

Classical Mechanics

June – 2014

(1) Consider a particle of mass m in simple harmonic oscillation about the origin with spring constant k ; then for the Lagrangian L and the Hamiltonian H of the system

(a) $L(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$

$H(x, \dot{x}) = \frac{p^2}{2m} + \frac{1}{2} k x^2$, p is generalized momentum.

(b) $(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} k x^2$, $p = m\dot{x}$

(c) $L(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$, $p = m\dot{x}$

(d) $L(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$

$H(x, \dot{x}) = \frac{m\dot{x}^2}{2} + \frac{1}{2} k x^2$

Answer: (a), (c) and (d)

Solution: $K.E = T = \frac{1}{2} m \dot{x}^2$

$F \propto x$

$F = -kx$

$F = -\text{grad } V = -\nabla V = -\frac{\partial V}{\partial x}$

$V = -\int F dx = \int kx dx = \frac{kx^2}{2}$

$\therefore L = T - V = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} kx^2$

$H = T + V = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} kx^2$

$P = \frac{\partial L}{\partial \dot{x}} = m \dot{x}$

Also $H = \frac{1}{2} m \cdot \frac{p^2}{m^2} + \frac{1}{2} kx^2 = \frac{p^2}{2m} + \frac{1}{2} kx^2$

So, the options (a), (c) and (d) are correct.

December – 2014

(1) Consider a body of unit mass falling freely from rest under gravity with velocity v . If the air resistance retards the acceleration by cv where c is a constant then

(a) $v = \frac{g}{c} (1 + e^{ct})$

(b) $v = \frac{g}{c} (1 + e^{-ct})$

(c) $v = \frac{g}{c} (1 - e^{-ct})$

(d) $v = \frac{g}{c} (1 - e^{ct})$

Answer: (c)

Solution: Equation of motion is

$m \cdot \frac{dv}{dt} = mg - mc v$

$$\text{or, } \frac{dv}{g\left(1-\frac{c}{g}v\right)} = dt$$

$$\text{Integrating, } \frac{\log\left(1-\frac{c}{g}v\right)}{g \times \left(-\frac{c}{g}\right)} = t + A \text{ (A being arbitrary constant)}$$

$$-\frac{1}{c} \log\left(1 - \frac{c}{g} v\right) = t + A$$

$$\text{Initially, } t = 0, v = 0 \Rightarrow A = 0$$

$$\therefore \log\left(1 - \frac{c}{g} v\right) = -ct$$

$$\Rightarrow 1 - \frac{c}{g} v = e^{-ct}$$

$$\Rightarrow v = \frac{g}{c} (1 - e^{-ct})$$

So, the option (c) is correct.

June-2015

1. Consider two weightless, inextensible roots AB and BC; Suspended at A and joined by a flexible joint at B. Then the degrees of freedom of the system is

- (a) 3
- (b) 4
- (c) 5
- (d) 6

Answer: (a)

Solution:

$$\text{Degrees of freedom is } = \frac{2(2+1)}{2} = 3$$

So, the option (a) is correct.

2. A particle of mass m is constrained to move on the surface of a cylinder $x^2 + y^2 = a^2$ under the influence of a force directed towards the origin and proportional to the distance of the particle from the origin. Then

- (a) The angular momentum about z-axis is constant.
- (b) The angular momentum about z-axis is not constant.
- (c) The motion is simple harmonic in z-direction.
- (d) The motion is not simple harmonic in z-direction.

Answer: (a), (c)

Solution:

Here the motion is simple harmonic in z-direction and the angular momentum about z-axis is constant.

So, the options (a) and (c) are correct.

December-2015

1. A force $5\hat{i} - 2\hat{j} + 3\hat{k}$ acts on a particle with position vector $2\hat{i} + \hat{j} - 2\hat{k}$. The torque of the force about the origin is

- (a) $\hat{i} + 16\hat{j} + 9\hat{k}$
- (b) $-\hat{i} - 16\hat{j} - 9\hat{k}$
- (c) $\hat{i} + 16\hat{j} - 9\hat{k}$
- (d) $\hat{i} - 16\hat{j} + 9\hat{k}$

Answer: (b)

Solution:

$$\text{Torque is } \vec{r} \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 2 & 1 & -2 \\ 5 & -2 & 3 \end{vmatrix} = -\hat{i} - 16\hat{j} - 9\hat{k}$$

So, the option (b) is correct.

2. Consider a mass m moving in an inverse square central force with characteristic coefficient μ and described by the Lagrangian:

$$L(r, \dot{r}, \theta, \dot{\theta}) = \frac{m}{2}(\dot{r}^2 + r^2\dot{\theta}^2) + \frac{\mu m}{r} \text{ then}$$

- (a) The generalized momenta of the system are $p_r = m\dot{r}$ and $p_\theta = mr^2\dot{\theta}$
- (b) The Hamiltonian of the system is $H = \frac{1}{2m} \left[p_r^2 + \frac{p_\theta^2}{r^2} \right] - \frac{1}{2} \frac{\mu m}{r}$
- (c) The Hamiltonian of the system is $H = \frac{1}{2m} \left[p_r^2 + \frac{p_\theta^2}{r^2} \right] - \frac{\mu m}{r}$
- (d) The generalized momenta of the system are $p_r = m\dot{r}$ and $p_\theta = -mr^2\dot{\theta}$

Answer: (a), (c)

Solution:

$$p_r = \frac{\partial L}{\partial \dot{r}} = m\dot{r}, \quad p_\theta = \frac{\partial L}{\partial \dot{\theta}} = mr^2\dot{\theta}$$

$$\text{Now, } \dot{r} = \frac{p_r}{m} \text{ and } \dot{\theta} = \frac{p_\theta}{mr^2}$$

$$\begin{aligned} H &= \sum p_i q_i - L = p_r \cdot \dot{r} + p_\theta \cdot \dot{\theta} - L \\ &= p_r \cdot \dot{r} + p_\theta \cdot \dot{\theta} - \frac{m}{2}(\dot{r}^2 + r^2\dot{\theta}^2) - \frac{\mu m}{r} \\ &= p_r \cdot \frac{p_r}{m} + p_\theta \cdot \frac{p_\theta}{mr^2} - \frac{m}{2} \left(\frac{p_r^2}{m^2} + \frac{r^2 \cdot p_\theta^2}{m^2 r^4} \right) - \frac{\mu m}{r} \\ &= \frac{p_r^2}{m} + \frac{p_\theta^2}{mr^2} - \frac{p_r^2}{2m} - \frac{p_\theta^2}{2mr^2} - \frac{\mu m}{r} \\ &= \frac{1}{2m} \left[p_r^2 + \frac{p_\theta^2}{r^2} \right] - \frac{\mu m}{r} \end{aligned}$$

So, the options (a) and (c) are correct.

3. Consider the Hamiltonian (H) and the Lagrangian (L) for a free particle of mass m and velocity v . Then

- (a) H and L are independent of each other.
- (b) H and L are related but have different dependence on v .
- (c) H and L are equal.
- (d) Both H and L are quadratic in v .

Answer: (d)

Solution:

For free particle potential energy $V = 0$

$$T = \frac{1}{2}mv^2$$

$$L = T - V, \quad H = T + V$$

$$\Rightarrow H = L = \frac{1}{2}mv^2 \text{ (quadratic in } v\text{)}$$

So, the options (c) and (d) are correct.

June – 2016

(1) Consider the equations motion for a system

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = 0, i = 1, 2, 3, \dots, n$$

Where $L = T - V$ [Where $T(t, q_i, \dot{q}_i)$ as kinetic energy and $V(t, q_i)$ as potential energy]
 q_i be the generalized coordinates, \dot{q}_i be the generalized velocities. Then the equations of motion in the form as above are

- (a) Necessarily restricted to a conservative system but there is no unique choice of L .
- (b) Not necessarily restricted to a conservative system and there is a unique choice of L .
- (c) Necessarily restricted to a conservative system and there is a unique choice of L .
- (d) Not Necessarily restricted to a conservative system and there is no unique choice of L .

Answer: (a)

Solution: Option (a) is correct.

(2) A particle of unit mass moves in the direction of x – axis such that it has the Lagrangian $L = \frac{1}{12}\dot{x}^4 + \frac{1}{2}x\dot{x}^2 - x^2$. Let $Q = \dot{x}^2\ddot{x}$ represent a force (not arising from a potential) acting on the particle in the x – direction. If $x(0) = 1$ and $\dot{x}(0) = 1$, then the value of \dot{x} is

- (a) Some non-zero finite value at $x = 0$.
- (b) 1 at $x = 1$.
- (c) $\sqrt{5}$ at $x = \frac{1}{2}$.
- (d) 0 at $x = \sqrt{\frac{3}{2}}$.

Answer:

Solution:

December – 2016

(1) A bead slides without friction on a frictionless wire in the shape of a cycloid with equation $x = a(\theta - \sin \theta)$, $y = a(1 + \cos \theta)$, $0 \leq \theta \leq 2\pi$.

The Lagrangian function is

- (a) $ma^2(1 + \cos \theta)\dot{\theta}^2 - mga(1 + \cos \theta)$
- (b) $ma^2(1 - \cos \theta)\dot{\theta}^2 - mga(1 + \cos \theta)$
- (c) $ma^2(1 - \cos \theta)\dot{\theta}^2 + mga(1 + \cos \theta)$
- (d) $ma^2(1 + \cos \theta)\dot{\theta}^2 - mga(1 - \cos \theta)$

Answer: (b)

Solution: $L = T - V$

$$\begin{aligned}\text{Where } T &= \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) \\ &= \frac{1}{2}m \left[\{a(\dot{\theta} - \cos \theta \dot{\theta})\}^2 + \{-a \sin \theta \dot{\theta}\}^2 \right] \\ &= \frac{1}{2}ma^2\dot{\theta}^2[1 - 2\cos \theta + \cos^2 \theta + \sin^2 \theta] \\ &= \frac{1}{2}ma^2\dot{\theta}^2 \cdot 2(1 - \cos \theta) \\ &= ma^2(1 - \cos \theta)\dot{\theta}^2\end{aligned}$$

$$V = mgy = mga(1 + \cos \theta)$$

$$\therefore L = ma^2(1 - \cos \theta)\dot{\theta}^2 - mga(1 + \cos \theta)$$

So, answer (b) is correct.

(2) Which of the following are canonical transformation? (where q, p represent generalized coordinate and generalized momentum respectively).

- (a) $P = \log \sin p$, $Q = q \tan p$
- (b) $P = qp^2$, $Q = \frac{1}{p}$
- (c) $P = q \cot p$, $Q = \log\left(\frac{1}{q} \sin p\right)$
- (d) $P = q^2 \sin 2p$, $Q = q^2 \cos 2p$

Answer: (a), (b), and (c)

Solution:

$$(i) P = \log \sin p, Q = q \tan p$$

$$\begin{aligned}\text{Now, } p dq - P dQ &= p dq - \log \sin p \cdot [q \sec^2 p dp + \tan p dq] \\ &= (-q \log \sin p \cdot \sec^2 p) dp + (p - \log \sin p \cdot \tan p) dq \\ &= d[q \cdot (p - \log \sin p \cdot \tan p)]\end{aligned}$$

Which is exact.

$$(ii) P = qp^2, Q = \frac{1}{p}$$

$$\begin{aligned}\text{Now, } p dq - P dQ &= p dq - qp^2 \cdot \left(-\frac{1}{p^2}\right) dp = p dq + q dp = d(pq) \text{ Which is exact.}\end{aligned}$$

$$(iii) P = q \cot p, Q = \log\left(\frac{1}{q} \sin p\right)$$

Now, $p dq - P dQ$

$$= p dq - q \cot p \cdot \left(\frac{\frac{1}{q} \cos p dp + \sin p \cdot \left(-\frac{1}{q^2}\right) dq}{\frac{1}{q} \sin p} \right)$$

$$= p dq - q \cot p \left[\cot p dp - \frac{1}{q} dq \right]$$

$$= p dq - q \cot^2 p dp + \cot p dq$$

$$= -q \cot^2 p dp + (p + \cot p) dq$$

$$= q(1 - \operatorname{cosec}^2 p) dp + (p + \cot p) dq$$

$$= d[q \cdot (p + \cot p)]$$

Which is exact.

$$(iv) P = q^2 \sin 2p, Q = q^2 \cos 2p$$

Now, $p dq - P dQ$

$$= p dq - q^2 \sin 2p \cdot (-q^2 \cdot 2 \sin 2p dp + 2q \cos 2p dq)$$

$$= p dq + 2q^4 \cdot \sin^2 2p dp - 2q^3 \sin 2p \cos 2p dq$$

$$= 2q^4 \sin^2 2p dp + (p - 2q^3 \sin 2p \cos 2p) dq$$

$$= q^4(1 - \cos 4p) dp + (p - 2q^3 \sin 4p) dq$$

$$= M dp + N dq (\text{say})$$

$$\frac{\partial M}{\partial q} = 4q^3(1 - \cos 4p), \frac{\partial N}{\partial p} = 1 - 4q^3 \cos p$$

$$\therefore \frac{\partial M}{\partial q} \neq \frac{\partial N}{\partial p}$$

So, it is not exact.

Hence the options (a), (b), and (c) are correct.

June – 2017

(1) A rigid body having one point 0 fixed and no external torque about 0 has equal principal moments of inertia. Then the body must rotate with

- (a) Angular velocity of variable magnitude.
- (b) Angular velocity with constant magnitude.
- (c) Constant angular momentum but varying angular velocity.
- (d) Varying angular momentum with varying angular velocity.

Answer: (b)

Solution: Option (b) is correct

(2) Consider a spherical pendulum consisting of a particle of mass m which moves under gravity on a smooth sphere of radius a . In terms of spherical polar angles θ, ϕ , with θ measured up from the downward vertical, the Lagrangian is given by

- (a) $ma \left[\frac{a}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) - g \cos \theta \right]$
- (b) $ma \left[\frac{a}{2} (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + g \cos \theta \right]$
- (c) $ma \left[\frac{a}{2} (\dot{\theta}^2 + \dot{\phi}^2 \cos^2 \theta) + g \sin \theta \right]$
- (d) $ma \left[\frac{a}{2} (\dot{\theta}^2 + \dot{\phi}^2 \cos^2 \theta) - g \sin \theta \right]$

Answer: (a)

Solution: $x = a \sin \theta \cos \phi, y = a \sin \theta \sin \phi, z = a \cos \theta$

Now $L = T - V$

$$= \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - mgz$$

$$\dot{x} = a [\cos \theta \cos \phi \dot{\theta} - \sin \theta \sin \phi \dot{\phi}]$$

$$\dot{y} = a [\cos \theta \sin \phi \dot{\theta} + \sin \theta \cos \phi \dot{\phi}]$$

$$\dot{z} = -a \sin \theta \dot{\theta}$$

$$\therefore L = \frac{1}{2} m \cdot a^2 [\cos^2 \theta \cos^2 \phi \dot{\theta}^2 + \sin^2 \theta \sin^2 \phi \dot{\phi}^2 - 2 \sin \theta \cos \theta \cos \phi \sin \phi \dot{\theta} \dot{\phi} + \cos^2 \theta \sin^2 \phi \dot{\theta}^2 + \sin^2 \theta \cos^2 \phi \dot{\phi}^2 + 2 \cos \theta \sin \theta \sin \phi \cos \phi \dot{\theta} \dot{\phi} + \sin^2 \theta \dot{\theta}^2] - mga \cos \theta$$

$$= \frac{1}{2} ma^2 [\cos^2 \theta \dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2 + \sin^2 \theta \dot{\theta}^2] - mga \cos \theta$$

$$= \frac{1}{2} ma^2 [\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2] - mga \cos \theta$$

$$= ma \left[\frac{a}{2} (\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2) - g \cos \theta \right]$$

So, the option (a) is correct.

(3) Let q_α and p_α ($\alpha = 1, 2, \dots, n$) be the generalized coordinates and the generalized momenta respectively. If H denotes the Hamiltonian and q_α (for some $\alpha = \alpha_0$) is an ignorable coordinate, then which of the following equations are satisfied?

(a) $\dot{p}_\alpha = -\frac{\partial H}{\partial q_\alpha}, \dot{q}_\alpha = \frac{\partial H}{\partial p_\alpha}, \forall \alpha$

(b) $\dot{p}_\alpha = \frac{\partial H}{\partial q_\alpha}, \dot{q}_\alpha = -\frac{\partial H}{\partial p_\alpha} \forall \alpha$

(c) $\dot{p}_{\alpha_0} = 0, \dot{q}_{\alpha_0} = \frac{\partial H}{\partial p_{\alpha_0}}$

(d) $\dot{p}_{\alpha_0} = \frac{\partial H}{\partial q_{\alpha_0}}, \dot{q}_{\alpha_0} = 0.$

Answer: (a), (c)

Solution: q_{α_0} is ignorable $\Rightarrow \frac{\partial L}{\partial q_{\alpha_0}} = 0$

Lagrange's equation is

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_{\alpha_0}} \right) - \frac{\partial L}{\partial q_{\alpha_0}} = 0 \Rightarrow \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_{\alpha_0}} \right) = 0$$

$$\therefore \frac{d}{dt} (\bar{p}_{\alpha_0}) = 0 \Rightarrow \dot{p}_{\alpha_0} = 0$$

$$\text{Also, } \dot{q}_{\alpha_0} = \frac{\partial H}{\partial p_{\alpha_0}}$$

Hamilton's equations are

$$\dot{p}_\alpha = -\frac{\partial H}{\partial q_\alpha} \text{ and } \dot{q}_\alpha = \frac{\partial H}{\partial p_\alpha} \forall \alpha.$$

So, the option (a) and (c) are correct.

(4) For a conservative system, the end configurations are fixed and the velocity in the varied motion is such that $T + V = E$. Here T, V and E represent, respectively the kinetic energy, the potential energy and the total energy. If $\delta(A)$ denotes the infinitesimal change in a variable A , and p_α and q_α ($\alpha = 1, 2, \dots, n$) represent the generalized momenta and generalized coordinates, respectively, then

(a) $\delta \int T dt = 0$

(b) $\delta \int \sum_{\alpha=1}^n p_\alpha dq_\alpha = 0$

(c) $\delta \int \sum_{\alpha=1}^n q_\alpha dp_\alpha = 0$

(d) $\delta \int \sum_{\alpha=1}^n (p_\alpha dq_\alpha + q_\alpha dp_\alpha) = 0$

Answer: (a), (b)

Solution: For conservative system,

$$\delta \int T dt = 0$$

$$\text{Also, } \delta H = 0$$

$$\delta \left(\sum_{\alpha=1}^n p_\alpha \dot{q}_\alpha - L \right) = 0$$

$$\Rightarrow \delta \int \sum_{\alpha=1}^n p_\alpha dq_\alpha = 0$$

So, the options (a) and (b) are correct.

December – 2017

(1) Let $I(m)$ denote the moment of inertia of a regular solid tetrahedron about an axis m passing through its centre of gravity. Which of the following is true?

- (a) If the axis l passes through a vertex and the axis l' does not pass through a vertex then $I(l) > I(l')$.
- (b) If the axis l passes through the midpoint of an edge and l' is any other axis then $I(l) < I(l')$.
- (c) $I(l)$ is the same for all axes l .
- (d) If the axis l passes through a vertex and the axis l' does not pass through a vertex then $I(l) < I(l')$.

Answer: (c)

Solution: Moment of inertia is same for all axes.

So, the option (c) is correct.

June – 2018

(1) Given that the Lagrangian for the motion of a simple pendulum is

$L = \frac{1}{2} m \dot{\theta}^2 + mgl \cos \theta$, where m is the mass of the pendulum *bob* suspended by a string of length l , g is the acceleration due to gravity and θ is the amplitude of the pendulum from the mean position, then a Hamiltonian corresponding to L is

- (a) $H(p, \theta) = \frac{p^2}{2ml^2} + mgl \cos \theta$
- (b) $H(p, \theta) = \frac{p^2}{2ml^2} - mgl \cos \theta$
- (c) $H(p, \theta) = \frac{p^2}{ml^2} - mgl \cos \theta$
- (d) $H(p, \theta) = \frac{3p^2}{2ml^2} + mgl \cos \theta$

Answer: (b)

Solution: $H(p, \theta) = p \cdot \dot{\theta} - L$

Where $p = \frac{\partial L}{\partial \dot{\theta}} = ml^2 \dot{\theta}$

$$\therefore H(p, \theta) = ml^2 \dot{\theta}^2 - \frac{1}{2} ml^2 \dot{\theta}^2 - mgl \cos \theta$$

$$= \frac{ml^2 \dot{\theta}^2}{2} - mgl \cos \theta = \frac{ml^2}{2} \cdot \frac{p^2}{m^2 l^4} - mgl \cos \theta \left[\because \dot{\theta} = \frac{p}{ml^2} \right] = \frac{p^2}{2ml^2} - mgl \cos \theta$$

So, the option (b) is correct.

(2) The Hamiltonian for a simple harmonic oscillator is $H(p, q) = \frac{p^2}{2m} + \frac{k}{2} q^2$. Then a possible Lagrangian corresponding to H can be

- (a) $L = \frac{1}{2} m \dot{q}^2 - \frac{k}{2} q^2$
- (b) $L = \frac{1}{2} m \dot{q}^2 - \frac{k}{2} (q^2 + 3q^2 \dot{q})$
- (c) $L = \frac{1}{2} m \dot{q}^2 + \frac{k}{2} q^2$
- (d) $L = \frac{1}{2} m \dot{q}^2 + \frac{k}{2} (q^2 + 3q^2 + \dot{q})$

Answer: (a)

$$\begin{aligned} \text{Solution: } H &= p \dot{q} - L \\ \text{or, } L &= p \dot{q} - \frac{p^2}{2m} - \frac{k}{2} q^2 \quad \left| \begin{array}{l} \dot{q} = \frac{\partial H}{\partial p} = \frac{p}{m} \\ \text{or, } p = m \dot{q} \end{array} \right. \\ &= m \dot{q}^2 - \frac{1}{2} m \dot{q}^2 - \frac{k}{2} q^2 \\ &= \frac{1}{2} m \dot{q}^2 - \frac{k}{2} q^2 \end{aligned}$$

This is option (a)

$$\text{For option (b) } L = \frac{1}{2} m \dot{q}^2 - \frac{k}{2} (q^2 + 3q^2 \dot{q})$$

$$p = \frac{\partial L}{\partial \dot{q}} = m \dot{q} - \frac{k}{2} 3q^2$$

$$\text{or, } \dot{q} = \frac{1}{m} \left(p + \frac{3}{2} k q^2 \right)$$

$$\text{Now, } H = p \dot{q} - L$$

$$\begin{aligned} &= \frac{p^2}{m} + \frac{3pkq^2}{2m} - \frac{1}{2} m \cdot \frac{1}{m^2} \left(p + \frac{3}{2} k q^2 \right)^2 + \frac{k}{2} (q^2 + 3q^2 \dot{q}) \\ &= \frac{p^2}{m} + \frac{3pkq^2}{2m} - \frac{p^2}{2m} - \frac{3pkq^2}{2m} - \frac{9}{8} \frac{k^2 q^4}{m} + \frac{k}{2} q^2 + \frac{3kq^2}{2m} \left(p + \frac{3}{2} k q^2 \right) \\ &= \frac{p^2}{2m} + \frac{k}{2} q^2 + \left(\frac{9}{8m} k^2 q^4 + \frac{3}{2m} p k q^2 \right) \end{aligned}$$

It is not true.

So, the option (a) is true.

December – 2018

(1) Consider the two-dimensional motion of a mass m attached to one end of a spring whose other end is fixed. Let k be the spring constant. The kinetic energy T and the potential energy V of the system are given by

$T = \frac{1}{2} m (\dot{r}^2 + (r\dot{\theta})^2)$ and $V = \frac{1}{2} k r^2$ where $\dot{r} = \frac{dr}{dt}$ and $\dot{\theta} = \frac{d\theta}{dt}$ with t as time. Then which of the following statement is correct?

- (a) r is an ignorable coordinate.
- (b) θ is not an ignorable coordinate.
- (c) $r^2 \dot{\theta}$ remains constant throughout the motion.
- (d) $r \dot{\theta}$ remains constant throughout the motion.

Answer:

Solution: r is not ignorable coordinate.

θ is ignorable coordinate.

$$L = T - V = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\theta}^2) - \frac{1}{2} k r^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\text{or, } \frac{d}{dt} (m r^2 \dot{\theta}) - 0 = 0$$

$$\Rightarrow m r^2 \dot{\theta} = \text{constant}$$

$$r^2 \dot{\theta} = \text{constant}$$

So, the option (c) is correct.

(2) Consider a point mass of mass m which is attached to a massless rigid rod of length ' a '. The other end of the rod is made to move vertically such that its downward displacement from the origin at time t is given by $z(t) = z_0 \cos(wt)$. The mass is moving in a fixed plane and its position vector at time t is given by $\vec{r}(t) = (a \sin \theta(t), z(t) + a \cos \theta(t))$. Then the equation of motion of the point mass is

(a) $a \frac{d^2 \theta}{dt^2} + (g + z_0 w^2 \cos(wt)) \sin \theta = 0$

(b) $a \frac{d^2 \theta}{dt^2} + (g - z_0 w^2 \cos(wt)) \sin \theta = 0$

(c) $a \frac{d^2 \theta}{dt^2} + (g + z_0^2 w^2 \cos(wt)) \cos \theta = 0$

(d) $a \frac{d^2 \theta}{dt^2} + (g - z_0 w^2 \cos(wt)) \cos \theta = 0$

Answer: (a)

Solution: $K.E = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$

$x = a \sin \theta,$

$\dot{x} = a \cos \theta \dot{\theta}$

$y = z(t) + a \cos \theta = z_0 \cos wt + a \cos \theta$

$\dot{y} = -z_0 w \sin wt - a \sin \theta \dot{\theta}$

$\therefore T = K.E = \frac{1}{2} m \left[a^2 \cos^2 \theta \dot{\theta}^2 + z_0^2 w^2 \sin^2 wt + a^2 \sin^2 \theta \dot{\theta}^2 + 2z_0 w a \sin wt \sin \theta \dot{\theta} \right]$

$= \frac{1}{2} m [a^2 \dot{\theta}^2 + z_0^2 w^2 \sin^2 wt + 2az_0 w \sin wt \sin \theta \dot{\theta}]$

$V = P.E \text{ at the pt } A \text{ is } = mg \cdot Oc$

$= mg(oB - Bc)$

$= mg(a - a \cos \theta)$

$= mga(1 - \cos \theta)$

$\therefore L = T - V$

$= \frac{1}{2} m [a^2 \dot{\theta}^2 + z_0^2 w^2 \sin^2 wt + 2a z_0 w \sin wt \sin \theta \dot{\theta}] - mga(1 - \cos \theta)$

Lagrange's equations are

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{d}{dt} \left[\frac{1}{2} m (2a^2 \dot{\theta} + 2a z_0 w \sin wt \sin \theta) \right] - \left[\frac{1}{2} m (2a z_0 w \sin wt \cdot \cos \theta \dot{\theta}) - mga \sin \theta \right] = 0$$

$$\text{or, } a \frac{d^2 \theta}{dt^2} + z_0 w^2 \cos wt \cdot \sin \theta + z_0 w \sin wt \cos \theta \dot{\theta} - z_0 w \sin wt \cos \theta \dot{\theta} + g \sin \theta = 0$$

$$\text{or, } a \frac{d^2 \theta}{dt^2} + (g + z_0 w^2 \cos wt) \sin \theta = 0$$

So, the option (a) is correct.

