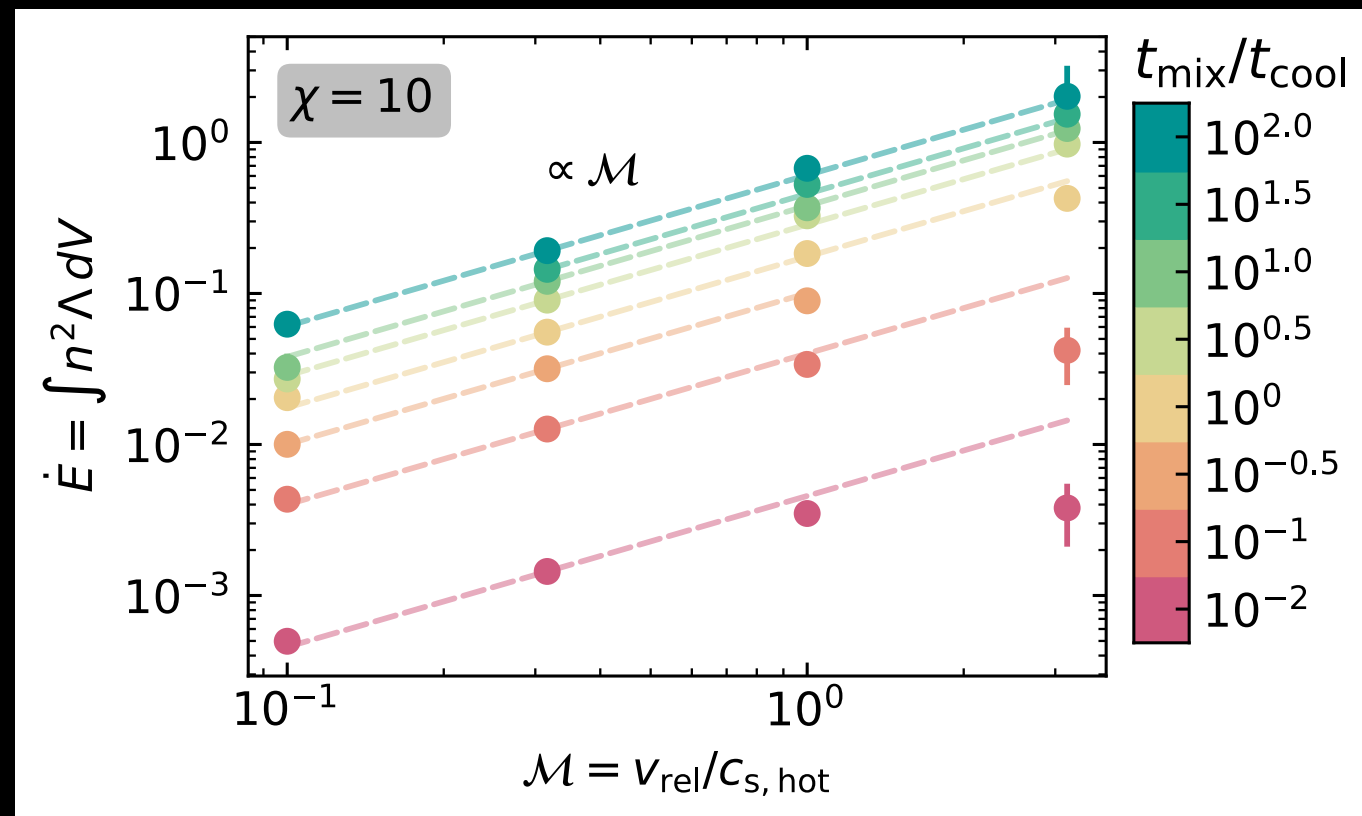
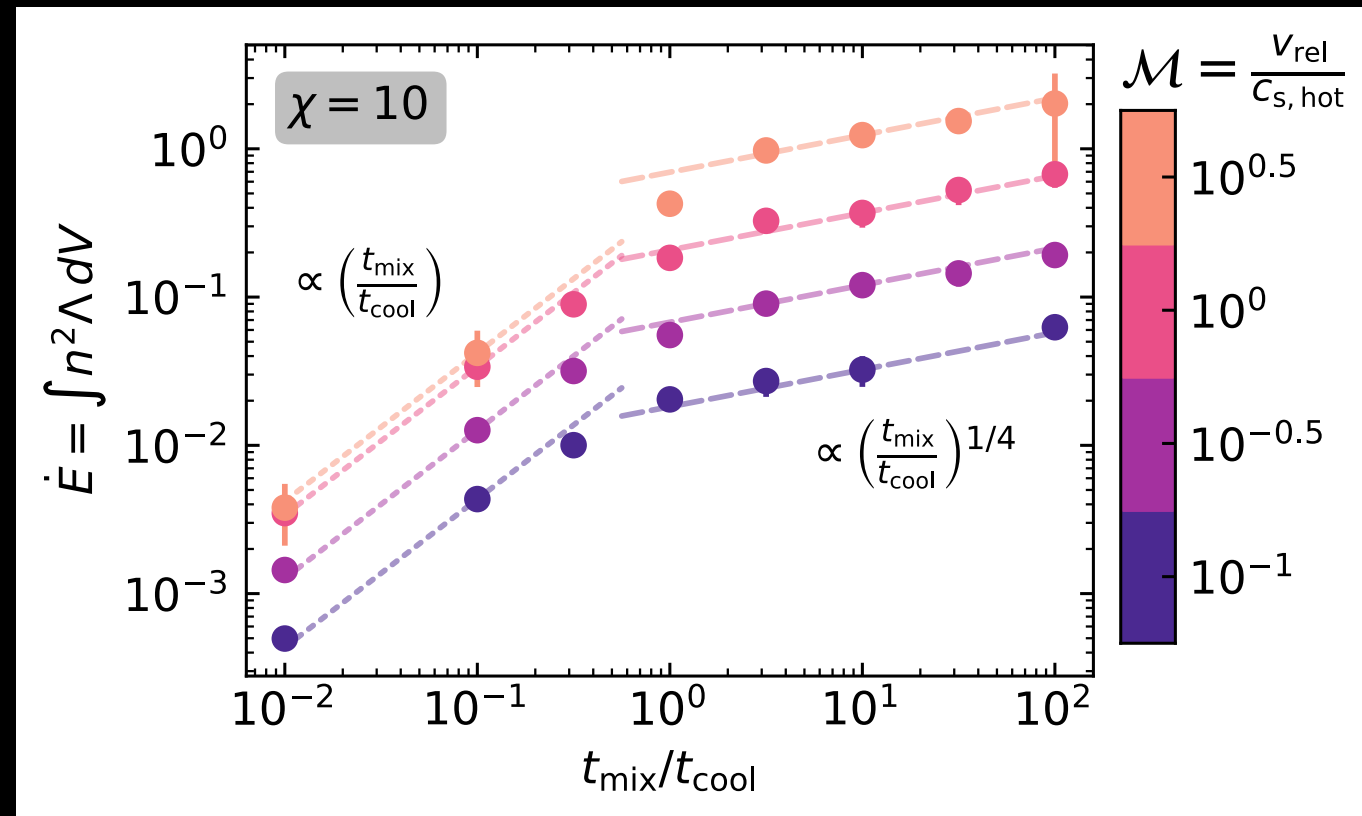


Entrainment, acceleration, & cooling

Inflow of high enthalpy gas replenishes cooling losses:

$$\dot{E}_{\text{cool}} \propto v_z$$



In the slowly cooling limit, simple interpretation of the energy equation gives the right scaling:

$$\frac{\partial}{\partial z} \frac{\gamma}{\gamma - 1} P v_z \left(1 + \frac{\text{Mach}^2}{2(\gamma - 1)} \right) = \dot{e}_{\text{cool}} = \frac{1}{\gamma - 1} \frac{P}{t_{\text{cool}}}$$

Approximately reduces to

$$v_z \approx \frac{h}{t_{\text{cool}}},$$

which, since $t_{\text{mix}} < t_{\text{cool}}$, $h \sim \sqrt{v_{\text{rel}} L t_{\text{mix}}} \sim L$, gives

$$\dot{E}_{\text{cool}} \propto 1/t_{\text{cool}}$$

Entrainment, acceleration, & cooling

In the rapid cooling limit the total cooling is set by the amount that can be refilled

The naive use of the energy equation gives the wrong scaling because

$$\dot{E}_{\text{cool}} = \int \frac{P}{t_{\text{cool}}} dV \bigg/ L^2 = \frac{P h}{t_{\text{cool}}} \frac{\text{Area}}{L^2} \neq \frac{P h}{t_{\text{cool}}}$$

which would if $\text{Area} = L^2$ predict $\dot{E}_{\text{cool}} \propto (v_{\text{rel}}/t_{\text{cool}})^{1/2}$ since $t_{\text{mix}} > t_{\text{cool}}$ and therefore $h \sim \sqrt{v_{\text{rel}} L t_{\text{cool}}}$

To recover the empirical result requires that

$$\frac{\text{Area}}{L^2} \propto (v_{\text{rel}} t_{\text{cool}})^{1/4}$$

