Modeling Ideas: Transit Modeling for Kepler Light Curves

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1 Notation & Preliminaries

Disclaimer: These are just preliminary ideas. There are likely to be typos etc. throughout the document. Use at your own risk. ©

Consider a random process Y defined as a function of time and depending on unknown parameter(s) θ i.e., $Y(t|\theta)$. For simplicity let $\theta = (\omega, \psi, \xi)$ and \mathcal{T} be the interval on which the process is defined. We assume that:

$$Y(t|\theta) = m(t|\omega, \xi) + \epsilon(t|\psi), \qquad \forall t \in \mathcal{T}$$

$$m(t|\omega, \xi) = [1 - q(t|\omega)] f(t|\xi),$$

where $\epsilon(t|\psi)$ is a mean-zero random process and $f(t|\omega)$ is the underlying intensity if there are no transits i.e., $q(t|\omega) \equiv 0$ if there are no transits. In the event of a transit we observe a version of $f(t|\xi)$ 'dampened' by a multiplicative factor $1 - q(t|\omega)$. There is a trade-off in the complexity we place in f and in ϵ , but for our purposes we assume a mean-zero Gaussian process for ϵ :

$$\epsilon_t \sim GP(0, \Sigma_{\epsilon}(\psi)),$$
(1)

where $\Sigma_{\epsilon}(t_i, t_j) = K(|t_i - t_j|; \psi)$ i.e., stationary covariance. Although the form in (1) is fairly rich, in practice we will often assume either independent noise (i.e., $\Sigma(t_i, t_j) = \sigma_{\epsilon}^2 \delta_{t_i t_j}$) or Ornstein-Uhlenbeck/AR(1)-type noise. Although we only observe a noisy version of m(t), the goal is to 'separate' m into two pieces: one representing 'transit-like variation', the other representing 'everything else'. Fortunately we are able to approximate the form of q parametrically. Let t_0 denote the time of the first transit, t_d the transit duration, α the multiplicative reduction in the signal and P the period. Then we assume:

$$q(t|\omega) = \begin{cases} 0 & \text{if } (t \bmod P) \notin [t_0, t_0 + t_d] \\ \alpha & \text{if } (t \bmod P) \in [t_0, t_0 + t_d] \end{cases},$$

where $\omega = \{t_0, t_D, \alpha, P\}$. More sophisticated models for the transit could also be used to allow for smooth transits. The remaining modeling task is to decide upon a structure for f, for which flexible models are required to capture all possible instrumental and astrophysical signals. We decide to model f in the wavelet domain. Let \mathbf{w} denote the DWT of f and let g denote a known thresholding function (hard or soft) applied to the wavelet coefficients then we could model f as:

$$f(\cdot|\mathbf{w}) = \mathcal{W}^{-1}\mathbf{w}, \quad \text{s.t.} \quad g(\mathbf{w}) = c.$$
 (2)

For example, g could be used to constrain the high-resolution wavelet coefficients to zero. This can be just be viewed as a modeling the DWT-transformed mean as a linear function of a subset of the wavelet coefficients (i.e., or more generally, just a flexible parametric form for f parametrized by the wavelet coefficients). This full model has parameters $\theta = \{\psi, \omega, \mathbf{w}\}$, corresponding to the noise

(co)variance parameters, the transit parameters and the (denoised) wavelet coefficients describing the noiseless and transit-removed signal. The full model is then just:

$$Y(t|\theta) = GP([1 - q(t|\omega)] f(t|\xi), \Sigma_{\epsilon}(\psi)),$$

$$L(\theta|y) = \mathcal{N}(y; [1 - q(\cdot|\omega)] f(\cdot|\xi), \Sigma_{\epsilon}(\cdot;\psi)).$$
(3)

In a likelihood framework direct maximization of (3) with respect to θ is likely to be very difficult (and given the high dimension of \mathbf{w} , probably not advisable). However, given there may be prior information available about certain parameters we will instead take a more flexible approach that actually leads to simpler computation. Thus, instead of applying hard constraints on the wavelet coefficients we instead propose to model:

$$W\mathbf{f} = \mathbf{w}, \quad \mathbf{w}|d, \Sigma_w \sim N(d, \Sigma_w), \quad \mathbf{d}|\Sigma_d \sim N(0, \Sigma_d).$$

where W represents the DWT in matrix form. Thus:

$$\mathbf{d}|\mathbf{w}, \Sigma_w, \Sigma_d \sim N\left(\left[\Sigma_d^{-1} + \Sigma_w^{-1}\right]^{-1} \left[\Sigma_w^{-1}\mathbf{w}\right], \left[\Sigma_d^{-1} + \Sigma_w^{-1}\right]^{-1}\right).$$

This illustrates the shrinkage of the wavelet coefficients toward zero by an amount controlled by Σ_w . Typically Σ_w would be specified to ensure a large amount of shrinkage for high-resolution coefficients, and minimal shrinkage for low-resolution coefficients (and would be diagonal). In practice as well, means and variances can be pooled across resolution levels to yield:

$$w_{km}|d_k, \sigma_{w,k}^2 \stackrel{ind}{\sim} N\left(d_k, \sigma_{w,k}^2\right), \qquad d_k \stackrel{ind}{\sim} N\left(0, \sigma_{d,k}^2\right), \qquad \sigma_{w,k}^2 \sim \text{Inv-}\chi^2(\nu_0, s^2).$$

Thus:

$$d_k | \{w_{k1}, \dots, w_{kn_k}\}, \sigma_{w,k}^2, \sigma_{d,k}^2 \sim N \left(\frac{\frac{\sum_m w_{km}}{\sigma_{w,k}^2}}{\frac{1}{\sigma_{d,k}^2} + \frac{n_k}{\sigma_{w,k}^2}}, \frac{1}{\frac{1}{\sigma_{d,k}^2} + \frac{n_k}{\sigma_{w,k}^2}} \right).$$

For levels with little or no shrinkage we obtain:

$$d_k | \{w_{k1} \dots, w_{kn_k}\}, \sigma_{w,k}^2, (\sigma_{d,k}^{-2} \approx 0) \stackrel{approx.}{\sim} N(\bar{w}_k, \sigma_{w,k}^2),$$

where $\bar{w}_k = \sum_m w_{km}$. For levels with a high degree of shrinkage, the posterior for d_k is concentrated around zero. For the time being, lets just consider Σ_d (and possibly Σ_w as well) to be specified by the analyst. Note that \mathbf{f} and \mathbf{w} are one-to-one so that:

This version of the full model either has parameters $\theta = \{\psi, \omega\}$ or $\theta = \{\psi, \omega, \Sigma_w\}$ depending upon how the wavelet coefficients are modeled. The likelihood is given by:

$$p(y, f, d|\theta) \propto p(y|f, \omega, \psi)p(f|d, \Sigma_w)p(d),$$
 (4)

$$\Rightarrow \qquad p(y|\theta) = \int p(y, f, d|\theta) df \, dd, \tag{5}$$

again, the full Bayesian version simply requires a prior on all components of θ and may be preferable.

$$p(f, d, \theta|y) \propto p(y|f, \omega, \psi)p(f|d, \Sigma_w^2)p(d)p(\psi, \omega, \Sigma_w), \tag{6}$$

$$\Rightarrow \qquad p(\theta|y) = \int p(f, d, \theta|y) df \, dd, \tag{7}$$

The form of (6) lends itself to a natural Gibbs sampler of the form:

$$\left[\psi|f,\omega\right],\quad \left[\omega|\psi,f\right],\quad \left[f|d,\omega,\psi,\Sigma_w^2\right],\quad \left[d|f,\Sigma_w^2\right],\quad \left[\Sigma_w^2|d,f\right],$$

where the conditioning on y throughout is dropped for brevity. Let us briefly describe each of the Gibbs updates:

- Sampling from $\psi|f,\omega$: This just amounts to sampling from the posterior distribution of the covariance parameters of a Gaussian process. If the covariance is diagonal then, with a conjugate prior, ψ can be sampled exactly (Inverse- χ^2). If ψ includes covariance parameters then this can be done using e.g., MH.
- Sampling from $\omega|\psi, f$: This will not be in closed form and other methods (e.g., MH) will be needed. Intuitively, varying ω given f varies the value of m(t), so values of ω that yield m to be an appropriate mean for g (given g) will be favored. Likely to be the most challenging step to effectively explore the posterior.
- Sampling from $f|d, \omega, \psi, \Sigma_w^2$: Since f is sandwiched within the multilevel model, the 'prior' on f is multivariate normal, and f appears linearly in the 'likelihood', this will also be multivariate normal (albeit with some linear algebra and an inverse DWT required to compute the mean and variance).
- Sampling from $d|f, \Sigma_w^2$: This is just multivariate normal. The current value of f is transformed to the wavelet domain, and then the mean of d is given by a shrunken version of the wavelet coefficients corresponding to the shrinkage induced by the relative weights of Σ_w^2 and Σ_d^2 .
- Sampling from $\Sigma_w^2 | d, f$: This just amounts to sampling from the posterior distribution of the covariance parameters of a multivariate normal. If the covariance is diagonal (which it will be) or fully unstructured then, with a conjugate prior this can be sampled exactly (Inverse- χ^2). For parametrized non-diagonal matrices it can be sampled using e.g., MH.

To make things more concrete, consider the special case with n evenly-spaced osbervations, with no missing data,

$$\Sigma_{\epsilon}(\psi) = \sigma_{\epsilon}^{2} I, \qquad \Sigma_{d} = \operatorname{diag}\left(\sigma_{d,1}^{2}, \dots, \sigma_{d,1}^{2}, \sigma_{d,2}^{2}, \dots, \sigma_{d,J}^{2}\right),$$
$$\Sigma_{w} = \operatorname{diag}\left(\sigma_{w,1}^{2}, \dots, \sigma_{w,1}^{2}, \sigma_{w,2}^{2}, \dots, \sigma_{w,J}^{2}\right),$$

where the elements of Σ_d are specified constants. Full details of the sampling procedure for this special case are given below:

• Sampling from $\sigma_{\epsilon}^2|f,\omega$: With a conjugate prior $\sigma^2 \sim \text{Inv-}\chi^2(\nu_0,s_0^2)$ on σ^2 we obtain the posterior:

$$\sigma^2 | f, \omega \sim \text{Inv-}\chi^2(\nu_n, s_n^2),$$

where

$$\nu_n = \nu_0 + n, \qquad s_n^2 = \frac{\nu_0 s_0^2 + (y - f)^T (y - f)}{\nu_n}.$$

- Sampling from $\omega | \sigma_{\epsilon}^2$, f: This will not be in closed form and other methods (e.g., MH) will be needed. Intuitively, varying ω given f varies the value of m(t), so values of ω that yield m to be an appropriate mean for g (given g) will be favored.
- Sampling from $f|d, \omega, \psi, \Sigma_w^2$: Since,

$$f|d, \Sigma_w \sim N\left(\mathcal{W}^T d, \mathcal{W}^T \Sigma_w \mathcal{W}\right), \qquad y|f \sim N\left((I-Q)f, \Sigma_\epsilon\right)$$

Therefore:

$$f|y, d, \Sigma_w \sim N(\mu_f, \Sigma_f),$$

$$\Sigma_f^{-1} = \mathcal{W}^T \Sigma_w^{-1} \mathcal{W} + \frac{1}{\sigma_\epsilon^2} (I - Q)(I - Q),$$

$$\mu_f = \Sigma_f \left[\mathcal{W}^T \Sigma_w^{-1} d + \frac{1}{\sigma_\epsilon^2} (I - Q) y \right].$$

Note that if $\Sigma_w = \sigma_w^2 I$ then Σ_f is diagonal and $\mathcal{W}^T d$ is the IDWT of d. For the more general case, $\Sigma_w^{-1} d$ is essentially a standardized vector of the mean wavelet coefficients at different scales. Equivalently, we can work with w:

$$w|d, \Sigma_w \sim N(d, \Sigma_w), \qquad y|w \sim N((I-Q)\mathcal{W}^T w, \Sigma_\epsilon).$$

Therefore:

$$w|d, \Sigma_w \sim N(d, \Sigma_w), \qquad y|w \sim N((I-Q)\mathcal{W}^T w, \Sigma_{\epsilon}),$$

therefore,

$$w|y, d, \Sigma_w \sim N(m_w, V_f),$$

$$V_f^{-1} = \Sigma_w^{-1} + \mathcal{W}(I - Q)\Sigma_{\epsilon}^{-1}(I - Q)\mathcal{W}^T,$$

$$m_f = V_f \left[\Sigma_w^{-1} d + \mathcal{W}(I - Q)\Sigma_{\epsilon}^{-1} y\right].$$

• Sampling from $d|f, \Sigma_w^2$: Let $w = \mathcal{W}f$, then:

$$d_k | \{w_{k1}, \dots, w_{kn_k}\}, \sigma_{w,k}^2, \sigma_{d,k}^2 \sim N\left(\frac{\frac{\sum_m w_{km}}{\sigma_{w,k}^2}}{\frac{1}{\sigma_{d,k}^2} + \frac{n_k}{\sigma_{w,k}^2}}, \frac{1}{\frac{1}{\sigma_{d,k}^2} + \frac{n_k}{\sigma_{w,k}^2}}\right).$$

Note that both Σ_w and Σ_d are diagonal, so there are no matrix inversions in this step.

• Sampling from $\Sigma_w^2 | d, f$: With independent conjugate priors on each $\sigma_{w,j}^2 \sim \text{Inv-}\chi^2(\nu_{w,j,0}, s_{w,j,0}^2)$ on each $\sigma_{w,j}^2$ for $j = 1, \ldots, J$ we obtain:

$$\sigma_{w,j}^2 | d, f \sim \text{Inv-}\chi^2(\nu_{w,j,n}, s_{w,j,n}^2),$$

where

$$\nu_{w,j,n} = \nu_{w,0} + n_j, \qquad s^2 = \frac{\nu_{w,0} s_{w,j,0}^2 + (w_j - d_j)^T (w_j - d_j)}{\nu_{w,j,n}}.$$

Here n_j is the number of coefficients in level j of the wavelet decomposition and w_j and d_j are the subcomponents of w and d corresponding to the wavelet coefficients at level j.

2 Missing Data

The model outlined in the previous sections can be extended to handle missing data in the time domain. Let $t_{mis} = \{t_{mis,1}, \dots, t_{mis,n_{mis}}\}$ denote the set of time points for which the value of Y(t) was not observed, similarly for t_{obs} . We introduce some new notation:

$$y_{com} = y,$$
 $y_{mis} = \begin{pmatrix} y(t_{mis,1}) \\ y(t_{mis,2}) \\ \dots \\ y(t_{mis,n_{mis}}) \end{pmatrix},$ $y_{obs} = \begin{pmatrix} y(t_{obs,1}) \\ y(t_{obs,2}) \\ \dots \\ y(t_{obs,n_{mis}}) \end{pmatrix}.$

Here Y_{mis} becomes another latent variable to be integrated out in the posterior distribution. The full posterior from (6) then becomes:

$$p(f, d, y_{mis}, \theta | y_{obs}) \propto p(y_{com} | f, \omega, \psi) p(f | d, \Sigma_w^2) p(d) p(\psi, \omega, \Sigma_w^2), \tag{8}$$

$$\Rightarrow \qquad p(\theta|y_{obs}) = \int p(f, d, y_{mis}, \theta|y) df \, dd \, dy_{mis}. \tag{9}$$

The form of (8) lends itself to a natural Gibbs sampler of the form:

$$\left[y_{mis} | y_{obs}, \psi, f \right], \quad \left[\psi | f, \omega, y_{com} \right], \quad \left[\omega | \psi, f, y_{com} \right], \quad \left[f | d, \omega, \psi, \Sigma_w^2, y_{com} \right], \quad \left[d | f, \Sigma_w^2 \right], \quad \left[\Sigma_w^2 | d, f \right].$$

where the conditioning on y_{obs} and y_{mis} is now made explicit throughout. Since all previous derivations hold for the complete data, the only additional sampling step is to sample y_{mis} from the conditional posterior distribution. Fortunately, by standard normal theory, Y_{mis} is just normal. Let $\Sigma_{\epsilon,oo}(\psi)$ and $\Sigma_{\epsilon,mm}(\psi)$ denote the submatrices of the covariance matrix $\Sigma_{\epsilon}(\psi)$ that correspond to the observed and missing matrices respectively. Similarly, let $\Sigma_{\epsilon,om} = \Sigma_{\epsilon,mo}^T$ denote the covariance matrix between the observed and missing portions of the data. Let:

$$\mu_{obs} = (I - Q_{obs})f_{obs}, \qquad \mu_{mis} = (I - Q_{mis})f_{mis},$$

where Q_{obs} , Q_{mis} , f_{obs} and f_{mis} are the subcomponents of Q and f corresponding to the observed and missing components respectively. The conditional distribution of Y_{mis} can then be seen to be:

$$Y_{mis}|y_{obs}, \psi, f \sim N\left(\mu_{mis} + \Sigma_{\epsilon,mo}\Sigma_{\epsilon,}^{-1}(y_{obs} - \mu_{obs}), \Sigma_{\epsilon,mm} - \Sigma_{\epsilon,mo}\Sigma_{\epsilon,oo}^{-1}\Sigma_{\epsilon,om}\right).$$

- 3 Searching for Transits: Finding Candidate Period/Phase Configurations
- 4 Using Metadata to Improve Performance

Any ideas?

5 Model Comparison

- Ideas to 'test' for presence of transits
- Stellar variability models.
- Spacecraft event models (e.g., sudden pixel dropout)