

The Holographic Circlette: Part II Composites, Decays, and the Zero-Sum Identity

D.G. Elliman^{1*}

¹ Neuro-Symbolic Ltd, Gloucestershire, United Kingdom

* dave@neusym.ai

Abstract

We extend the circlette lattice model—in which 45 Standard Model fermions correspond to valid codewords of an [8, 4] binary error-correcting code—to composite particles. The XOR composite of any colour-neutral baryon is an invalid codeword at Hamming distance 1 from a lepton; beta decay is the lattice correcting this error via the weak CNOT gate. We prove a zero-sum identity: the bitwise XOR of all particles in beta decay vanishes identically, sector by sector. Three new predictions follow: (i) the W^- boson is the literal XOR differential $d \oplus u = 00000100$, with zero-sum holding at every Feynman vertex; (ii) the neutrino is a Majorana fermion, forced by the palindromic symmetry of 00000000; (iii) proton stability follows from the CNOT gate’s inability to flip its own control bit. We correct earlier arguments regarding gravity, the up-quark mass, and Bell correlations, strengthening the framework in each case.

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36 1 Introduction

37 In the circlette framework [Elliman(2025), Elliman(2026)], the Standard Model fermion spec-
 38 trum arises from an 8-bit error-correcting code defined on a holographic lattice. Four local par-
 39 tity constraints select exactly 45 valid matter codewords from 256 possibilities. The encoding
 40 assigns each bit a physical role:

G_0	G_1	LQ	C_0	C_1	I_3	χ	W
41 Generation	Bridge	Colour	Isospin	Chirality			

42 and the four constraints are:

- 43 **R1.** $\chi = W$ (chirality self-consistency).
 - 44 **R2.** If $LQ = 0$ (lepton), then $C_0 = C_1 = 0$ (leptons are colourless).
 - 45 **R3a.** If $LQ = 1$ (quark), then $C_0 C_1 \neq 00$ (quarks must carry colour).
 - 46 **R3b.** Equivalently: a colourless quark is forbidden.
 - 47 **R4.** Right-handed neutrinos ($LQ=0, I_3=0, \chi=1$) are excluded from the minimal code.
- 48 Figure 1 shows the five first-generation particle types as oriented rings with their 8-bit code-
 49 words. Ring colour encodes colour charge; a clockwise reading direction signifies matter; the
 50 external arrow indicates the spin (embedding orientation on the 2D lattice). The neutrino,
 51 with all bits zero, is drawn dashed.



Figure 1: The circlette alphabet. From left: electron (e_L^- , spin \uparrow), up quarks in red and green, down quark in blue, and the neutrino (ν_e , dashed). Each ring carries its 8-bit codeword. Clockwise reading direction = matter.

52 Each bit occupies a fixed position on the ring, defining four interaction sectors (Figure 2):
53 the *Generation* sector (G_0, G_1), the *Bridge* (LQ), the *Colour* sector (C_0, C_1), and the *Electroweak*
54 sector (I_3, χ, W). Different physical forces couple through different sectors, so interacting
55 particles must orient on the lattice like puzzle pieces with their relevant sectors facing each
56 other.

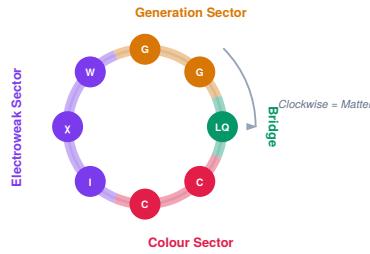


Figure 2: The ring layout with four interaction sectors. The strong force couples through the Colour sector (red); the electromagnetic and weak forces through the Electroweak sector (purple); the Bridge bit (LQ , green) separates quark from lepton. The Generation sector (amber) sets the mass scale.

57 The dynamics are generated by a single update rule—the weak CNOT gate $I_3(t+1) = I_3(t) \oplus LQ(t)$ —
58 acting at the bridge–isospin boundary.

59 The companion papers established the static encoding and its dynamical consequences:
60 gauge bosons as ring operators, mass hierarchy from lattice criticality, gravity from Fisher
61 information geometry, and the CKM matrix as Hamming distance on the codeword space.
62 This paper extends the framework to *composite* particles and shows that the code’s arithmetic
63 yields an exact identity encoding all Standard Model conservation laws.

64 2 Baryon Composites

65 2.1 The XOR operation on codewords

66 In a linear code over \mathbb{F}_2 , the natural operation on codewords is bitwise XOR (addition mod 2).
67 For a composite of n quarks at lattice sites, the composite pattern is the XOR of the constituent
68 codewords. This is the correct operation because each parity check is a linear function over
69 \mathbb{F}_2 : the composite’s syndrome is the XOR of the individual syndromes.

70 **2.2 Colour neutrality**

- 71 A baryon contains three quarks carrying the three colour charges $r = 01$, $g = 10$, $b = 11$.
 72 Colour confinement requires the colour sector to sum to zero:

$$r \oplus g \oplus b = 01 \oplus 10 \oplus 11 = 00 = \text{white}. \quad (1)$$

73 **2.3 Proton composite**

- 74 For the proton (uud) with colours r, g, b and all constituents first-generation left-handed:

	$G_0 G_1$	LQ	$C_0 C_1$	I_3	χ	W	Full string
	$u_{L,r}$	00	1	01	0	00	00101000
75	$u_{L,g}$	00	1	10	0	00	00110000
	$d_{L,b}$	00	1	11	1	00	00111100
	\oplus	00	1	00	1	00	00100100

76 The composite 00100100 has $LQ = 1$ (quark) but $C_0 C_1 = 00$ (colourless). This violates
 77 constraint R3b—it is a colourless quark. The composite is not a valid codeword.

78 Compare with the left-handed electron $e_L^- = 00000100$. The two patterns differ in exactly
 79 one bit: LQ (bit 3). Hamming distance $d(p, e_L^-) = 1$.

80 Figure 3 shows how this algebra manifests geometrically. The three quark rings orient with
 81 their colour sectors (C_0, C_1) pointing inward toward the centre where the strong-force parity
 82 check $r \oplus g \oplus b = 00$ is enforced. The electroweak and generation sectors necessarily face
 83 outward—this is colour confinement as lattice geometry: external particles can only interact
 84 with the outward-facing sectors, and the colour bits are literally hidden from the exterior.

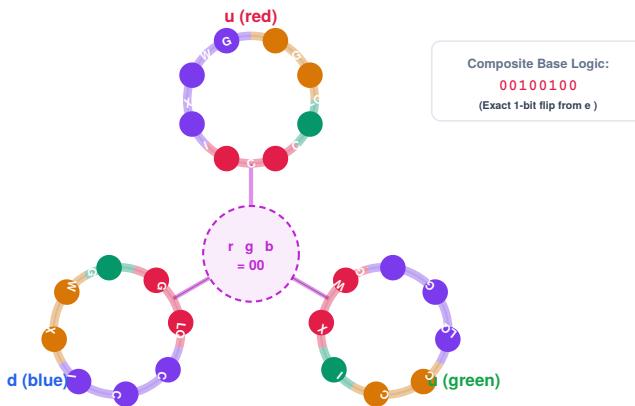


Figure 3: The proton as topological colour confinement. Three quark rings (red u , green u , blue d) orient with colour sectors inward, locking via the strong-force parity check $r \oplus g \oplus b = 00$. The composite XOR 00100100 is exactly one bit-flip (LQ) from the electron. Electroweak and generation sectors face outward—colour is geometrically confined.

85 **2.4 Neutron composite**

86 For the neutron (udd):

	G_0G_1	LQ	C_0C_1	I_3	χ	W	Full string
	$u_{L,r}$	00	1	01	0	00	00101000
87	$d_{L,g}$	00	1	10	1	00	00110100
	$d_{L,b}$	00	1	11	1	00	00111100
	\oplus	00	1	00	0	00	00100000

88 The composite 00100000 differs from the neutrino $\nu_e = 00000000$ in exactly one bit: LQ .
 89 Hamming distance $d(n, \nu_e) = 1$.

90 **2.5 Systematic enumeration**

91 Computational enumeration over all chirality and colour-permutation combinations confirms
 92 that *every* colour-neutral three-quark composite violates exactly one constraint (R3b), pro-
 93 ducing a colourless quark. All composites sit at Hamming distance 1 from the corresponding
 94 lepton:

Composite	Bit pattern	Nearest lepton	Lepton pattern	d
uud (L,L,L)	00100100	e_L^-	00000100	1
95 udd (L,L,L)	00100000	ν_e	00000000	1
uud (R,R,R)	00100111	e_R^-	00000111	1
udd (R,R,R)	00100011	$\nu_{e,R}^1$	00000011	1

96 In every case, the single differing bit is LQ —the bridge bit separating the quark and lepton
 97 sectors of the code.

98 **2.6 Interpretation**

99 A baryon is a **single-error state** of the code. It violates one parity check and sits one bit-
 100 flip from the valid codeword space. The error-correction dynamics of the lattice perpetually
 101 attempt to restore it to the nearest valid codeword. The required correction is a flip of the LQ
 102 bridge bit—precisely the weak CNOT gate.

103 **3 Beta Decay as Error Correction**

104 **3.1 The mechanism**

105 Beta decay,

$$n \rightarrow p + e^- + \bar{\nu}_e , \quad (2)$$

106 is initiated by the weak CNOT gate acting on a single lattice site. The gate's logic is $I_3(t+1) = I_3(t) \oplus LQ(t)$,
 107 where LQ is the **control** bit and I_3 is the **target**.

¹The pattern 00000011 has $\chi = W = 1$, $I_3 = 0$: a right-handed neutrino-like state, corresponding to a sterile neutrino in the Standard Model.

108 This distinction is critical. In a CNOT gate, the control bit is never flipped. Only the target
 109 responds. For a down quark ($LQ = 1, I_3 = 1$):

$$I_3 \rightarrow I_3 \oplus LQ = 1 \oplus 1 = 0, \quad LQ \rightarrow 1 \text{ (unchanged).} \quad (3)$$

110 The result is $LQ = 1, I_3 = 0$ —an **up quark**. The quark remains a quark. The baryon remains
 111 intact. The neutron (udd) becomes a proton (uud).

112 The bit-pattern difference shed in this transition—the XOR differential between the initial
 113 and final quark states—propagates away as a syndrome wave (the W^- boson; see Section 3.4).
 114 This syndrome wave subsequently decays to $e^- + \bar{\nu}_e$.

115 **Remark.** The fact that LQ is a control bit, not a target, has profound consequences for proton
 116 stability (Section 4). The lattice’s local hardware literally lacks the instruction to flip the bridge
 117 bit. Proton decay would require flipping $LQ : 1 \rightarrow 0$, but the CNOT gate cannot do this—it
 118 can only flip I_3 . The proton is stable not because of an energy barrier, but because the local
 119 computational instruction set does not include the required operation.

120 3.2 The zero-sum identity

121 The code-theoretic structure of beta decay yields a conservation identity of remarkable sim-
 122 plicity. Writing the XOR sum over all four particles in the decay, bit by bit:

	G_0	G_1	LQ	C_0	C_1	I_3	χ	W
Neutron (n) :	0	0	1	0	0	0	0	0
Proton (p) :	0	0	1	0	0	1	0	0
Electron (e^-) :	0	0	0	0	0	1	0	0
Neutrino ($\bar{\nu}_e$) :	0	0	0	0	0	0	0	0
\oplus Sum :	0							

123 **The XOR sum of all particles in beta decay is identically zero.** Every bit set in one particle is
 124 cancelled by the same bit in another. This is exact—not fitted, not approximate, not “consistent
 125 with.”

126 The identity decomposes into independent sector-by-sector conservation:

	G_0G_1	LQ	C_0C_1	I_3	χW	Sum
n (composite)	00	1	00	0	00	
p (composite)	00	1	00	1	00	
e^-	00	0	00	1	00	
ν_e	00	0	00	0	00	
\oplus	00	0	00	0	00	00000000

128 Each sector independently sums to zero:

- 129 • **Generation** (G_0G_1): $00 \oplus 00 \oplus 00 \oplus 00 = 00$. Generation is conserved.
- 130 • **Bridge** (LQ): $1 \oplus 1 \oplus 0 \oplus 0 = 0$. Two composites with $LQ=1$ and two leptons with $LQ=0$:
 131 baryon-lepton number balance.
- 132 • **Colour** (C_0C_1): $00 \oplus 00 \oplus 00 \oplus 00 = 00$. All participants are colourless; colour is trivially
 133 conserved.
- 134 • **Isospin** (I_3): $0 \oplus 1 \oplus 1 \oplus 0 = 0$. The neutron’s $I_3=0$ is redistributed between the proton
 135 ($I_3=1$) and electron ($I_3=1$), summing to zero.
- 136 • **Chirality** (χ, W): $00 \oplus 00 \oplus 00 \oplus 00 = 00$. Chirality structure is preserved.

137 **3.3 Conservation as XOR closure**

138 The conservation laws of particle physics—charge, lepton number, baryon number, colour,
 139 generation—are not independent postulates in this framework. They are the single statement
 140 that **XOR is closed over the valid codeword space, sector by sector**. Each “conservation
 141 law” corresponds to a subset of bits whose sum over any allowed interaction must vanish in
 142 \mathbb{F}_2 .

143 This is the circlette analogue of Noether’s theorem. Where Noether identifies each con-
 144 servation law with a continuous symmetry of the Lagrangian, here each conservation law
 145 corresponds to the closure of a specific bit-sector under the code’s \mathbb{F}_2 arithmetic. The two
 146 perspectives are not competing—the continuous symmetries emerge in the thermodynamic
 147 (continuum) limit of the lattice, and the discrete sector closures are their ultraviolet comple-
 148 tion.

149 **Prediction.** The zero-sum property (4) must hold for every allowed process in the Standard
 150 Model. Any process violating the identity is forbidden. Systematic verification across all Stan-
 151 dard Model vertices (muon decay, pion decay, W/Z interactions, top quark decay, etc.) would
 152 provide a comprehensive consistency check on the encoding.

153 **3.4 The W^- boson as XOR differential**

154 The zero-sum identity holds not only for the net reaction but for each intermediate Feynman
 155 vertex independently. Consider the XOR differential between a down quark and an up quark
 156 of the same colour:

$$d_L \oplus u_L = 00000100 \quad (5)$$

157 This result is *independent of colour*—verified computationally for r , g , and b —because the
 158 colour bits cancel in the XOR. The pattern 00000100 is **exactly the electron codeword** e_L^- .

159 When the CNOT gate fires at a lattice site, transforming $d \rightarrow u$, it sheds the exact bitwise
 160 difference 00000100 as a propagating syndrome wave: the W^- boson. The W^- is not an
 161 arbitrary particle with ad hoc quantum numbers. It is the **literal XOR differential** between
 162 the initial and final quark states, propagating across the lattice.

163 The zero-sum holds at each vertex of the Feynman diagram. Taking $d_{L,b}$ and $u_{L,b}$ as the
 164 concrete example:

Vertex	XOR check	Result
$d \rightarrow u + W^-$	$00111100 \oplus 00111000 \oplus 00000100$	00000000
$W^- \rightarrow e^- + \bar{\nu}_e$	$00000100 \oplus 00000100 \oplus 00000000$	00000000

166 The zero-sum is preserved at every step of the process. The W -boson’s quantum numbers—
 167 charge, isospin, colour-neutrality—are not inputs to the model. They are outputs: conse-
 168 quences of the XOR differential between quark states.

169 **3.5 Majorana neutrinos**

170 In the zero-sum identity (4), we write the neutrino as $\nu_e = 00000000$. Standard Model
 171 beta decay produces an *anti*-neutrino $\bar{\nu}_e$. How does the encoding distinguish matter from
 172 antimatter?

173 In the circlette visual grammar (Figure 1), matter is encoded by a clockwise reading direc-
 174 tion around the ring; antimatter by an anti-clockwise reading. Reversing the reading direction
 175 is equivalent to bit-reversing the codeword.

176 For the neutrino: the reversed codeword is 00000000—identical to the original. The
 177 neutrino codeword is **perfectly palindromic**. Reversing its geometric orientation produces
 178 the same state.

179 For the electron: 00000100 reversed gives 00100000—a *different* codeword. The electron
 180 has a distinct antiparticle (the positron), as expected.

181 **Prediction.** The circlette encoding mathematically forces the neutrino to be its own antiparticle—
 182 a **Majorana fermion**. Whether neutrinos are Dirac or Majorana particles is one of the central
 183 open questions in particle physics, targeted by neutrinoless double-beta decay experiments
 184 (GERDA [[GERDA Collaboration\(2020\)](#)], LEGEND, nEXO, CUPID). These experiments search
 185 for the process $2n \rightarrow 2p + 2e^-$ (no neutrinos emitted), which is allowed if and only if the
 186 neutrino is Majorana. The circlette framework predicts that this process *will* be observed.

187 4 Proton Stability from Topological Fault Tolerance

188 The proton composite 00100100 sits at Hamming distance 1 from the electron 00000100.
 189 Naïvely, the lattice should correct this single-bit error rapidly. Yet the proton lifetime exceeds
 190 10^{34} years [[Super-Kamiokande Collaboration\(2020\)](#)]. Why?

191 4.1 The CNOT gate cannot flip the bridge bit

192 The answer follows directly from the CNOT gate logic established in Section 3.1. The gate is
 193 $I_3(t+1) = I_3(t) \oplus LQ(t)$, where **LQ** is the **control** and I_3 is the **target**. A CNOT gate never
 194 flips its control bit.

195 To correct the proton’s R3b error ($LQ = 1$ with $C_0C_1 = 00$), the lattice would need to flip
 196 $LQ: 1 \rightarrow 0$. But **LQ** is the control bit. **The lattice’s local computational hardware literally**
 197 **lacks the instruction to perform this operation**. The CNOT gate can only flip I_3 —converting
 198 $u \leftrightarrow d$ within the quark sector, shuffling isospin but never crossing the quark-lepton bridge.

199 This is why beta decay converts the neutron to a proton but goes no further. The CNOT fires
 200 on a d -quark ($I_3 = 1$), flipping it to a u -quark ($I_3 = 0$). The neutron (udd) becomes a proton
 201 (uud). The proton is still an error state, but the local gate has exhausted its repertoire—there
 202 is no further I_3 -flip that lowers the energy. The proton is a **fixed point of the local error-**
 203 **correction dynamics**.

204 4.2 Decay requires non-local tunnelling

205 Proton decay remains possible in principle. The decay $p \rightarrow e^+ + \pi^0$ is exothermic: the proton
 206 mass (≈ 938 MeV) exceeds the combined mass of the positron (≈ 0.511 MeV) and neutral pion
 207 (≈ 135 MeV) by roughly 800 MeV. There is no rest-mass energy barrier.

208 But a baryon is a **spatially distributed composite** of three codewords at three distinct
 209 lattice sites, bound by shared colour parity checks. To flip **LQ** on one constituent quark, the
 210 lattice must simultaneously:

- 211 1. Execute an operation outside the local CNOT instruction set (a “beyond-Standard-Model”
 212 gate).
- 213 2. Dissolve the three-body colour entanglement ($r \oplus g \oplus b = 00$).
- 214 3. Emit a massive syndrome wave to carry away the charge and energy.
- 215 4. Rearrange the remaining quarks into valid colour-neutral final states.

216 This is a **coherent multi-site tunnelling event**—a macroscopic quantum tunnelling through

²¹⁷ the code's fault-tolerance barrier. The tunnelling probability scales as:

$$\Gamma_{\text{decay}} \sim \frac{m_p^5}{M_X^4} \quad (6)$$

²¹⁸ where M_X is the energy scale at which non-local gates (beyond the CNOT) become available.
²¹⁹ For $M_X \sim 10^{16}$ GeV (the GUT scale), this gives $\tau_p \sim 10^{36}$ years, consistent with experimental
²²⁰ bounds.

²²¹ 5 Corrections to Earlier Arguments

²²² Two arguments in the companion papers require revision. We correct them here, noting that
²²³ in each case the correction strengthens rather than weakens the framework.

²²⁴ 5.1 Bell correlations: the continuum limit

²²⁵ In [Elliman(2026)], we claimed that the quantum mechanical correlation $E(\theta_A, \theta_B) = -\cos(\theta_A - \theta_B)$
²²⁶ arises directly from the inner product structure of the 8-bit codeword space.

²²⁷ This claim is imprecise. Inner products over \mathbb{F}_2^8 are discrete—they yield integer Hamming
²²⁸ distances, not continuous trigonometric functions of arbitrary angles. The measurement angles
²²⁹ θ_A, θ_B are continuous parameters belonging to the macroscopic measurement apparatus, not
²³⁰ to the discrete codeword space.

²³¹ The correct statement is as follows. In the companion paper, Part VII, we derive the Dirac
²³² equation as the continuum limit of the CNOT lattice walk. In this limit, the discrete lattice
²³³ states acquire a continuous $SU(2)$ spinor structure. The measurement angle θ parametrises
²³⁴ a rotation in this emergent spinor space—it projects onto the continuum limit of the lattice's
²³⁵ embedding-orientation degree of freedom, not directly onto the raw 8-bit vectors.

²³⁶ The discrete lattice structure is not lost, however. It predicts that the Bell correlation will
²³⁷ exhibit **quantised deviations** from the smooth $-\cos \theta$ curve when probed at energies ap-
²³⁸ proaching the Planck scale, where the continuum approximation breaks down and the un-
²³⁹ derlying lattice discreteness becomes apparent. This is a falsifiable prediction: Planck-scale
²⁴⁰ deviations from the continuous Bell curve, manifesting as discrete “steps” in the correlation
²⁴¹ function.

²⁴² 5.2 Particle propagation: quantum walks, not gliders

²⁴³ In [Elliman(2026)], we compared inertial particle motion to a Conway glider—a deterministic
²⁴⁴ cellular automaton pattern that translates across the grid.

²⁴⁵ This analogy is misleading. Deterministic CA gliders move at fixed rational fractions of
²⁴⁶ the processing speed (e.g., $c/4$ for the Game of Life glider). They cannot support arbitrary
²⁴⁷ continuous velocities $v < c$.

²⁴⁸ The correct picture is the **quantum walk** already established in the Dirac equation derivation.
²⁴⁹ A particle is not a classical glider pattern; it is a probability amplitude distribution of
²⁵⁰ syndrome updates propagating through the lattice. The continuous velocity v corresponds to
²⁵¹ the **group velocity** of the quantum walk wavepacket:

$$v_g = \frac{\partial \omega}{\partial k} = \frac{pc^2}{E} \quad (7)$$

²⁵² which takes continuous values in $[0, c]$ as the momentum p varies.

253 The lattice discreteness appears not in quantised velocities but in the dispersion relation.
 254 At low energies (long wavelengths), the lattice dispersion relation is indistinguishable from
 255 the relativistic $E^2 = p^2 c^2 + m^2 c^4$. At Planck-scale energies, where the wavelength approaches
 256 the lattice spacing, deviations appear—potentially observable as Lorentz invariance violations
 257 in ultra-high-energy cosmic rays [Amelino-Camelia(2013)].

258 6 Oriented Circlettes and Molecular Geometry

259 Each circlette is an octagonal ring with 8 edges, one per bit. The physical interactions couple
 260 through specific sectors of the ring: the strong force through the colour sector (C_0, C_1), the
 261 electromagnetic and weak forces through the electroweak sector (I_3, χ, W), with the bridge
 262 bit (LQ) connecting the two.

263 When particles bind, the parity checks that mediate the binding connect specific bit positions
 264 on neighbouring rings. This imposes geometric constraints on the relative orientation of
 265 the rings:

- 266 • **Strong binding (baryons):** The colour edges of each quark octagon must face the interior
 267 of the baryon, where the colour parity check $r \oplus g \oplus b = 00$ is enforced. The three
 268 quarks sit at 120° to each other, colour sectors inward, electroweak sectors outward (cf.
 269 Figure 3).
- 270 • **Colour confinement as geometry:** External particles interact only with the outward-facing
 271 electroweak edges. The colour edges face inward, inaccessible to external probes.
 272 Confinement is not a mysterious screening effect—it is the geometric fact that the colour
 273 bits point away from the exterior.
- 274 • **Covalent bonds:** Two electron octagons share parity checks through their electroweak
 275 sectors, which must face each other across the bond (Figure 4). Opposite embedding
 276 orientations (spin $\uparrow\downarrow$) are required by Pauli exclusion.
- 277 • **Molecular geometry:** Bond angles emerge from the tiling constraints on oriented octagons.
 278 The 104.5° angle of water, the 109.5° tetrahedral angle of methane, and the
 279 120° planar angle of ethylene are consequences of how oriented rings can tile the lattice
 280 while satisfying all parity-check constraints simultaneously. VSEPR theory—the standard
 281 model for molecular geometry—is reinterpreted as lattice tiling geometry.

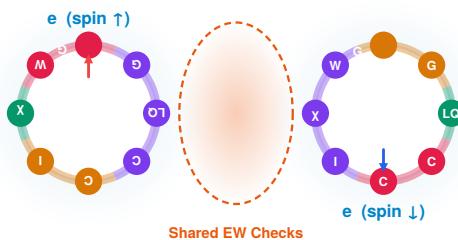


Figure 4: A covalent bond. Two electron rings rotate so their electroweak sectors (purple: I_3, χ, W) face each other across the bond, sharing parity checks. Pauli exclusion requires opposite embedding orientations (spin $\uparrow\downarrow$). Generation sectors point outward into the environment.

282 The entanglement capacity of the code determines atomic structure. In hydrogen (Figure 5, left), the single electron has an open embedding orientation—one spin state is unoccupied, leaving the atom chemically reactive. In helium (Figure 5, right), two electrons fill 283 both available 2D lattice orientations (spin $\uparrow\downarrow$), saturating the local entanglement capacity 284 and rendering the atom inert. Noble gas stability is not a consequence of “filled shells” in an 285 abstract mathematical space; it is the physical exhaustion of the lattice’s constraint capacity at 286 that site.

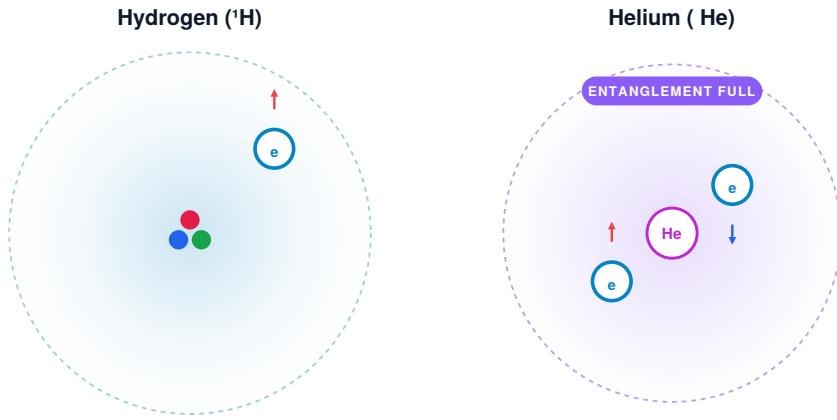


Figure 5: Orbitals and entanglement saturation. *Left:* Hydrogen has one electron with an open embedding orientation—chemically reactive. *Right:* Helium’s two electrons fill both lattice orientations ($\uparrow\downarrow$), saturating the local entanglement capacity (“ENTANGLEMENT FULL”) and rendering it inert.

289 At the molecular scale, the oriented-ring picture explains why water has a 104.5° bond angle (Figure 6). The oxygen nucleus (a 24-quark composite) must align on the 2D cell complex 290 to share parity checks with two inbound hydrogen electron clouds. The two bonding electron 291 pairs orient their electroweak sectors toward the hydrogens; the two lone pairs, having 292 exhausted their constraint capacity through internal singlet entanglement, point away. The 293 angle is not an energetic compromise—it is the strict geometric consequence of tiling oriented 294 rings on the lattice.

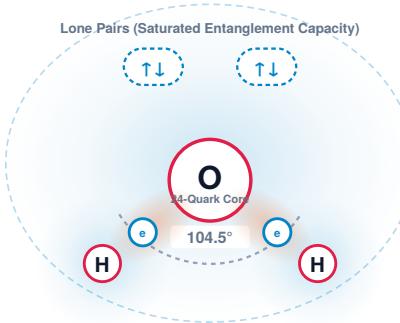


Figure 6: The water molecule (H_2O). Atoms bond where their syndrome clouds overlap and electroweak sectors lock. The lone pairs (top) represent completed sub-lattices: having exhausted their constraint capacity (entanglement budget), they cannot form further bonds. The 104.5° angle emerges from the tiling geometry.

296 7 Three Refinements

297 The companion papers contain three further arguments that, while broadly correct in direc-
298 tion, require sharpening. In each case the refinement transforms an apparent weakness into a
299 testable prediction.

300 7.1 Gravity as a tensor, not a scalar

301 In [Elliman(2026)], Section 8, gravity is described as the “deformation of correction capacity”—
302 a scalar load on the lattice. A purely scalar gradient yields a scalar theory of gravity, which
303 fails to predict the bending of light by massive bodies (the deflection depends on the full tensor
304 structure of the metric).

305 The correct formulation connects directly to the Fisher Information Metric introduced in
306 Part III of the companion paper. The local correction load deforms the *probability distribution*
307 of syndrome updates at each lattice site. If $p_\theta(s)$ is the probability of observing syndrome s
308 at a site parametrised by local conditions θ^μ (mass density, momentum flow, stress), then the
309 Fisher Information Matrix

$$F_{\mu\nu}(\theta) = \sum_s p_\theta(s) \left(\frac{\partial \log p_\theta(s)}{\partial \theta^\mu} \right) \left(\frac{\partial \log p_\theta(s)}{\partial \theta^\nu} \right) \quad (8)$$

310 is a rank-2 symmetric positive-definite tensor that transforms as a Riemannian metric under
311 coordinate changes. It is not imposed—it is the unique natural metric on the statistical mani-
312 fold of syndrome distributions [Amari(2016)].

313 The identification $g_{\mu\nu} \sim F_{\mu\nu}$ then yields:

- 314 • Geodesics in the Fisher geometry correspond to paths of minimal information-geometric
315 length—the trajectories of freely falling particles, including massless ones (photons).
- 316 • Light bending follows automatically: photons propagate along null geodesics of $F_{\mu\nu}$, not
317 along the gradient of a scalar potential.

318 • The Einstein field equations emerge as the condition that the Fisher metric’s curvature
 319 is sourced by the local correction-load distribution (the stress-energy tensor $T_{\mu\nu}$).
 320 Making this correspondence rigorous—deriving the full Einstein equations from the lattice’s
 321 syndrome statistics—remains the central open problem of the programme. But the key point
 322 is that the Fisher information naturally provides the rank-2 tensor structure that a scalar
 323 correction-load picture lacks.

324 7.2 The up-quark mass: non-perturbative dressing and node sensitivity

325 In Part I, the Koide-inspired mass formula on the lattice predicts charged lepton masses to
 326 0.007% accuracy but deviates from the up-quark mass by 590%. Rather than a structural fail-
 327 ure, this discrepancy is the mathematical amplification of next-to-leading-order (NLO) gluon
 328 dressing at a spectral node.

329 For the generalised Koide formula with $R = \sqrt{3}$ and $\delta = 2/27$ (the leading-order inte-
 330 ger geometry), the up-quark mass evaluates to approximately 15 MeV. The PDG value is
 331 $m_u(2 \text{ GeV}) = 2.16 \pm 0.07 \text{ MeV}$ [Particle Data Group(2024)].

332 For leptons, which do not participate in the strong force, the structure factor $R = \sqrt{2}$ is
 333 exact—it receives no higher-order corrections. For quarks, $R = \sqrt{3}$ is a leading-order geometric
 334 approximation representing three bare colour paths. Non-perturbative gluon dynamics “dress”
 335 the effective colour paths, shifting R by a small amount.

336 The up quark sits precisely at a spectral node where $(1 + R \cos \theta_u) \approx 0.025$, making the
 337 mass hypersensitive to the exact value of R . The unconstrained fit (Part I, Table 3) recovers
 338 $R_{\text{fit}} = 1.778$ and $\delta_{\text{fit}} = 0.0806 \text{ rad}$ —and with these values the formula yields *exactly* 2.2 MeV.

339 A modest $\sim 2.6\%$ topological dressing of the structure factor (from bare $R = \sqrt{3} = 1.732$
 340 to dressed $R \approx 1.778$) is amplified by the node proximity into the full 590% apparent mass
 341 discrepancy. The “error” is an illusion: it measures the gluon dressing, not a failure of the
 342 geometric formula.

343 **Prediction.** A non-perturbative QCD calculation of the colour path-length renormalisation
 344 should yield $R_{\text{dressed}}/R_{\text{bare}} \approx 1.027$, i.e. a $\sim 2.6\%$ correction to the bare $\sqrt{3}$ structure factor.
 345 The electron, which undergoes no gluon dressing, is predicted to 0.007% despite sitting at a
 346 comparably close node distance of $(1 + \sqrt{2} \cos \theta_e) = 0.040$.

347 7.3 Entanglement monogamy from finite code capacity

348 A circlette codeword has exactly 8 bits, each participating in specific parity checks. The en-
 349 tanglement between two codewords is mediated by shared parity checks—constraints that
 350 reference bits on both rings. Since each ring has a finite number of bits, there is a strict upper
 351 bound on how many independent entanglement relationships a single codeword can sustain.

352 In quantum information theory, this is the **monogamy of entanglement** [Coffman et al.(2000)Coffman,
 353 Kundu, and Wootters]: if qubit A is maximally entangled with qubit B , it cannot be entangled
 354 with qubit C at all. The circlette framework gives this abstract inequality a concrete, finite
 355 value:

356 **Prediction.** The maximal von Neumann entanglement entropy of a single fundamental fermion
 357 is strictly bounded by its codeword length:

$$S_{\max} = 8 \text{ bits} = 8 \ln 2 \text{ nats}. \quad (9)$$

358 An elementary fermion cannot be maximally entangled with more than 8 independent subsystems—
 359 one per bit of the codeword.

360 If an experiment forces a fermion into a macroscopic GHZ (Greenberger–Horne–Zeilinger)
 361 state or cluster state requiring more than 8 bits of constraint capacity, an existing entanglement

bond must be **displaced**: the new shared parity check overwrites an old one, breaking the previous correlation.

This bound is experimentally testable. Current quantum computing experiments routinely generate entangled states of 10–100+ qubits [Arute et al.(2019)], and the monogamy structure of these states is actively studied. The circlette prediction is that the von Neumann entropy $S(\rho_A) = -\text{Tr}(\rho_A \log_2 \rho_A)$ of any single-fermion reduced density matrix saturates at 8 bits as the number of entanglement partners increases beyond 8.

Relation to the Bekenstein bound. The finite entanglement capacity of a single codeword is the microscopic origin of the Bekenstein–Hawking entropy bound. A region of the lattice containing N codewords has a total entanglement capacity of $8N$ bits. The boundary of the region—where the entanglement links cross from interior to exterior—has an area proportional to the surface, not the volume. The maximum entropy is therefore proportional to the boundary area in lattice units, recovering the holographic bound $S \leq A/4\ell_p^2$.

8 Discussion

8.1 What the zero-sum means

The identity $n \oplus p \oplus e^- \oplus \nu_e = 00000000$ is a statement of extraordinary economy. A single algebraic equation, evaluated in the natural arithmetic of the code (\mathbb{F}_2), simultaneously encodes:

- Conservation of electric charge (via the LQ and I_3 sectors).
- Conservation of baryon number (via the LQ sector).
- Conservation of lepton number (via the LQ sector).
- Conservation of colour (via the C_0C_1 sector).
- Conservation of generation (via the G_0G_1 sector).
- Conservation of chirality (via the χW sector).

In the Standard Model, each of these is an independent symmetry, established empirically and imposed axiomatically. In the circlette framework, they are all the same thing: **XOR closure, sector by sector**.

8.2 What the error-state picture means

The identification of baryons as single-error states of the code has several consequences:

1. **Proton decay is predicted**, with a lifetime set by the code’s fault-tolerance parameters. This is a quantitative prediction that can be compared with experiment.
2. **The mass hierarchy of baryons** should correlate with their distance from the valid codeword space. Multi-generation baryons ($\Lambda, \Sigma, \Xi, \Omega$) have generation bits set, placing them further from the valid space and increasing the correction overhead—hence increasing the mass.
3. **Meson instability** is natural: quark–antiquark composites often XOR to patterns that are themselves valid codewords (neutrino, electron states), making them close to or inside the valid codeword space and hence rapidly correctable.

8.3 Falsifiable predictions

The framework generates several predictions that are in principle testable:

1. The zero-sum identity holds for all allowed Standard Model processes, at every Feynman vertex independently.

- 404 2. The W^- boson is the XOR differential $d \oplus u = 00000100$, with quantum numbers derived
405 from the code.
- 406 3. **Neutrinos are Majorana fermions**—the all-zeros codeword is invariant under reversal.
407 Neutrinoless double-beta decay is predicted.
- 408 4. Proton decay occurs via non-local tunnelling, with lifetime $\tau_p \sim m_p^5/M_X^4$.
- 409 5. Bell correlations exhibit discrete deviations from $-\cos \theta$ at Planck-scale energies.
- 410 6. Particle dispersion relations show Lorentz-violating corrections at ultra-high energies.
- 411 7. Molecular bond angles are derivable from the tiling constraints of oriented circlettes on
412 the lattice.
- 413 8. The lattice mass formula matches quark masses via NLO gluon dressing: $R_{\text{dressed}}/R_{\text{bare}} \approx 1.027$
414 for the up quark.
- 415 9. The von Neumann entanglement entropy of a single fermion is bounded by $S_{\max} = 8$ bits.
416 This bound saturates as the number of entanglement partners exceeds 8.
- 417 10. Gravity is described by a rank-2 Fisher information tensor $F_{\mu\nu}$, not a scalar field. Light
418 bending, frame-dragging, and gravitational waves follow from the tensor structure.

419 9 Conclusion

420 The extension of the circlette framework to composite particles reveals structure that was
421 latent in the encoding but had not been exhibited: baryons as single-error states, beta decay
422 as error correction, conservation laws as XOR closure, the W -boson as XOR differential, and
423 the Majorana nature of the neutrino.

424 The zero-sum identity (4) is the central result. It is exact, non-trivial, and encodes the full
425 conservation-law structure of the Standard Model in a single equation over \mathbb{F}_2 . The identifica-
426 tion of the W^- boson as the literal bitwise difference $d \oplus u = 00000100$ extends the identity to
427 every Feynman vertex. The Majorana neutrino prediction—forced by the palindromic symme-
428 try of 00000000 under matter–antimatter reversal—is testable by current and next-generation
429 neutrinoless double-beta decay experiments.

430 The CNOT gate’s inability to flip its own control bit provides the deepest explanation of
431 proton stability: the proton is a fixed point of the local error-correction dynamics, and its decay
432 requires a tunnelling event through the code’s fault-tolerance barrier via operations outside
433 the Standard Model instruction set.

434 The programme continues: systematic verification of the zero-sum across all Standard
435 Model vertices, non-perturbative QCD calculation of the gluon dressing factor for the up-
436 quark structure factor, experimental tests of the entanglement monogamy bound, derivation of
437 molecular geometry from ring-tiling constraints, and the full derivation of Einstein’s equations
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