogonal on [-L, L].

The Fourier series of $f \in L^2([-L, L])$ w.r.t. orthogonal set B is

$$a_0 + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right)$$

$$||1||^2 = 2L$$
, $||\cos \frac{n\pi x}{L}||^2 = L$, $||\sin \frac{n\pi x}{L}||^2 = L$

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So the Fourier coefficients are given by

$$a_0 = \frac{1}{2L} \int_{-L}^{L} f(x) \, dx$$

$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos \frac{n\pi x}{L} dx, \quad n > 0$$

$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin \frac{n\pi x}{L} dx, \quad n > 0$$

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Convergence of Fourier Series in norm

For $f \in L^2([-L, L])$, the partial sum of Fourier series of f converges to f in norm, i.e.

$$||f - a_0 - \sum_{n=1}^{N} \left(a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right)||$$

conveges to 0 as $N \to \infty$.

What about pointwise convergence of Fourier series.

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Definition. Recall that a function f is said to be piecewise smooth if

- f has atmost finitely many discontinuity.
- f' exists and is continuous except at finitely many points.
- $f(x_0+) = \lim_{x \to x_0^+} f(x) \text{ and } f'(x_0+) = \lim_{x \to x_0^+} f'(x) \text{ exists if } a \le x_0 < b.$
- $f(x_0-) = \lim_{x \to x_0^-} f(x)$ and $f'(x_0-) = \lim_{x \to x_0^-} f'(x)$ exists if $a < x_0 \le b$.

Hence f is piecewise smooth if and only if f, f' have atmost finitely many jump discontinuity and f' is piecewise continuous.

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Pointwise Convergence of Fourier series

Let f(x) be a piecewise smooth on [-L, L]. Extend f(x) to $\mathbb R$ by periodicity f(x+2L)=f(x). Then the Fourier series of f on [-L, L],

$$F(x) = a_0 + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right)$$

converges to $\frac{1}{2}[f(x^+)+f(x^-)]$ for all $x\in\mathbb{R}$. Therefore, F(x)=f(x) if f is continuous at x.

If we change f(x) to $\frac{1}{2}[f(x^+) + f(x^-)]$ at discontinuous points x, then Fourier series of f converges to f(x). Thus two functions can have same Fourier series.

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The Fourier series converges to F(x) for all $x \in [-L, L]$, hence for a fixed x, the error

$$E_N(x) = \left| F(x) - a_0 - \sum_{n=1}^N \left(a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right) \right|$$

is small if N is large.

But we can not make the error $E_N(x)$ arbitrary small "uniformly" for all $x \in [-L, L]$, by choosing N sufficiently large if

- f is discontinuous at some point $\alpha \in (-L, L)$ or
- $\bullet \ f(-L+) \neq f(L-)$

The next result explains this.

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Gibbs Phenomenon.

• If f has a jump discontinuity at $\alpha \in (-L, L)$, then there exist a sequence $u_N \in (-L, \alpha)$ and $v_N \in (\alpha, L)$ s.t.

$$\lim_{N \to \infty} u_N = \alpha, \quad E_N(u_N) \simeq .09 |f(\alpha -) - f(\alpha +)|$$

$$\lim_{N \to \infty} v_N = \alpha, \quad E_N(v_N) \simeq .09 |f(\alpha -) - f(\alpha +)|$$

Maximum of error $E_N(x) \not\to 0$ near α as $N \to \infty$.

• If $f(-L+) \neq f(L-)$, $\exists u_N, v_N \text{ in } (-L, L) \text{ s.t.}$

$$\lim_{N \to \infty} u_N = -L, \quad E_N(u_N) \simeq .09 |f(-L+) - f(L-)|$$

$$\lim_{N \to \infty} v_N = L, \quad E_N(v_N) \simeq .09 |f(-L+) - f(L-)|$$

This is called Gibbs phenomenon.

Example. Find the Fourier series of piecewise smooth function on [-2, 2]

$$f(x) = \begin{cases} -x, & -2 < x < 0 \\ 1/2, & 0 < x < 2 \end{cases}$$

and determine the sum of the Fourier series.

We need not define f at points of discontinuities, i.e. at points -2,0 and 2, since the coefficients in the Fourier series

$$F(x) = a_0 + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{2} + b_n \sin \frac{n\pi x}{2} \right)$$

are not affected by them.

Note f is continuous for $x \in (-2,0) \cup (0,2)$.

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So F(x) = f(x) for x in $(-2,0) \cup (0,2)$. So let us compute F(x) at discontinuous points.

$$F(-2) = F(2) = \frac{1}{2} \left(f(-2+) + f(2-) \right) = \frac{1}{2} \left(2 + \frac{1}{2} \right) = \frac{5}{4}$$

$$F(0) = \frac{1}{2} \left(f(0-) + f(0+) \right) = \frac{1}{2} \left(0 + \frac{1}{2} \right) = \frac{1}{4}$$

To summarize,

$$F(x) = \begin{cases} 5/4, & x = \pm 2 \\ -x, & -2 < x < 0 \\ 1/4, & x = 0 \\ 1/2, & 0 < x < 2 \end{cases}$$

Let us compute the Fourier coefficients now.

$$a_0 = \frac{1}{4} \int_{-2}^{2} f(x) \, dx = \frac{1}{4} \left[\int_{-2}^{0} (-x) \, dx + \int_{0}^{2} \frac{1}{2} \, dx \right] = \frac{3}{4}$$

If $n \geq 1$, then

$$a_n = \frac{1}{2} \int_{-2}^{2} f(x) \cos \frac{n\pi x}{2} dx$$

$$= \frac{1}{2} \left[\int_{-2}^{0} (-x) \cos \frac{n\pi x}{2} dx + \int_{0}^{2} \frac{1}{2} \cos \frac{n\pi x}{2} dx \right]$$

$$= \frac{1}{2} \left[-x \frac{2}{n\pi} \sin \frac{n\pi x}{2} \Big|_{-2}^{0} + \int_{-2}^{0} \frac{2}{n\pi} \sin \frac{n\pi x}{2} dx + 0 \right]$$

$$= \frac{1}{2} \frac{4}{n^2 \pi^2} \left(-\cos \frac{n\pi x}{2} \right) \Big|_{-2}^{0} = \frac{2}{n^2 \pi^2} (\cos n\pi - 1)$$

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$$b_n = \frac{1}{2} \int_{-2}^{2} f(x) \sin \frac{n\pi x}{2} dx$$

$$= \frac{1}{2} \left[\int_{-2}^{0} (-x) \sin \frac{n\pi x}{2} dx + \int_{0}^{2} \frac{1}{2} \sin \frac{n\pi x}{2} dx \right]$$

$$= \frac{1}{2n\pi} (1 + 3\cos n\pi)$$

Thus, the Fourier series of f(x) is

$$F(x) = \frac{3}{4} + \frac{2}{\pi^2} \sum_{n=1}^{\infty} \frac{\cos n\pi - 1}{n^2} \cos \frac{n\pi x}{2} + \frac{1}{2\pi} \sum_{n=1}^{\infty} \frac{1 + 3\cos n\pi}{n} \sin \frac{n\pi x}{2}$$

Definition. Recall a function f is odd if f(-x) = -f(x) and even if f(-x) = f(x).

Product of odd and even functions:

- (odd) (odd)=(even)
- (odd) (even)=(odd)
- (even) (even)=(even)

If f is an odd, then $\int_{-L}^{L} f(x) \, dx = 0$, if f is an even, then $\int_{-L}^{L} f(x) \, dx = 2 \int_{0}^{L} f(x) \, dx$ Note that $\sin \frac{n\pi x}{L}$ is odd and $\cos \frac{n\pi x}{L}$ is even.

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Let $f \in L^2([-L, L])$ be piecewise smooth.

• If f is odd, then $f(x) \cos \frac{n\pi x}{L}$ is odd. Hence $a_n = 0$ and the Fourier series of f is

$$F(x) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$$

which has only sine terms.

• If f is <u>even</u>, then $f(x) \sin \frac{n\pi x}{L}$ is odd. Hence $b_n = 0$ and the Fourier series of f is

$$F(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L}$$

which has only cosine terms.

Fourier sine series.

If $f \in L^2([0,L])$, then extend f to [-L,L] as an odd function by defining f(-x) = -f(x) for $x \in (-L,0)$. Then Fourier series of f on [-L,L] is

$$F(x) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}, \ b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

called Fourier sine series of f on [0, L].

This is the Fourier series of f on [0, L] with respect to orthogonal system of eigen-functions

$$B = \{ \sin \frac{\pi x}{L}, \sin \frac{2\pi x}{L}, \dots, \sin \frac{n\pi x}{L}, \dots \}$$
 of EVP 1: $y'' + \lambda y = 0, \quad y(0) = 0 = y(L).$

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Example. Consider the periodic square wave with $L=\pi$.

$$f(x) = \begin{cases} 1 & \text{if } 0 < x < \pi \\ -1 & \text{if } -\pi < x < 0 \end{cases}, \qquad f(x+2\pi) = f(x)$$

Since f is an odd function, Fourier series of f is Fourier sine series

$$F(x) = \sum_{n=1}^{\infty} b_n \sin n\pi$$

Show that

$$b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin nx \, dx = \begin{cases} \frac{4}{n\pi} & \text{if } n \text{ is odd} \\ 0 & \text{if } n \text{ is even} \end{cases}$$

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Thus the Fourier series of f(x) is

$$\frac{4}{\pi} \left(\sin x + \frac{\sin 3x}{3} + \frac{\sin 5x}{5} + \dots \right)$$

Since F(x) = f(x) at all points of continuity of f and

$$F(0) = \frac{1}{2}[f(0+) + f(0-)] = 0$$

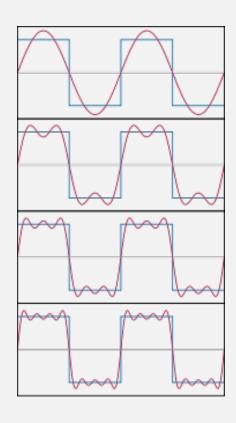
In particular, evaluating at $x = \pi/2$

$$f\left(\frac{\pi}{2}\right) = 1 = \frac{4}{\pi} \left(1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots\right)$$

$$\implies 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots = \frac{\pi}{4}$$

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The <u>partial sums</u> of the Fourier series wiggle around the square wave.



Fourier cosine series

If $f \in L^2([0,L])$, then extend f to [-L,L] as even function by f(-x) = f(x) for $x \in (-L,0)$. Then Fourier series of f on [-L,L] is

$$F(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L}, \quad a_0 = \frac{1}{L} \int_0^L f(x) \, dx$$
$$a_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi x}{L} \, dx, \quad n \ge 1$$

called Fourier cosine series of f on [0, L].

This is the Fourier series of f on [0,L] with respect to orthogonal system of eigen-functions of EVP 2

$$B = \{1, \cos\frac{\pi x}{L}, \cos\frac{2\pi x}{L}, \dots, \cos\frac{n\pi x}{L}, \dots\}.$$

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Example. Find the Fourier series of the function

$$f(x) = x^2 \qquad -\pi \le x \le \pi$$

Since f is an even, $b_n = 0$.

$$a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} x^2 \, dx = \frac{\pi^2}{3}$$

$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} x^2 \cos nx \, dx = \frac{4}{n^2} \cos n\pi$$

$$= \begin{cases} \frac{4}{n^2} & n : \text{is even} \\ -\frac{4}{n^2} & n : \text{is odd} \end{cases}$$

Thus the Fourier series of f(x) is

$$F(x) = \frac{\pi^2}{3} - 4\left(\cos x - \frac{\cos 2x}{4} + \frac{\cos 3x}{9} - \dots\right)$$

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Since f is continuous with $f(\pi) = f(-\pi)$, F(x) = f(x) for all $x \in \mathbb{R}$. Hence

$$x^{2} = \frac{\pi^{2}}{3} - 4\left(\cos x - \frac{\cos 2x}{4} + \frac{\cos 3x}{9} - \dots\right)$$

Evaluating at $x = \pi$,

$$\pi^2 = \frac{\pi^2}{3} + 4\left(1 + \frac{1}{4} + \frac{1}{9} + \dots\right)$$

This yields the identity

$$1 + \frac{1}{2^2} + \frac{1}{3^2} + \dots = \frac{1}{4} \cdot \frac{2\pi^2}{3} = \frac{\pi^2}{6}$$

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