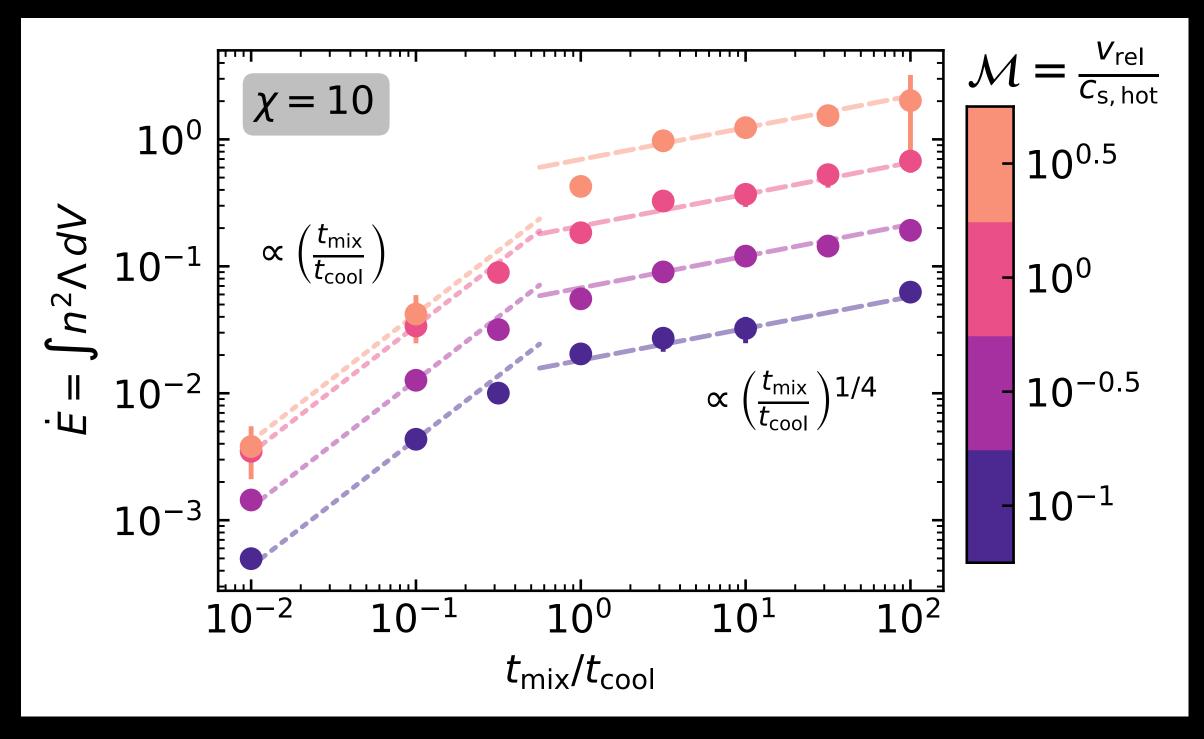
## Entrainment, acceleration, & cooling

Rapid cooling limit, i.e. when  $t_{mix}/t_{cool} \gtrsim 1$ :

$$\dot{E}_{cool} \propto \left(t_{mix}/t_{cool}\right)^{1/4} \; \text{Mach} \sim \left(\frac{\chi^{1/2} \, L}{v_{rel} \, t_{cool}}\right)^{1/2} \left(\frac{v_{rel}}{c_{s}}\right) \propto \frac{L^{1/4} \, v_{rel}^{3/4}}{t_{cool}^{1/4}}$$



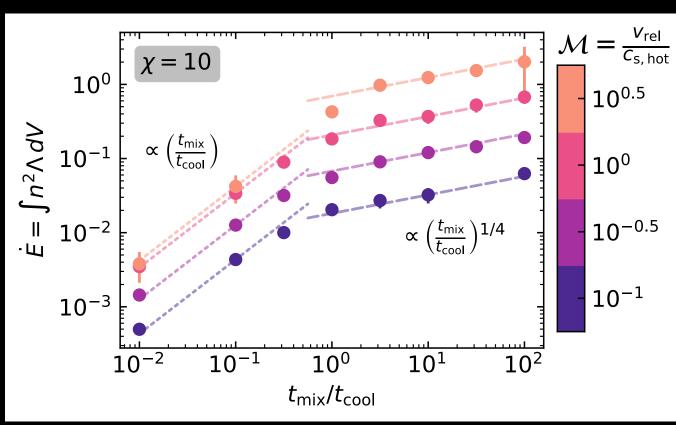
Does this really depend on L in this way To test this I should do strong cooling simulations With fixed  $\Delta x$  and vrel and toool, but I need to get at Least a range of 16 in L to get a factor 2 difference In Edot\_cool, but if I do a L=1 sim with 128, then going 4x larger gives me 512, which is probably too big to afford. So first I need to see if using a L=1 with 64 is Converged (enough) to use. However then going the Other way to a L = 1/4, means I have only 16 cells Across which is probably too small.....
Maybe I can get a way with a range of 8 in L, which Corresponds to a 1.68 change in Edot\_cool.

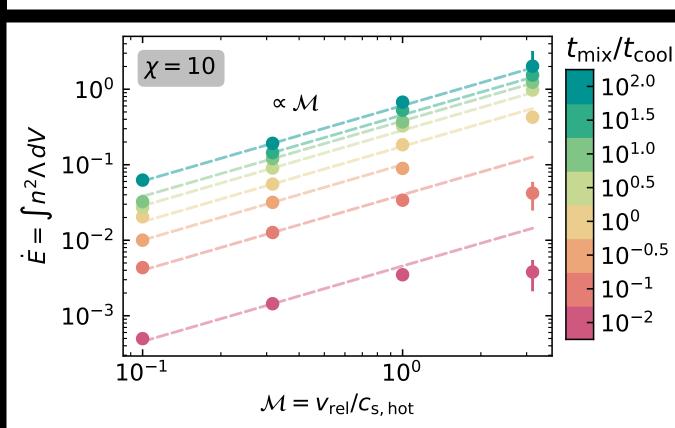
Drummond Fielding

## Entrainment, acceleration, & cooling

Inflow of high enthalpy gas replenishes cooling losses:

$$\dot{E}_{cool} \propto v_z$$





In the slowly cooling limit, simple interpretation of the energy equation gives the right scaling:

$$\frac{\partial}{\partial z} \frac{\gamma}{\gamma - 1} Pv_z \left( 1 + \frac{Mach^2}{2(\gamma - 1)} \right) = \dot{\varepsilon}_{cool} = \frac{1}{\gamma - 1} \frac{P}{t_{cool}}$$

Approximately reduces to

$$v_z \approx \frac{h}{t_{cool}}$$

which, since  $t_{mix} < t_{cool}$ ,  $h \sim \sqrt{v_{rel} L t_{mix}} \sim L$ , gives

$$\dot{E}_{cool} \propto 1/t_{cool}$$